# Robust Computational Models for Water Waves

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August 31, 2018

### Outline of the Presentation

- Motivation
- History
- Contribution
  - Method
  - Validation

# Physical Phenomena: Water Waves

#### Water wave hazards:

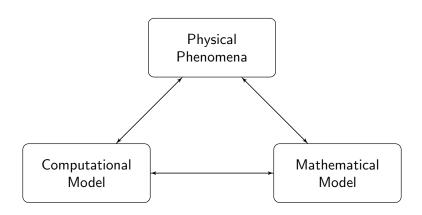
- Tsunamis
- Storm Surges
- Rogue Waves

### Phenomena caused by water waves:

- Nutrient Transport
- Beach Erosion
- Breakup of Sea Ice

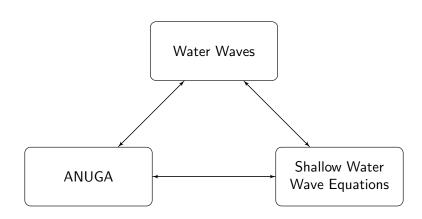
## Computational Modelling

Goal: Model Physics On Computers



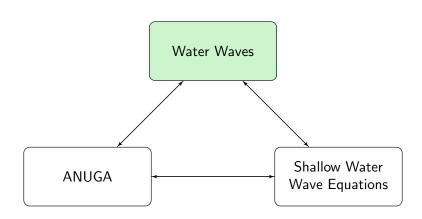
Introduction

## **ANUGA**

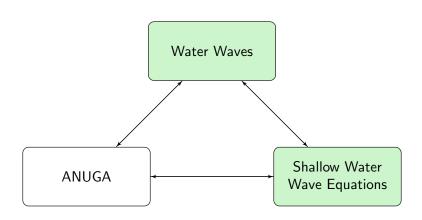


└─ Introduction

### ANUGA: Water Waves

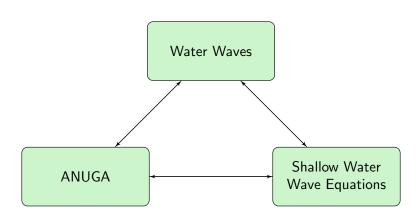


## ANUGA: Shallow Water Wave Equations



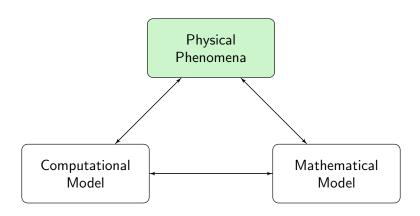
Introduction

## **ANUGA**



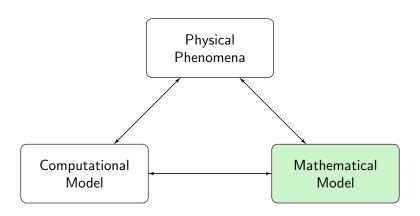
Physical Phenomena

# Physical Phenomena



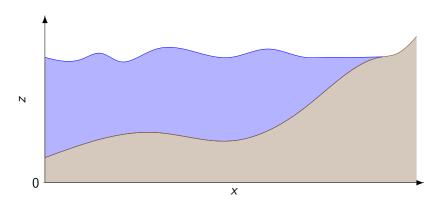
Free Surface Flows

### Mathematical Model



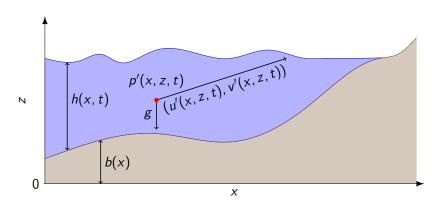
Free Surface Flows

# Typical Scenario



Free Surface Flows

### Full Water Model



Shallow Water Wave Model

## Shallow Water Wave Model

 $./{\tt Pics/WaterModelDiagrams/SWWE.pdf}$ 

# Assumptions

- ▶ particle : u'(x, z, t) constant in z
- particle : v'(x, z, t) = 0
- particle :  $p'(x, z, t) = g\xi$

with

$$\xi(x,z,t)=(h(x,t)+b(x))-z$$

# **Equations**

Mass: 
$$\frac{\partial h}{\partial t} + \frac{\partial}{\partial x} (uh) = 0$$

Momentum: 
$$\frac{\partial uh}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{1}{2} g h^2 \right) + g h \frac{\partial b}{\partial x} = 0$$

