

Variational methods for finding ground states of quantum systems

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Outline

1. Physical task
2. Numerical task
 - 2.1. Simple sampling
 - 2.2. Importance sampling
3. Examples



Physical task

Stationary Schrödinger equation

- given a Hilbert space \mathcal{H} and a Hamiltonian H the stationary Schrödinger equation reads:

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- rarely analytically solvable

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- no physical intuition is used

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- solve the integral and find the minimizing variational parameter λ_{min}



Numerical task

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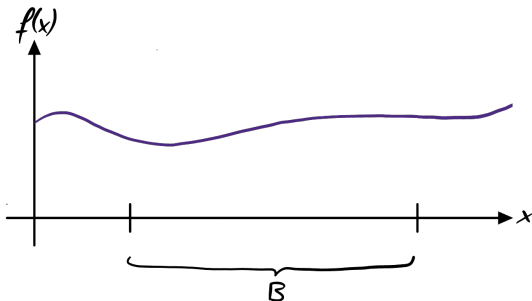
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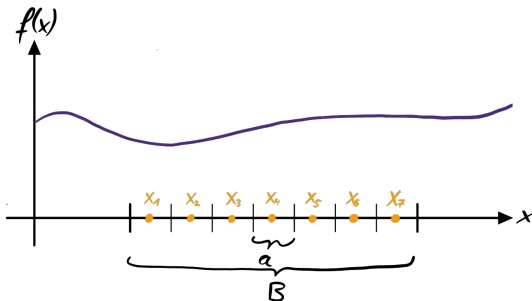
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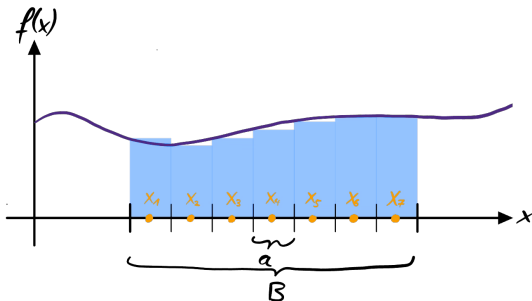
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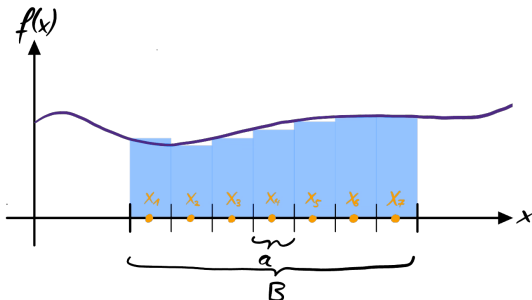
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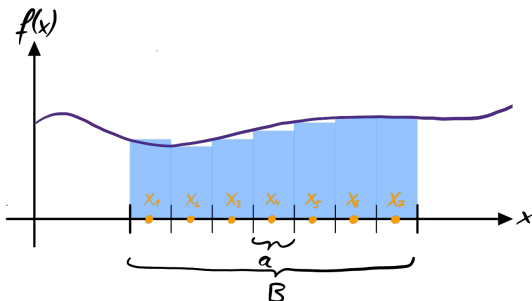
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- problem: slow convergence for high dimension D

