

Positron Converter Model

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1 Background

The positron converter provides CESR with its positrons. The converter is a slab of heavy metal (usually tungsten), which is bombarded with electrons whose energy is on the order of ~ 100 MeV. As the incident electrons pass through the converter, they emit photons via Bremsstrahlung, which in turn decay to e^+e^- pairs:

$$e^- + Z \rightarrow e^- + Z + \gamma \rightarrow e^- + Z + e^+ + e^- \quad (1)$$

The production of positrons in the converter is a stochastic process, the details of which are computationally expensive to simulate. As such, it is desirable to have a model for the properties of the produced positrons (their energy, radial displacement, and direction of motion) in terms of probability distributions.

2 Coordinate System

Figure 1 illustrates the coordinate system we use to describe the converter. The incoming electron beam is taken to be along the z -axis. Outgoing positrons will exit the target with some displacement r off of the z -axis, and have some energy E_+ and momentum \mathbf{p}_+ . It is convenient to define a coordinate system (x', y', z) with $\hat{\mathbf{x}}'$ pointing along the direction of $\hat{\mathbf{r}}$, and $\hat{\mathbf{y}}'$ taken perpendicular to $\hat{\mathbf{x}}'$ so that (x', y', z) is a right-handed coordinate system. This defines p'_x and p'_y , the components of the outgoing positron momentum in the primed coordinate system. We then define the “transverse momenta” $\frac{dx'}{ds}$ and $\frac{dy'}{ds}$ by

