

# A Search for Resonant and Non-Resonant Di-Higgs Production in the $\gamma\gamma bb$ Channel Using the ATLAS Detector

Dissertation Defense  
Tyler James Burch  
March 26, 2020

**Committee:**

Jahred Adelman (Chair)  
Dhiman Chakraborty  
Michael Eads



# Welcome!

Thanks for staying flexible!

## **How this defense will proceed:**

- Will start with this presentation (~1 hour) in this room
  - Open Zoom meeting, plus cross-streaming to YouTube
- After, committee and I will transition to a separate room to conclude
- Once we've finished I'll return to this room, and the committee will deliberate and return to this room afterwards

## **Some general guidelines:**

- Please mute if not speaking
- Given Brianna Stamas room host privileges to manage low-level Zoom concerns.

Thanks Bri!

**Thank you for attending!**



# Outline

## Introduction

- The Standard Model
- The LHC and ATLAS Detector
- Di-Higgs Production

$\text{HH} \rightarrow \gamma\gamma\text{bb}$  search using 2015+2016 data

## VBF $\text{HH} \rightarrow \gamma\gamma\text{bb}$

- Motivation
- Integrating into  $\gamma\gamma\text{bb}$  analysis

## Photon Identification Optimization

- Background and current methods
- Topological Cluster inputs
- MVA Approaches

## Conclusion



# Introduction



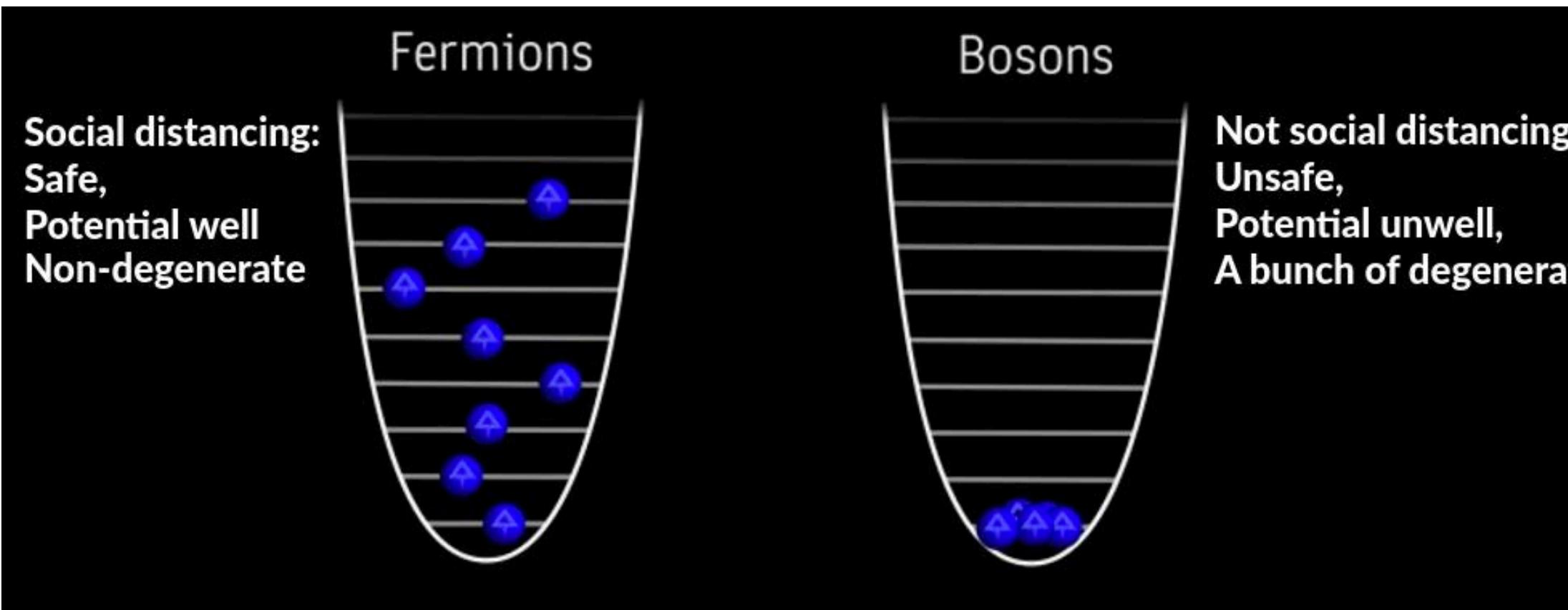
# The Standard Model (SM)

Best working description of the fundamental particles and their interactions

**Fermions** - particles that “take up space.”  
(Precisely, they obey the Pauli exclusion principle)

**Bosons** - force carriers, explain interactions

- Photon → Electromagnetic force
- $W^\pm$  and  $Z$  → Weak nuclear force
- Gluon → Strong nuclear force

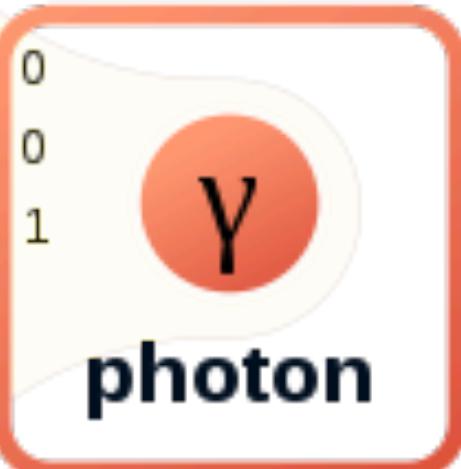
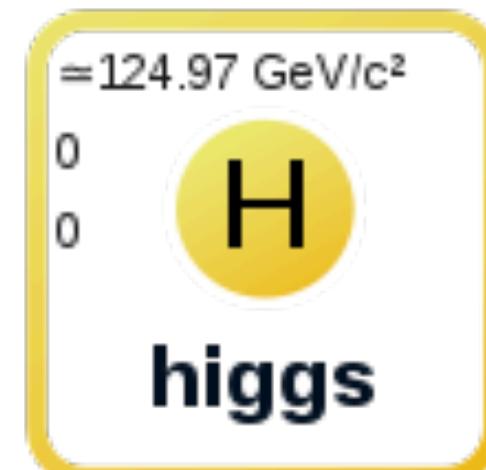


three generations of matter (fermions)				interactions / force carriers (bosons)
I	II	III		
mass charge spin				
=2.2 MeV/c <sup>2</sup> $\frac{2}{3}$ $\frac{1}{2}$ u up	=1.28 GeV/c <sup>2</sup> $\frac{2}{3}$ $\frac{1}{2}$ c charm	=173.1 GeV/c <sup>2</sup> $\frac{2}{3}$ $\frac{1}{2}$ t top	0 0 1 g gluon	
QUARKS				
=4.7 MeV/c <sup>2</sup> $-\frac{1}{3}$ $\frac{1}{2}$ d down	=96 MeV/c <sup>2</sup> $-\frac{1}{3}$ $\frac{1}{2}$ s strange	=4.18 GeV/c <sup>2</sup> $-\frac{1}{3}$ $\frac{1}{2}$ b bottom	0 0 1 $\gamma$ photon	
LEPTONS				
=0.511 MeV/c <sup>2</sup> -1 $\frac{1}{2}$ e electron	=105.66 MeV/c <sup>2</sup> -1 $\frac{1}{2}$ $\mu$ muon	=1.7768 GeV/c <sup>2</sup> -1 $\frac{1}{2}$ $\tau$ tau	0 1 Z Z boson	
$<1.0 \text{ eV}/c^2$ 0 $\frac{1}{2}$ $\nu_e$ electron neutrino	$<0.17 \text{ MeV}/c^2$ 0 $\frac{1}{2}$ $\nu_\mu$ muon neutrino	$<18.2 \text{ MeV}/c^2$ 0 $\frac{1}{2}$ $\nu_\tau$ tau neutrino	$<80.39 \text{ GeV}/c^2$ $\pm 1$ 1 W W boson	GAUGE BOSONS VECTOR BOSONS

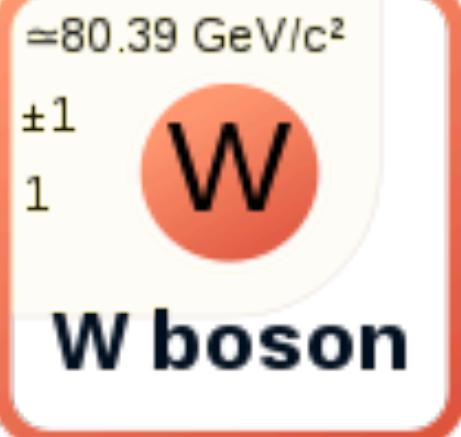


# The missing piece: the Higgs Boson

interactions / force carriers  
(bosons)



**GAUGE BOSONS**  
**VECTOR BOSONS**



**SCALAR BOSONS**

## Electroweak Theory

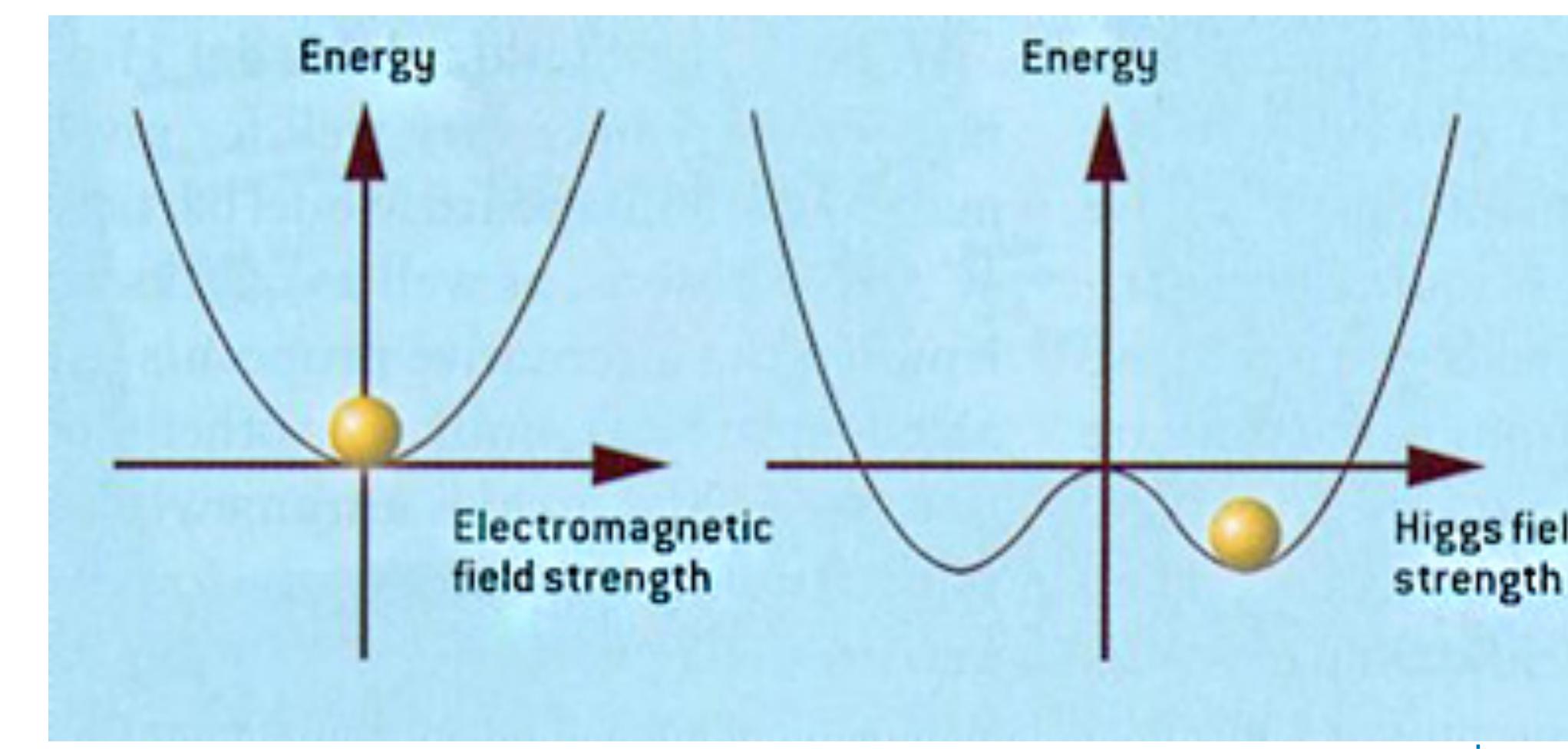
Weak and electromagnetic force are manifestations of the same “electroweak” force, which split as universe cooled (under  $\sim 160$  GeV)

- W and Z gained mass while photon remained massless... but how?

## Brout-Englert-Higgs

Interactions with a field present everywhere in space leads to mass

- Non-zero vacuum expectation value!



[enigma](#)



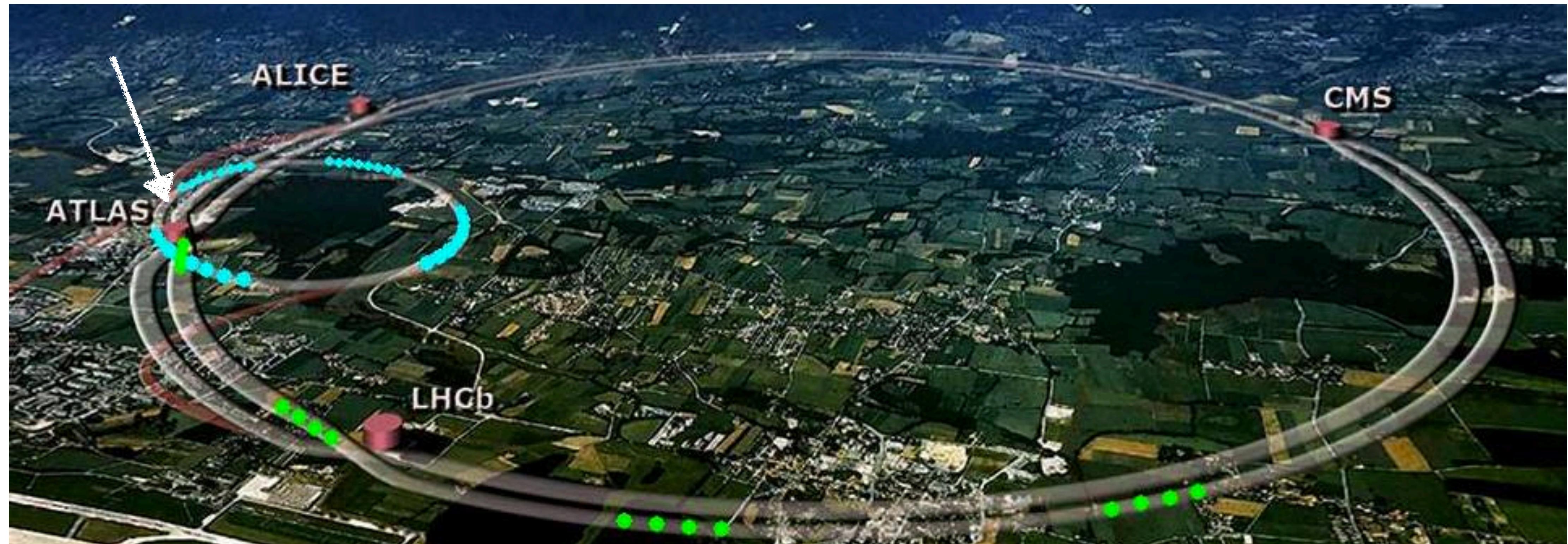
# Small particles require large solutions...

## The Large Hadron Collider

27 kilometers in circumference, situated on the Franco-Swiss border

Primarily  $p\bar{p}$  collisions, accelerated to 99.999990% the speed of light, collide at 4 primary points where detectors are situated

Center-of-mass energy of 13 TeV (6.5 TeV per beam)



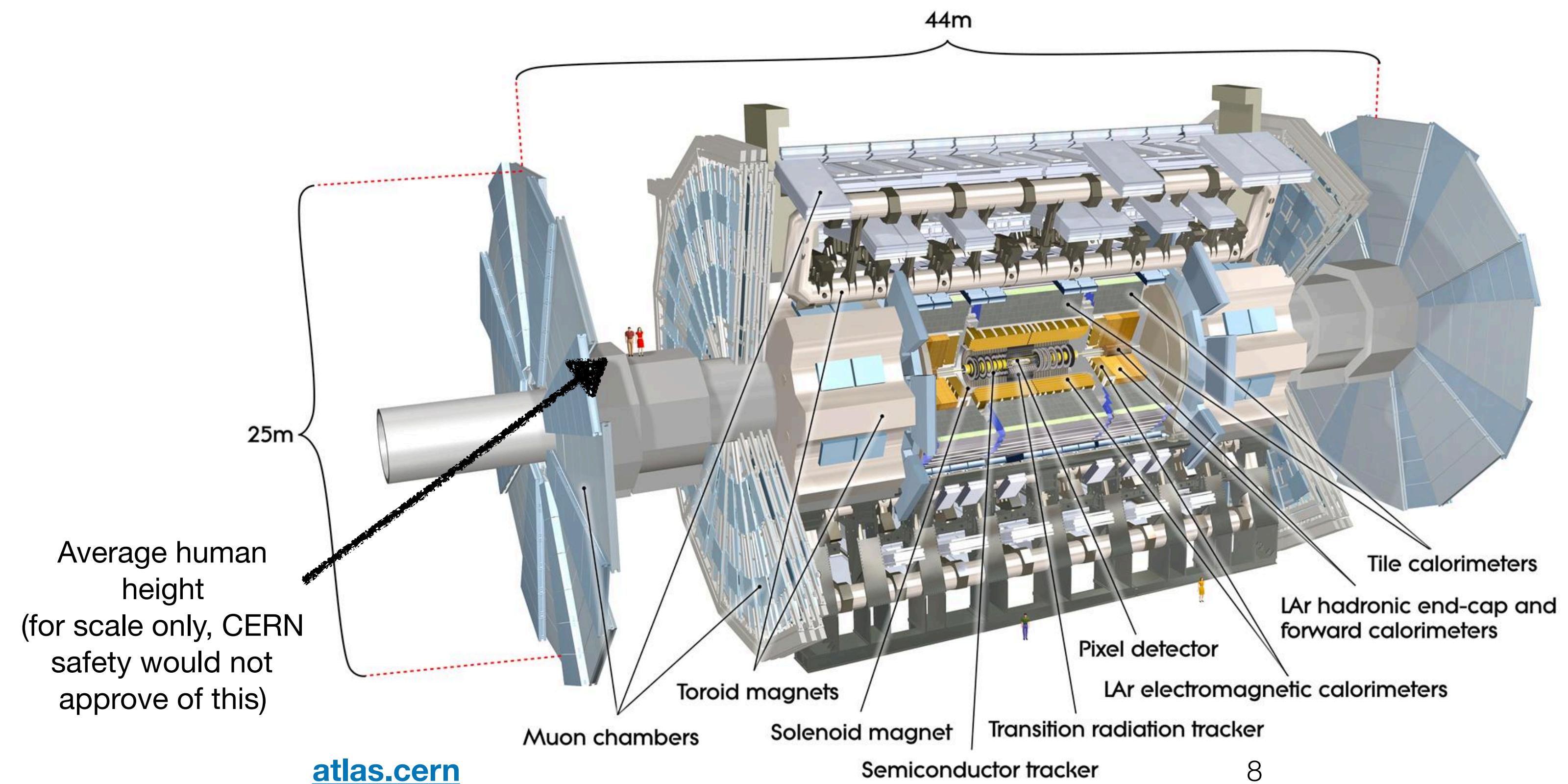
# Small particles require large solutions...

## The ATLAS Detector

General purpose detector, aiming to measure signals resulting from  $pp$  collisions to cover a vast range of analyses  
- Beyond just the Higgs program, searches for supersymmetry, dark matter, precision measurements, etc.