### Pros and Cons

#### Pros:

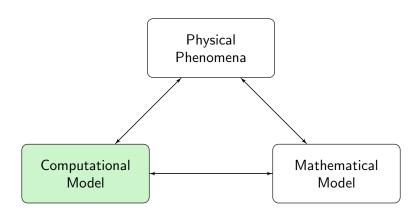
- ► Far simpler than the full water wave model
- Models waves with long wavelengths very well
- Shows good agreement with experimental results

#### Cons:

- No dispersion
- Poor model for short waves
- Cannot model breaking waves

Shallow Water Wave Model

## Computational Model



Shallow Water Wave Model

### **ANUGA**

- ▶ 1999 : Stephen Roberts and Chris Zoppou Paper solving SWWE with a Finite Volume Method
- ▶ 2004 : ANUGA development begins originally focusing on storm surges
- ▶ 2005 : ANUGA refocused to tsunamis
- ▶ 2006 : ANUGA has first public release

Shallow Water Wave Model

### Pros and Cons

#### Pros

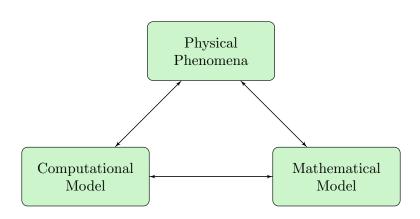
 Robust computational model for water waves based on the Shallow Water Wave Equations

#### Cons

- ▶ Limited by the Shallow Water Wave Equations:
  - no dispersion of waves (recent papers suggest dispersion important for tsunami modelling)
  - not appropriate for shorter waves

Shallow Water Wave Model

### **ANUGA**

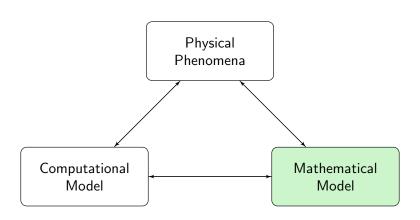


#### Robust Computational Models for Water Waves Water Wave Modelling

### Outcome

New project at the ANU to build robust computational model from dispersive mathematical models

# Mathematical Model



Serre Model

## Serre Model

 $./{\tt Pics/WaterModelDiagrams/SWWE.pdf}$ 

# Assumptions

▶ Particle: u'(x, z, t) constant in z

▶ Particle: 
$$v'(x, z, t) = u \frac{\partial b}{\partial x} - (z - b) \frac{\partial b}{\partial x}$$

► Particle: 
$$p'(x, z, t) = g\xi + \xi\Psi + \frac{1}{2}\xi(2h - \xi)\Phi$$

with

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2},$$

$$\Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

# **Equations**

Mass: 
$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$
 
$$\text{Momentum:} \qquad \frac{\partial (uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right)$$

$$+\frac{\partial b}{\partial x}\left(gh+h\Psi+\frac{h^2}{2}\Phi\right)=0.$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

### Pros and Cons

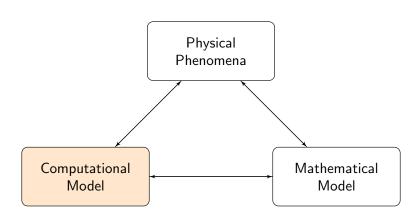
#### Pro:

- ▶ Far simpler than the Euler equations
- Includes dispersive effects
- Still a good model for long wavelength waves and also a good model for shorter wavelengths
- Considered one of the best models for water waves up to wave breaking

#### Cons:

- More complicated than the Shallow Water Wave Equations
- Cannot model breaking waves

# Computational Model



### Previous Work at the ANU

- 2014: Chris Zoppou's PhD thesis
  Demonstrated computational model based on Finite Volume
  Method for the Serre equations with varying bathymetry in
  2D.
- 2014: My Honours thesis Independent reproduction of Chris Zoppou's computational model

Open problems:

3D: Extension of the method to 3D flows

Robust: Validation of model with steep gradients in free surface

Robust: Validation of model in the presence of dry beds

### Thesis Goals

Solve these open problems:

3D: Extension of the method to 3D flows

Robust: Validation of model with steep gradients in free surface

Robust: Validation of model in the presence of dry beds

Technique: Develop a robust computational model from the 2D

Serre equations that can be easily extended to 3D.