$$\frac{dx'}{ds} = \frac{p'_x}{p_z} \quad (2)$$

$$\frac{dy'}{ds} = \frac{p'_y}{p_z} \quad (3)$$

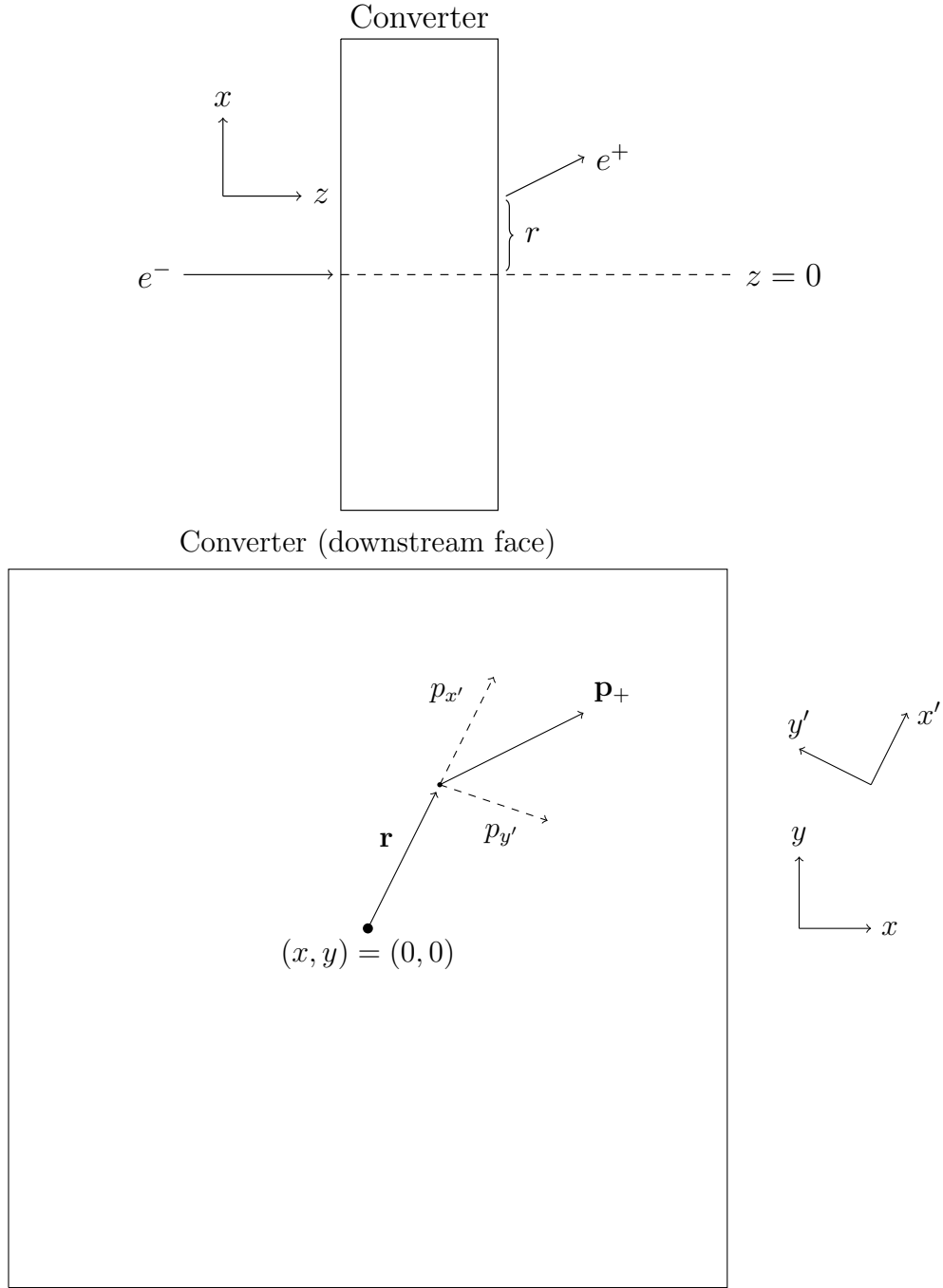


Figure 1: Coordinates used to describe the positrons exiting the converter.

3 The Model

Given a positron converter of thickness T and incoming electrons of energy E_- , we wish to predict E_+ , r , $\frac{dx'}{ds}$, and $\frac{dy'}{ds}$ for the outgoing positrons. We do this in two steps:

- a. First, we pull E_+ and r from a two-dimensional probability distribution $P_1(E_+, r)$. P_1 is determined by interpolation of the simulation data (see Section 4).
- b. For outgoing positrons of any given (E_+, r) , $\frac{dx'}{ds}$ and $\frac{dy'}{ds}$ appear to be distributed as

$$P_2\left(\frac{dx'}{ds}, \frac{dy'}{ds}; E_+, r\right) = A \frac{1 + \beta x}{1 + \alpha_x \left(\frac{dx'}{ds} - c_x\right)^2 + \alpha_y \left(\frac{dy'}{ds}\right)^2}. \quad (4)$$

Note that this functional form is empirically derived, and is not heavily motivated by the underlying physical processes that occur in the converter. In our testing, we have found that this form does the best job of capturing the peak of the transverse momentum distribution, which is of greatest importance when the user cares about the positron capture efficiency of downstream linac elements.

4 Obtaining the Model Coefficients

Using the Geant4[1] software package developed at CERN, we have developed a program for simulating the production of positrons in the converter. This program simulates the results of sending electrons of a fixed energy into a target of fixed thickness, and records the E_+ , r , $\frac{dx'}{ds}$, and $\frac{dy'}{ds}$ values of the positrons that emerge at the downstream face of the converter. The number of electrons used is determined dynamically so that good statistics on the properties of the produced positrons can be obtained.

With a large sample of positron data in hand, the program first bins the E_+ and r data into a 2D histogram. This binned data can then be interpolated to approximate the distribution of E_+ and r values, $P_1(E_+, r)$. After the (E_+, r) bins are chosen, the $\frac{dx'}{ds}$ and $\frac{dy'}{ds}$ values for positrons in each bin are then themselves binned into 2D histograms. We then fit the functional form described in Section 3 to the binned $\frac{dx'}{ds}$, $\frac{dy'}{ds}$ data in each of the (E_+, r) bins to obtain the fit parameters A , β , α_x , α_y and c_x . These fits give us an approximation of $P_2\left(\frac{dx'}{ds}, \frac{dy'}{ds}; E_+, r\right)$.

Note that the distributions P_1 and P_2 will vary with incoming electron energy E_- and target thickness T . The program can be used to sample the behavior of the converter at various values of E_- and T , and performs the above procedure once per (E_-, T) pair. *Bmad* will then interpolate between the discrete set of (E_-, T) values simulated to model the converter at any electron energy and target thickness in the range of interest.

