

Variational methods for finding ground states of quantum systems

Jonathan Schümann

Dezember 12, 2024

Outline

1. Physical task
2. Numerical task
 - 2.1. Simple sampling
 - 2.2. Importance sampling
3. Examples

Physical task

[Pottorf et al; Eur. J. Phys. 20 (1999) 205]

[Gould et al; An Introduction to Computer Simulation Methods (Addition Wesley, 2006); Chapt. 16.7]

Stationary Schrödinger equation

- given a Hilbert space \mathcal{H} and a Hamiltonian H the stationary Schrödinger equation reads:

$$H\psi_{i,\alpha}(\mathbf{x}) = E_i\psi_{i,\alpha}(\mathbf{x}), \psi_{i,\alpha} \in \mathcal{H}, E_i \in \mathbb{R}$$

Stationary Schrödinger equation

- given a Hilbert space \mathcal{H} and a Hamiltonian H the stationary Schrödinger equation reads:

$$H\psi_{i,\alpha}(\mathbf{x}) = E_i\psi_{i,\alpha}(\mathbf{x}), \psi_{i,\alpha} \in \mathcal{H}, E_i \in \mathbb{R}$$

- task: find the groundstate energy $E_0 \leq E_i \forall i$

Stationary Schrödinger equation

- given a Hilbert space \mathcal{H} and a Hamiltonian H the stationary Schrödinger equation reads:

$$H\psi_{i,\alpha}(\mathbf{x}) = E_i\psi_{i,\alpha}(\mathbf{x}), \psi_{i,\alpha} \in \mathcal{H}, E_i \in \mathbb{R}$$

- task: find the groundstate energy $E_0 \leq E_i \forall i$
- rarely analytically solvable

Disadvantages of diagonalizing methods

- if H is discretized by a $(n \times n)$ -matrix the time complexity is $\mathcal{O}(n^3)$
 \Rightarrow high spatial resolution is computationally expensive

Disadvantages of diagonalizing methods

- if H is discretized by a $(n \times n)$ -matrix the time complexity is $\mathcal{O}(n^3)$
 \Rightarrow high spatial resolution is computationally expensive
- difficult to implement boundary conditions

Disadvantages of diagonalizing methods

- if H is discretized by a $(n \times n)$ -matrix the time complexity is $\mathcal{O}(n^3)$
 \Rightarrow high spatial resolution is computationally expensive
- difficult to implement boundary conditions
- we're not interested in any other eigenenergies $E_i, i \neq 0$

Disadvantages of diagonalizing methods

- if H is discretized by a $(n \times n)$ -matrix the time complexity is $\mathcal{O}(n^3)$
 \Rightarrow high spatial resolution is computationally expensive
- difficult to implement boundary conditions
- we're not interested in any other eigenenergies $E_i, i \neq 0$
- we're not interested in the wavefunctions $\psi_i(\mathbf{x})$

Disadvantages of diagonalizing methods

- if H is discretized by a $(n \times n)$ -matrix the time complexity is $\mathcal{O}(n^3)$
 \Rightarrow high spatial resolution is computationally expensive
- difficult to implement boundary conditions
- we're not interested in any other eigenenergies $E_i, i \neq 0$
- we're not interested in the wavefunctions $\psi_i(\mathbf{x})$
- no physical intuition is used

Variational method

- consider only one spatial dimension ($\mathbf{x} = x$)

Variational method

- consider only one spatial dimension ($\mathbf{x} = x$)
- we're interested in bound states $\Rightarrow \psi_0(x)$ can be assumed to be real function

Variational method

- consider only one spatial dimension ($\mathbf{x} = x$)
- we're interested in bound states $\Rightarrow \psi_0(x)$ can be assumed to be real function
- choose a physically reasonable variational ansatz $\psi(x; \lambda)$

Variational method

- consider only one spatial dimension ($\mathbf{x} = x$)
- we're interested in bound states $\Rightarrow \psi_0(x)$ can be assumed to be real function
- choose a physically reasonable variational ansatz $\psi(x; \lambda)$
- one can find an upper bound for the groundstate energy:

$$E_0 \leq E[\psi(x; \lambda)] = E(\lambda) = \frac{\int_{-\infty}^{\infty} dx \psi(x; \lambda) H \psi(x; \lambda)}{\int_{-\infty}^{\infty} dx \psi(x; \lambda) \psi(x; \lambda)}$$

Variational method

- consider only one spatial dimension ($\mathbf{x} = x$)
- we're interested in bound states $\Rightarrow \psi_0(x)$ can be assumed to be real function
- choose a physically reasonable variational ansatz $\psi(x; \lambda)$
- one can find an upper bound for the groundstate energy:

$$E_0 \leq E[\psi(x; \lambda)] = E(\lambda) = \frac{\int_{-\infty}^{\infty} dx \psi(x; \lambda) H \psi(x; \lambda)}{\int_{-\infty}^{\infty} dx \psi(x; \lambda) \psi(x; \lambda)}$$

- solve the integral and find the minimizing variational parameter λ_{min}

Numerical task

[Jeckelmann; Lecture 07: Monte-Carlo Integration (Lecture on Computerphysik, 2022)]

[Gould et al.; An Introduction to Computer Simulation Methods (Addition Wesley, 2006); Chapt. 11]

Quadrature rules

- let B be a domain $B \subset \mathbb{R}^D$ and f a function $f: \mathbb{R}^D \rightarrow \mathbb{R}$

Quadrature rules

- let B be a domain $B \subset \mathbb{R}^D$ and f a function $f: \mathbb{R}^D \rightarrow \mathbb{R}$
- task: approximate the integral $I = \int_B dx^D f(\mathbf{x})$

Quadrature rules

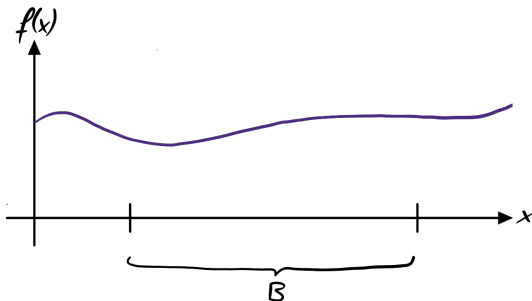
- let B be a domain $B \subset \mathbb{R}^D$ and f a function $f: \mathbb{R}^D \rightarrow \mathbb{R}$
- task: approximate the integral $I = \int_B dx^D f(\mathbf{x})$
- discretize B into M cells of volume $a^D = \frac{\text{vol}(B)}{M}$

Quadrature rules

- let B be a domain $B \subset \mathbb{R}^D$ and f a function $f: \mathbb{R}^D \rightarrow \mathbb{R}$
- task: approximate the integral $I = \int_B dx^D f(\mathbf{x})$
- discretize B into M cells of volume $a^D = \frac{\text{vol}(B)}{M}$
- general quadrature rule: $S_M^Q = \sum_{j=1}^M w_j f(\mathbf{x}_j)$