Collaboration of over 3,000 scientists from all over the world

Composed of several subsystems, each aiming to collect specific information about collisions



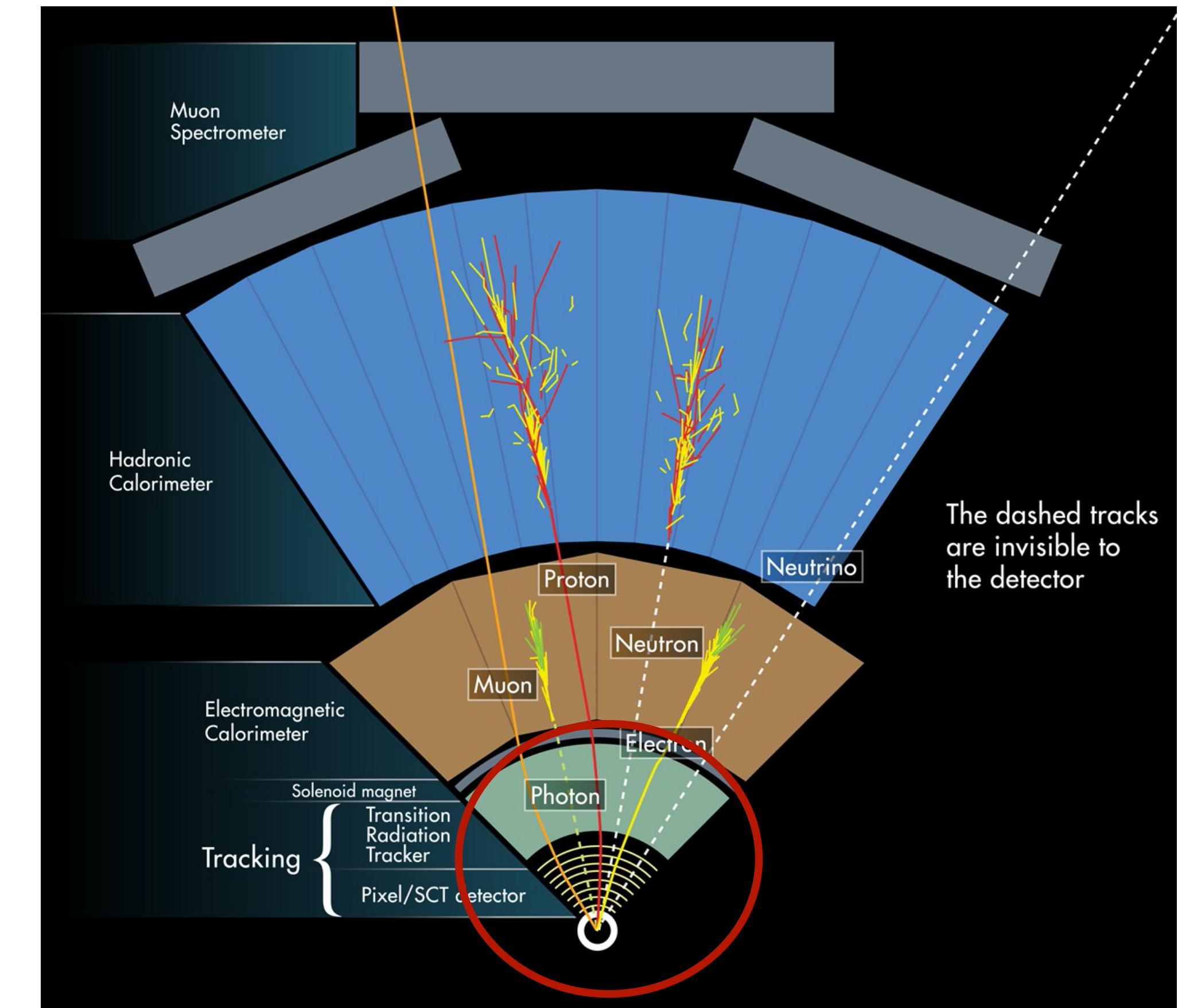
# Particles in...

## ATLAS is a many-layered detector

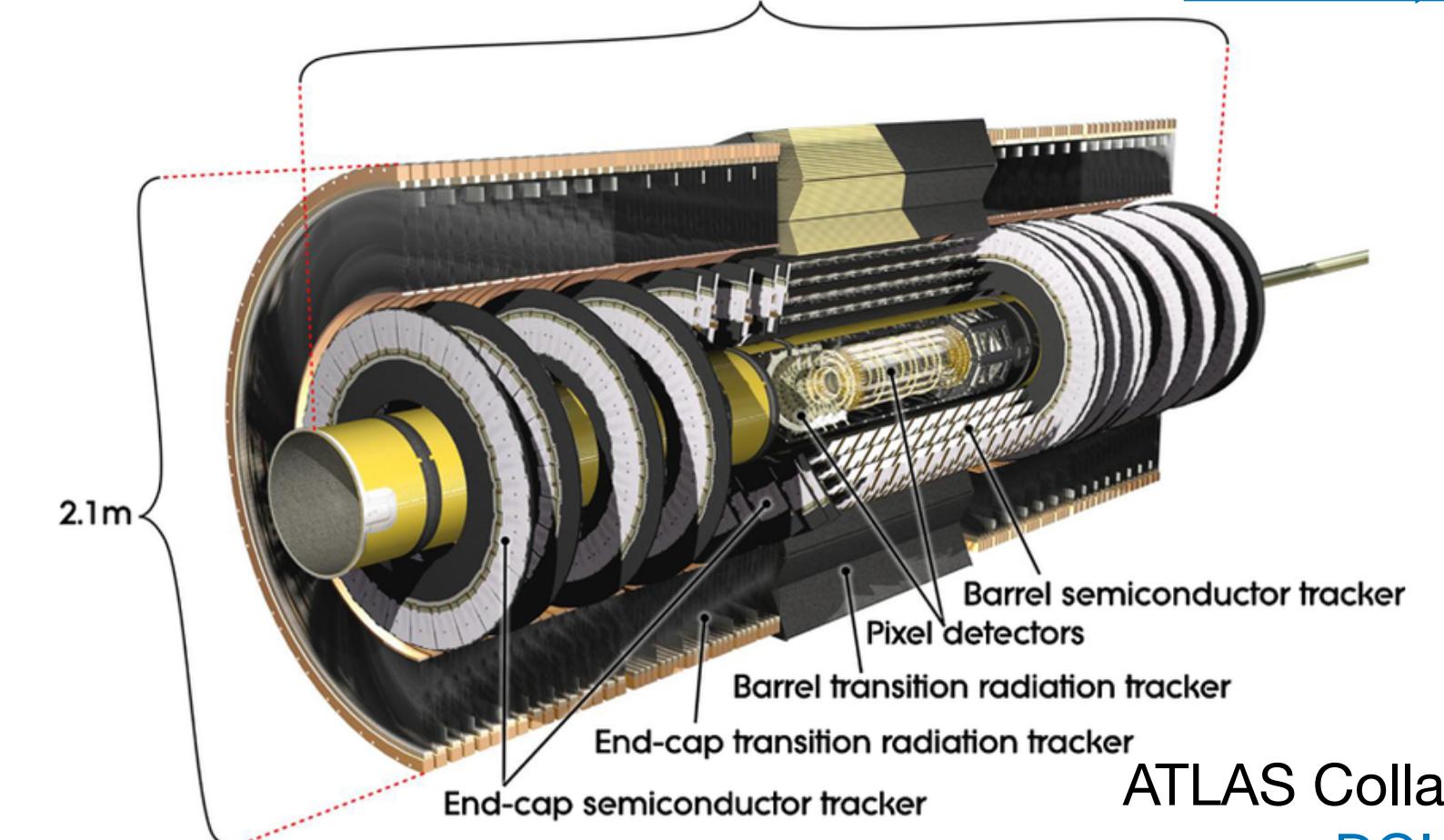
Each layer targets specific information from collision decay products

### Inner detector - Tracking information

- Describes charged particle trajectory through the detector and magnetic field



[Castillo, 2015](#)



ATLAS Collaboration

[DOI 10.1088](#)

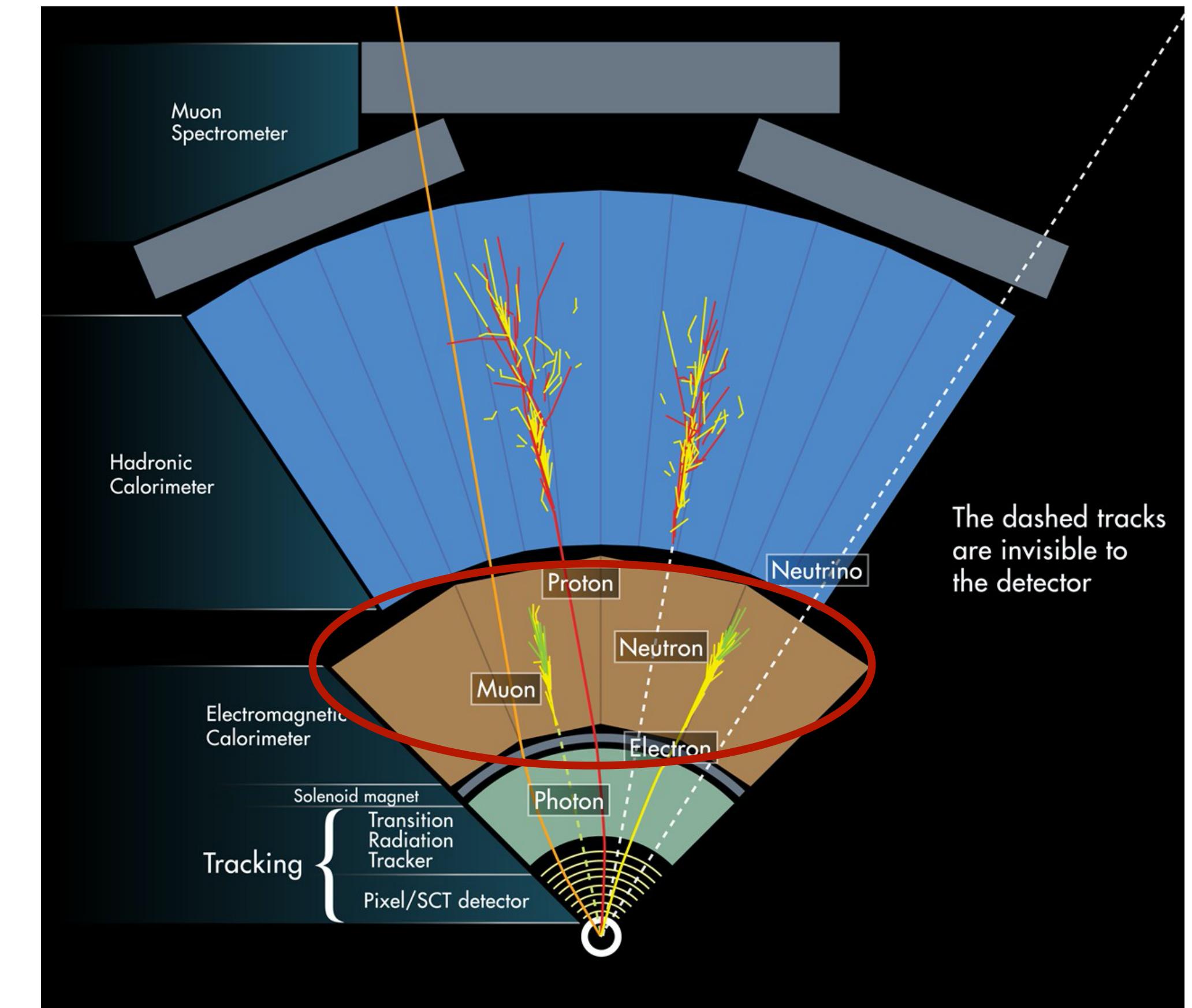
# Particles in...

**ATLAS is a many-layered detector**

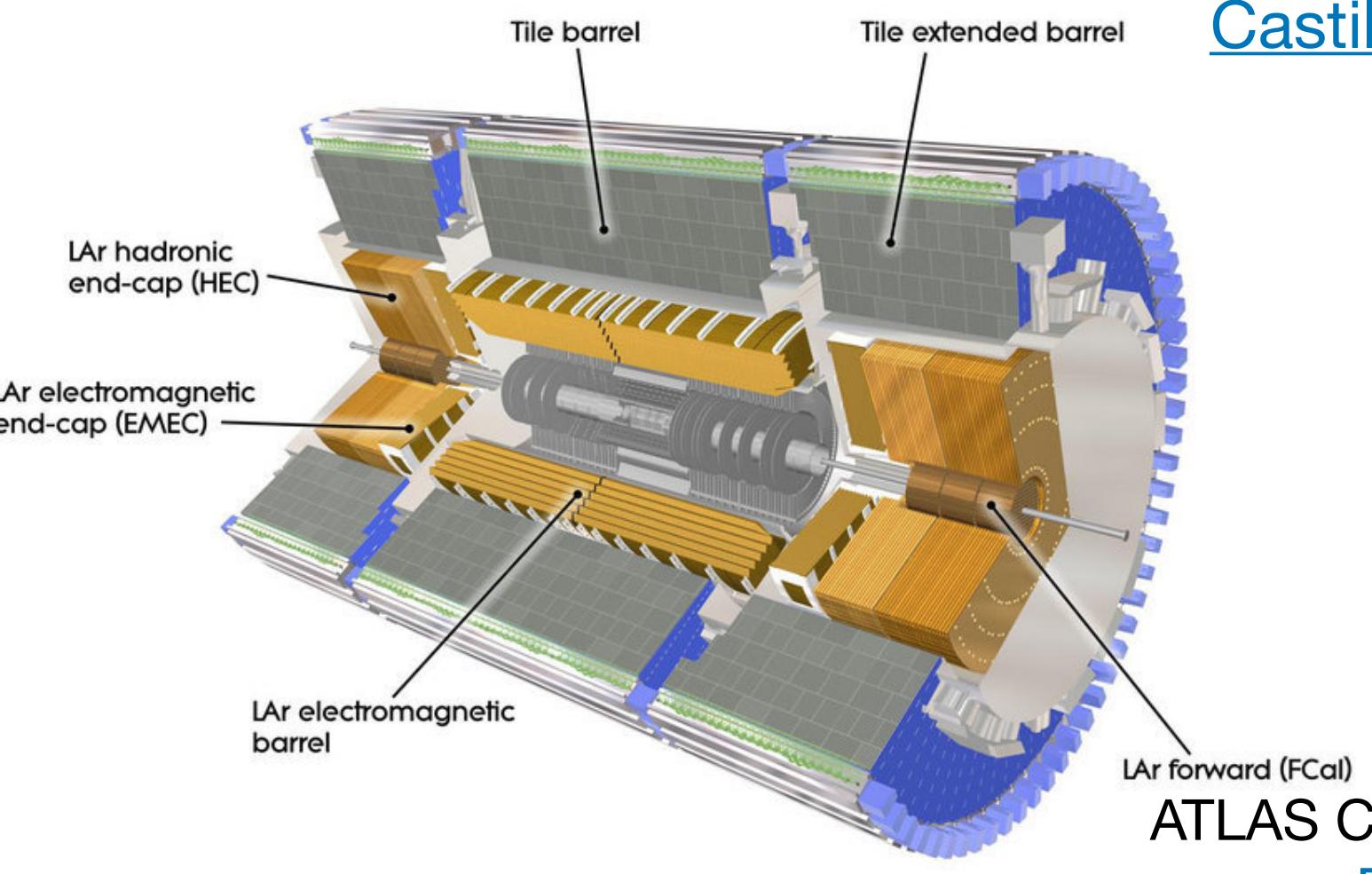
Each layer targets specific information from collision decay products

Inner detector - Tracking information

Electromagnetic Calorimeter - Electromagnetic (EM) signatures ( $\gamma, e$ )



[Castillo, 2015](#)



ATLAS Collaboration

[DOI 10.1088](#)

# Particles in...

## ATLAS is a many-layered detector

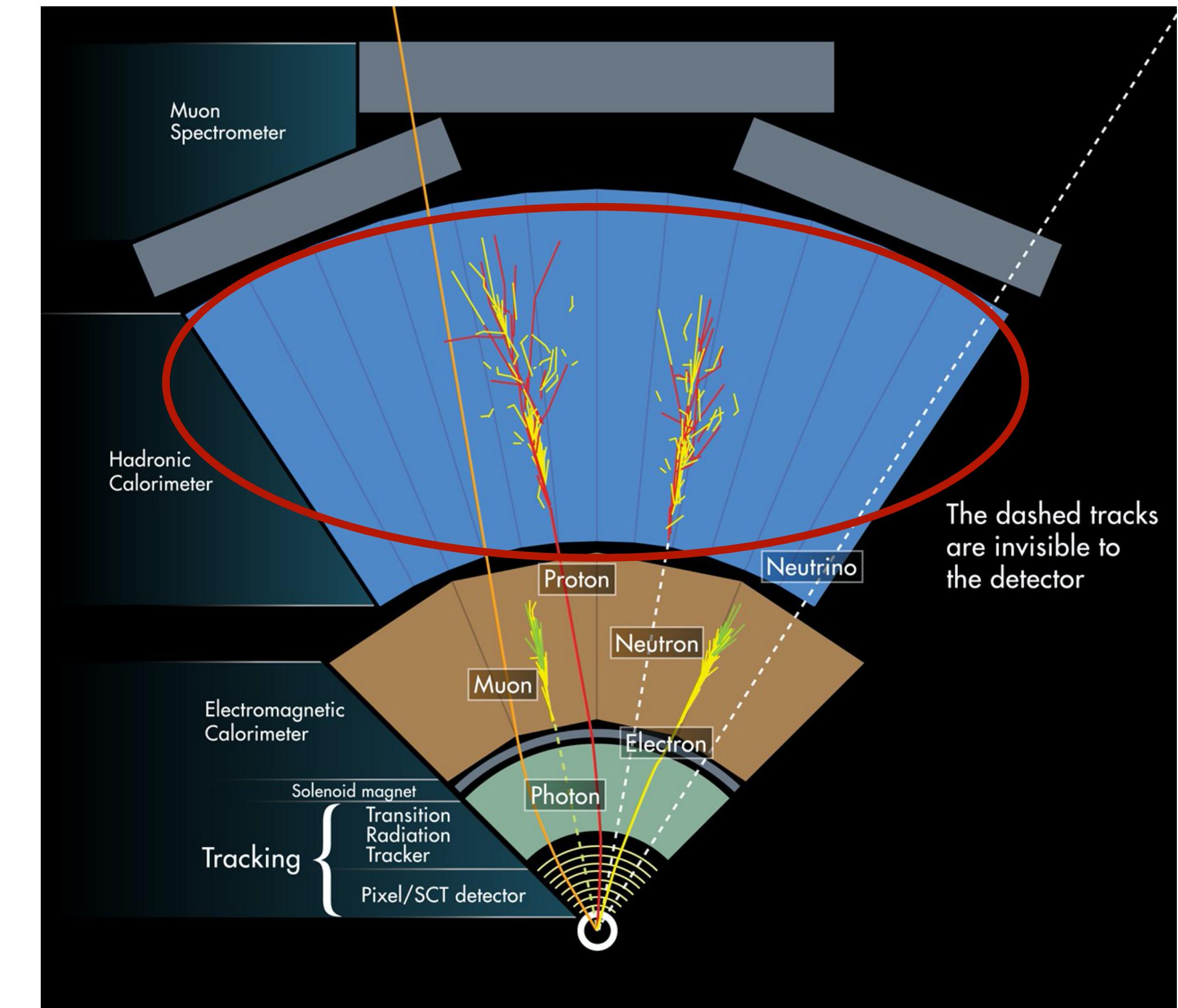
Each layer targets specific information from collision decay products