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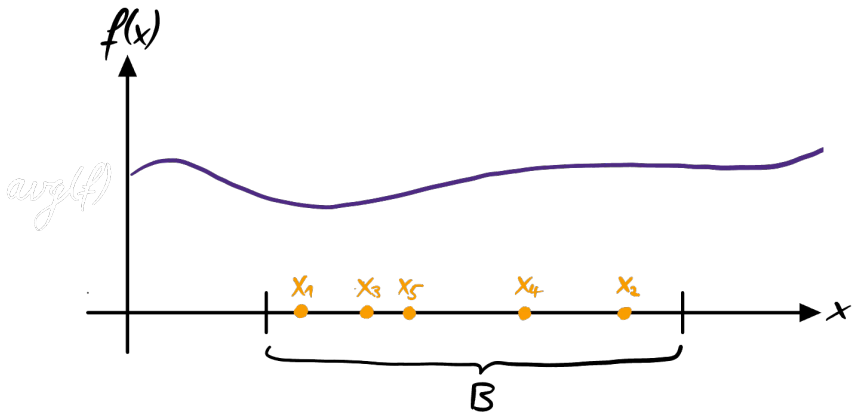
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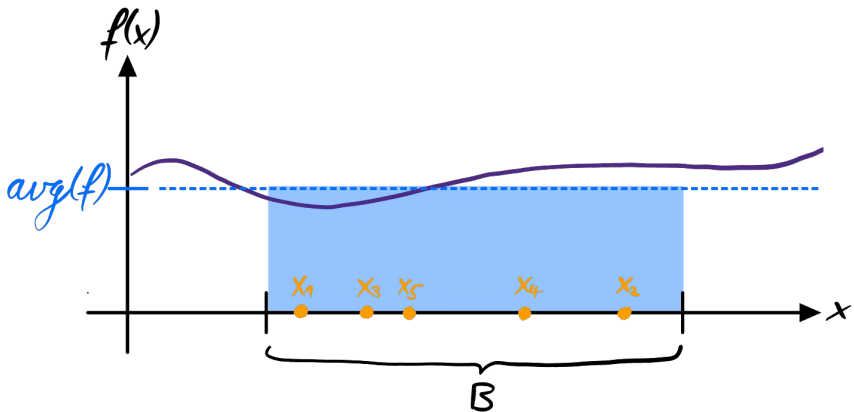
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- the Monte-Carlo sum is a normal distributed random variable

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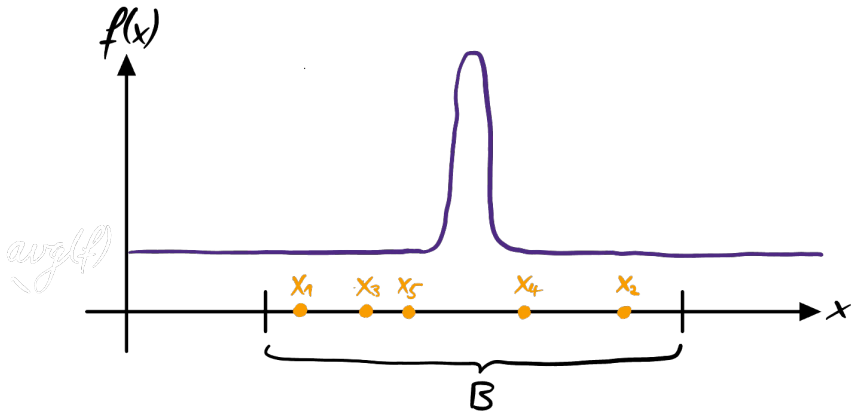
- problem: $\text{vol}(B)$ might be unknown and σ_M might be very large

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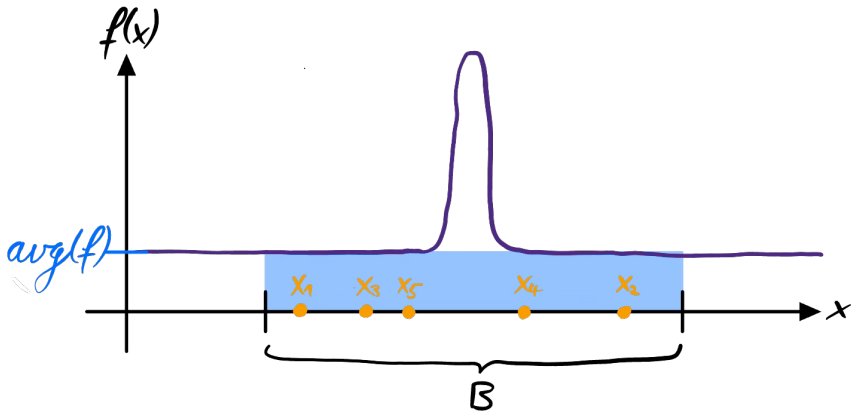
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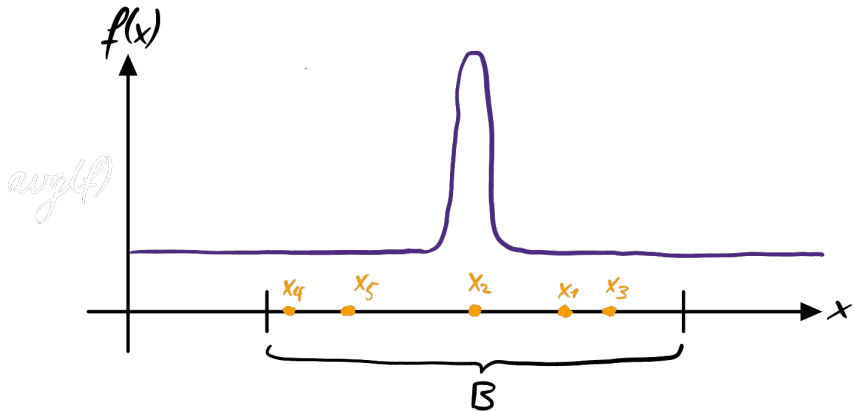
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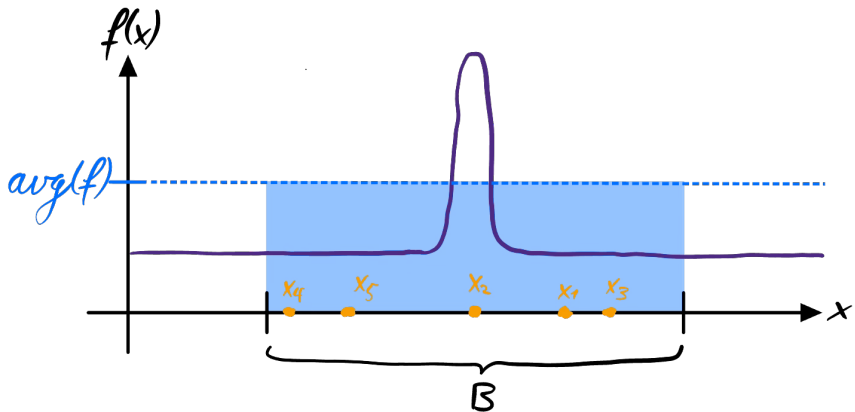
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Importance sampling for variational methods