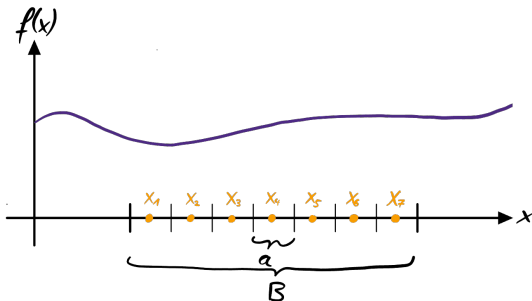
Quadrature rules

- let B be a domain $B \subset \mathbb{R}^D$ and f a function $f: \mathbb{R}^D \rightarrow \mathbb{R}$
- task: approximate the integral $I = \int_B dx^D f(\mathbf{x})$
- discretize B into M cells of volume $a^D = \frac{\text{vol}(B)}{M}$
- general quadrature rule: $S_M^Q = \sum_{j=1}^M w_j f(\mathbf{x}_j)$



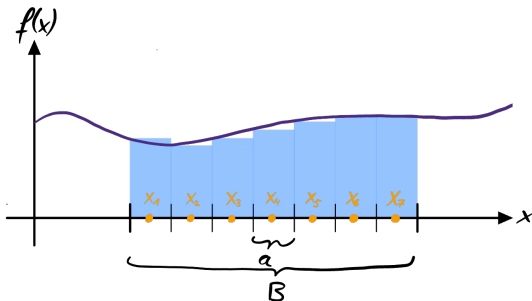
Quadrature rules

- let B be a domain $B \subset \mathbb{R}^D$ and f a function $f: \mathbb{R}^D \rightarrow \mathbb{R}$
- task: approximate the integral $I = \int_B dx^D f(\mathbf{x})$
- discretize B into M cells of volume $a^D = \frac{\text{vol}(B)}{M}$
- general quadrature rule: $S_M^Q = \sum_{j=1}^M w_j f(\mathbf{x}_j)$



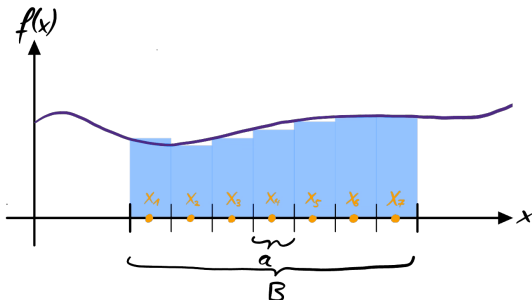
Quadrature rules

- let B be a domain $B \subset \mathbb{R}^D$ and f a function $f: \mathbb{R}^D \rightarrow \mathbb{R}$
- task: approximate the integral $I = \int_B dx^D f(\mathbf{x})$
- discretize B into M cells of volume $a^D = \frac{\text{vol}(B)}{M}$
- general quadrature rule: $S_M^Q = \sum_{j=1}^M w_j f(\mathbf{x}_j)$



Quadrature rules

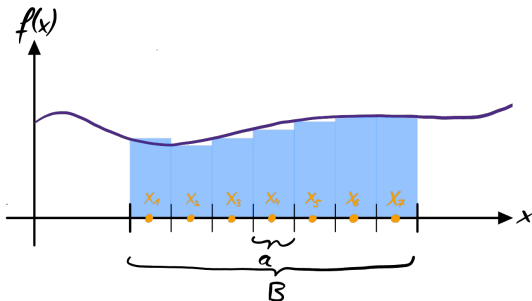
- let B be a domain $B \subset \mathbb{R}^D$ and f a function $f: \mathbb{R}^D \rightarrow \mathbb{R}$
- task: approximate the integral $I = \int_B dx^D f(\mathbf{x})$
- discretize B into M cells of volume $a^D = \frac{\text{vol}(B)}{M}$
- general quadrature rule: $S_M^Q = \sum_{j=1}^M w_j f(\mathbf{x}_j)$



- error of quadrature rule of order k : $|S_M^Q - I| \propto a^k \propto M^{-\frac{k}{D}}$

Quadrature rules

- let B be a domain $B \subset \mathbb{R}^D$ and f a function $f: \mathbb{R}^D \rightarrow \mathbb{R}$
- task: approximate the integral $I = \int_B dx^D f(\mathbf{x})$
- discretize B into M cells of volume $a^D = \frac{\text{vol}(B)}{M}$
- general quadrature rule: $S_M^Q = \sum_{j=1}^M w_j f(\mathbf{x}_j)$



- error of quadrature rule of order k : $|S_M^Q - I| \propto a^k \propto M^{-\frac{k}{D}}$
- problem: slow convergence for high dimension D

Simple sampling

- interpret X as a random variable with probability distribution $\rho(\mathbf{x}) = \frac{1}{\text{vol}(B)}$

Simple sampling

- interpret X as a random variable with probability distribution $\rho(\mathbf{x}) = \frac{1}{\text{vol}(B)}$
- rewrite the integral:

$$I = \int_B dx^D f(\mathbf{x}) = \int_B dx^D f(\mathbf{x}) \frac{\text{vol}(B)}{\text{vol}(B)} = \left(\int_B dx^D f(\mathbf{x}) \rho(\mathbf{x}) \right) \text{vol}(B) = \mathbb{E}[f(X)] \text{vol}(B)$$

Simple sampling

- interpret X as a random variable with probability distribution $\rho(\mathbf{x}) = \frac{1}{\text{vol}(B)}$
- rewrite the integral:

$$I = \int_B dx^D f(\mathbf{x}) = \int_B dx^D f(\mathbf{x}) \frac{\text{vol}(B)}{\text{vol}(B)} = \left(\int_B dx^D f(\mathbf{x}) \rho(\mathbf{x}) \right) \text{vol}(B) = \mathbb{E}[f(X)] \text{vol}(B)$$

- generate M random sample points $\mathbf{x}_j \in B$ with distribution $\rho(\mathbf{x})$

Simple sampling

- interpret X as a random variable with probability distribution $\rho(\mathbf{x}) = \frac{1}{\text{vol}(B)}$
- rewrite the integral:

$$I = \int_B dx^D f(\mathbf{x}) = \int_B dx^D f(\mathbf{x}) \frac{\text{vol}(B)}{\text{vol}(B)} = \left(\int_B dx^D f(\mathbf{x}) \rho(\mathbf{x}) \right) \text{vol}(B) = \mathbb{E}[f(X)] \text{vol}(B)$$

- generate M random sample points $\mathbf{x}_j \in B$ with distribution $\rho(\mathbf{x})$
- approximate the expectation value by an average:

$$\mathbb{E}[f(X)] \approx \text{avg}(f) = \frac{1}{M} \sum_{j=1}^M f(\mathbf{x}_j)$$

Simple sampling

- interpret X as a random variable with probability distribution $\rho(\mathbf{x}) = \frac{1}{\text{vol}(B)}$
- rewrite the integral:

$$I = \int_B dx^D f(\mathbf{x}) = \int_B dx^D f(\mathbf{x}) \frac{\text{vol}(B)}{\text{vol}(B)} = \left(\int_B dx^D f(\mathbf{x}) \rho(\mathbf{x}) \right) \text{vol}(B) = \mathbb{E}[f(X)] \text{vol}(B)$$

- generate M random sample points $\mathbf{x}_j \in B$ with distribution $\rho(\mathbf{x})$
- approximate the expectation value by an average:

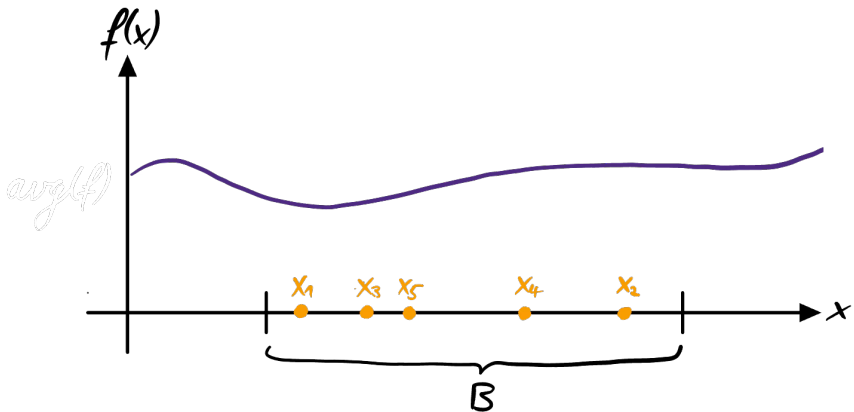
$$\mathbb{E}[f(X)] \approx \text{avg}(f) = \frac{1}{M} \sum_{j=1}^M f(\mathbf{x}_j)$$

- Monte-Carlo sum:

$$S_M^{MC} = \text{vol}(B) \text{avg}(f) = \frac{\text{vol}(B)}{M} \sum_{j=1}^M f(\mathbf{x}_j)$$

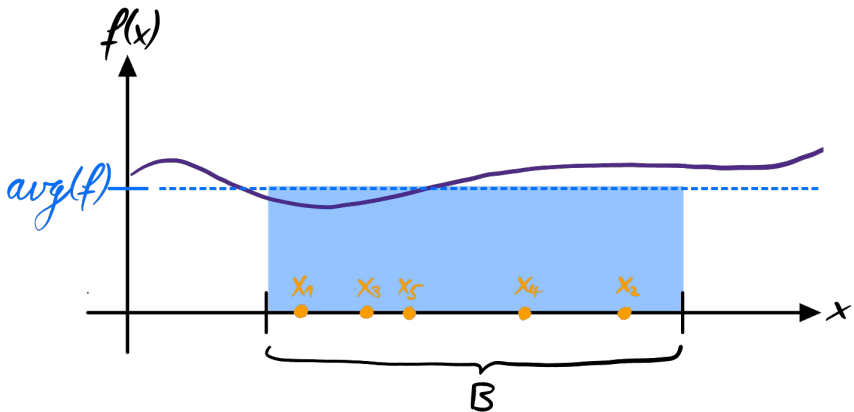
Simple sampling

- Monte-Carlo sum: $S_M^{MC} = \text{vol}(B) \text{avg}(f) = \frac{\text{vol}(B)}{M} \sum_{j=1}^M f(\mathbf{x}_j)$



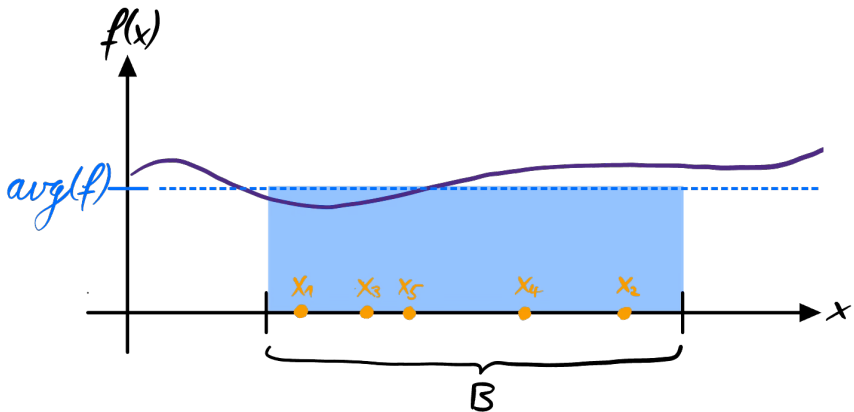
Simple sampling

- Monte-Carlo sum: $S_M^{MC} = \text{vol}(B) \text{avg}(f) = \frac{\text{vol}(B)}{M} \sum_{j=1}^M f(\mathbf{x}_j)$



Simple sampling

- Monte-Carlo sum: $S_M^{MC} = \text{vol}(B) \text{avg}(f) = \frac{\text{vol}(B)}{M} \sum_{j=1}^M f(\mathbf{x}_j)$



- error: $|I - S_M^{MC}| \propto M^{-\frac{1}{2}}$

Simple sampling

- the Monte-Carlo sum is a normal distributed random variable

$$P(S_M^{MC}) = \frac{1}{\sqrt{2\pi}\sigma_M} \exp\left(-\frac{(S_M^{MC} - I)^2}{2\sigma_M^2}\right)$$

Simple sampling

- the Monte-Carlo sum is a normal distributed random variable

$$P(S_M^{MC}) = \frac{1}{\sqrt{2\pi}\sigma_M} \exp\left(-\frac{(S_M^{MC} - I)^2}{2\sigma_M^2}\right)$$

- approximate the integral by averaging N Monte-Carlo sums:

$$I \approx \frac{1}{N} \sum_{i=1}^N (S_M^{MC})^{(i)} = \frac{1}{N} \frac{\text{vol}(B)}{M} \sum_{i=1}^N \sum_{j=1}^M f(\mathbf{x}_j^{(i)})$$

Simple sampling

- the Monte-Carlo sum is a normal distributed random variable

$$P(S_M^{MC}) = \frac{1}{\sqrt{2\pi}\sigma_M} \exp\left(-\frac{(S_M^{MC} - I)^2}{2\sigma_M^2}\right)$$

- approximate the integral by averaging N Monte-Carlo sums:

$$I \approx \frac{1}{N} \sum_{i=1}^N (S_M^{MC})^{(i)} = \frac{1}{N} \frac{\text{vol}(B)}{M} \sum_{i=1}^N \sum_{j=1}^M f(\mathbf{x}_j^{(i)})$$

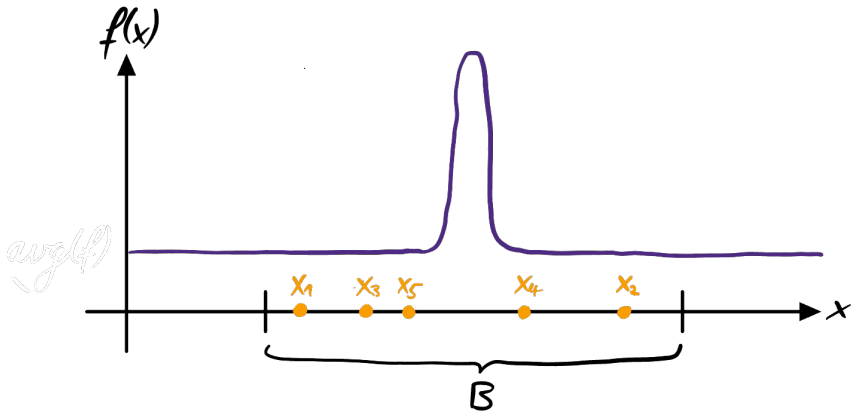
- problem: $\text{vol}(B)$ might be unknown and σ_M might be very large

Simple sampling

- high σ_M for rapidly changing $f(x)$:

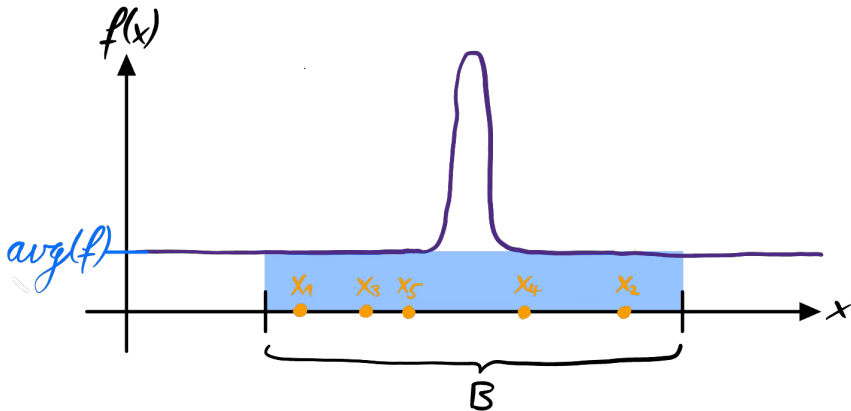
Simple sampling

- high σ_M for rapidly changing $f(x)$:



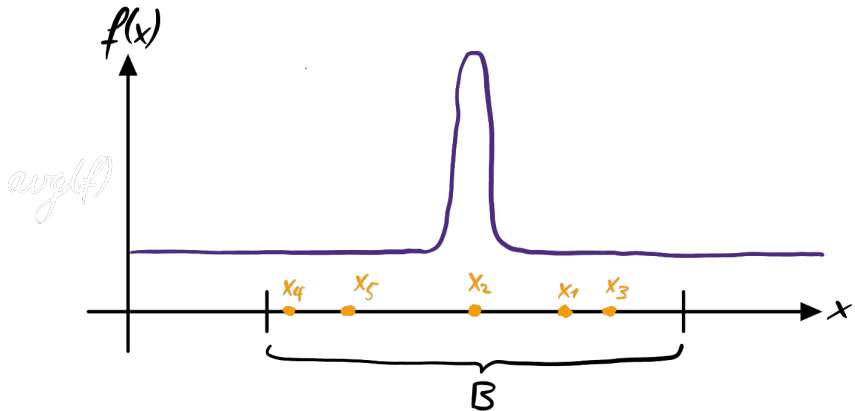
Simple sampling

- high σ_M for rapidly changing $f(x)$:



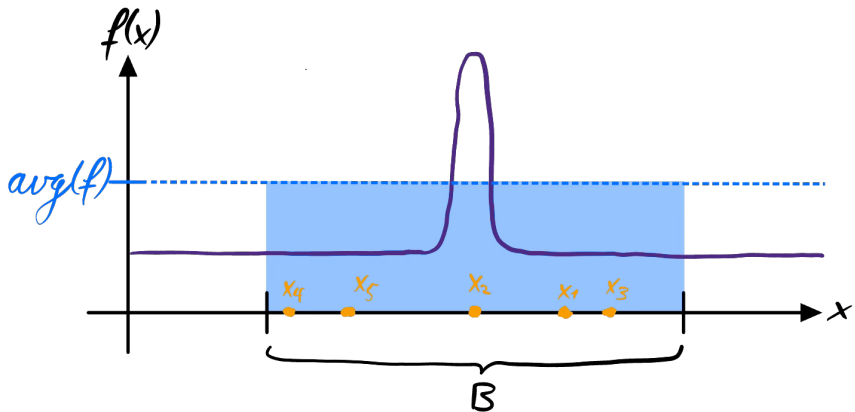
Simple sampling

- high σ_M for rapidly changing $f(x)$:



Simple sampling

- high σ_M for rapidly changing $f(x)$:



Importance sampling

- choose a probability distribution $\rho(\mathbf{x})$ such that $g(\mathbf{x}) := \frac{f(\mathbf{x})}{\rho(\mathbf{x})}$ varies slowly

Importance sampling

- choose a probability distribution $\rho(\mathbf{x})$ such that $g(\mathbf{x}) := \frac{f(\mathbf{x})}{\rho(\mathbf{x})}$ varies slowly
- rewrite the integral:

$$I = \int_B dx^D f(\mathbf{x}) = \int_B dx^D \frac{\rho(\mathbf{x})}{\rho(\mathbf{x})} f(\mathbf{x}) = \int_B dx^D \rho(\mathbf{x}) g(\mathbf{x}) = \mathbb{E}[g(X)]$$

Importance sampling

- choose a probability distribution $\rho(\mathbf{x})$ such that $g(\mathbf{x}) := \frac{f(\mathbf{x})}{\rho(\mathbf{x})}$ varies slowly
- rewrite the integral:

$$I = \int_B dx^D f(\mathbf{x}) = \int_B dx^D \frac{\rho(\mathbf{x})}{\rho(\mathbf{x})} f(\mathbf{x}) = \int_B dx^D \rho(\mathbf{x}) g(\mathbf{x}) = \mathbb{E}[g(X)]$$

- generate M random sample points $\mathbf{x}_j \in B$ with distribution $\rho(\mathbf{x})$

Importance sampling

- choose a probability distribution $\rho(\mathbf{x})$ such that $g(\mathbf{x}) := \frac{f(\mathbf{x})}{\rho(\mathbf{x})}$ varies slowly
- rewrite the integral:

$$I = \int_B dx^D f(\mathbf{x}) = \int_B dx^D \frac{\rho(\mathbf{x})}{\rho(\mathbf{x})} f(\mathbf{x}) = \int_B dx^D \rho(\mathbf{x}) g(\mathbf{x}) = \mathbb{E}[g(X)]$$

- generate M random sample points $\mathbf{x}_j \in B$ with distribution $\rho(\mathbf{x})$
- Monte-carlo sum:

$$S_M^{MC} = \frac{1}{M} \sum_{j=1}^M g(\mathbf{x}_j)$$

Importance sampling

- choose a probability distribution $\rho(\mathbf{x})$ such that $g(\mathbf{x}) := \frac{f(\mathbf{x})}{\rho(\mathbf{x})}$ varies slowly
- rewrite the integral:

$$I = \int_B dx^D f(\mathbf{x}) = \int_B dx^D \frac{\rho(\mathbf{x})}{\rho(\mathbf{x})} f(\mathbf{x}) = \int_B dx^D \rho(\mathbf{x}) g(\mathbf{x}) = \mathbb{E}[g(X)]$$

- generate M random sample points $\mathbf{x}_j \in B$ with distribution $\rho(\mathbf{x})$
- Monte-carlo sum:

$$S_M^{MC} = \frac{1}{M} \sum_{j=1}^M g(\mathbf{x}_j)$$

- approximate the integral by averaging N Monte-Carlo sums:

$$I \approx \frac{1}{N} \sum_{i=1}^N (S_M^{MC})^{(i)} = \frac{1}{N} \frac{1}{M} \sum_{i=1}^N \sum_{j=1}^M g(\mathbf{x}_j^{(i)})$$

Importance sampling for variational methods

- our goal is to calculate the variational energy $E(\lambda)$:

$$E(\lambda) = \frac{\int_{-\infty}^{\infty} dx \psi(x; \lambda) H \psi(x; \lambda)}{\int_{-\infty}^{\infty} dx \psi(x; \lambda) \psi(x; \lambda)} = \int_{-\infty}^{\infty} dx \left(\frac{\psi^2(x; \lambda)}{\int_{-\infty}^{\infty} dy \psi^2(y; \lambda)} \frac{H \psi(x; \lambda)}{\psi(x; \lambda)} \right)$$