Inner detector - Tracking information

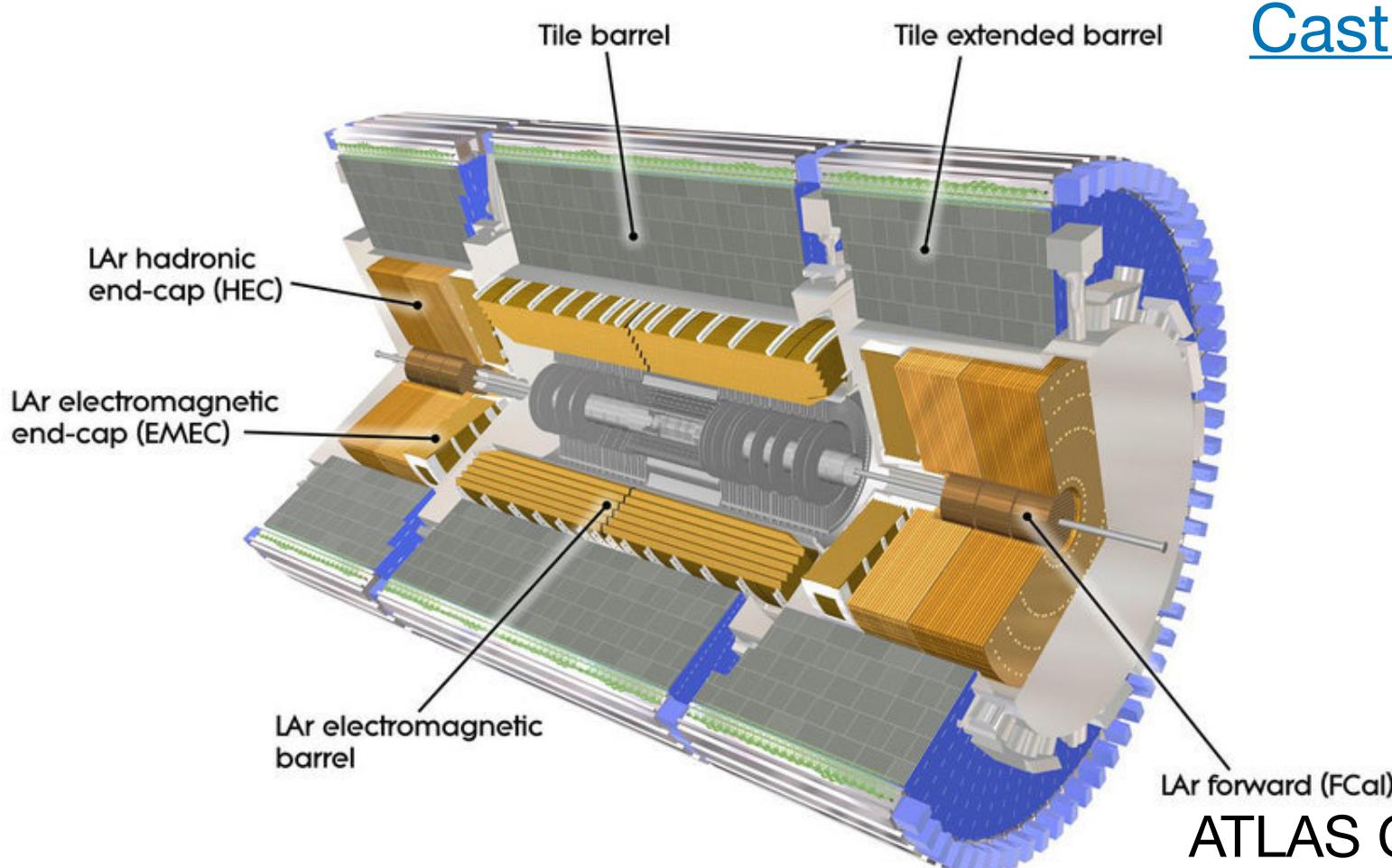
Electromagnetic Calorimeter - Electromagnetic (EM) signatures ( $\gamma, e$ )

Hadronic Calorimeter - Particles that interact via the strong force (quarks, gluons)

- Initiating strongly interacting particles may radiate additional quarks and gluons
- Must all be in colorless bound states due to *confinement*, so eventually turn into hadrons ( $\pi, K, p, n$ , etc.)



[Castillo, 2015](#)



ATLAS Collaboration

[DOI 10.1088](#)

# Particles in...

# ATLAS is a many-layered detector

Each layer targets specific information from collision decay products

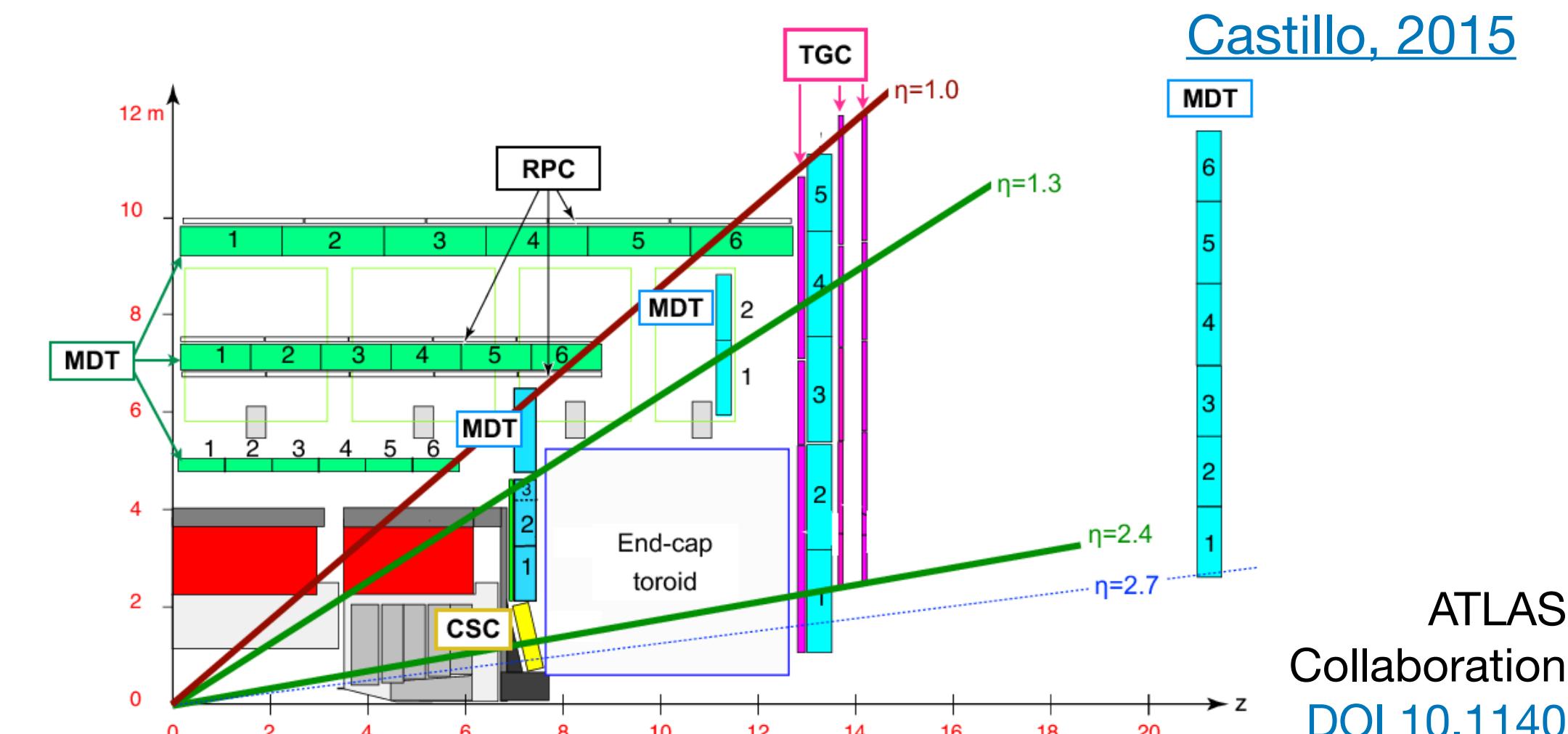
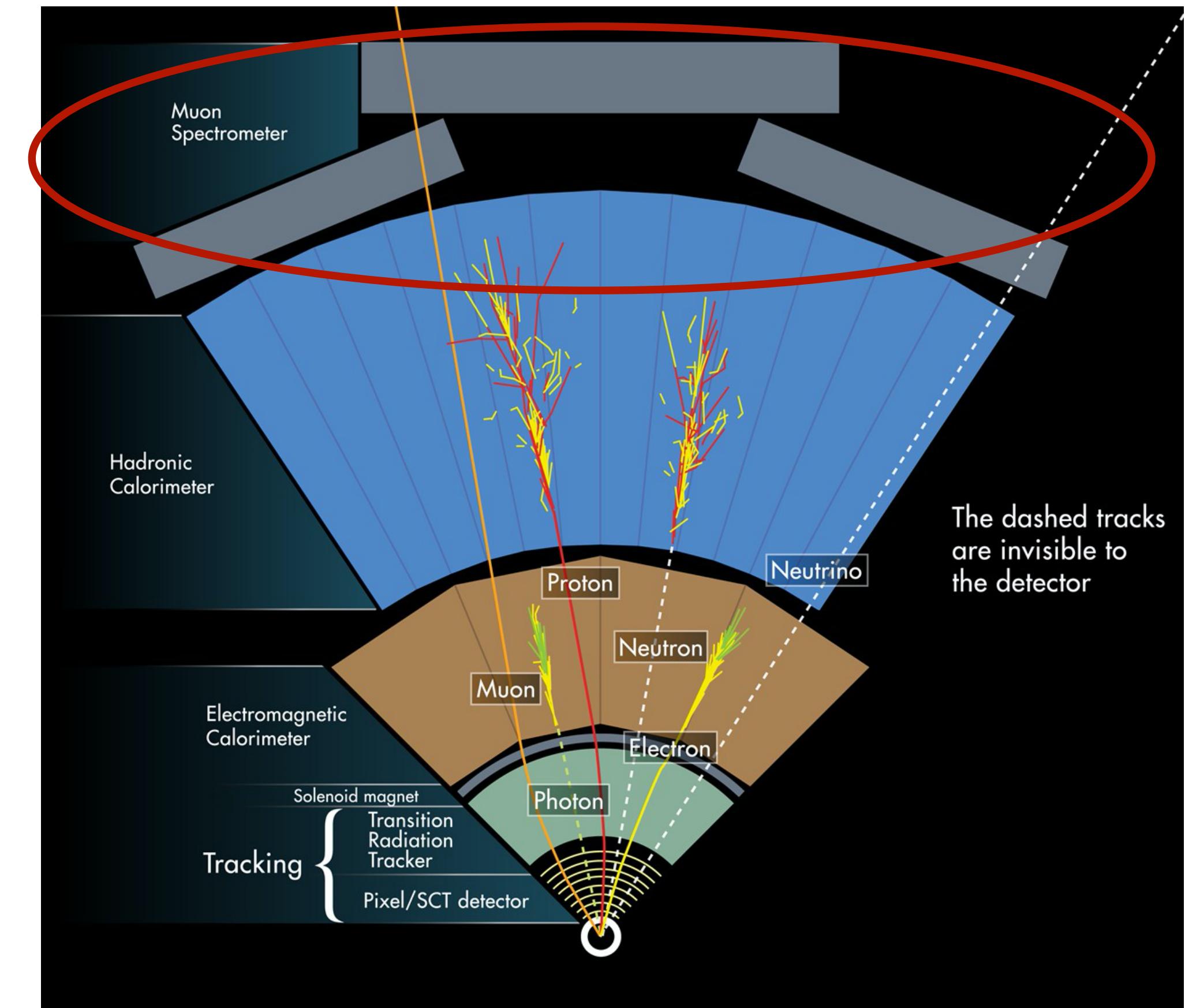
# Inner detector - Tracking information

# Electromagnetic Calorimeter - Electromagnetic (EM) signatures ( $\gamma, e$ )

Hadronic Calorimeter - Particles that interact via the strong force (quarks, gluons)

# Muon Spectrometer - Muons

- Don't interact in upstream material, so dedicated subsystem for detecting



# Particles in...

## ATLAS is a many-layered detector

Each layer targets specific information from collision decay products

Inner detector - Tracking information

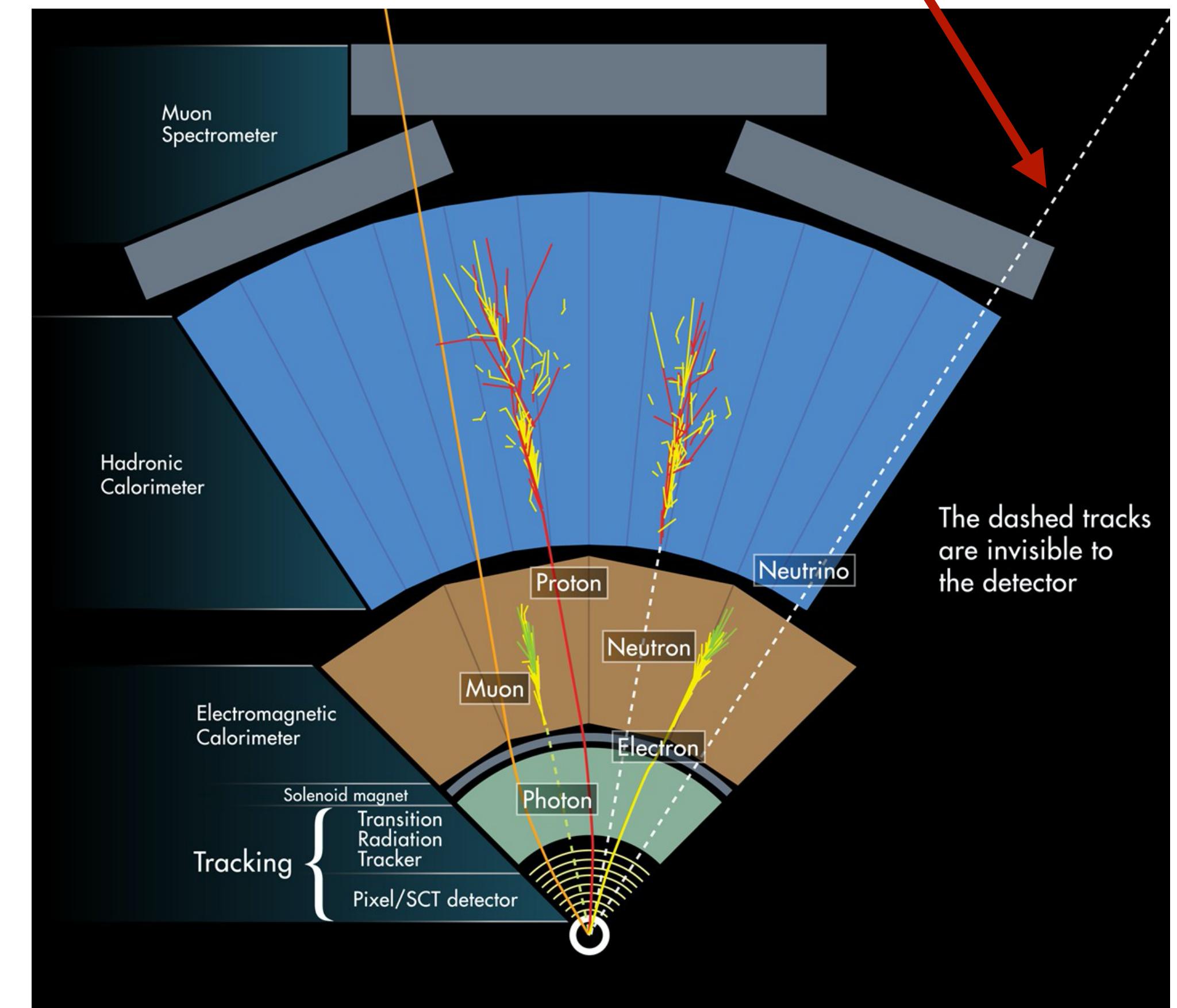
Electromagnetic Calorimeter - Electromagnetic (EM) signatures ( $\gamma, e$ )

Hadronic Calorimeter - Particles that interact via the strong force (quarks, gluons)

Muon Spectrometer - Muons

Missing Transverse Energy (MET) - Neutrinos

- Measurement for particles that don't interact with any portion of the detector, using conservation of momentum



[Castillo, 2015](#)

# ...objects defined...

Use signatures in various subsystems to define various physics objects

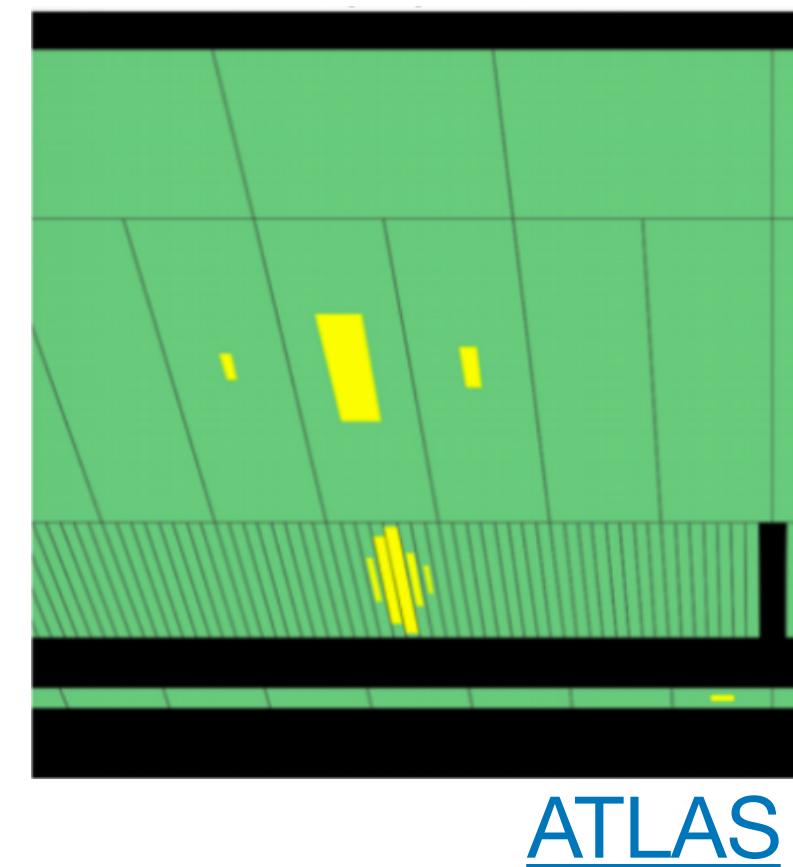
Most relevant to this analysis are photons and jets

## Photons

Signatures in the electromagnetic calorimeter used

- Measures electromagnetic showers based on 2 processes
  - Bremsstrahlung photon emission
  - Electron-positron conversion

Topological clusters of dynamic size constructed moving out from high signal to noise cells, then grouped into superclusters, ultimately defined as photons



ATLAS

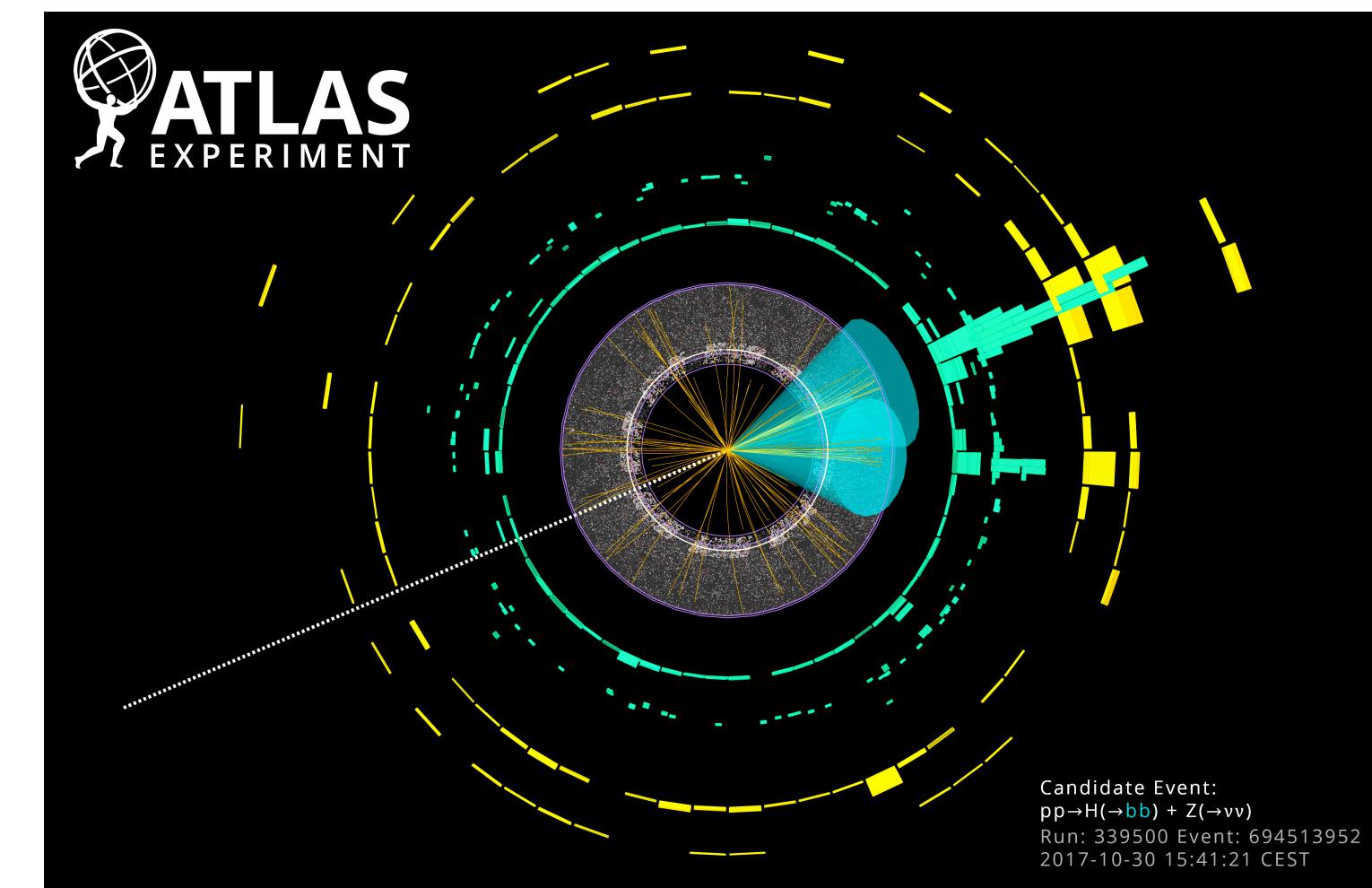
## Jets

Final signature of strongly interacting particles is a collimated spray of hadrons called “jets.” Information in both calorimeters used