- our goal is to calculate the variational energy $E(\lambda)$:

$$E(\lambda) = \frac{\int_{-\infty}^{\infty} dx \psi(x; \lambda) H \psi(x; \lambda)}{\int_{-\infty}^{\infty} dx \psi(x; \lambda) \psi(x; \lambda)} = \int_{-\infty}^{\infty} dx \left(\frac{\psi^2(x; \lambda)}{\int_{-\infty}^{\infty} dx \psi^2(x; \lambda)} \frac{H \psi(x; \lambda)}{\psi(x; \lambda)} \right)$$

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- define:

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- using importance sampling we find:

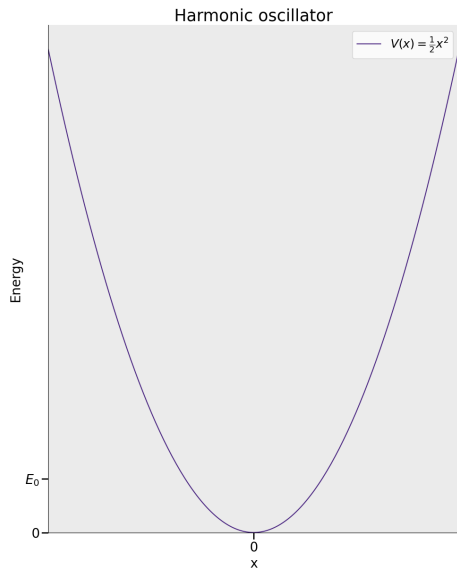
$$E(\lambda) = \int_{-\infty}^{\infty} dx \rho(x; \lambda) E_{local}(x; \lambda) \approx \frac{1}{N} \frac{1}{M} \sum_{i=1}^N \sum_{j=1}^M E_{local}(x; \lambda) (x_j^{(i)}; \lambda)$$