Importance sampling for variational methods

- our goal is to calculate the variational energy $E(\lambda)$:

$$E(\lambda) = \frac{\int_{-\infty}^{\infty} dx \psi(x; \lambda) H \psi(x; \lambda)}{\int_{-\infty}^{\infty} dx \psi(x; \lambda) \psi(x; \lambda)} = \int_{-\infty}^{\infty} dx \left(\frac{\psi^2(x; \lambda)}{\int_{-\infty}^{\infty} dy \psi^2(y; \lambda)} \frac{H \psi(x; \lambda)}{\psi(x; \lambda)} \right)$$

- define:

$$\rho(x; \lambda) = \frac{\psi^2(x; \lambda)}{\int_{-\infty}^{\infty} dy \psi^2(y; \lambda)} \quad \text{and} \quad E_{\text{local}}(x; \lambda) = \frac{H \psi(x; \lambda)}{\psi(x; \lambda)}$$

Importance sampling for variational methods

- our goal is to calculate the variational energy $E(\lambda)$:

$$E(\lambda) = \frac{\int_{-\infty}^{\infty} dx \psi(x; \lambda) H \psi(x; \lambda)}{\int_{-\infty}^{\infty} dx \psi(x; \lambda) \psi(x; \lambda)} = \int_{-\infty}^{\infty} dx \left(\frac{\psi^2(x; \lambda)}{\int_{-\infty}^{\infty} dy \psi^2(y; \lambda)} \frac{H \psi(x; \lambda)}{\psi(x; \lambda)} \right)$$

- define:

$$\rho(x; \lambda) = \frac{\psi^2(x; \lambda)}{\int_{-\infty}^{\infty} dy \psi^2(y; \lambda)} \quad \text{and} \quad E_{local}(x; \lambda) = \frac{H \psi(x; \lambda)}{\psi(x; \lambda)}$$

- using importance sampling we find:

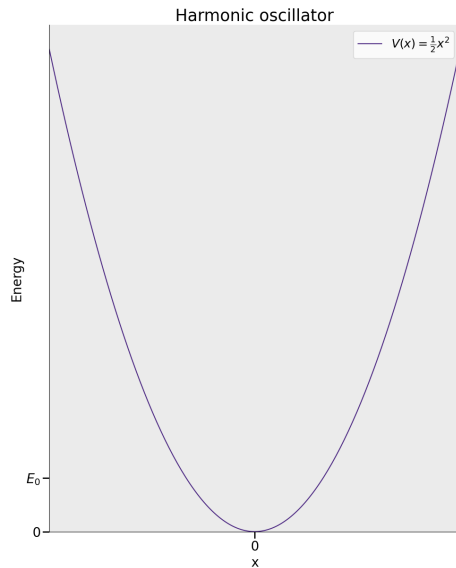
$$E(\lambda) = \int_{-\infty}^{\infty} dx \rho(x; \lambda) E_{local}(x; \lambda) \approx \frac{1}{N} \frac{1}{M} \sum_{i=1}^N \sum_{j=1}^M E_{local}(x_j^{(i)}; \lambda)$$

Examples

[<https://github.com/JSchuemann-physics/VariationalMethods>]

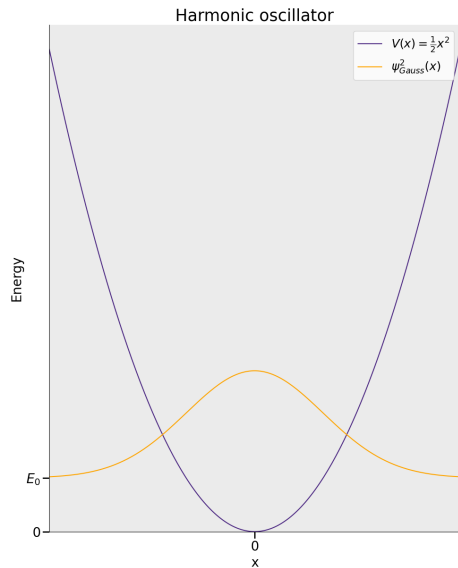
Harmonic oscillator

- Hamiltonian: $H = -\frac{1}{2}\frac{\partial^2}{\partial x^2} + \frac{1}{2}x^2$
($\hbar = m = 1$)



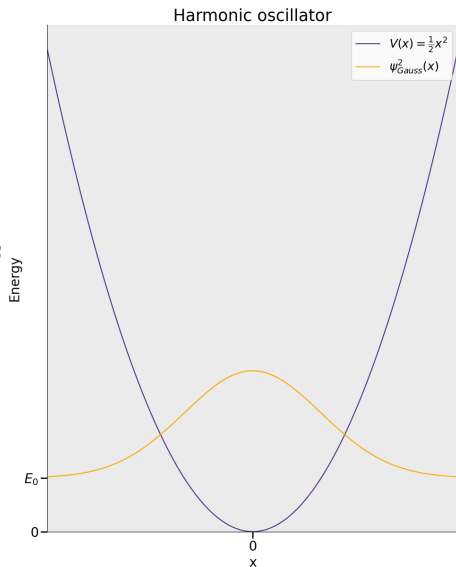
Harmonic oscillator

- Hamiltonian: $H = -\frac{1}{2}\frac{\partial^2}{\partial x^2} + \frac{1}{2}x^2$
($\hbar = m = 1$)
- variational ansatz: $\psi_{Gauss}(x; \lambda) = e^{-\lambda x^2}$



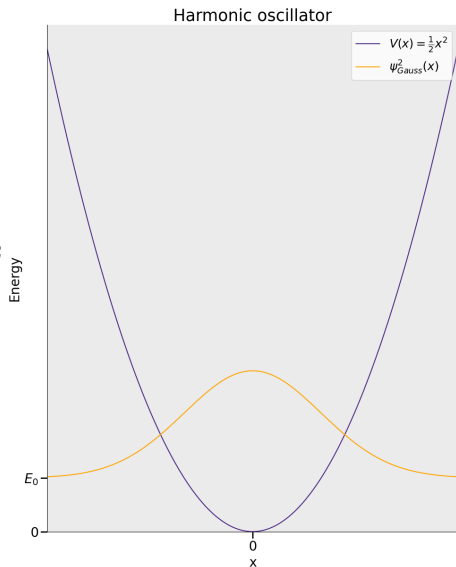
Harmonic oscillator

- Hamiltonian: $H = -\frac{1}{2}\frac{\partial^2}{\partial x^2} + \frac{1}{2}x^2$
($\hbar = m = 1$)
- variational ansatz: $\psi_{Gauss}(x; \lambda) = e^{-\lambda x^2}$
- local energy: $E_{local} = \frac{H\psi(x;\lambda)}{\psi(x;\lambda)} = \lambda - 2\lambda^2 x^2 + \frac{1}{2}x^2$