- As hadrons interact with matter they scatter and produce additional hadrons, cascade known as hadron shower
- Cone-based algorithm ( $\text{Anti-}k_t$ ) used

Algorithms used to discern the source of the jet, specifically, b-jets

- For this work a BDT-based algorithm
- Several calibrated working points defined based on signal purity



ATLAS

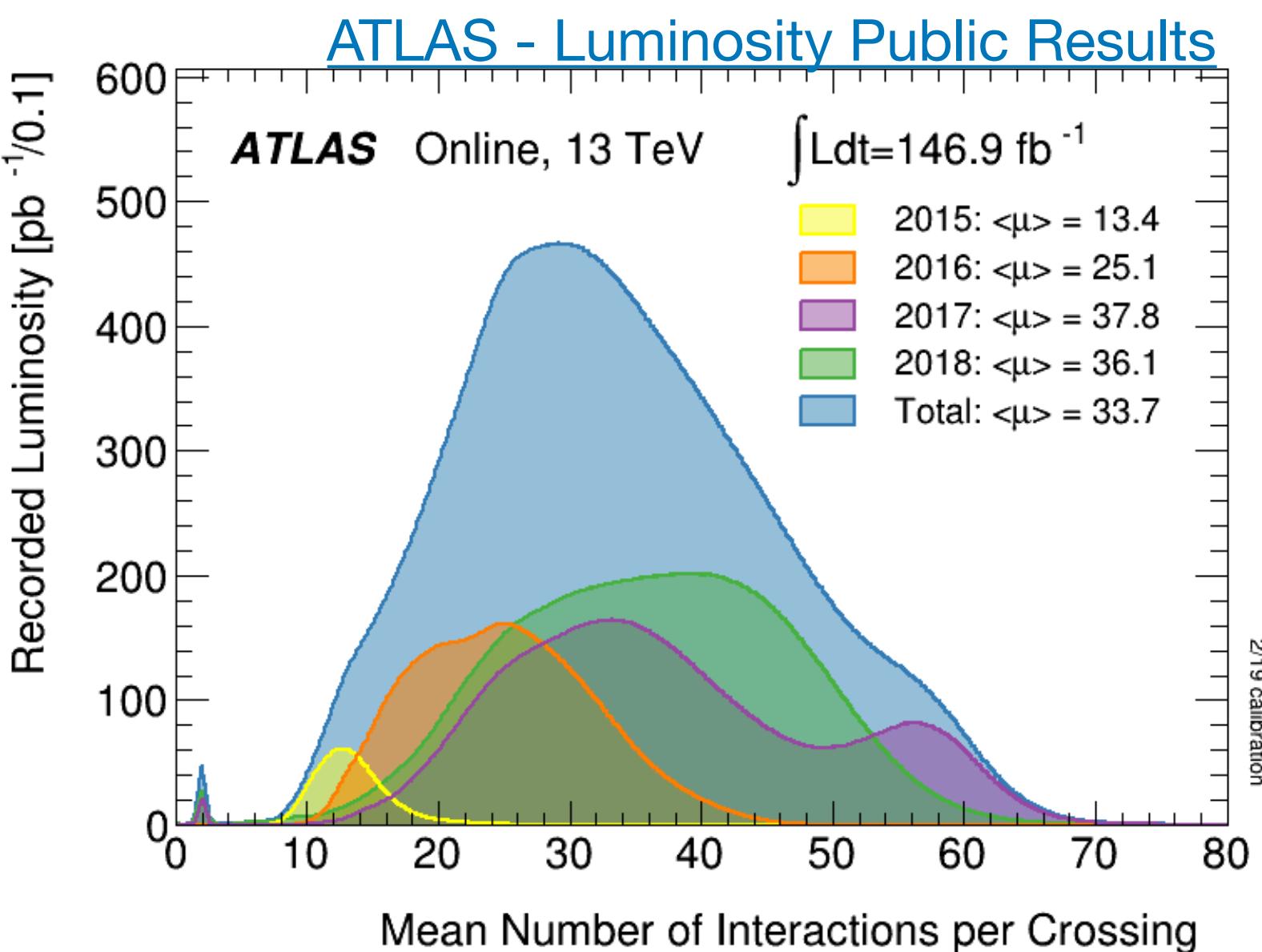
# ...data out

Bunches of protons collide at a rate of 40 MHz (every 25 ns)

- On average, 33.7 collisions per crossing in Run 2
  - 23.7 on average through 2015-2016 (presented analysis)

Each raw event ~1.6 MB of information, so around 64 TB/s data created

**No possible way to save all of this data!**



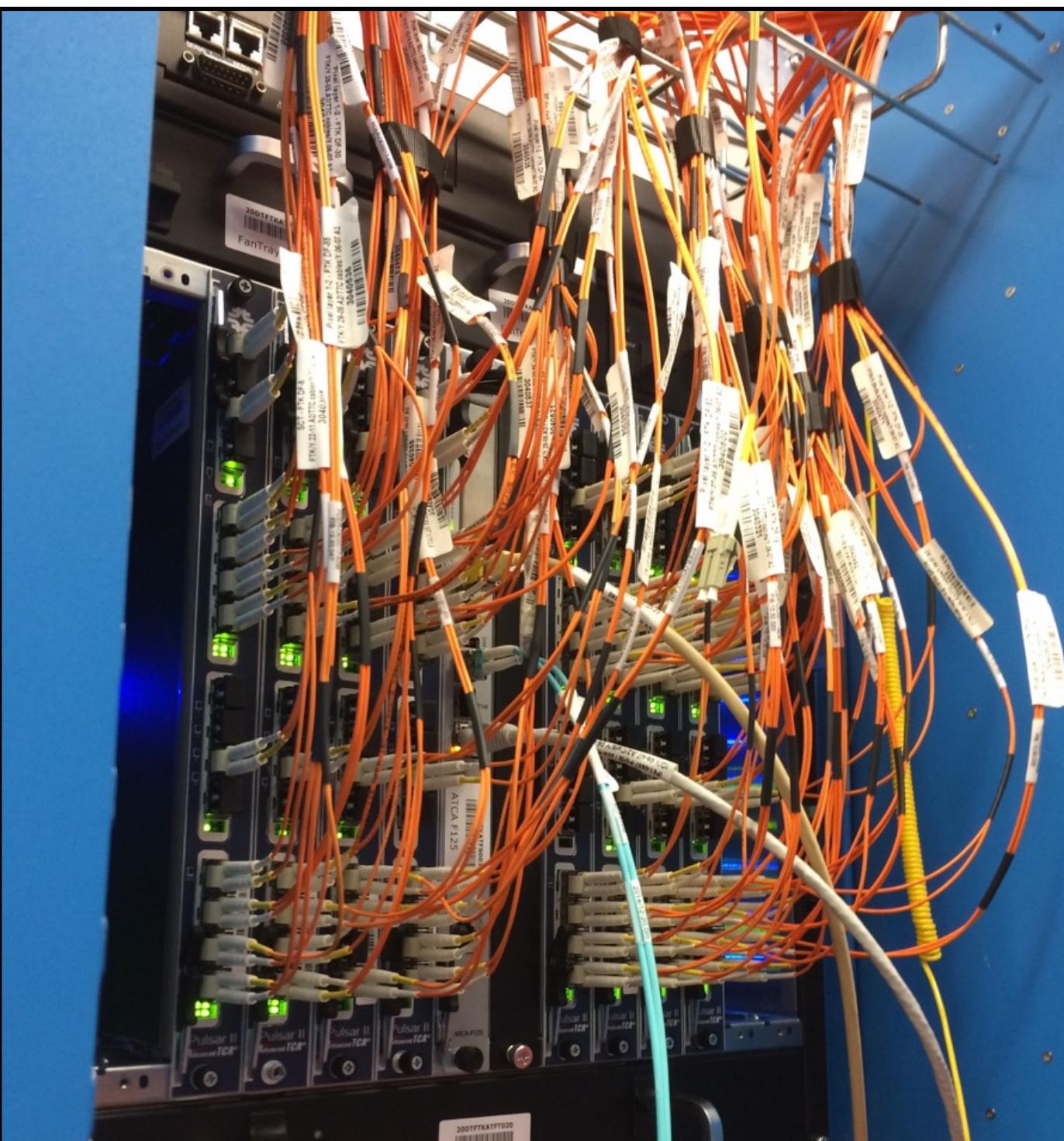
ATLAS employs a 2 stage trigger system

- Level 1: hardware based, performing coarse object reconstruction.  
Brings rate to under 100 kHz
- High-Level Trigger (HLT): Software based, performing object reconstruction as close to offline as possible. Brings rate to 1 kHz

**Still, incredible amounts of data**

200 petabytes milestone passed in June 2017

**and with this data....**



# Discovery of the Higgs Boson

**July 4, 2012** - Standard Model-like Higgs boson discovered by ATLAS and CMS!

- Mass reconstructed at  $\sim 125$  GeV
- Observed in multiple channels with rates consistent with the Standard Model (SM) prediction
- Only fundamental scalar boson, which motivates further study beyond just the discovery

[NYTimes](#)



**Plenty of outstanding questions! Among these...**

- Does the discovered Higgs boson couple as predicted by the SM?
- Are there any additional Higgs bosons?
- Do particles beyond those predicted by the SM exist? If so, how do they couple to the Higgs boson?

**Studying di-Higgs production gives insight to these questions**

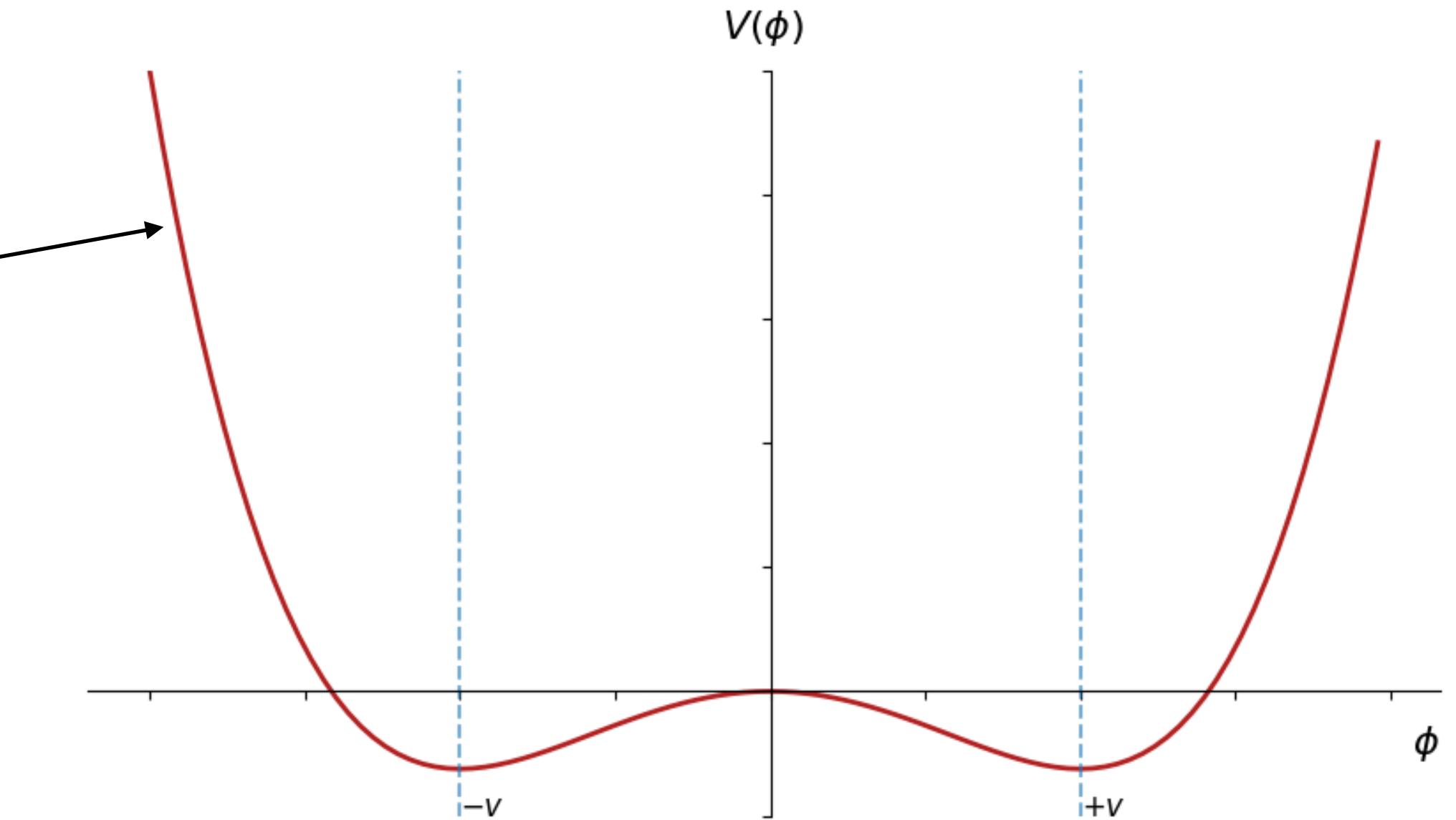
# Di-Higgs Production

Higgs potential function is given by

$$V(\phi) = \mu^2 |\phi|^2 + \lambda(|\phi|^2)^2$$

for scalar field  $\phi$  with mass  $\mu$ , and Higgs self coupling strength  $\lambda$ . For  $\mu^2 < 0$ , potential shown, with minima of found at

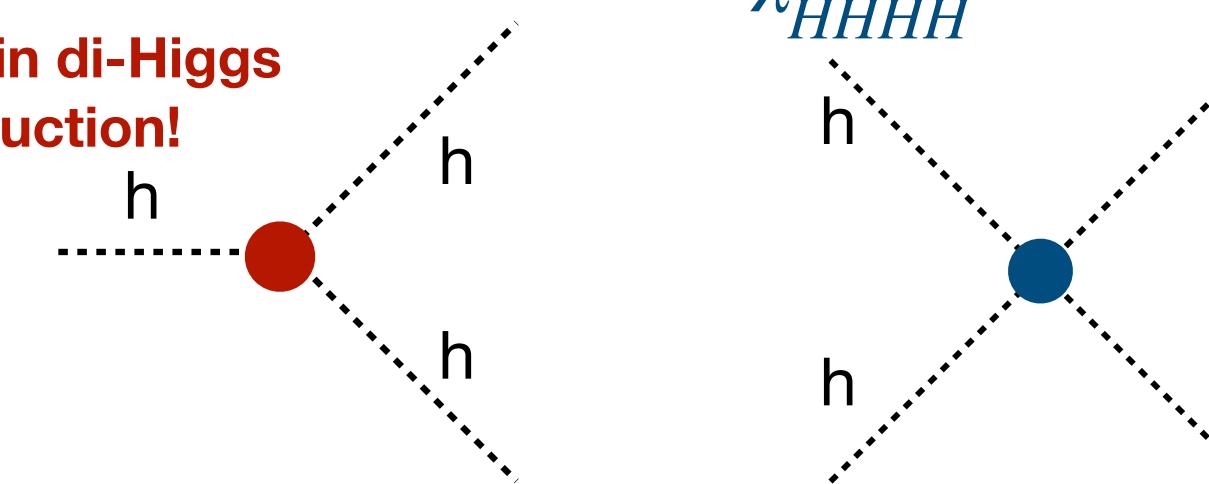
$$v = \sqrt{-\mu^2/\lambda} = 246 \text{ GeV (from experiment)},$$



which meets earlier criteria of **non-zero values in empty space**. Pick a minima and expand:

$$V(h) = \lambda v^2 h^2 + \underbrace{\lambda v h^3}_{\lambda_{HHH}^{SM}} + \frac{1}{4} \underbrace{\lambda h^4}_{\lambda_{HHHH}^{SM}}$$

Results in di-Higgs production!



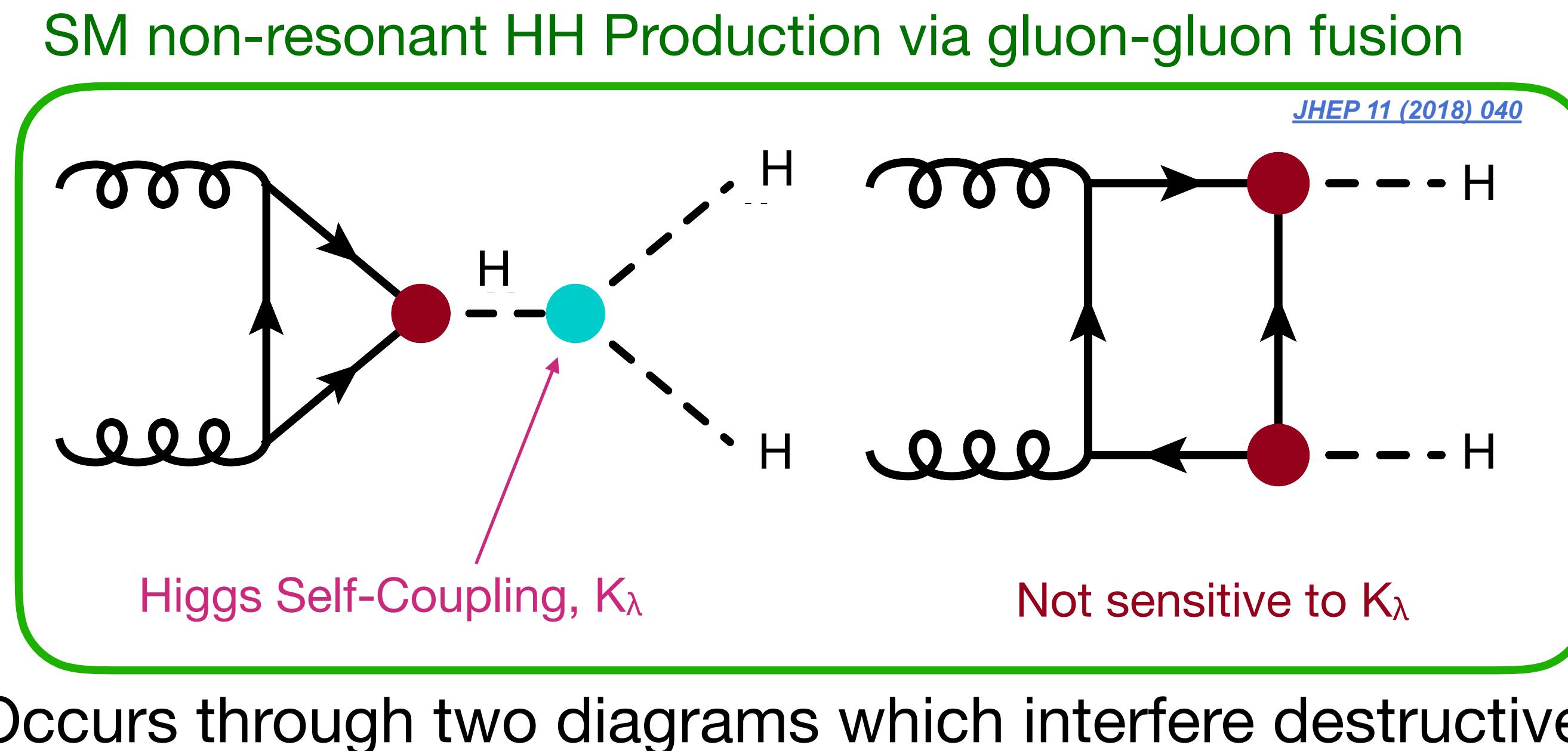
Quartic coupling - extraordinarily rare, out of reach of LHC

Picking this minima is known as **electroweak symmetry breaking**, and implies the Higgs **self-couples**. Measurement of the Higgs self-coupling provides a direct probe on the Higgs potential.