Examples

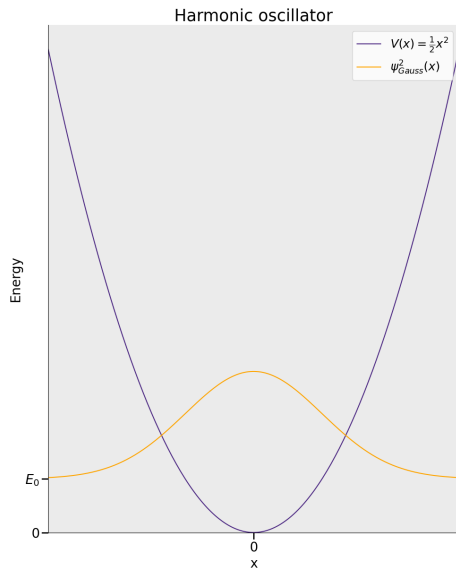
Harmonic oscillator

- Hamiltonian: $H = -\frac{1}{2}\frac{\partial^2}{\partial x^2} + \frac{1}{2}x^2$
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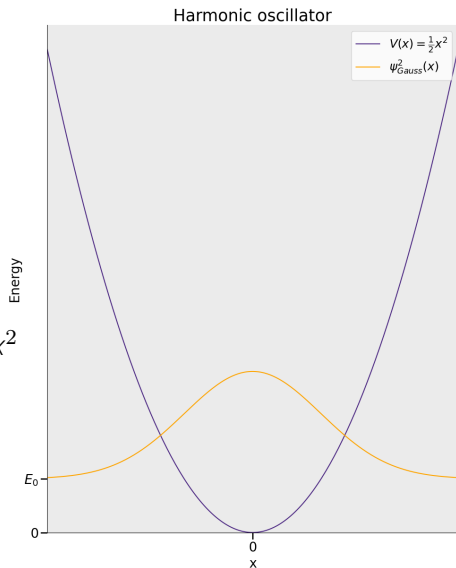
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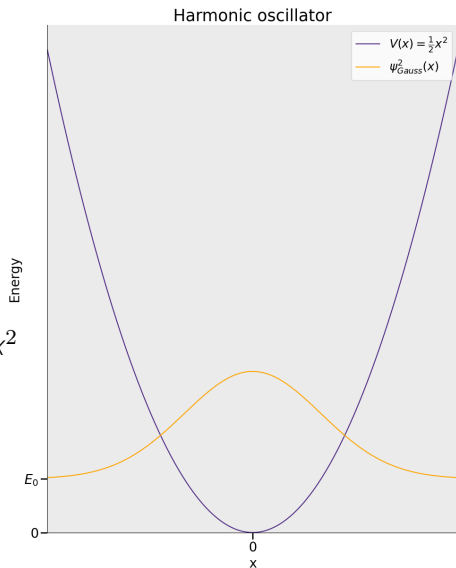
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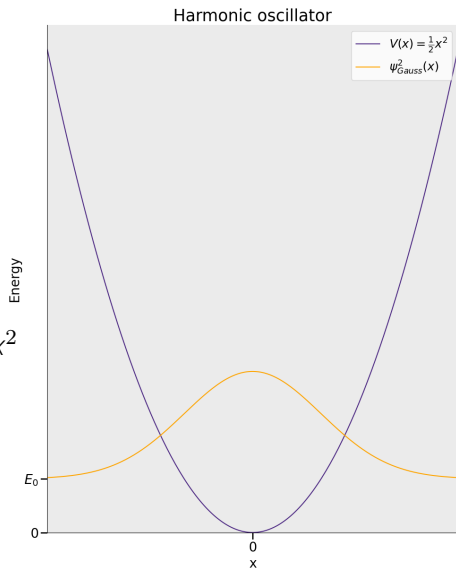
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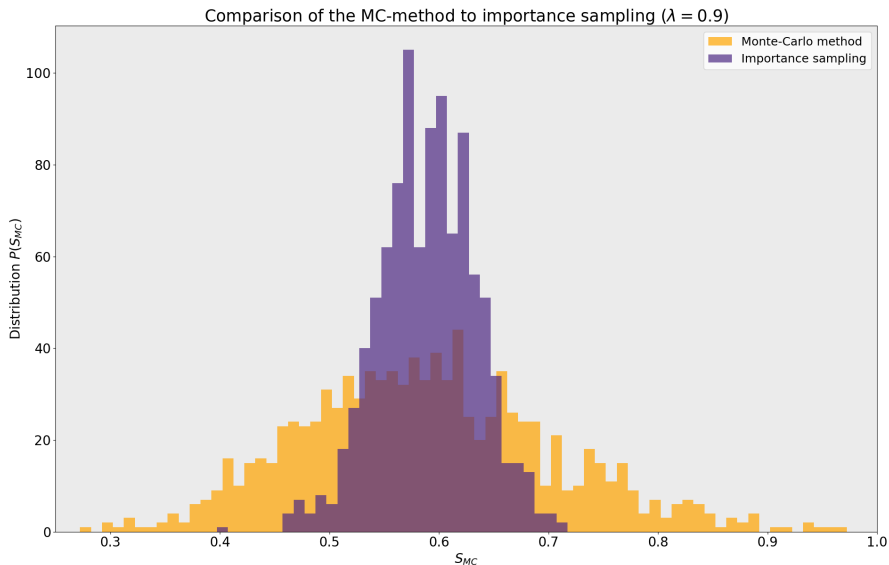