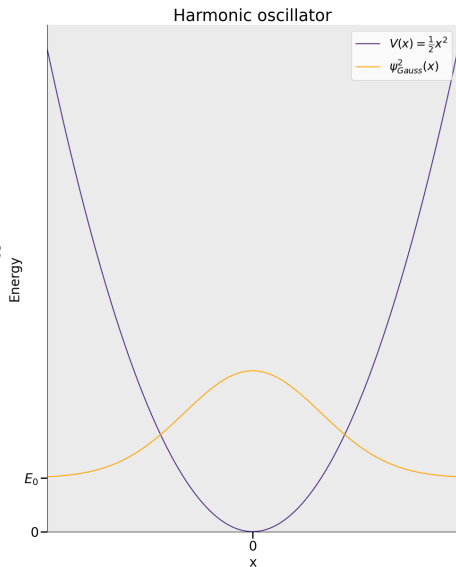
Harmonic oscillator

- Hamiltonian: $H = -\frac{1}{2}\frac{\partial^2}{\partial x^2} + \frac{1}{2}x^2$
($\hbar = m = 1$)
- variational ansatz: $\psi_{Gauss}(x; \lambda) = e^{-\lambda x^2}$
- local energy: $E_{local} = \frac{H\psi(x;\lambda)}{\psi(x;\lambda)} = \lambda - 2\lambda^2 x^2 + \frac{1}{2}x^2$
- exact solution: $E_0 = \frac{1}{2}$

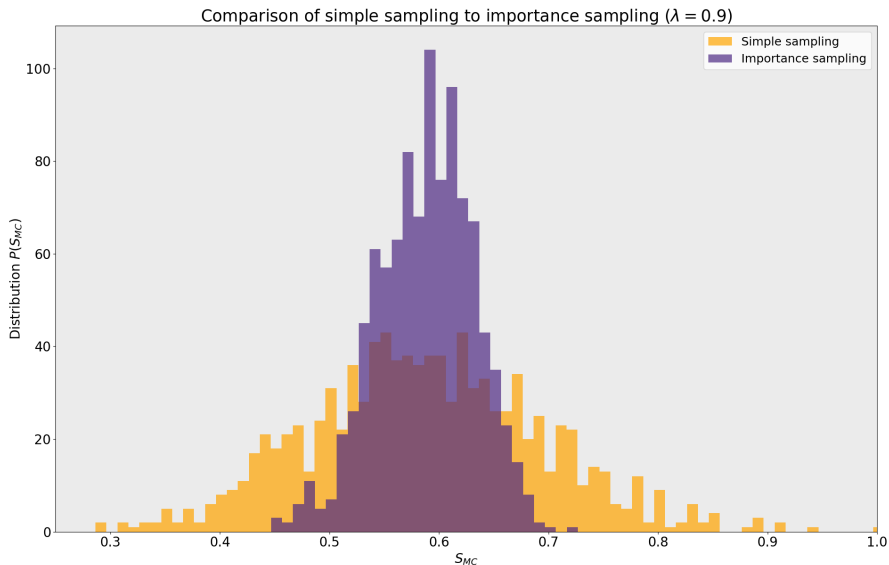


Harmonic oscillator

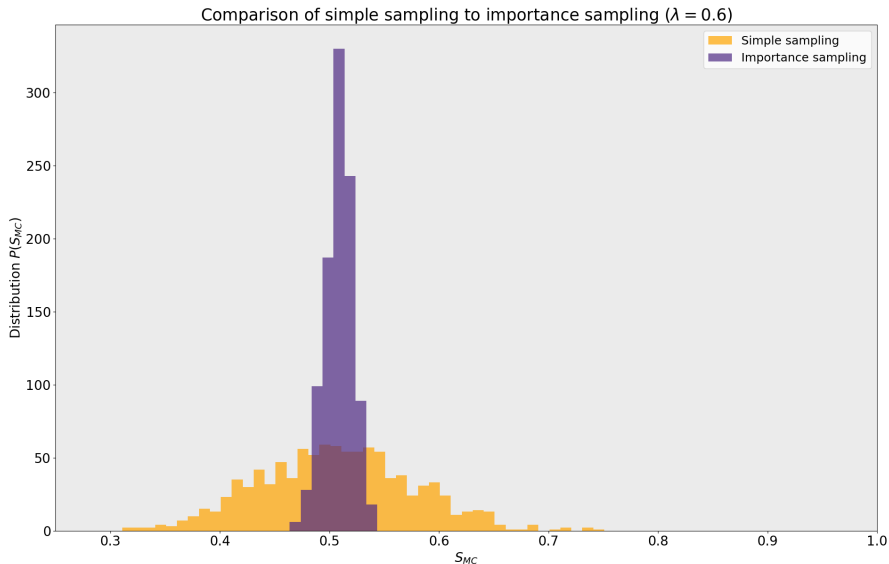
- Hamiltonian: $H = -\frac{1}{2}\frac{\partial^2}{\partial x^2} + \frac{1}{2}x^2$
($\hbar = m = 1$)
- variational ansatz: $\psi_{\text{Gauss}}(x; \lambda) = e^{-\lambda x^2}$
- local energy: $E_{\text{local}} = \frac{H\psi(x; \lambda)}{\psi(x; \lambda)} = \lambda - 2\lambda^2 x^2 + \frac{1}{2}x^2$
- exact solution: $E_0 = \frac{1}{2}$
- for $\lambda = \frac{1}{2}$ the variational ansatz is the exact solution
 $\Rightarrow E_{\text{local}}(x; \lambda = \frac{1}{2}) \equiv E_0 = \frac{1}{2}$