# Di-Higgs Production in the SM

Gluon-gluon fusion (ggF) is the dominant di-Higgs production mode, accounting for ~87% of di-Higgs events at 13 TeV, most published analyses target this production mode



Very small cross section!

$$\sigma(gg \rightarrow HH)_{SM} \approx 33 \text{ fb} @ 13 \text{ TeV}$$

di-Higgs production is a rare process, which will test the limits of the LHC



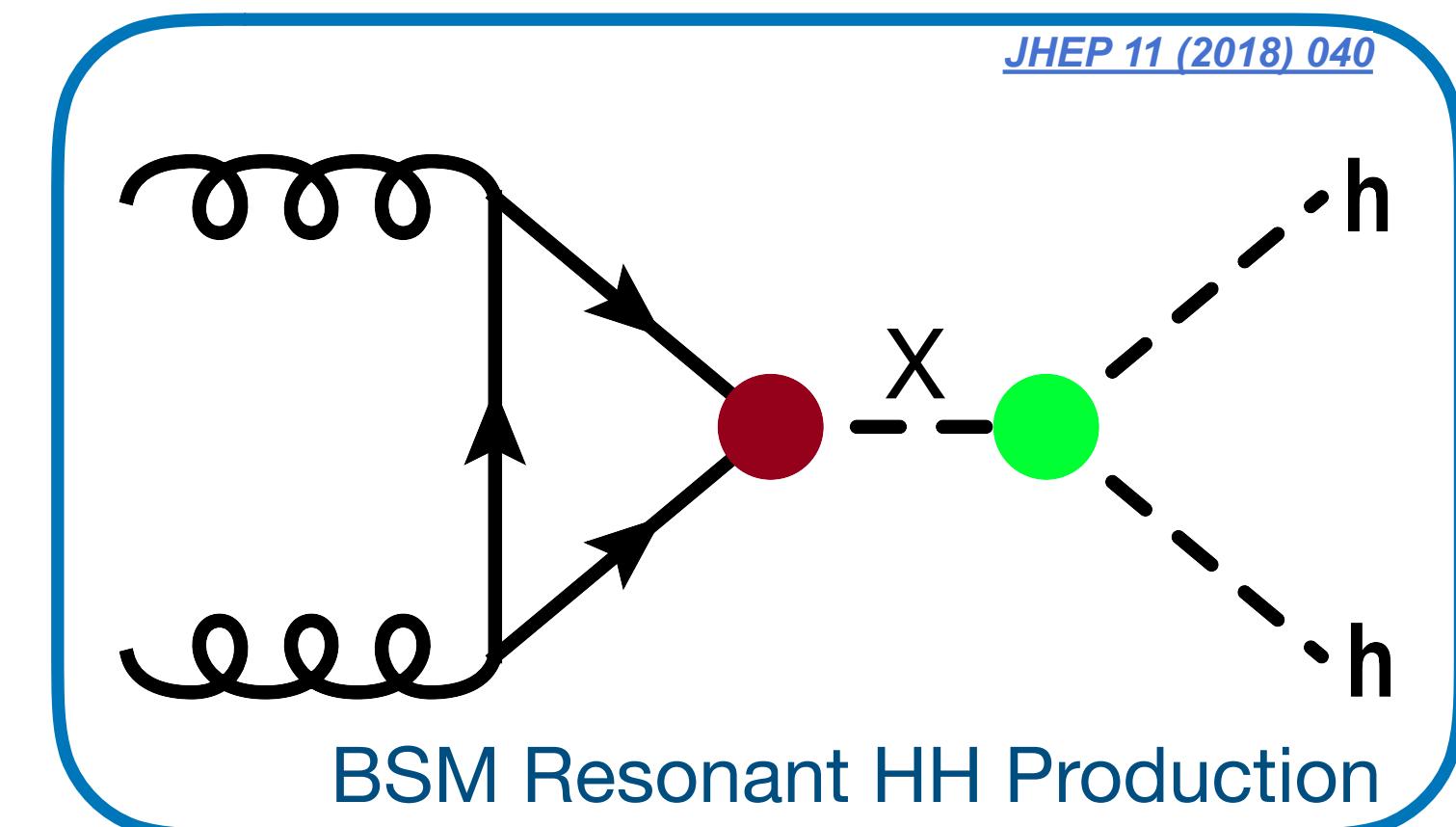
# BSM Di-Higgs Production

[JHEP 11 \(2018\) 040](#)

## Resonant enhancements:

Enhancements can occur through a theoretical particle beyond the Standard Model (BSM) that decays to two Higgs bosons

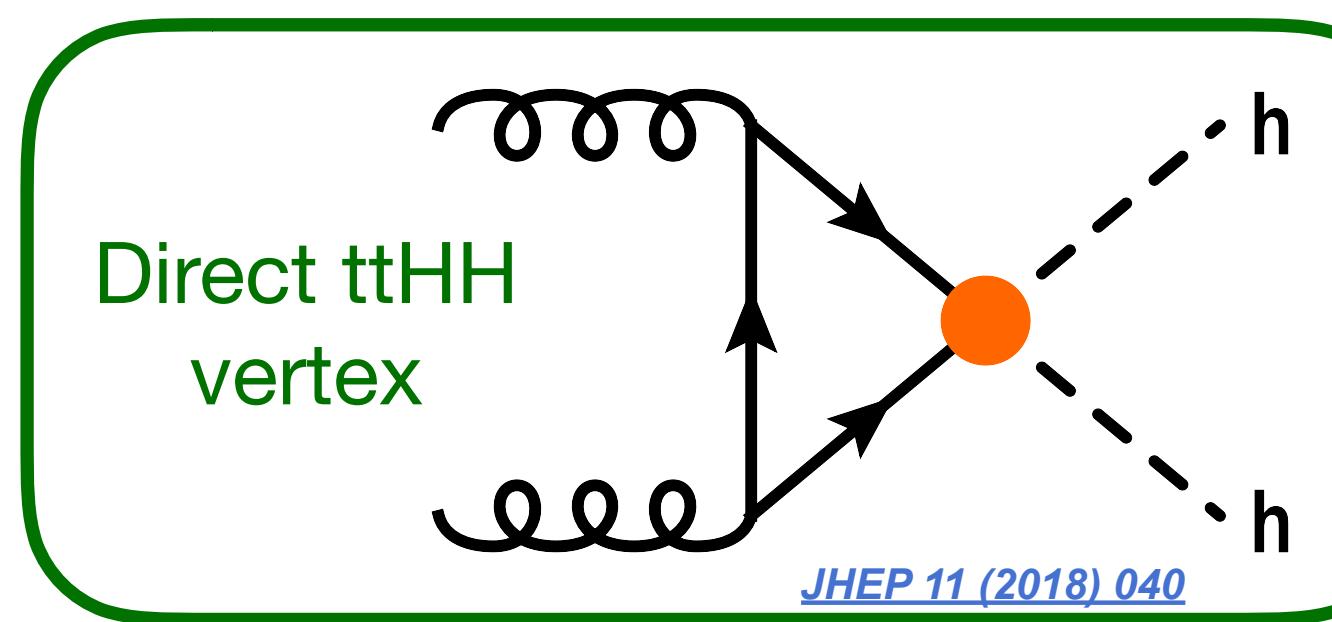
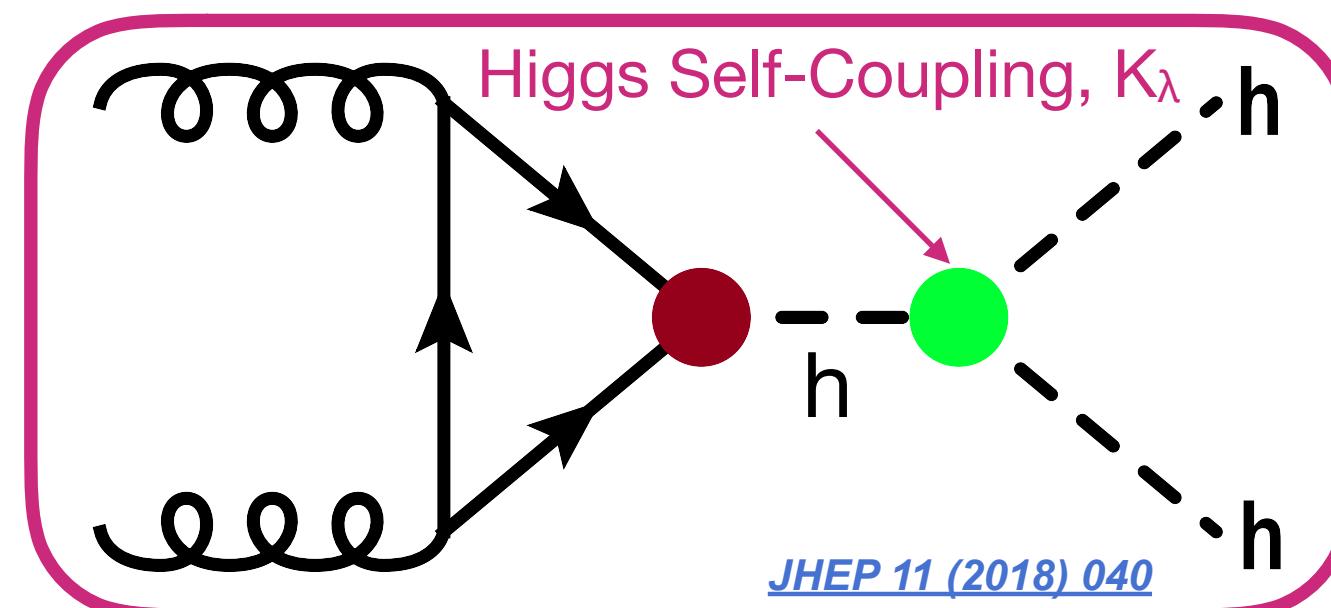
- Two-Higgs-doublet model (THDM) - Resonance is a heavy scalar



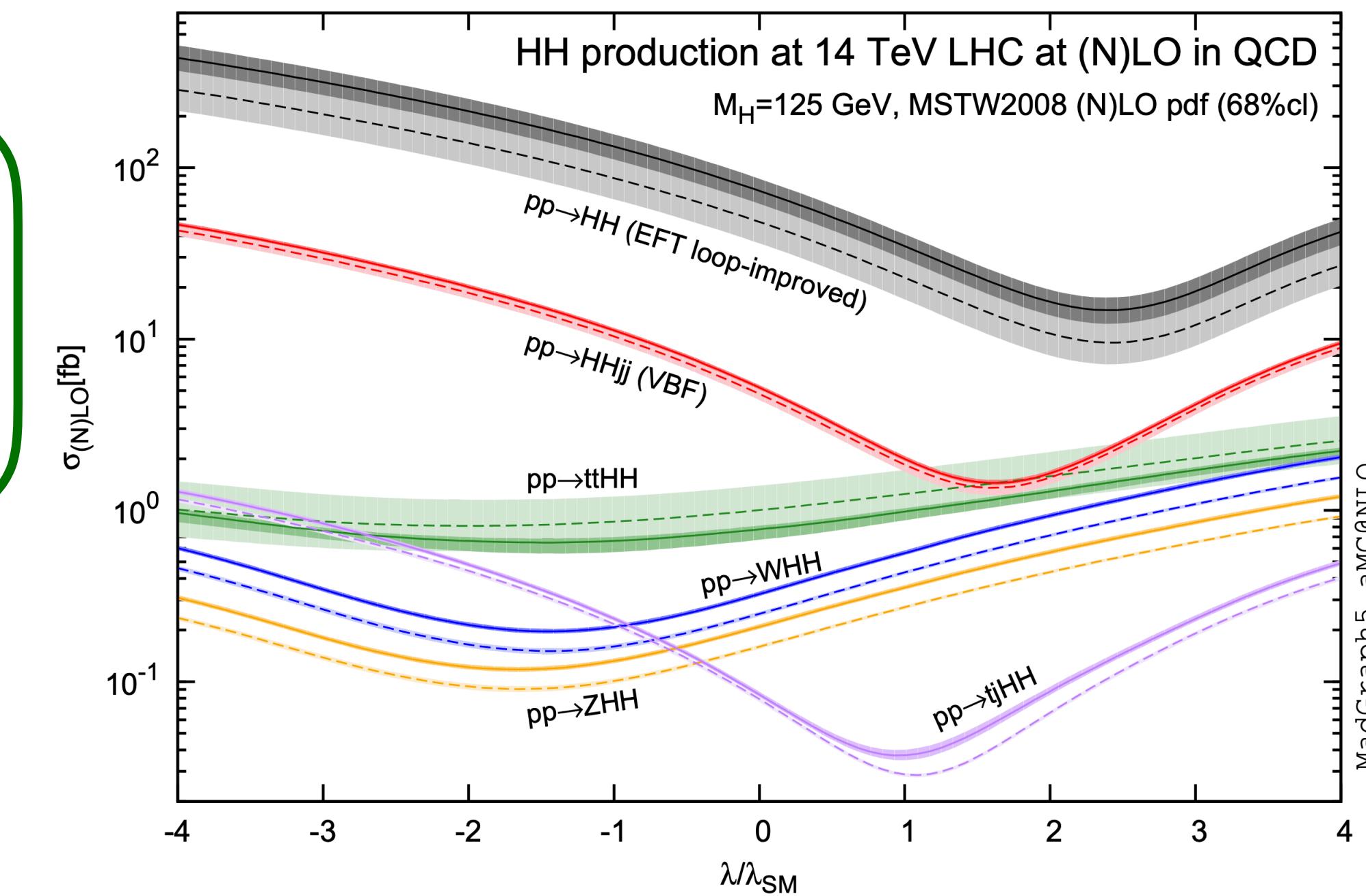
## Non-Resonant enhancements:

Can also look toward enhanced production caused by couplings that differ from the SM

- Higgs Self-Coupling ( $K_\lambda$ ) - strength of coupling deviating from SM value can lead to enhanced production
- Additional couplings not predicted by the SM, e.g. direct ttHH vertex



**Models with enhanced production make di-Higgs already interesting to study with Run-2 data**

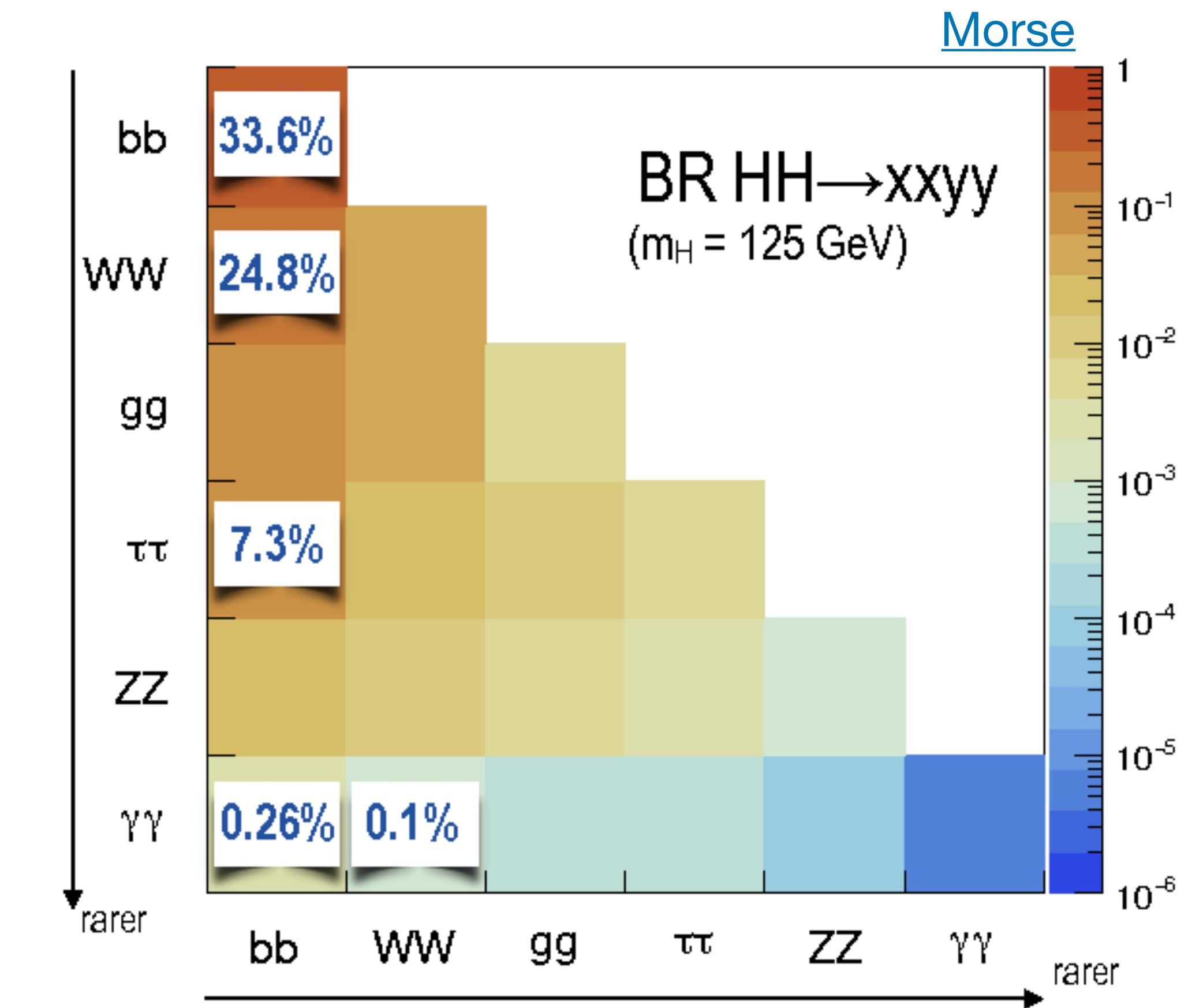


# Decay Modes

Strong motivation to look at certain channels based on properties of decay products

Strongest contribution to current limits:

- $b\bar{b}b\bar{b}$ : Fully takes advantage of high  $bb$  branching ratio, but suffers from large multijet background
- $\gamma\gamma b\bar{b}$ : Excellent trigger and mass resolution for photons, high  $bb$  branching ratio
- $b\bar{b}\tau\tau$ : Taus are relatively clean while still having a large branching ratio, high  $bb$  branching ratio