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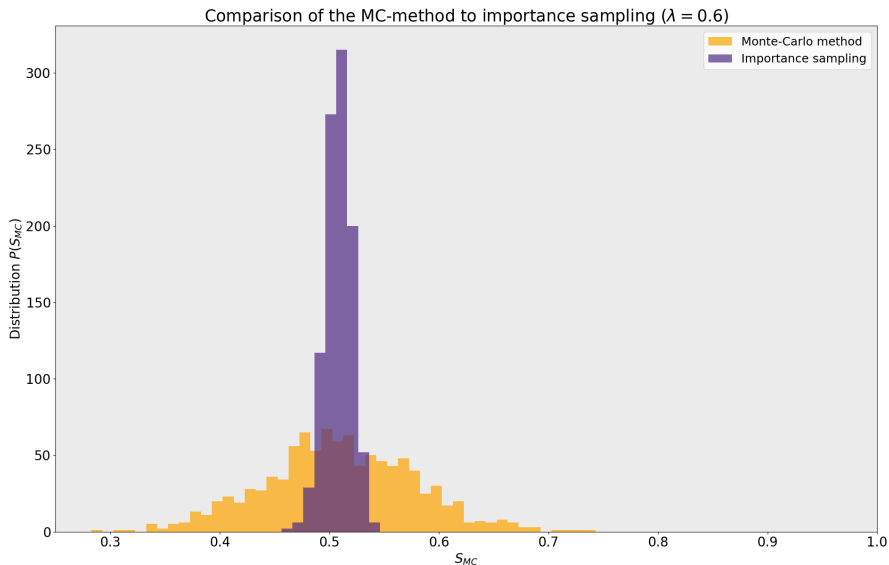
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- for $\lambda = \frac{1}{2}$ the variational ansatz is the exact solution
 $\Rightarrow E(x; \lambda = \frac{1}{2}) \equiv E_0 = \frac{1}{2}$



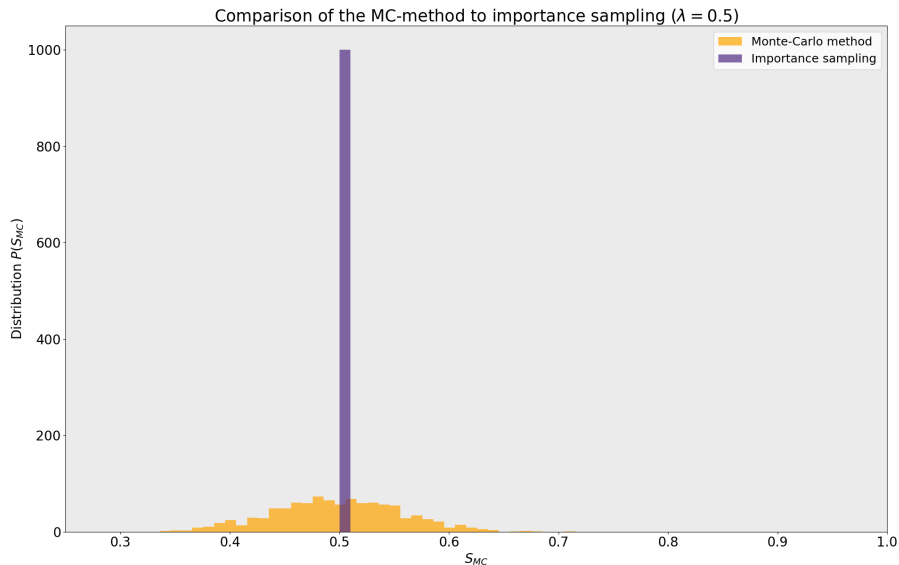
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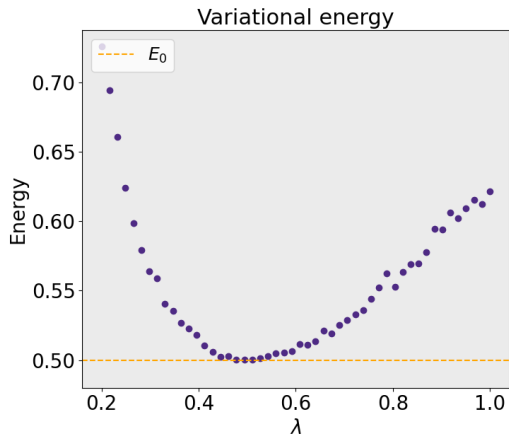