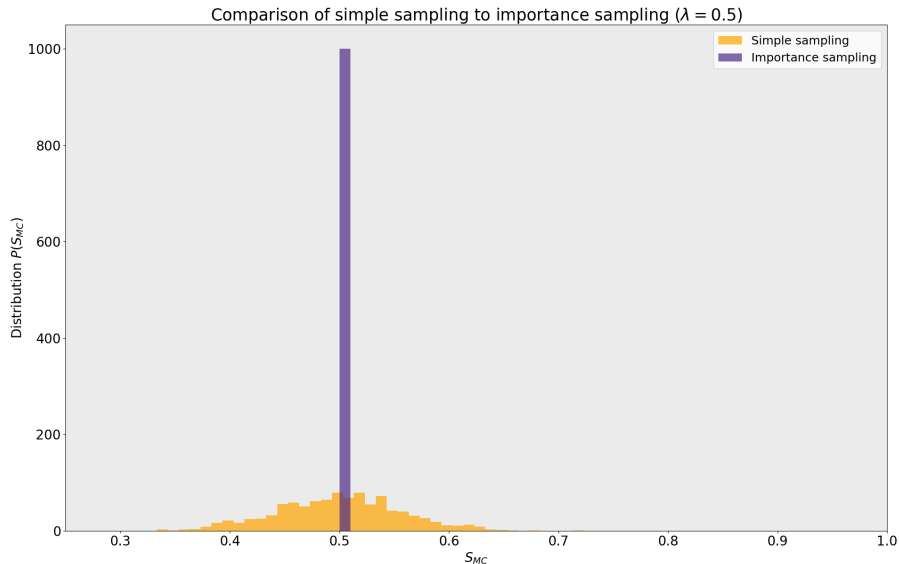
Simple sampling vs. importance sampling



Simple sampling vs. importance sampling

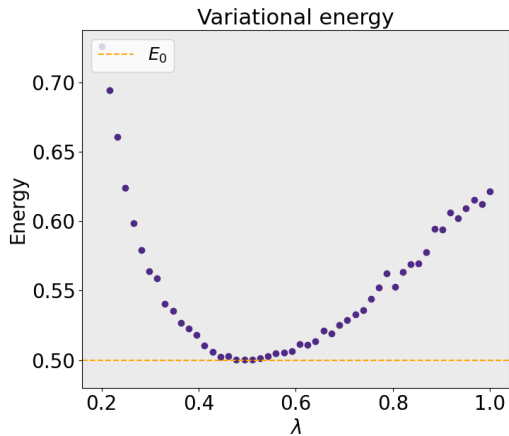


Simple sampling vs. importance sampling



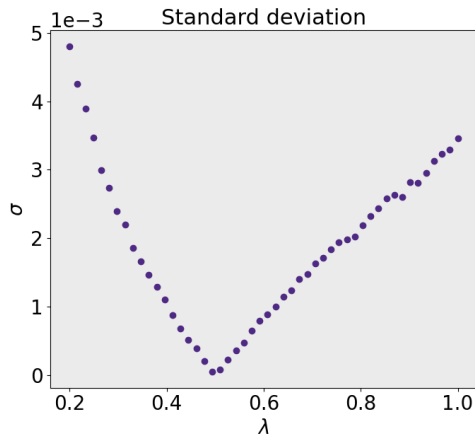
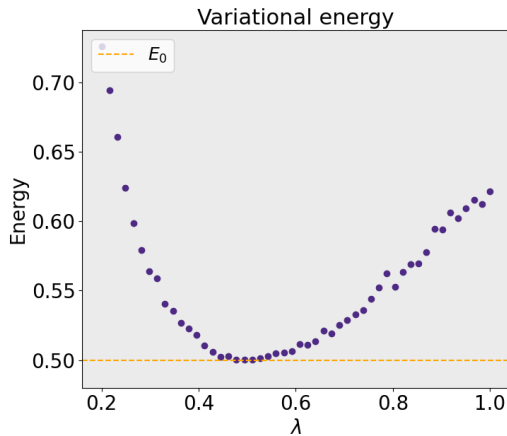
Harmonic Oscillator - Groundstate

Harmonic oscillator with Gaussian trial function



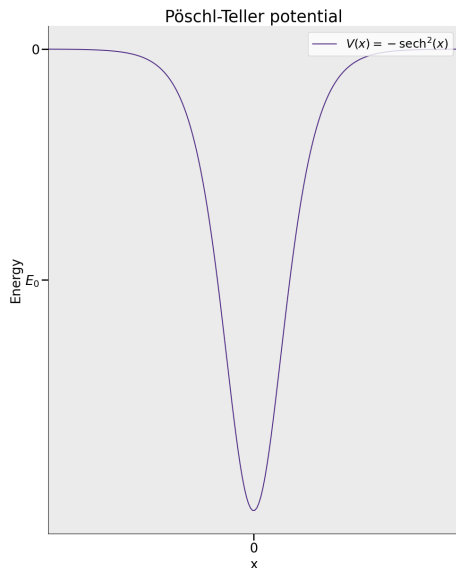
Harmonic Oscillator - Groundstate

Harmonic oscillator with Gaussian trial function



Pöschl-Teller potential

- Hamiltonian: $H = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \text{sech}^2(x)$
($\hbar = m = 1$)



Pöschl-Teller potential

- Hamiltonian: $H = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \text{sech}^2(x)$
($\hbar = m = 1$)
- variational ansatz: $\psi_{\text{Gauss}}(x; \lambda) = e^{-\lambda x^2}$



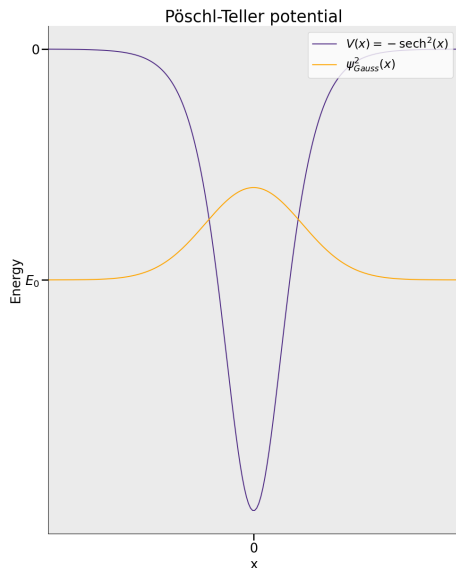
Pöschl-Teller potential

- Hamiltonian: $H = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \text{sech}^2(x)$
($\hbar = m = 1$)
- variational ansatz: $\psi_{\text{Gauss}}(x; \lambda) = e^{-\lambda x^2}$
- local energy: $E_{\text{local}} = \lambda - 2\lambda^2 x^2 - \text{sech}^2(x)$



Pöschl-Teller potential

- Hamiltonian: $H = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \text{sech}^2(x)$
($\hbar = m = 1$)
- variational ansatz: $\psi_{\text{Gauss}}(x; \lambda) = e^{-\lambda x^2}$
- local energy: $E_{\text{local}} = \lambda - 2\lambda^2 x^2 - \text{sech}^2(x)$
- exact solution: $E_0 = -\frac{1}{2}$



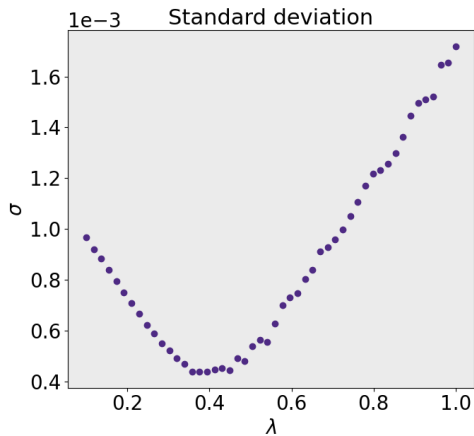
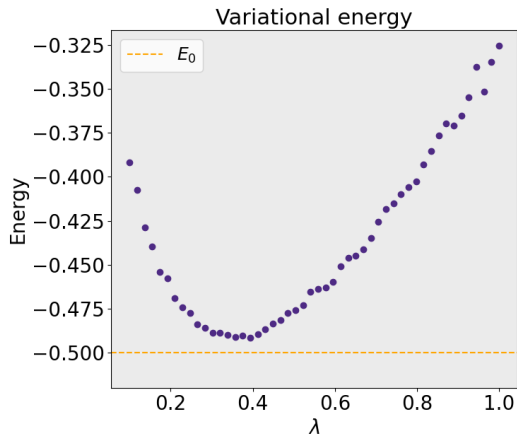
Pöschl-Teller potential

- Hamiltonian: $H = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \text{sech}^2(x)$
($\hbar = m = 1$)
- variational ansatz: $\psi_{\text{Gauss}}(x; \lambda) = e^{-\lambda x^2}$
- local energy: $E_{\text{local}} = \lambda - 2\lambda^2 x^2 - \text{sech}^2(x)$
- exact solution: $E_0 = -\frac{1}{2}$
- the variational ansatz is not the exact solution
 $\Rightarrow E_0 < E(\lambda_{\min})$



Pöschl-Teller-Potential - Groundstate

Pöschl-Teller potential with Gaussian trial function



Hydrogen atom

- assume no angular momentum ($l = 0$)