# $\text{HH} \rightarrow \gamma\gamma \text{bb}$ search using 2015+2016 data

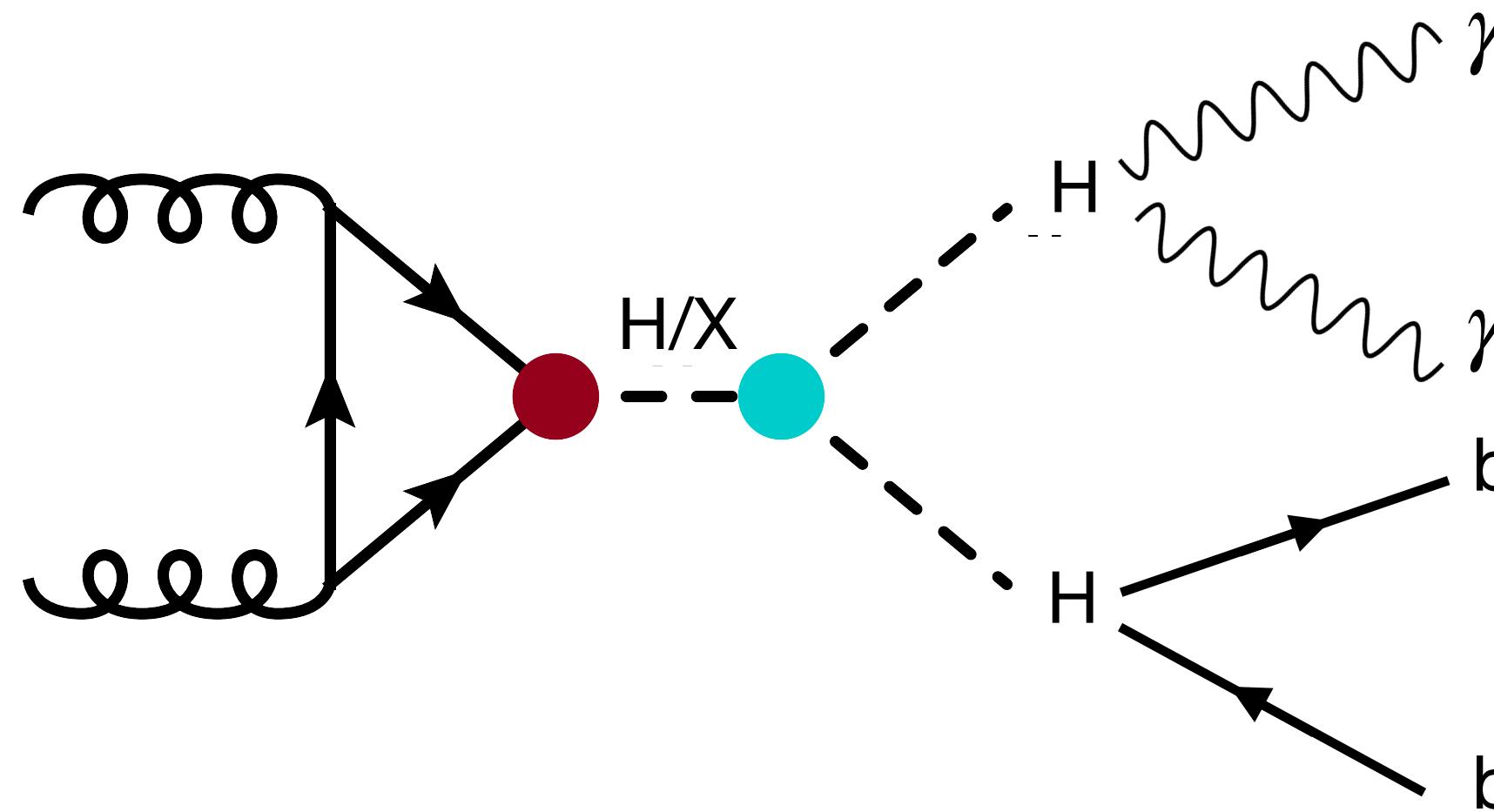


# Analysis Overview

## Searching for resonant and non-resonant di-Higgs production

Final state involves one Higgs decaying to photons, the other to a b quark pair

- Search for this final state in the detector - 2 photons and 2 b-jets



Full decay shown (similar for box diagram)

Analysis provides sensitivity to 3 values:

- The SM di-Higgs cross-section
- The Higgs trilinear coupling strength
- The existence of a generic scalar particle (resonance) which decays to Higgs boson pairs
- Generic scalar chosen to remain model independent



# Analysis Selection

## Preselection - target decay products

Matches other  $H \rightarrow \gamma\gamma$  analyses - 2 photons passing tight criteria, within the mass window of [105, 160] GeV

- Diphoton trigger used - one photon with  $p_{T,\gamma 1} > 35$  GeV,  $p_{T,\gamma 2} > 25$  GeV

2 jets within b-tagging region for  $H \rightarrow bb$

Split into 3 categories using b-tagging information

### 2-tag

2 jets passing medium b-tagging  
(70% signal efficiency)

Used as a signal region

### 1-tag

1 jet passing tight b-tagging  
(60% signal efficiency)

Used as a signal region, second jet chosen  
using a NN

### 0-tag

Passes preselection, but fails other  
categories

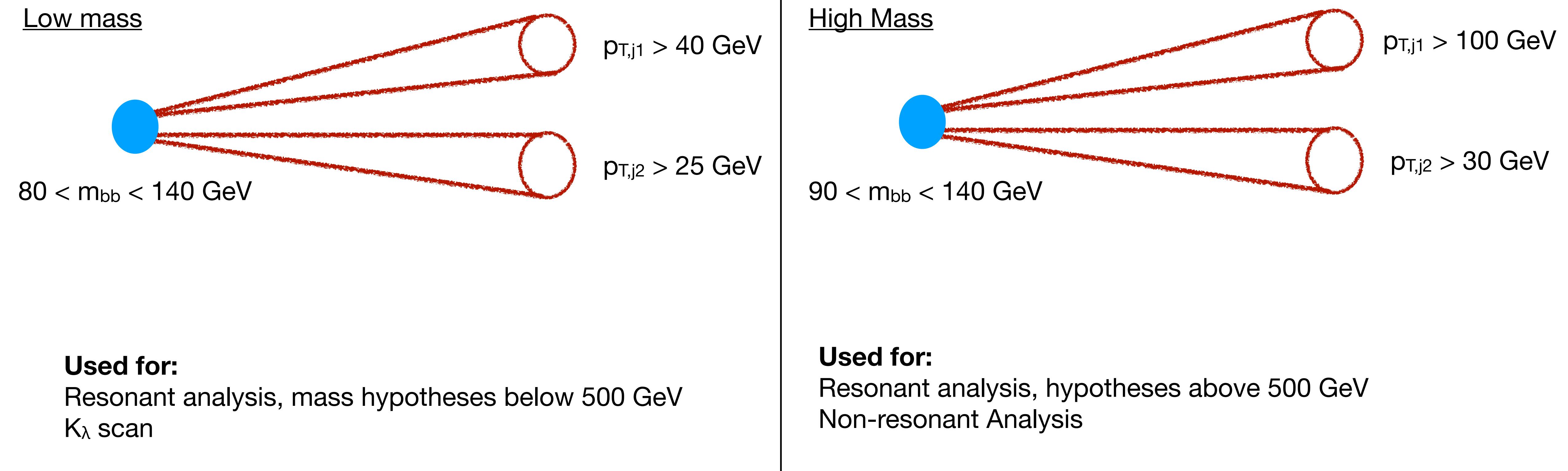
Used for shape information



# Analysis Selection

## 2 Overlapping Analysis Categories

Built to target different portions of the analysis, applying constraints to jets and bb system



**Used for:**  
Resonant analysis, mass hypotheses below 500 GeV  
 $K_\lambda$  scan

**Used for:**  
Resonant analysis, hypotheses above 500 GeV  
Non-resonant Analysis

## Approach - Perform unbinned likelihood fits to signal and background

Non-resonant analysis - fit  $m_{\gamma\gamma}$  spectrum

Resonant analysis - cut in  $m_{\gamma\gamma}$  and fit  $m_{\gamma\gamma bb}$  spectrum



# Background Estimation

Largest background in analysis is  $\gamma\gamma$ -continuum processes, subdominant is single Higgs boson production

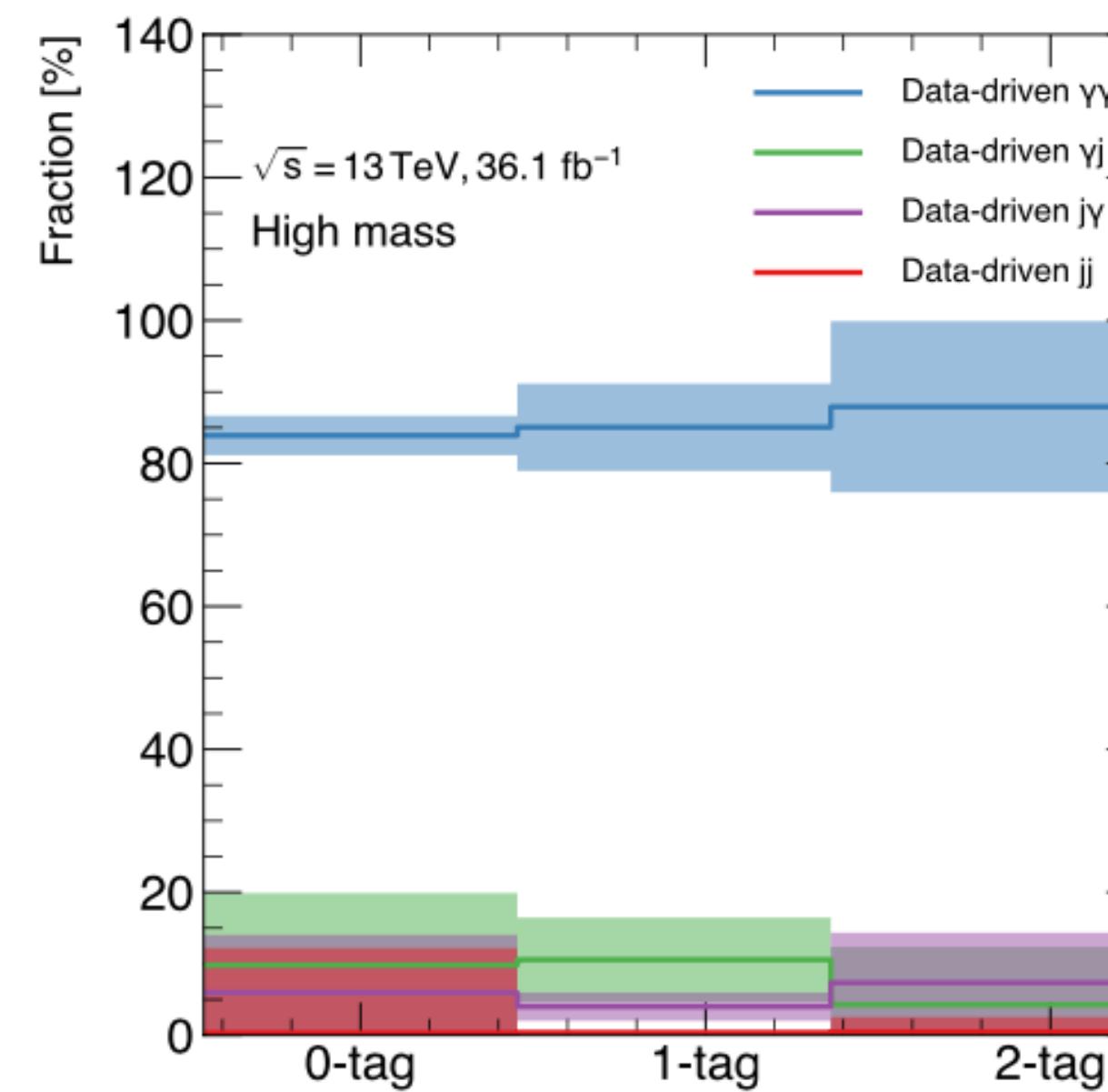
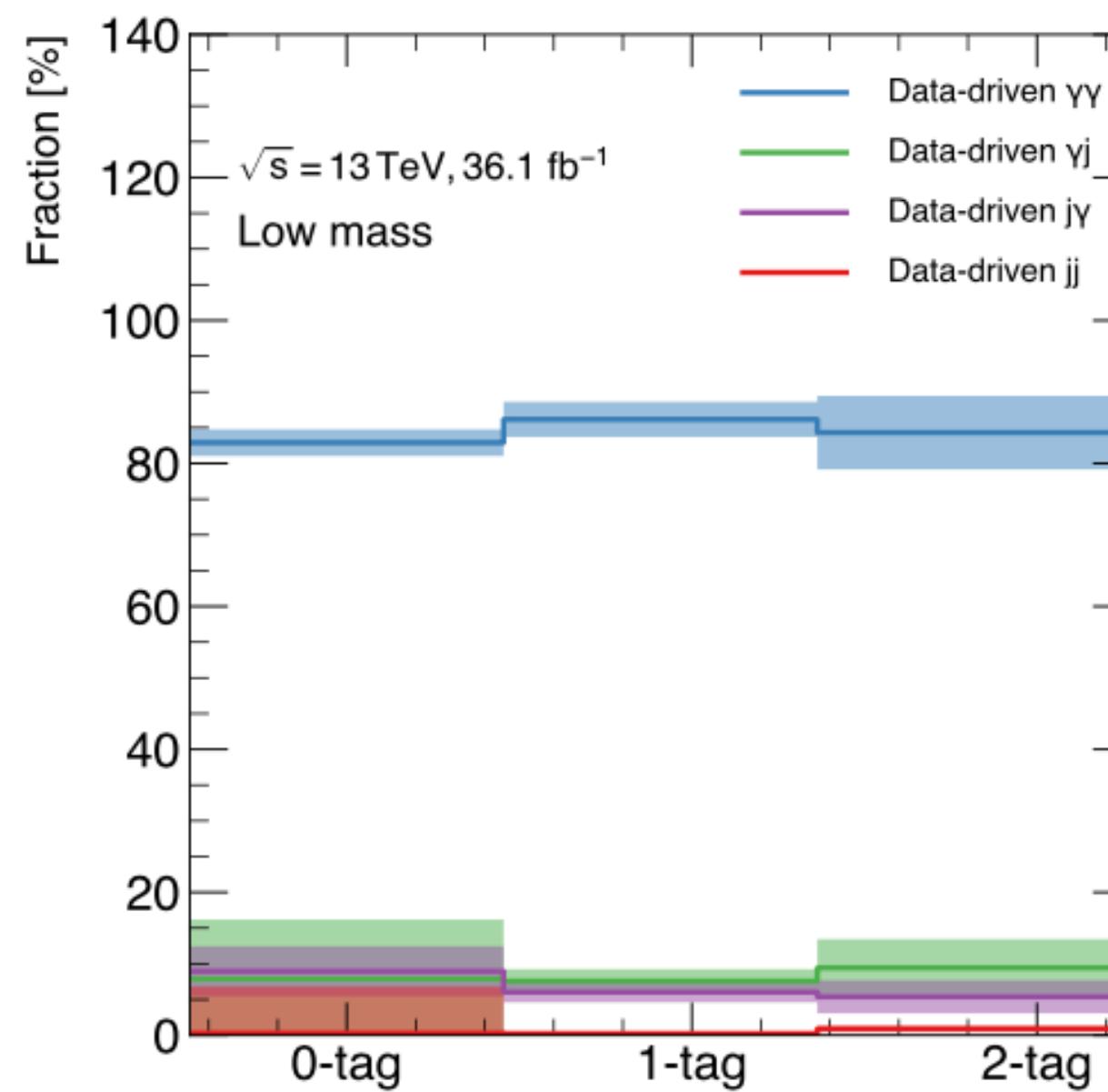
**Use data-driven method at estimating  $\gamma\gamma$ -continuum contribution from real photons and jet fakes ( $\gamma\gamma$ ,  $\gamma j$ ,  $j\gamma$ ,  $jj$ )**

- Different kinematics not well modeled in MC

## Use 2x2D sideband method

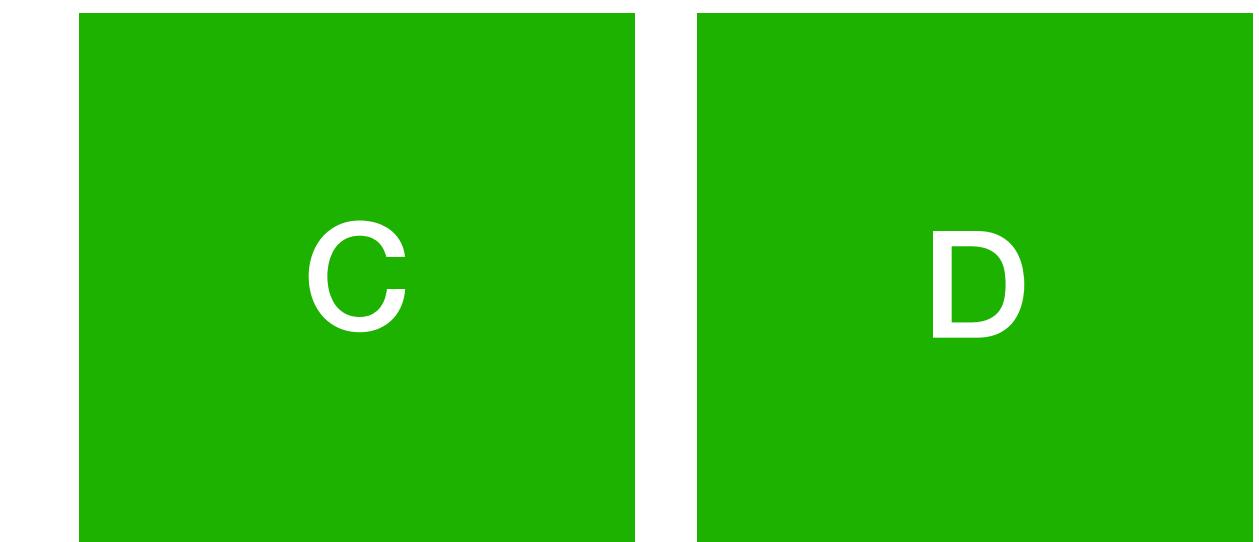
Method constructs 4 regions for each photon, using orthogonal cuts on photon identification and isolation criteria.

- 4 regions \*\* 2 photons = 16 total regions
- Use to establish a system of 16 equations on yield with 19 unknowns, constrain 4 from MC, 2 are assumed to be 1.0
- Perform bin-wise in  $m_{\gamma\gamma}$  0-tag (statistics), to establish a smooth reweighting function, apply across categories

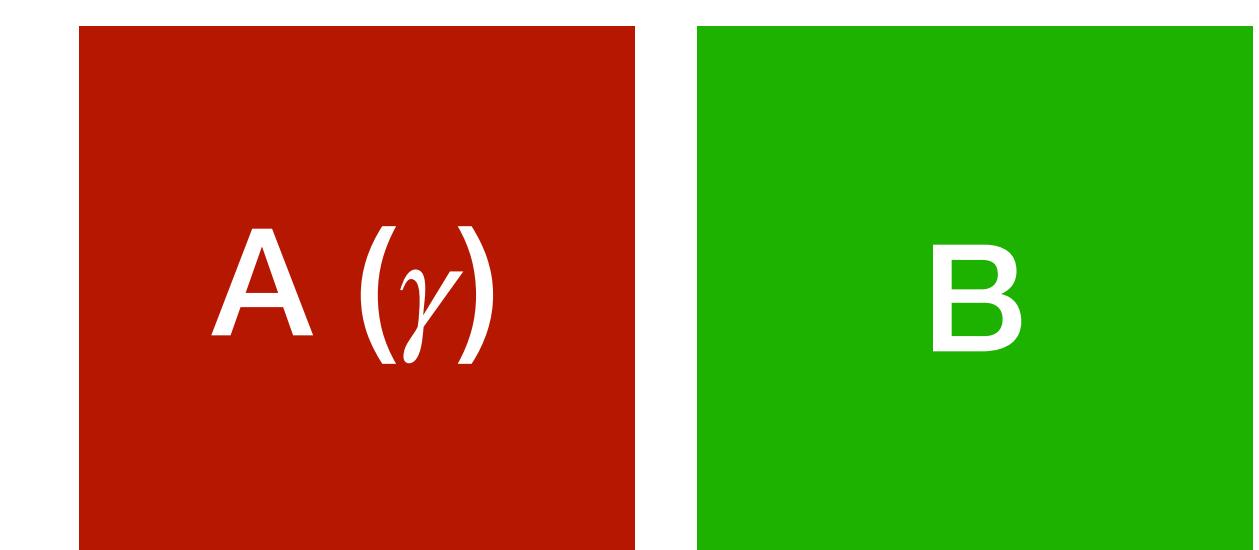


### Identification

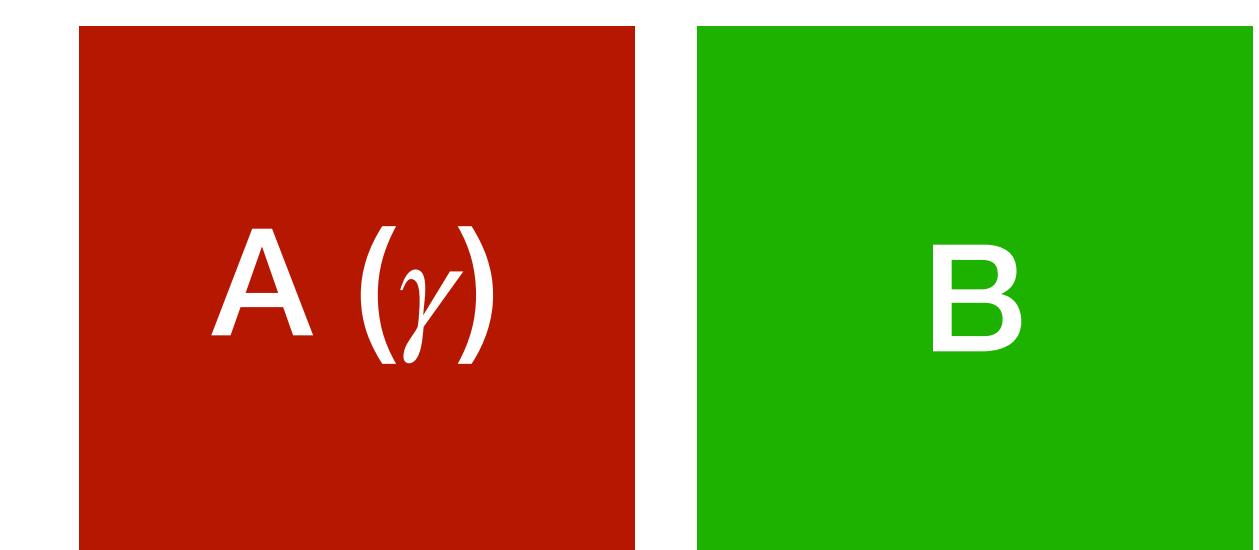
Loose criteria  
(but not tight)