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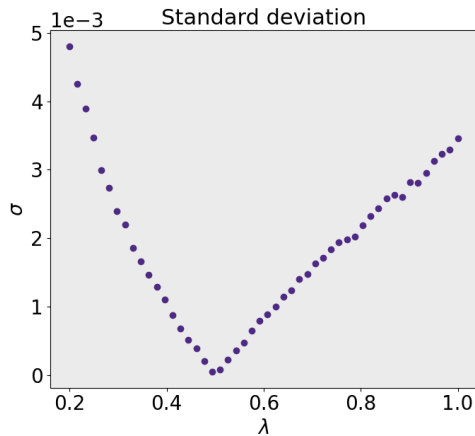
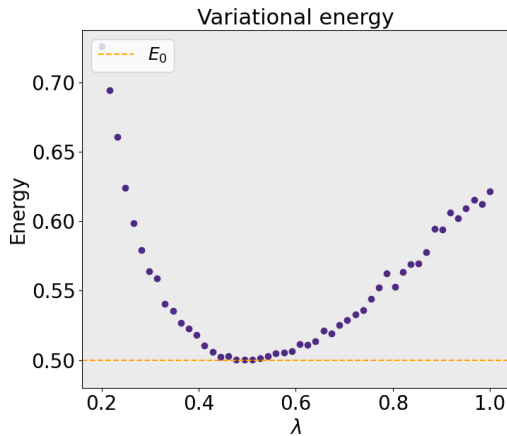
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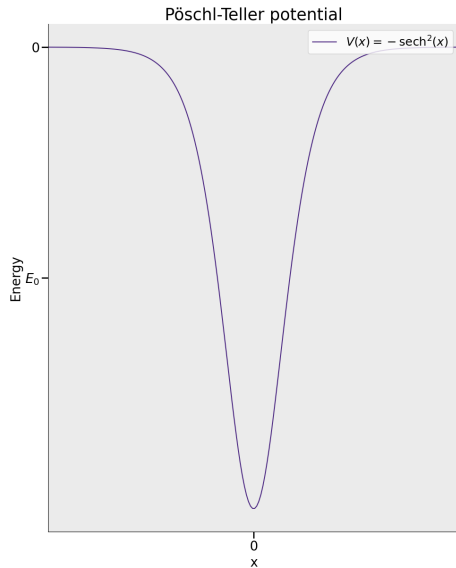
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- local energy: $E_{\text{local}} = \lambda - 2\lambda^2 x^2 - \text{sech}^2(x)$