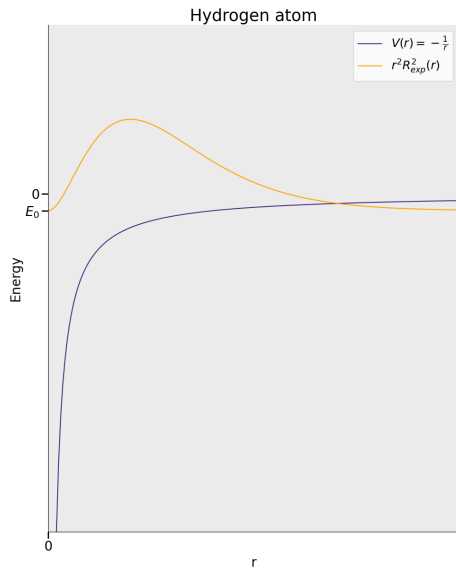
Hydrogen atom

- assume no angular momentum ($l = 0$)
- radial Hamiltonian: $H = -\frac{1}{r} \frac{\partial}{\partial r} - \frac{1}{2} \frac{\partial^2}{\partial r^2} - \frac{1}{r}$
($\hbar = e = m_e = 4\pi\epsilon_0 = 1$)



Hydrogen atom

- assume no angular momentum ($l = 0$)
- radial Hamiltonian: $H = -\frac{1}{r} \frac{\partial}{\partial r} - \frac{1}{2} \frac{\partial^2}{\partial r^2} - \frac{1}{r}$
($\hbar = e = m_e = 4\pi\epsilon_0 = 1$)
- variational ansatz: $R_{\text{exp}}(r; \lambda) = e^{-\lambda r}$



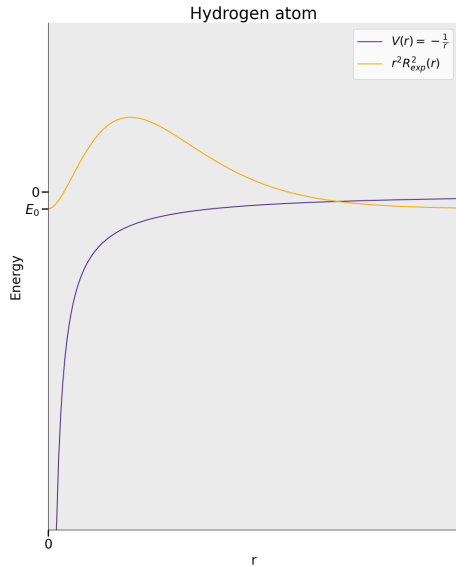
Hydrogen atom

- assume no angular momentum ($l = 0$)
- radial Hamiltonian: $H = -\frac{1}{r} \frac{\partial}{\partial r} - \frac{1}{2} \frac{\partial^2}{\partial r^2} - \frac{1}{r}$
($\hbar = e = m_e = 4\pi\epsilon_0 = 1$)
- variational ansatz: $R_{\text{exp}}(r; \lambda) = e^{-\lambda r}$
- local energy: $E_{\text{local}} = \frac{\lambda}{r} - \frac{\lambda^2}{2} - \frac{1}{r}$



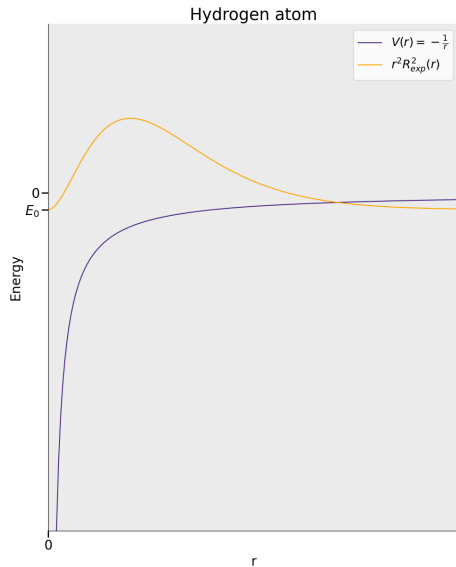
Hydrogen atom

- assume no angular momentum ($l = 0$)
- radial Hamiltonian: $H = -\frac{1}{r} \frac{\partial}{\partial r} - \frac{1}{2} \frac{\partial^2}{\partial r^2} - \frac{1}{r}$
($\hbar = e = m_e = 4\pi\epsilon_0 = 1$)
- variational ansatz: $R_{\text{exp}}(r; \lambda) = e^{-\lambda r}$
- local energy: $E_{\text{local}} = \frac{\lambda}{r} - \frac{\lambda^2}{2} - \frac{1}{r}$
- exact solution: $E_0 = -\frac{1}{2}$



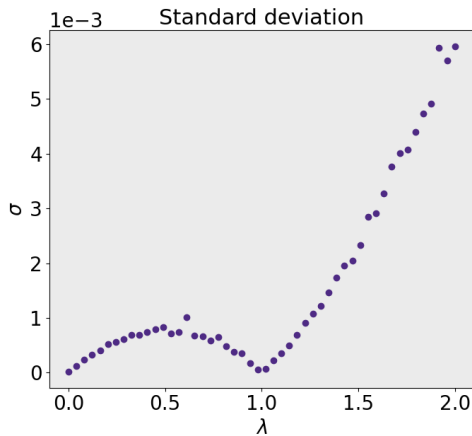
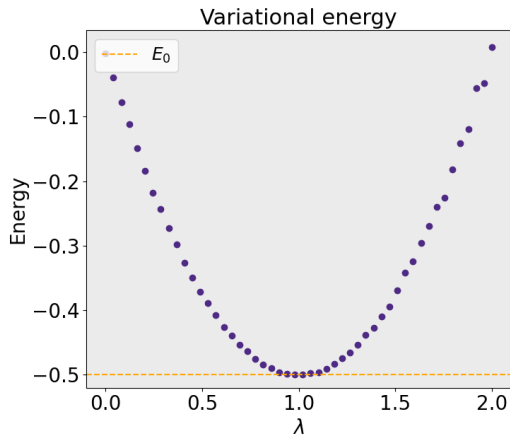
Hydrogen atom

- assume no angular momentum ($l = 0$)
- radial Hamiltonian: $H = -\frac{1}{r} \frac{\partial}{\partial r} - \frac{1}{2} \frac{\partial^2}{\partial r^2} - \frac{1}{r}$
($\hbar = e = m_e = 4\pi\epsilon_0 = 1$)
- variational ansatz: $R_{\text{exp}}(r; \lambda) = e^{-\lambda r}$
- local energy: $E_{\text{local}} = \frac{\lambda}{r} - \frac{\lambda^2}{2} - \frac{1}{r}$
- exact solution: $E_0 = -\frac{1}{2}$
- for $\lambda = 1$ the variational ansatz is the exact solution
 $\Rightarrow E(x; \lambda = 1) \equiv E_0 = -\frac{1}{2}$



Hydrogen atom - Groundstate

Hydrogen atom with exponentially decaying trial function.



Conclusion

1. using variational methods we can find an upper bound for the groundstate energy $E_0 \leq E(\lambda)$ of a quantum system if we have a good guess for the wavefunction

Conclusion

1. using variational methods we can find an upper bound for the groundstate energy $E_0 \leq E(\lambda)$ of a quantum system if we have a good guess for the wavefunction
2. for many particles this involves solving high dimensional integrals for which stochastical integration methods are superior to quadrature rules

Conclusion

1. using variational methods we can find an upper bound for the groundstate energy $E_0 \leq E(\lambda)$ of a quantum system if we have a good guess for the wavefunction
2. for many particles this involves solving high dimensional integrals for which stochastical integration methods are superior to quadrature rules
3. using importance sampling rather than simple sampling reduces the error and therefore the needed computation time significantly