### Method

Brief description of the method which is a combination of a Finite Volume Method and a Finite Element method.

### Finite Volume Method

3D: Extends well to 3D

Robust: Stable in the presence of steep gradients

Robust: Stable in the presence of dry beds

Maintains conservation properties of the equations

Chris Zoppou's thesis demonstrated how an adaptation of the Finite Volume Method to the Serre Equations.

## **Equations**

$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x}\left(u^2h + \frac{gh^2}{2} + \frac{h^2}{2}\Psi + \frac{h^3}{3}\Phi\right) + \frac{\partial b}{\partial x}\left(gh + h\Psi + \frac{h^2}{2}\Phi\right) = 0$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

For a Finite Volume Method we require equations in the form

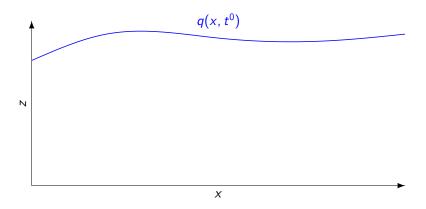
$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0$$

where f(q) and s(q) do not contain temporal derivatives

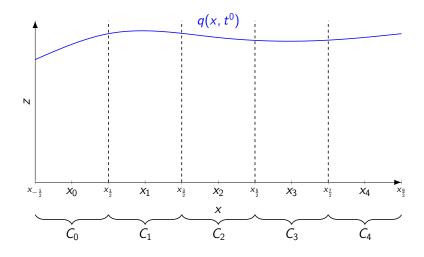
# Finite Volume Method Example

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0$$

## Function at $t = t_0$

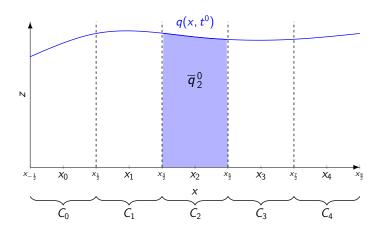


### Cell Discretisation



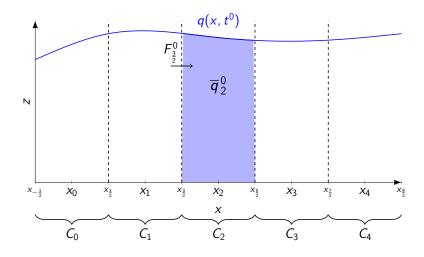
Finite Volume Method

### **Total Amount**



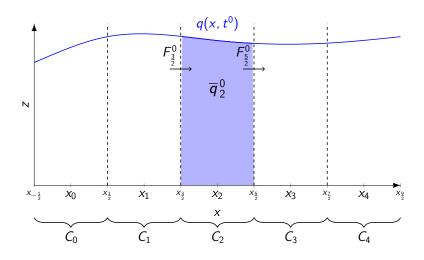
Finite Volume Method

### Flux Left



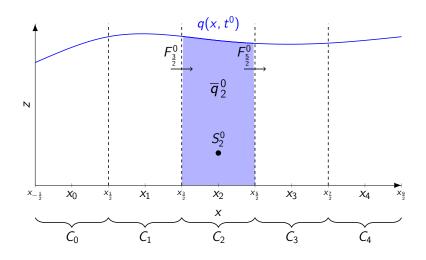
Finite Volume Method

# Flux Right



Finite Volume Method

### Source



# Finite Volume Update

$$\overline{q}_{2}^{1} = \overline{q}_{2}^{0} - \left(F_{\frac{5}{2}}^{0} - F_{\frac{3}{2}}^{0}\right) - \left(S_{2}^{0}\right)$$
$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0$$

# Finite Volume Update

$$\overline{q}_{2}^{1} = \overline{q}_{2}^{0} - \left(F_{\frac{5}{2}}^{0} - F_{\frac{3}{2}}^{0}\right) - \left(S_{2}^{0}\right)$$
$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0$$

$$\int_{C_{2}} q(x, t^{1}) dx = \int_{C_{2}} q(x, t^{0}) dx - \left( \int_{t^{0}}^{t^{1}} f(q(x_{5/2}, t)) dt - \int_{t^{0}}^{F_{\frac{3}{2}}^{0}} \int_{C_{2}}^{F_{\frac{3}{2}}^{0}} f(q(x_{3/2}, t)) dt \right) - \int_{t^{0}}^{t^{1}} \int_{C_{2}} s(q(x, t)) dt$$