Tight criteria



Tight criteria



Loose criteria  
(but not tight)

### Isolation



# Analysis Approach - Signal Modeling

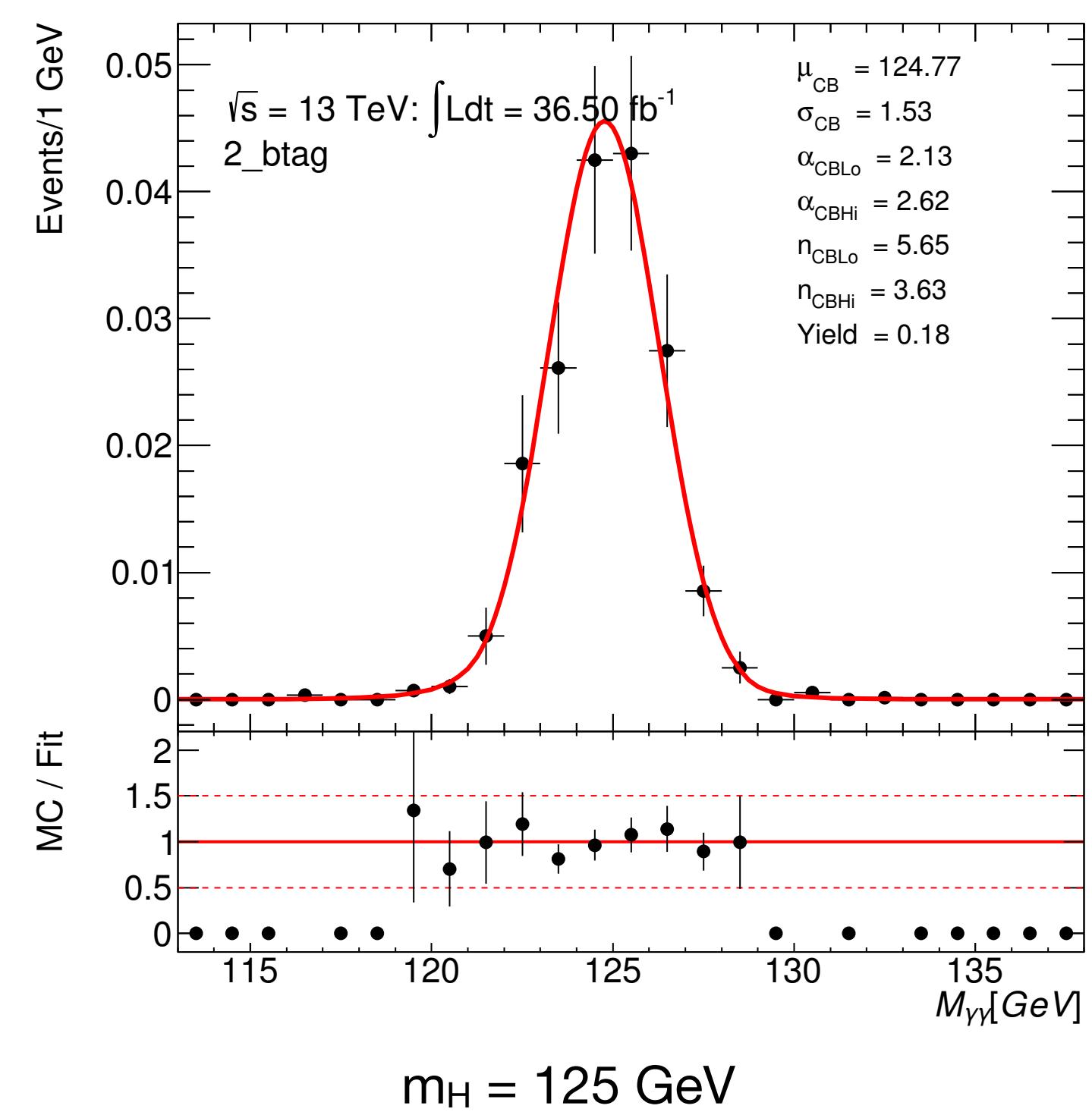
Signal modeling - parametrize Higgs peak (non-resonant) or peak for generic scalar particle decaying to Higgs pairs (resonant)

## Non-resonant Analysis:

Fit  $m_{\gamma\gamma}$  spectrum

Double-sided crystal ball

- Gaussian with power law tails

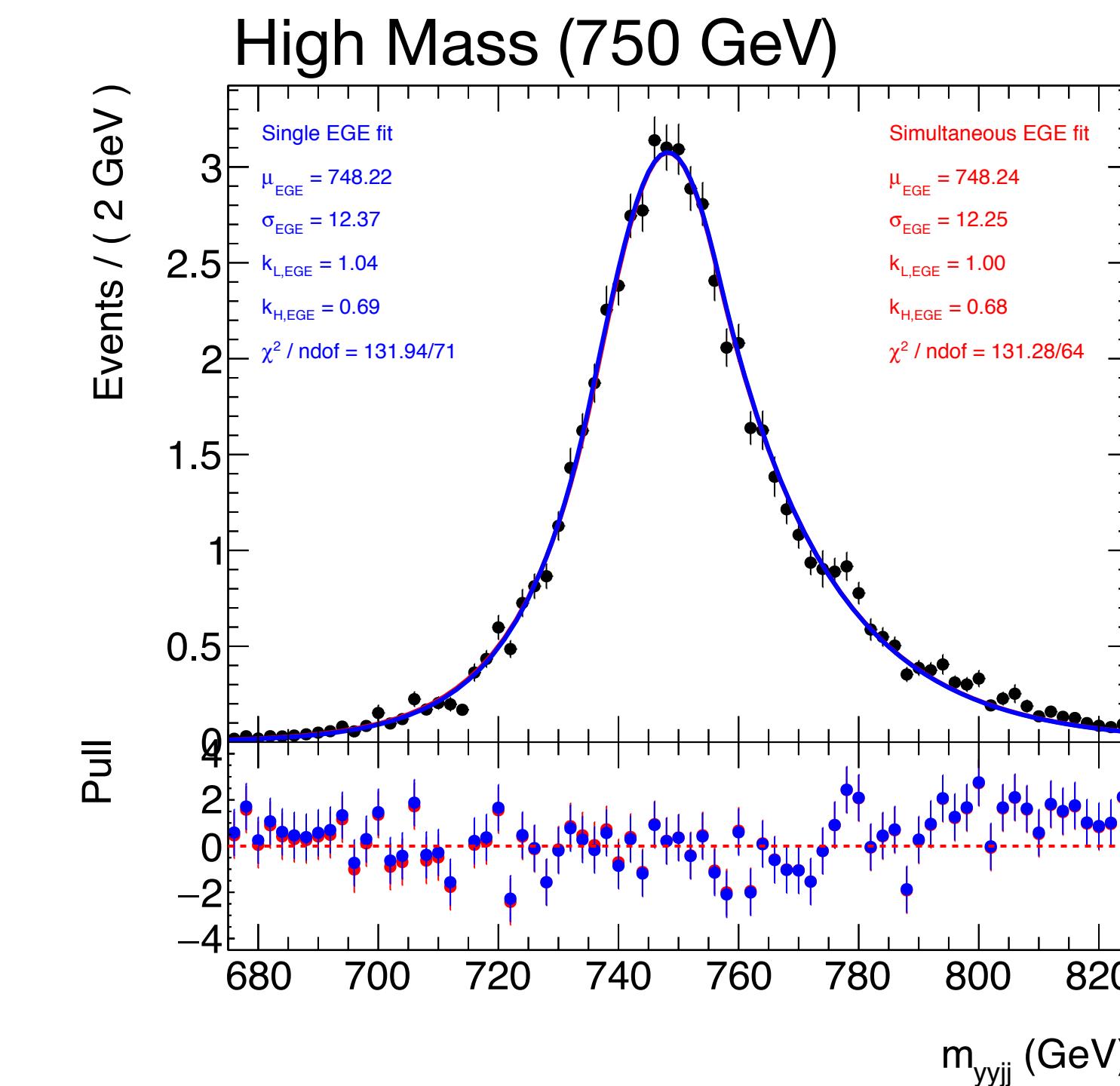
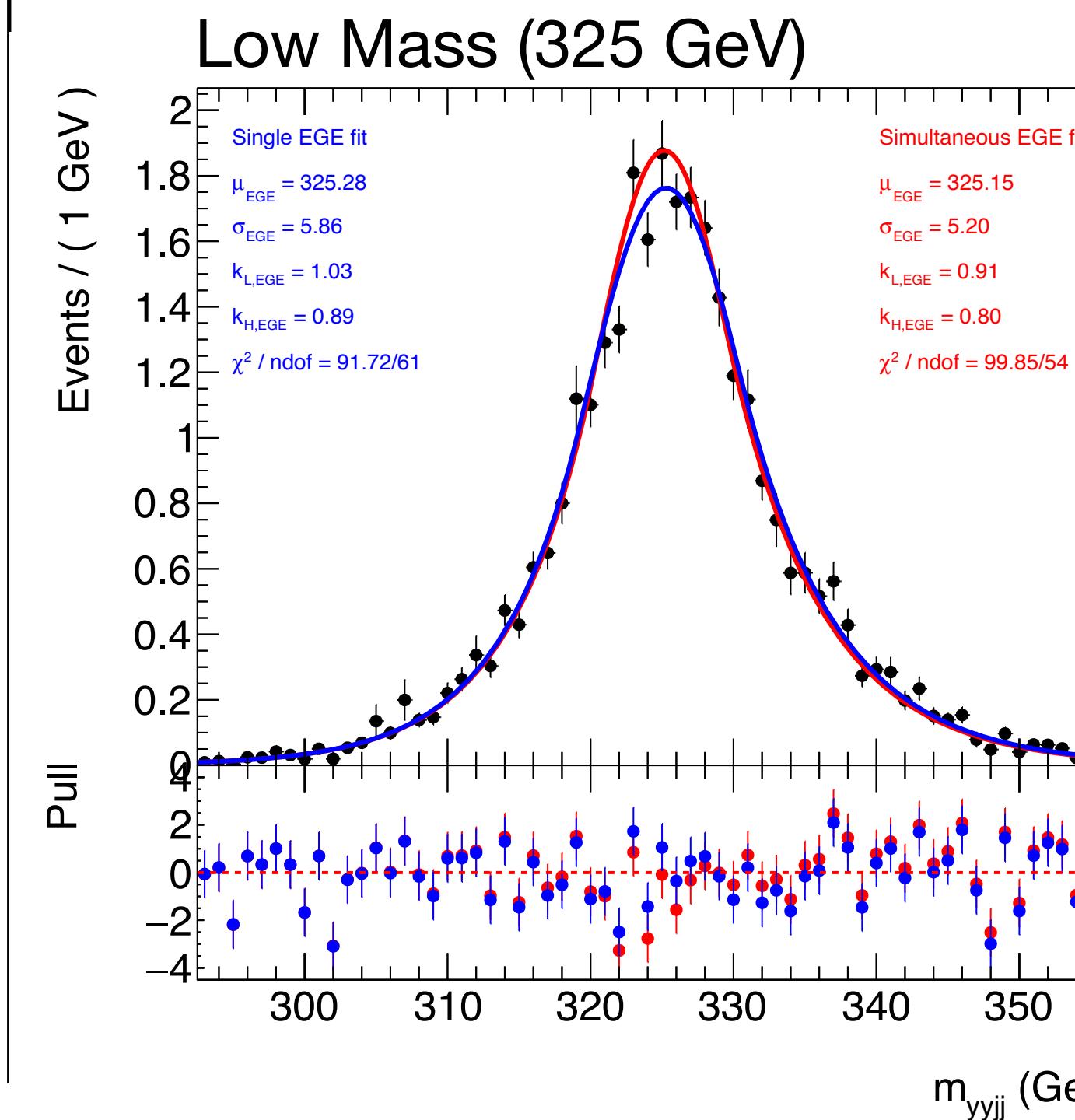


## Resonant Analysis:

Cut  $120 < m_{\gamma\gamma}/\text{GeV} < 130$  Fit  $m_{\gamma\gamma\text{bb}}$  spectrum, low and high mass categories separately

Gaussian with exponential tails

- Selected converging fit with minimal number of parameters



# Analysis Approach - Background Modeling

Perform unbinned profile likelihood fit

Models selected through “spurious signal” test -  
perform signal+background fits to just the  
background, take as bias. Select function that  
minimizes this bias

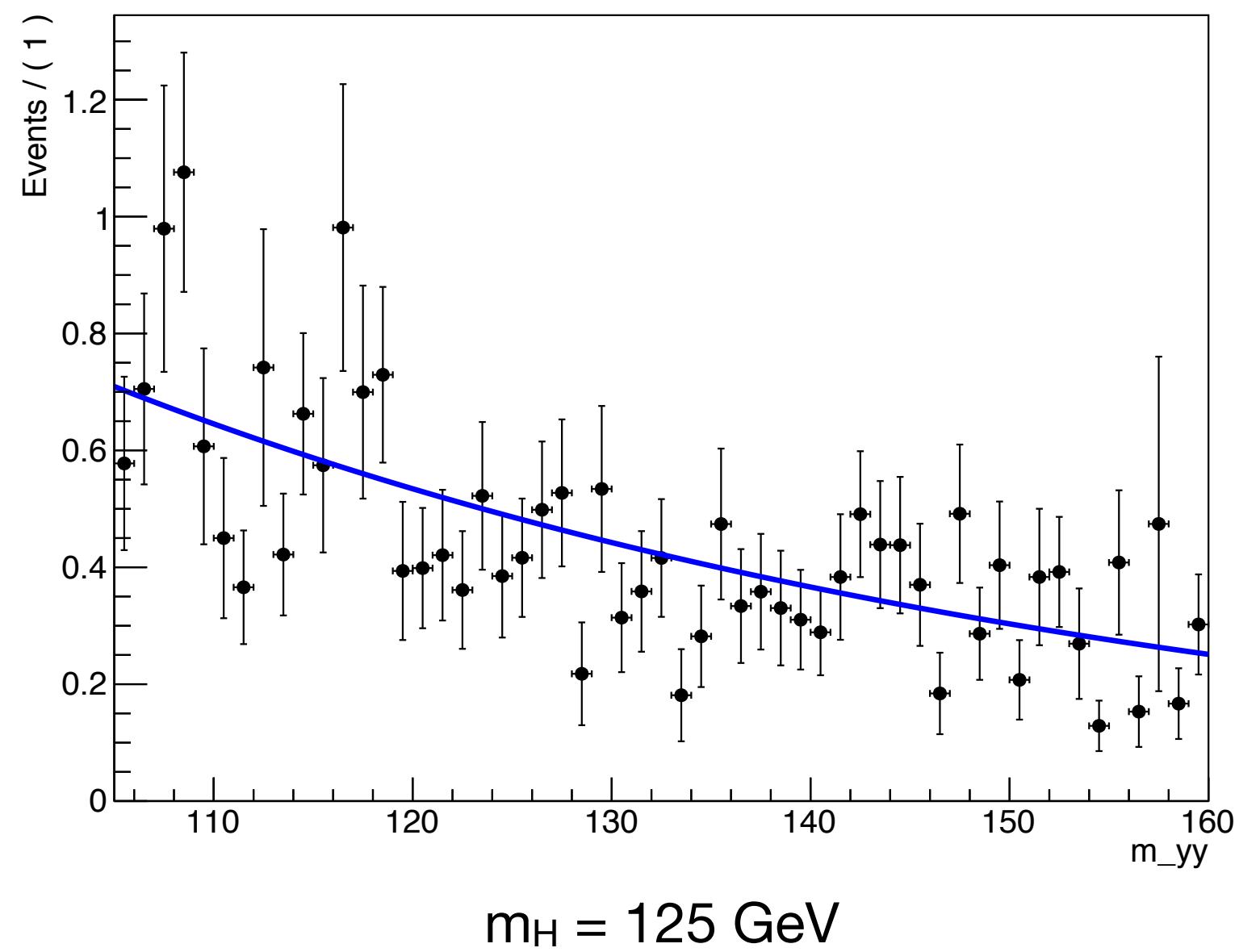
**Non-resonant Analysis:**

Fit  $m_{\gamma\gamma}$  spectrum

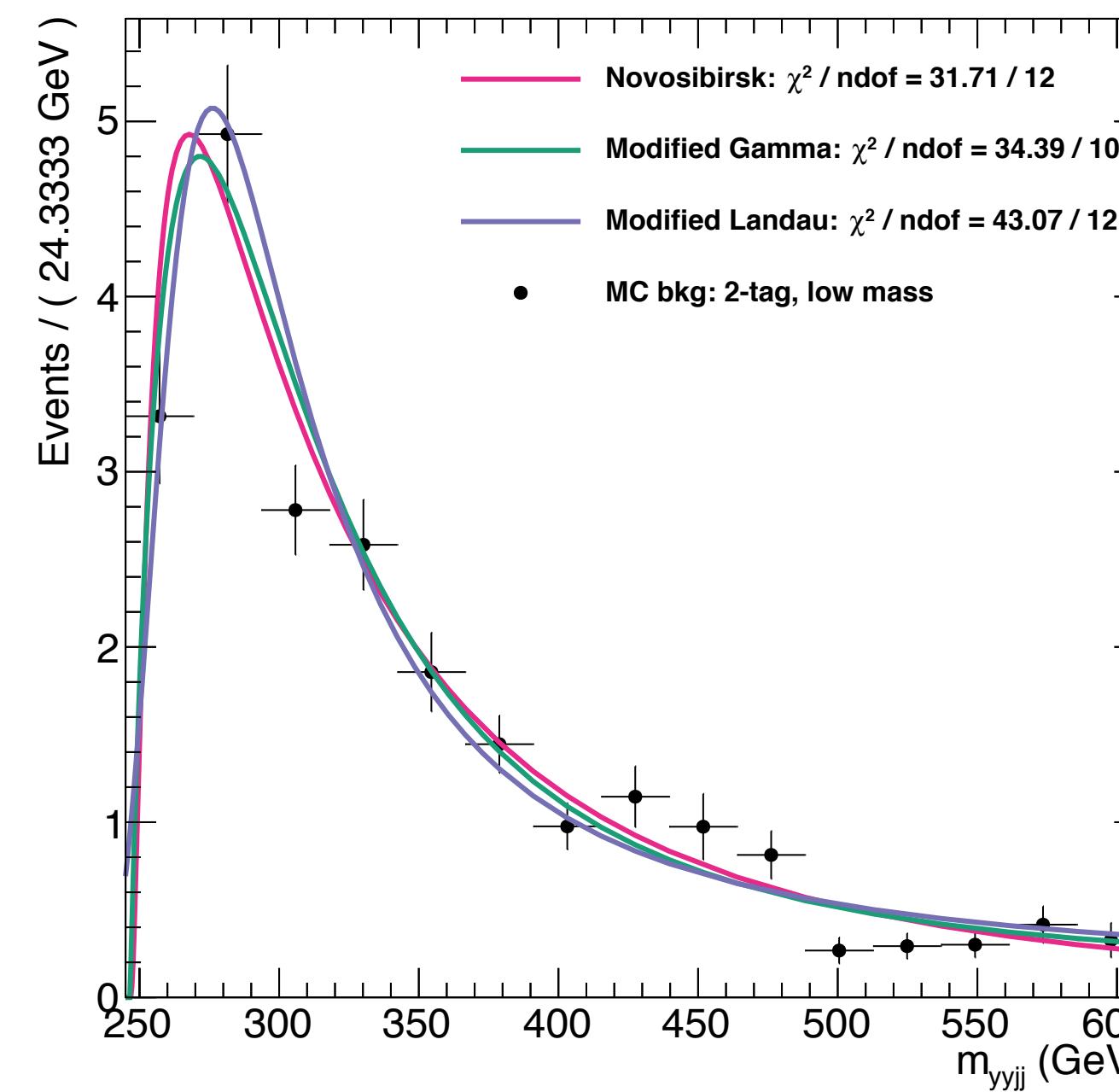
**Resonant Analysis:**

Cut  $120 < m_{\gamma\gamma}/\text{GeV} < 130$  Fit  $m_{\gamma\gamma\text{bb}}$  spectrum, low and high mass categories separately

