

Serre Notes

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1 Regularised Serre Equation

Clamond and Dutykh[?] derived the following regularised Shallow Water Wave equations

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0 \quad (1a)$$

$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left(u^2 h + \frac{gh^2}{2} + \epsilon(x, t) \mathcal{R} h^2 \right) = 0 \quad (1b)$$

where

$$\mathcal{R} \stackrel{\text{def}}{=} h \left(\frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - \frac{\partial^2 u}{\partial x \partial t} - u \frac{\partial^2 u}{\partial x^2} \right) - g \left(h \frac{\partial^2 h}{\partial x^2} + \frac{1}{2} \frac{\partial h}{\partial x} \frac{\partial h}{\partial x} \right).$$

In this context, regularisation means adding additional terms to an equation to control or eliminate fluctuations or oscillations in the solution.

If $\epsilon = 0$ the non-linear shallow water wave equation are recovered. For $\epsilon \neq 0$, \mathcal{R} is a regularisation term that prevents the formation of shocks. It consists of dispersive term that characterises the Serre equation and additional regularisation terms.

We want to make the same transformation as usual, but now with ϵ not being constant we will pick up some extra terms to make sure that

We usually solve the equivalent Serre equations

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0 \quad (2a)$$

$$\frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left[uG + \frac{gh^2}{2} - \epsilon(x, t) h^2 \left(2h \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} + gh \frac{\partial^2 h}{\partial x^2} + \frac{g}{2} \frac{\partial h}{\partial x} \frac{\partial h}{\partial x} \right) \right] = 0 \quad (2b)$$

where

$$G \stackrel{\text{def}}{=} uh - \epsilon(x, t) \frac{\partial}{\partial x} \left(h^3 \frac{\partial u}{\partial x} \right)$$

However, (2) is not equivalent to (1), which can be seen since $\frac{\partial G}{\partial t}$ will result in a derivative of η w.r.t time.

We will settle for modified version of (2) that includes some source terms.

$$\begin{aligned}
\frac{\partial G}{\partial t} &= \frac{\partial(uh)}{\partial t} - \frac{\partial\epsilon}{\partial t} \frac{\partial}{\partial x} \left(h^3 \frac{\partial u}{\partial x} \right) - \epsilon \frac{\partial}{\partial x \partial t} \left(h^3 \frac{\partial u}{\partial x} \right) \\
&= \frac{\partial(uh)}{\partial t} - \frac{\partial\epsilon}{\partial t} \frac{\partial}{\partial x} \left(h^3 \frac{\partial u}{\partial x} \right) - \epsilon \frac{\partial}{\partial x} \left(3 \frac{\partial h}{\partial t} h^2 \frac{\partial u}{\partial x} + h^3 \frac{\partial u}{\partial x \partial t} \right) \\
&= \frac{\partial(uh)}{\partial t} - \frac{\partial\epsilon}{\partial t} \frac{\partial}{\partial x} \left(h^3 \frac{\partial u}{\partial x} \right) - \epsilon \frac{\partial}{\partial x} \left(-3 \frac{\partial(uh)}{\partial x} h^2 \frac{\partial u}{\partial x} + h^3 \frac{\partial u}{\partial x \partial t} \right) \\
&= \frac{\partial(uh)}{\partial t} - \frac{\partial\epsilon}{\partial t} \frac{\partial}{\partial x} \left(h^3 \frac{\partial u}{\partial x} \right) - \epsilon \frac{\partial}{\partial x} \left(-3h^3 \left[\frac{\partial u}{\partial x} \right]^2 - 3uh^2 \frac{\partial h}{\partial x} \frac{\partial u}{\partial x} + h^3 \frac{\partial u}{\partial x \partial t} \right) \\
&= \frac{\partial(uh)}{\partial t} - \frac{\partial\epsilon}{\partial t} \frac{\partial}{\partial x} \left(h^3 \frac{\partial u}{\partial x} \right) - \epsilon \frac{\partial}{\partial x} \left(-3h^3 \left[\frac{\partial u}{\partial x} \right]^2 - 3uh^2 \frac{\partial h}{\partial x} \frac{\partial u}{\partial x} + h^3 \frac{\partial u}{\partial x \partial t} \right) \quad (3)
\end{aligned}$$

Flux term

$$\begin{aligned} \frac{\partial}{\partial x} \left[uG - 2\epsilon h^3 \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} \right] \\ = \frac{\partial}{\partial x} \left[u^2 h - u\epsilon \frac{\partial}{\partial x} \left(h^3 \frac{\partial u}{\partial x} \right) - 2\epsilon h^3 \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} \right] \\ = \frac{\partial}{\partial x} \left[u^2 h - u\epsilon \left(3h^2 \frac{\partial h}{\partial x} \frac{\partial u}{\partial x} \right) - u\epsilon \left(h^3 \frac{\partial^2 u}{\partial x^2} \right) - 2\epsilon h^3 \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} \right] \quad (4) \end{aligned}$$

Add them

[illegible]

Excellent so we get the additional terms

$$S_a = \frac{\partial \epsilon}{\partial t} \frac{\partial}{\partial x} \left(h^3 \frac{\partial u}{\partial x} \right) - \frac{\partial \epsilon}{\partial x} \left[-u \left(3h^2 \frac{\partial h}{\partial x} \frac{\partial u}{\partial x} \right) - u \left(h^3 \frac{\partial^2 u}{\partial x^2} \right) - 2h^3 \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} \right]$$

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0 \quad (6a)$$

$$\frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left[uG + \frac{gh^2}{2} - \epsilon(x, t)h^2 \left(2h \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} + gh \frac{\partial^2 h}{\partial x^2} + \frac{g}{2} \frac{\partial h}{\partial x} \frac{\partial h}{\partial x} \right) \right] \quad (6b)$$

$$= \frac{\partial \epsilon}{\partial t} \frac{\partial}{\partial x} \left(h^3 \frac{\partial u}{\partial x} \right) - \frac{\partial \epsilon}{\partial x} \left[-u \left(3h^2 \frac{\partial h}{\partial x} \frac{\partial u}{\partial x} \right) - u \left(h^3 \frac{\partial^2 u}{\partial x^2} \right) - 2h^3 \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} \right] \quad (6c)$$