# Update Formula for Serre Equations

$$\overline{h}_{j}^{n+1} = \overline{h}_{j}^{n} - \left[ F_{j+1/2}^{n} - F_{j-1/2}^{n} \right]$$

$$\overline{G}_{j}^{n+1} = \overline{G}_{j}^{n} - \left[ F_{j+1/2}^{n} - F_{j-1/2}^{n} \right] - S_{j}^{n}$$

- ▶ All the fluxes  $F_{j+1/2}^n$  and  $F_{j-1/2}^n$  and the source term  $S_j^n$  require u
- require a method to obtain u from  $\overline{h}_j^n$ ,  $\overline{G}_j^n$  and b

# Calculate Velocity

$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2} h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3} h^3 \frac{\partial u}{\partial x} \right).$$

- Chris Zoppou's Thesis used a Finite Difference Method
- Goal: Solve using a Finite Element Method

#### Finite Element Method

3D: Extends well to 3D

Robust: Stable in the presence of steep gradients

▶ Maintains conservation properties for conservation equations

# Finite Element Method Example

Example:

$$-\frac{\partial^2 u}{\partial x^2} = f.$$

Weak Form

$$-\int_{\Omega} \frac{\partial^2 u}{\partial x^2} v = \int_{\Omega} f v \ dx$$

Integrate by parts (Dirichlet boundary conditions):

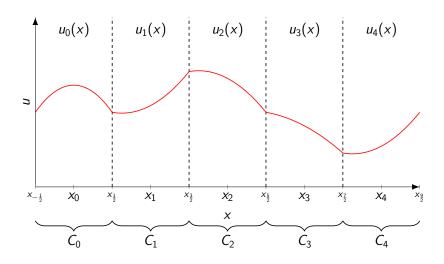
$$\int_{\Omega} \frac{\partial u}{\partial x} \frac{\partial v}{\partial x} \ dx = \int_{\Omega} fv \ dx$$

### Finite Element Method

$$\int_{\Omega} \frac{\partial u}{\partial x} \frac{\partial v}{\partial x} dx = \int_{\Omega} fv dx$$

$$\sum_{j} \left[ \int_{C_{j}} \frac{\partial u}{\partial x} \frac{\partial v}{\partial x} dx \right] = \sum_{j} \left[ \int_{C_{j}} fv dx \right]$$

# Piecewise Polynomial Representation



### Finite Element Method

$$\sum_{j} \left[ \int_{C_{j}} \frac{\partial u_{j}}{\partial x} \frac{\partial v_{j}}{\partial x} dx \right] = \sum_{j} \left[ \int_{C_{j}} f_{j} v_{j} dx \right]$$

$$\mathbf{A}\vec{u} = \vec{c}$$

#### where

- ▶ **A** depends on the polynomial representation of *v*
- $ightharpoonup \vec{u}$  determines the polynomial representation of u
- $ightharpoonup ec{c}$  depends on polynomial representation of f and v

### Finite Element Method for Serre Equations

$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2} h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3} h^3 \frac{\partial u}{\partial x} \right).$$

$$\mathbf{A}\vec{u}=\vec{c}$$

#### where

- ► **A** depends on the polynomial representation of *h*, *b* and test function
- $ightharpoonup \vec{u}$  determines the polynomial representation of u
- $ightharpoonup \vec{c}$  depends on polynomial representation of G and test function

Finite Flement Method

#### Method

- ▶ Reconstruction: Calculate the representations of h and G over the cells from the averages  $\bar{h}$  and  $\bar{G}$
- ▶ Finite Element: use the representations of *h* and *G* over the cells to calculate the representation of *u* over the cell
- ► Finite Volume Method: Update *h* and *G* to the next time using the Finite Volume Method update

### **Progress**

3D: Extension of the method to 3D flows ✓

Robust: Validation of model with steep gradients in free surface

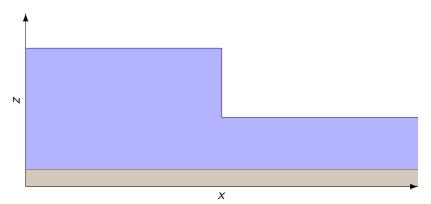
Robust: Validation of model in the presence of dry beds

Thesis

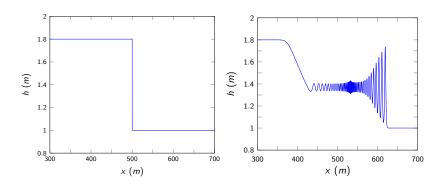
☐ Model Validation for Steep Gradients in the Flow

### Statement of Problem

How does this initially still body of water evolve?