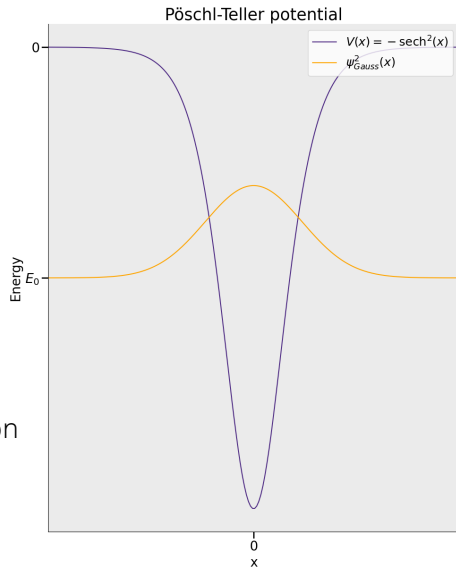
Pöschl-Teller potential

- Hamiltonian: $H = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \text{sech}^2(x)$
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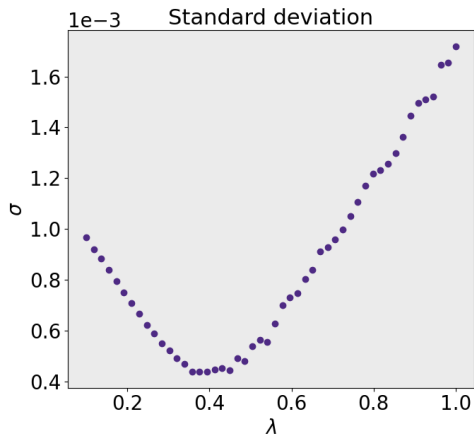
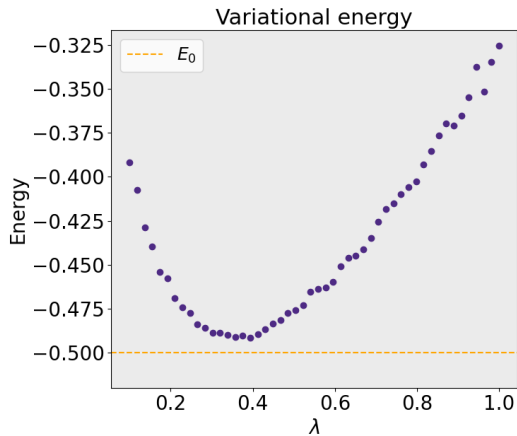
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- exact solution: $E_0 = -\frac{1}{2}$
- the variational ansatz is not the exact solution
 $\Rightarrow E_0 < E(x; \lambda_{\min})$



Pöschl-Teller-Potential - Groundstate

Pöschl-Teller potential with Gaussian trial function



Hydrogen atom

- assume no angular momentum ($l = 0$)

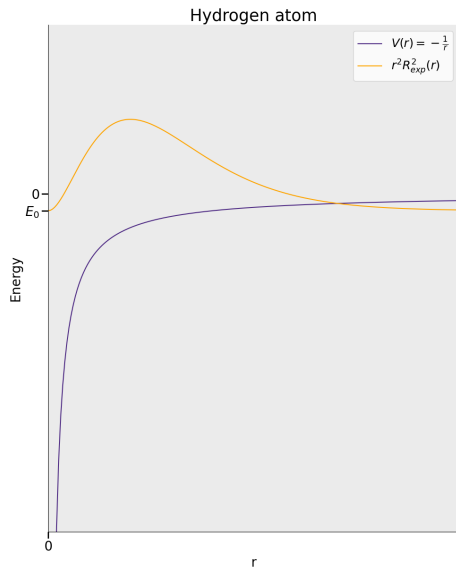
Hydrogen atom

- assume no angular momentum ($l = 0$)
- radial Hamiltonian: $H = -\frac{1}{r} \frac{\partial}{\partial r} - \frac{1}{2} \frac{\partial^2}{\partial r^2} - \frac{1}{r}$
($\hbar = e = m_e = 4\pi\epsilon_0 = 1$)



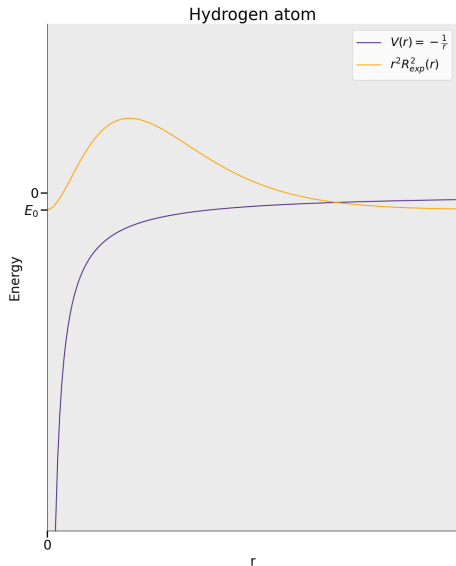
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- variational ansatz: $R_{\text{exp}}(r; \lambda) = e^{-\lambda r}$



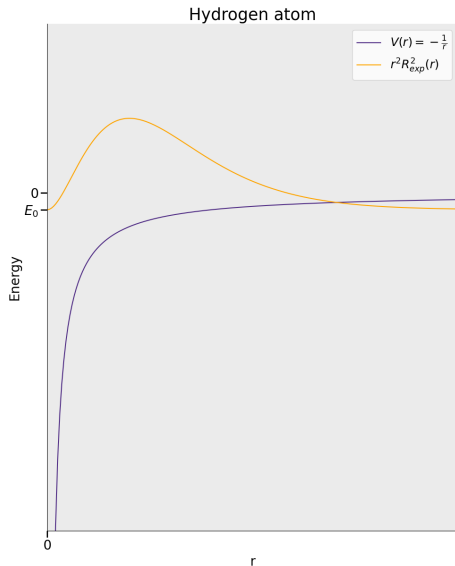
Hydrogen atom

- assume no angular momentum ($l = 0$)
- radial Hamiltonian: $H = -\frac{1}{r} \frac{\partial}{\partial r} - \frac{1}{2} \frac{\partial^2}{\partial r^2} - \frac{1}{r}$
($\hbar = e = m_e = 4\pi\epsilon_0 = 1$)
- variational ansatz: $R_{\text{exp}}(r; \lambda) = e^{-\lambda r}$
- local energy: $E_{\text{local}} = \frac{\lambda}{r} - \frac{\lambda^2}{2} - \frac{1}{r}$