5 Input and Output of the Simulation

To configure the simulation, the user edits the file `config.txt` in the same directory as the simulation program. In this file, the user specifies the target material, a list of incoming

electron energies E_- to be tested, and a list of target thicknesses T to be tested. For each electron energy and target thickness pair, the program simulates the behavior of the converter for many incoming electrons, and outputs the following:

- A table of the probability densities $P_1(E_+, r)$, automatically binned over the range of outgoing positron energies E_+ and radial displacements r where positrons are produced. This table is used by Bmad to approximate the probability distribution $P_1(E_+, r)$.
- Fits to the parameters c_x , α_x , α_y , and β which describe the distribution of positron $\frac{dx'}{ds}$ and $\frac{dy'}{ds}$ values. The program first bins the $\left(\frac{dx'}{ds}, \frac{dy'}{ds}\right)$ values of the positrons produced in each (E_+, r) bin and fits the form of Equation 4 to this binned data. This yields values for c_x , α_x , α_y , and β in each (E_+, r) bin. The program then performs the following fits to $c_x(E_+, r)$, $\alpha_x(E_+, r)$, $\alpha_y(E_+, r)$, and $\beta(E_+, r)$:

- At all values of E_+ and r ,

$$c_x(E_+, r) = A_c E_+^{k_E} r^{k_r} \quad (5)$$

- For α_x , α_y , and β , their volatile behavior at low E_+ make obtaining a single 2D fit at all E_+ and r values too difficult. We therefore adopt a hybrid approach: at low E_+ , a series of 1D fits are performed at each value of E_+ to obtain $\alpha_x(r)$, $\alpha_y(r)$, and $\beta(r)$. The values of α_x , α_y , and β at values of E_+ in between the binned values are then obtained by linear interpolation. At high E_+ , a single 2D fit is performed to obtain $\alpha_x(E_+, r)$, $\alpha_y(E_+, r)$, and $\beta(E_+, r)$.
- For α_x and α_y , the form

$$\alpha_x(r), \alpha_y(r) = (a + br + cr^2 + dr^3)e^{-kr} \quad (6)$$

is used at low E_+ , while the form

$$\alpha_x(E_+, r), \alpha_y(E_+, r) = \frac{1}{E_+} e^{-(k_E E_+ + k_r r)} (a_E + b_E E_+ + c_E E_+^2 + d_E E_+^3) \quad (7)$$

$$(a_r + b_r r + c_r r^2 + d_r r^3) \quad (8)$$

is used at high E_+ .

- For β , the form

$$\beta(r) = a_0 + a_1 r + a_2 r^2 + a_3 r^3 + a_4 r^4 \quad (9)$$

is used at low E_+ , while the form

$$\beta(E_+, r) = A_\beta E_+^{k_E} r^{k_r} \quad (10)$$

is used at high E_+ .

These coefficients (A_c , k_E , k_r , etc) are obtained once for each electron energy and target thickness pair, and can be used to characterize the distribution $P_2\left(\frac{dx'}{ds}, \frac{dy'}{ds}; E_+, r\right)$.

Using this output, *Bmad* can simulate the converter’s outgoing positrons with the probability distribution $P\left(E_+, r, \frac{dx'}{ds}, \frac{dy'}{ds}\right) = P_1(E_+, r)P_2\left(\frac{dx'}{ds}, \frac{dy'}{ds}; E_+, r\right)$ at each of the discrete E_- and T values specified in the setup. For electron energies at target thicknesses between the discrete set sampled by the Geant simulation, interpolation is used between the distributions at the four surrounding (E_-, T) points sampled.

References

- [1] S. Agostinelli et al. “Geant4—a simulation toolkit”. In: *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 506.3 (2003), pp. 250–303. ISSN: 0168-9002. DOI: [https://doi.org/10.1016/S0168-9002\(03\)01368-8](https://doi.org/10.1016/S0168-9002(03)01368-8). URL: <http://www.sciencedirect.com/science/article/pii/S0168900203013688>.