### Our New Numerical Solution



Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. Wave Motion, 76(1):6177.

#### What was known

- ► No analytic solutions
- Asymptotic results for step gradient problems as  $t \to \infty$
- Some experimental comparisons <sup>1</sup>
- Other numerical solutions from the literature; some solving the actual steep gradient problem and some solving a smoothed initial water profile.

<sup>&</sup>lt;sup>1</sup>Zoppou, C. (2014). Numerical Solution of the One-dimensional and Cylindrical Serre Equations for Rapidly Varying Free Surface Flows. PhD thesis, Australian National University.

### Solution

- Demonstrate convergence to one solution for many numerical methods
- Demonstrate good agreement with asymptotic results
- Comprehensively review many numerical methods and smoothing techniques from the literature

Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. Wave Motion, 76(1):6177.

└─ Model Validation for Steep Gradients in the Flow

# Convergence

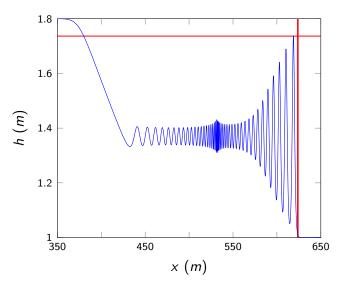
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obust Computational Models for Water Waves				
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Model Validation for Steep Gradients in the Flow				
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# Asymptotic Results



# Review of Smoothing and Methods

- Demonstrated that behaviour is consistent across many numerical methods
- Were able to explain why the behaviour had not previously been observed

#### **Result**

Our numerical solutions for the steep gradient problems are well validated  $^{2}$ 

 $<sup>^2</sup>$ Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. Wave Motion, 76(1):6177.

Thesis

└ Model Validation for Steep Gradients in the Flow

### **Progress**

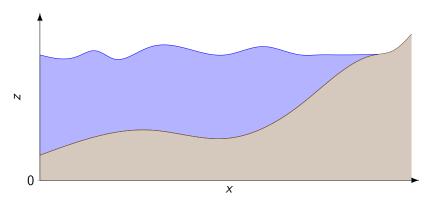
3D: Extension of the method to 3D flows ✓

Robust: Validation of model with steep gradients in free surface  $\checkmark$ 

Robust: Validation of model in the presence of dry beds

### Statement of Problem

Properly handle interaction of waves and the dry bed



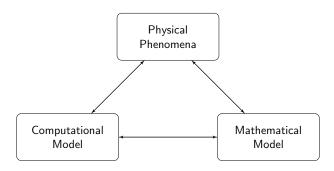
Solution in the Presence of Dry Beds

#### What was known

- No analytic solutions
- A variety of numerical techniques only validated against experimental data

### Solution

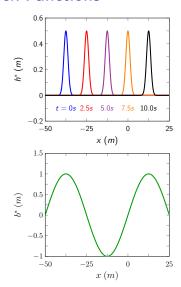
- Solved modified equations that did possess analytic solutions
- Compared with experimental data

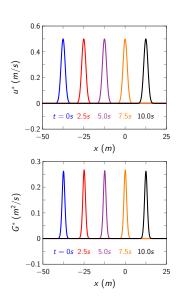


# **Constructing Modified Equations**

- ▶ Pick functions for height, velocity and bed:  $h^*$ ,  $u^*$  and  $b^*$
- Add Source terms to Serre equations that force a solution for h\*, u\* and b\*
- Validation tests