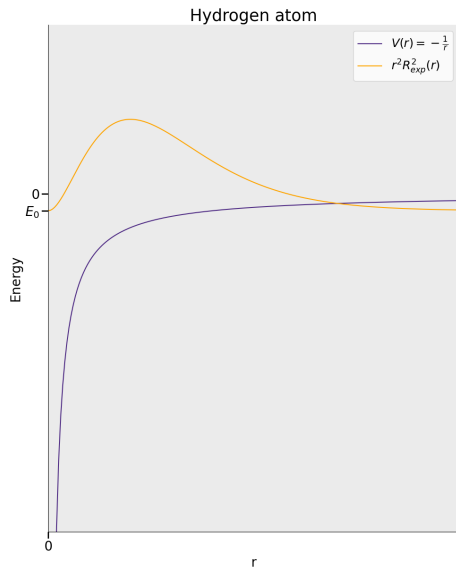
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- assume no angular momentum ($l = 0$)
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- local energy: $E_{\text{local}} = \frac{\lambda}{r} - \frac{\lambda^2}{2} - \frac{1}{r}$
- exact solution: $E_0 = -\frac{1}{2}$



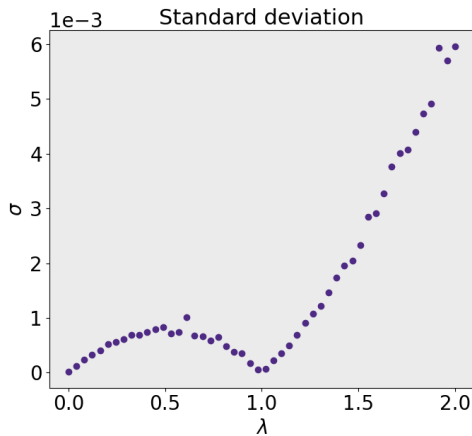
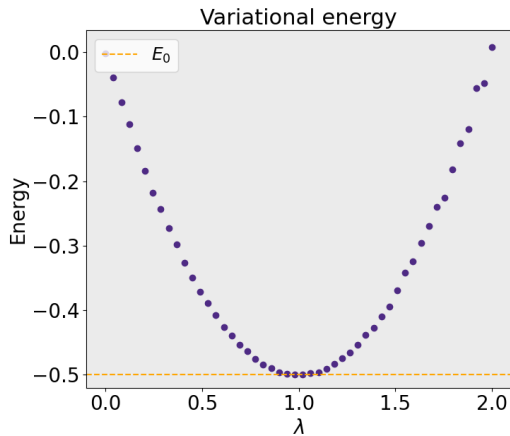
Hydrogen atom

- assume no angular momentum ($l = 0$)
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- variational ansatz: $R_{\text{exp}}(r; \lambda) = e^{-\lambda r}$
- local energy: $E_{\text{local}} = \frac{\lambda}{r} - \frac{\lambda^2}{2} - \frac{1}{r}$
- exact solution: $E_0 = -\frac{1}{2}$
- for $\lambda = 1$ the variational ansatz is the exact solution
 $\Rightarrow E(x; \lambda = 1) \equiv E_0 = -\frac{1}{2}$



Hydrogen atom - Groundstate

Hydrogen atom with exponentially decaying trial function.



Conclusion

1. using variational methods we can find an upper bound for the groundstate energy $E_0 \leq E(\lambda)$ of a quantum system if we have a good guess for the wavefunction

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Conclusion

1. using variational methods we can find an upper bound for the groundstate energy $E_0 \leq E(\lambda)$ of a quantum system if we have a good guess for the wavefunction
2. for many particles this involves solving high dimensional integrals for which stochastic integration methods are superior to quadrature rules
3. using importance sampling rather than simple sampling reduces the error and therefore the needed computation time significantly