### **Pick Functions**



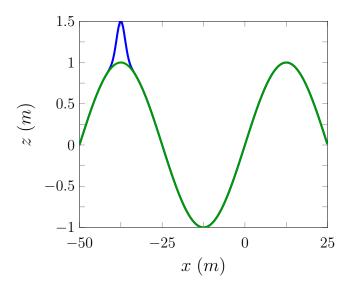


# **Modify Equations**

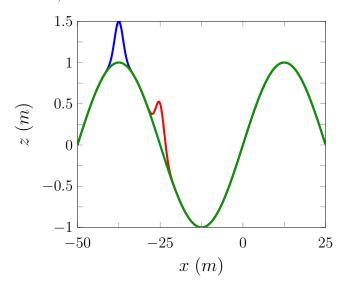
$$\begin{split} \frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} &= S_h^*, \\ \frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left( uG + \frac{gh^2}{2} - \frac{2}{3}h^3 \left[ \frac{\partial u}{\partial x} \right]^2 + h^2 u \frac{\partial u}{\partial x} \frac{\partial b}{\partial x} \right) \\ &+ \frac{1}{2}h^2 u \frac{\partial u}{\partial x} \frac{\partial^2 b}{\partial x^2} - hu^2 \frac{\partial b}{\partial x} \frac{\partial^2 b}{\partial x^2} + gh \frac{\partial b}{\partial x} &= S_G^*. \end{split}$$

 $S_h^*$  and  $S_G^*$  are just the LHS with the quantities replaced by their associated chosen function. We solve the LHS using our method and add in the source terms on the RHS analytically.

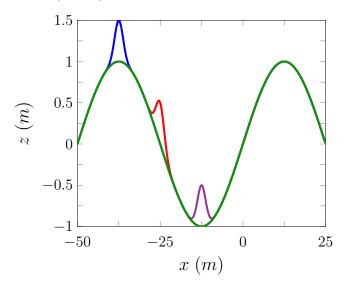
### Results t = 0s



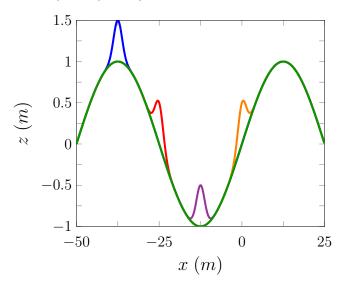
### Results t = 0s, 2.5s



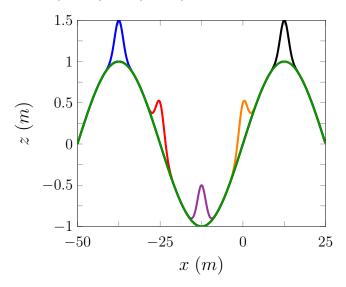
# Results t = 0s, 2.5s, 5.0s



# Results t = 0s, 2.5s, 5.0s, 7.5s



### Results t = 0s, 2.5s, 5.0s, 7.5s, 10.0s



### Modified Equations Validation Conclusions

- Very strong test that we are actually solving the Serre equations accurately as all terms must be accurately approximated
- Can measure the convergence of numerical solutions to the force solutions

# **Experimental Data**

 $./{\tt Pics/DryBed/Syn/WavetankArtifical.pdf}$ 

t = 30s

 $./{\tt Pics/DryBed/Syn/30s.pdf}$ 

t = 40s

 $./{\tt Pics/DryBed/Syn/40s.pdf}$ 

t = 50s

./Pics/DryBed/Syn/50s.pdf

t = 60s

./Pics/DryBed/Syn/60s.pdf

t = 70s

 $./{\tt Pics/DryBed/Syn/70s.pdf}$ 

### **Experimental Validation Conclusions**

- Demonstrates that our computational model agrees with the physical process
- not a very stringent test as there are many source of errors
- few experimental results for non breaking waves

#### **Progress**

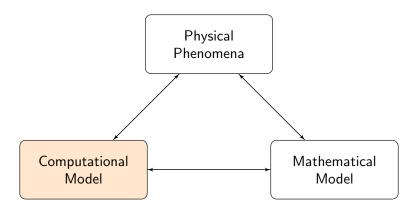
3D: Extension of the method to 3D flows ✓

Robust: Validation of model with steep gradients in free surface  $\checkmark$ 

Robust: Validation of model in the presence of dry beds ✓

#### Conclusions

 Developed a Robust Computational Model from the Serre equations for the 2D water wave problem



Thesis

└Solution in the Presence of Dry Beds

#### References I