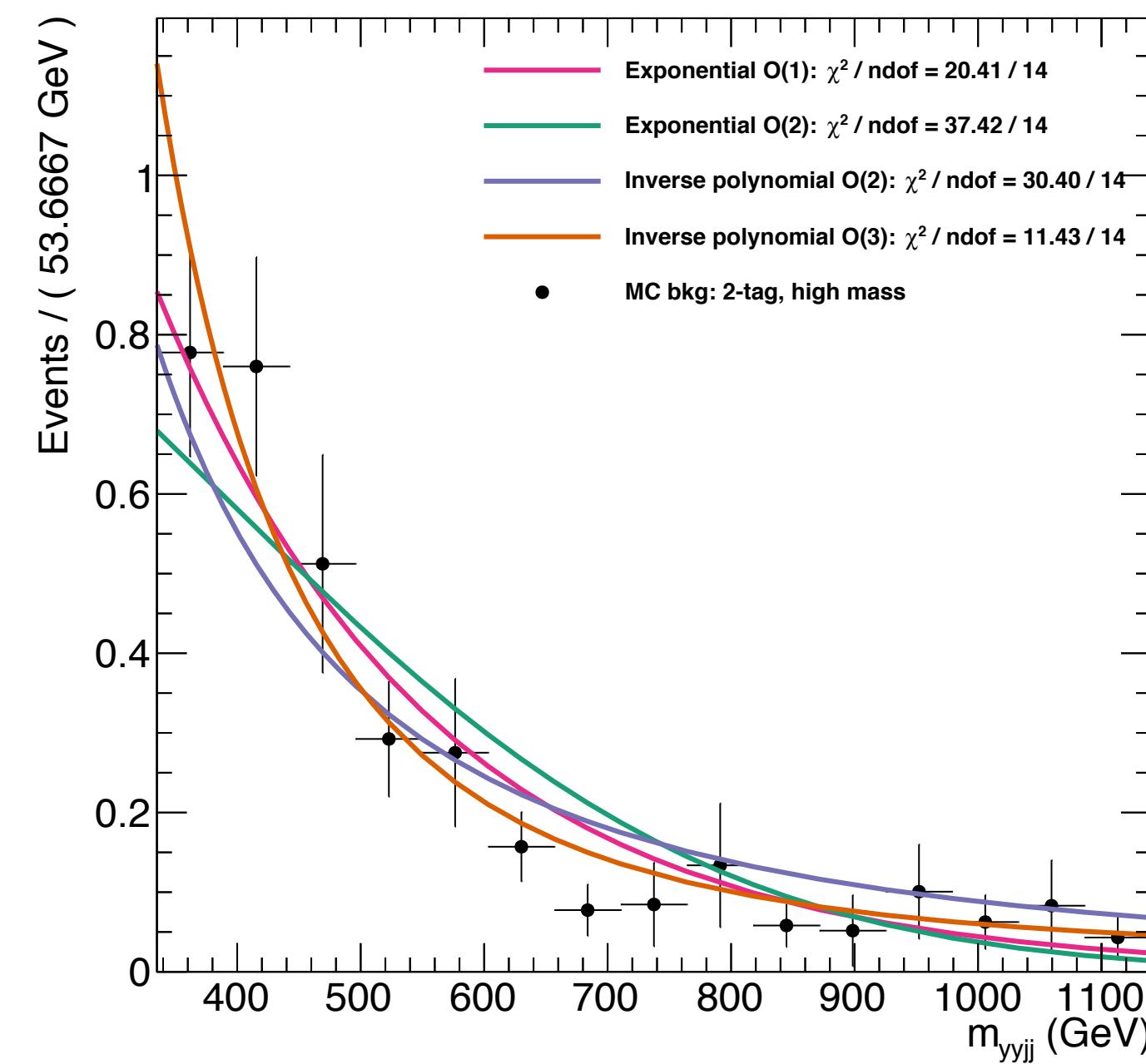
First order exponential



Low mass - Novosibirsk  
(Used for  $m_x < 500$ )



High mass - Exponential  
(Used for  $m_x \geq 500$ )



\* Additionally single Higgs background contributions modeled with double sided crystal ball functions

**Perform signal plus background fits using functions** - No excesses seen, so set limits

# Uncertainty

- **Non-resonant Signal and single Higgs:** Shape and yields extracted from fit to MC, uncertainties propagated to yield, peak, and width
- **Resonant Signal:** Shape and yield evaluated at each mass hypothesis, maximum taken for all
- **Background Parameters:** Unconstrained nuisance parameters

Source of systematic uncertainty	% effect relative to nominal in the 2-tag (1-tag) category					
	Non-resonant analysis		Resonant analysis: BSM $HH$			
	SM $HH$ signal	Single- $H$ bkg	Loose selection	Tight selection		
Luminosity	$\pm 2.1$ ( $\pm 2.1$ )	$\pm 2.1$ ( $\pm 2.1$ )	$\pm 2.1$ ( $\pm 2.1$ )	$\pm 2.1$ ( $\pm 2.1$ )	$\pm 2.1$ ( $\pm 2.1$ )	
Trigger	$\pm 0.4$ ( $\pm 0.4$ )	$\pm 0.4$ ( $\pm 0.4$ )	$\pm 0.4$ ( $\pm 0.4$ )	$\pm 0.4$ ( $\pm 0.4$ )	$\pm 0.4$ ( $\pm 0.4$ )	
Pile-up modelling	$\pm 3.2$ ( $\pm 1.3$ )	$\pm 2.0$ ( $\pm 0.8$ )	$\pm 4.0$ ( $\pm 4.2$ )	$\pm 4.0$ ( $\pm 3.8$ )		
Photon	identification	$\pm 2.5$ ( $\pm 2.4$ )	$\pm 1.7$ ( $\pm 1.8$ )	$\pm 2.6$ ( $\pm 2.6$ )	$\pm 2.5$ ( $\pm 2.5$ )	
	isolation	$\pm 0.8$ ( $\pm 0.8$ )	$\pm 0.8$ ( $\pm 0.8$ )	$\pm 0.8$ ( $\pm 0.8$ )	$\pm 0.9$ ( $\pm 0.9$ )	
	energy resolution	-	-	$\pm 1.0$ ( $\pm 1.3$ )	$\pm 1.8$ ( $\pm 1.2$ )	
	energy scale	-	-	$\pm 0.9$ ( $\pm 3.0$ )	$\pm 0.9$ ( $\pm 2.4$ )	
Jet	energy resolution	$\pm 1.5$ ( $\pm 2.2$ )	$\pm 2.9$ ( $\pm 6.4$ )	$\pm 7.5$ ( $\pm 8.5$ )	$\pm 6.4$ ( $\pm 6.4$ )	
	energy scale	$\pm 2.9$ ( $\pm 2.7$ )	$\pm 7.8$ ( $\pm 5.6$ )	$\pm 3.0$ ( $\pm 3.3$ )	$\pm 2.3$ ( $\pm 3.4$ )	
Flavor tagging	$b$ -jets	$\pm 2.4$ ( $\pm 2.5$ )	$\pm 2.3$ ( $\pm 1.4$ )	$\pm 3.4$ ( $\pm 2.6$ )	$\pm 2.5$ ( $\pm 2.6$ )	
	$c$ -jets	$\pm 0.1$ ( $\pm 1.0$ )	$\pm 1.8$ ( $\pm 11.6$ )	-	-	
	light-jets	$<0.1$ ( $\pm 5.0$ )	$\pm 1.6$ ( $\pm 2.2$ )	-	-	
Theory	PDF+ $\alpha_S$	$\pm 2.3$ ( $\pm 2.3$ )	$\pm 3.1$ ( $\pm 3.3$ )	n/a	n/a	
	Scale	$+4.3$ ( $+4.3$ )	$+4.9$ ( $+5.3$ )	n/a	n/a	
	EFT	$-6.0$ ( $-6.0$ )	$+7.0$ ( $+8.0$ )	n/a	n/a	

**Dominant:**

Yield: Flavor tagging, JES, Photon ID, Pile-up

Scale: e/gamma Scale

Width: e/gamma Resolution, e/gamma Scale, JES, Pile-up, flavor tagging

# Results

No significant excesses seen, so exclusion limits set

Limits on non-resonant HH production

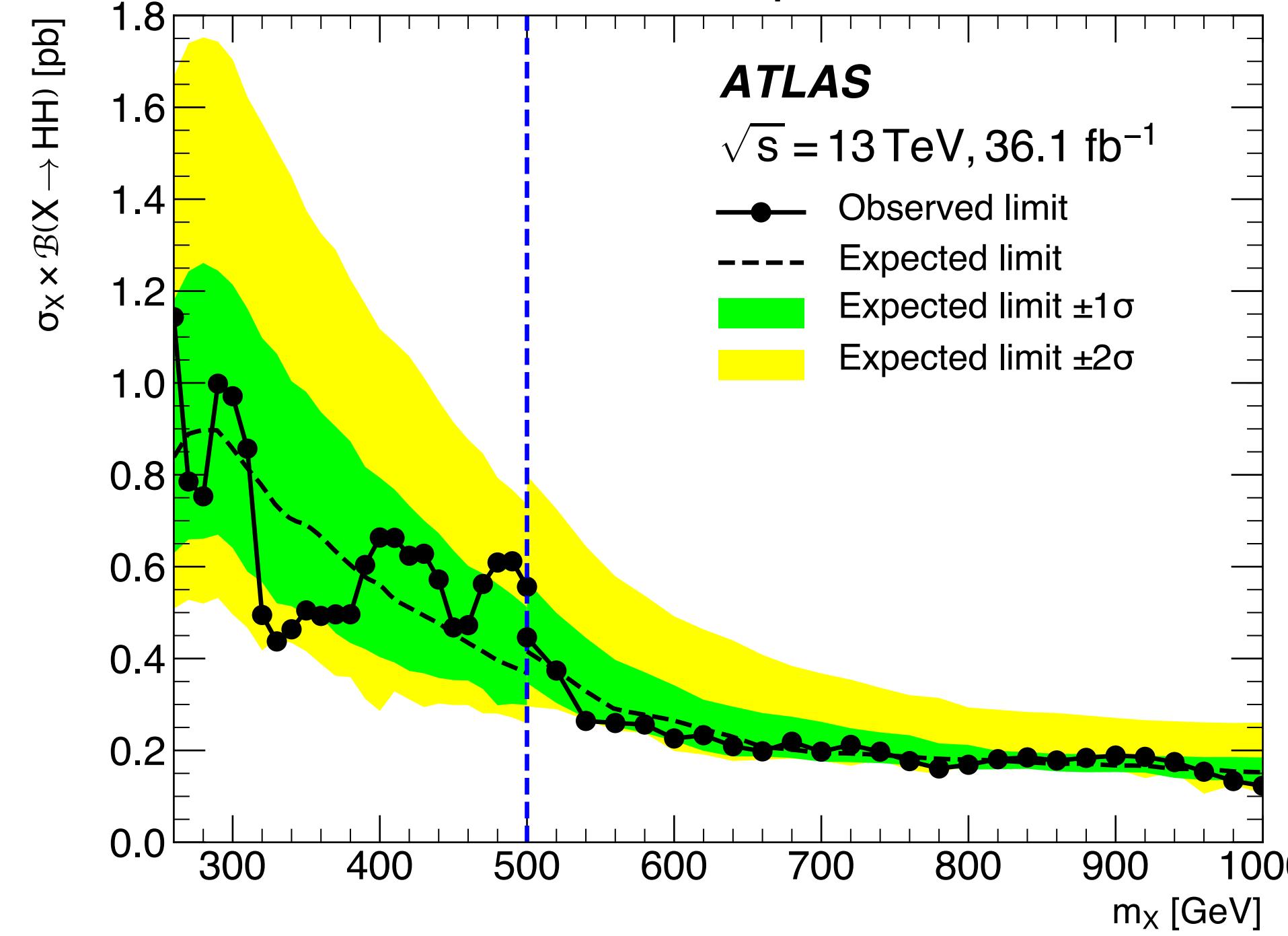
	Observed	Expected	$-1\sigma$	$+1\sigma$
$\sigma_{gg \rightarrow HH}$ [pb]	0.73	0.93	0.66	1.3
Multiple of $\sigma_{SM}$	22	28	20	40

Use profile likelihood ratio test statistic on  $m_{\gamma\gamma bb}$  ( $m_{\gamma\gamma}$ ) for the (non-)resonant analysis to derive signal yields

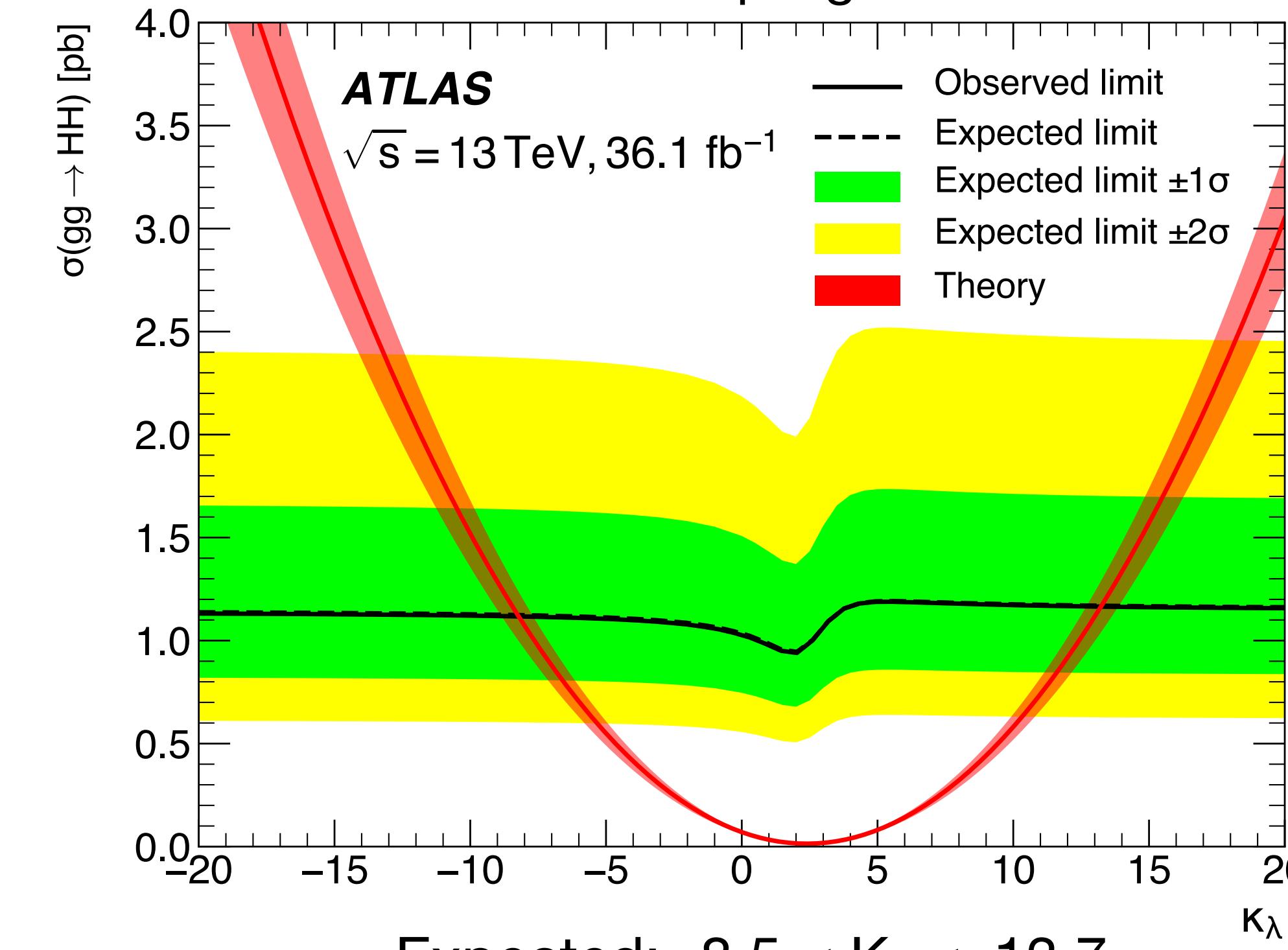
CLs prescription used on POI to set 95% CL

Limits are generated by using “toy” pseudo-experiments

Limits on resonant HH production



Limits on trilinear coupling



Expected:  $-8.5 < K_\lambda < 13.7$   
Observed:  $-8.2 < K_\lambda < 13.3$

But also looking forward to future di-Higgs analyses!  
- ~4x more data when using full Run 2 dataset  
- New analysis techniques and approaches



Northern Illinois  
University

# Vector Boson Fusion HH Production

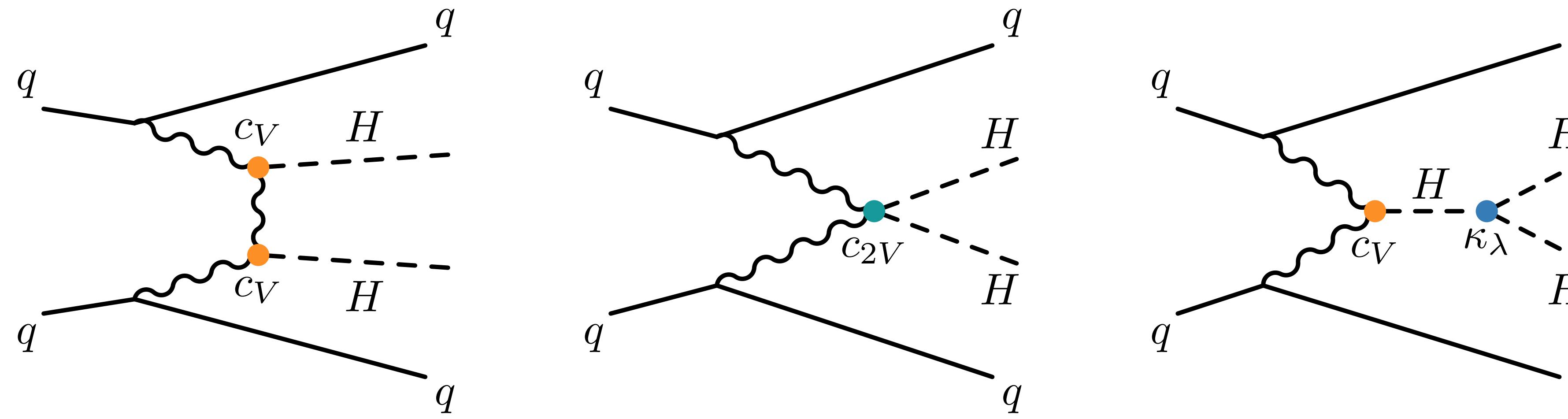


# SM Vector Boson Fusion Production

ggF is the dominant HH production mode, but what if we consider subdominant modes?

## Vector Boson Fusion

- Subdominant production mode, cross-section factor of 20 smaller than ggF



$$\sigma_{HH}^{ggF} = 33.4 \text{ fb}^{-1}$$

$$\sigma_{HH}^{vbf} = 1.64 \text{ fb}^{-1}$$

Small cross-section motivates use  
of machine learning (ML) to get the  
most pure possible signal  
selection

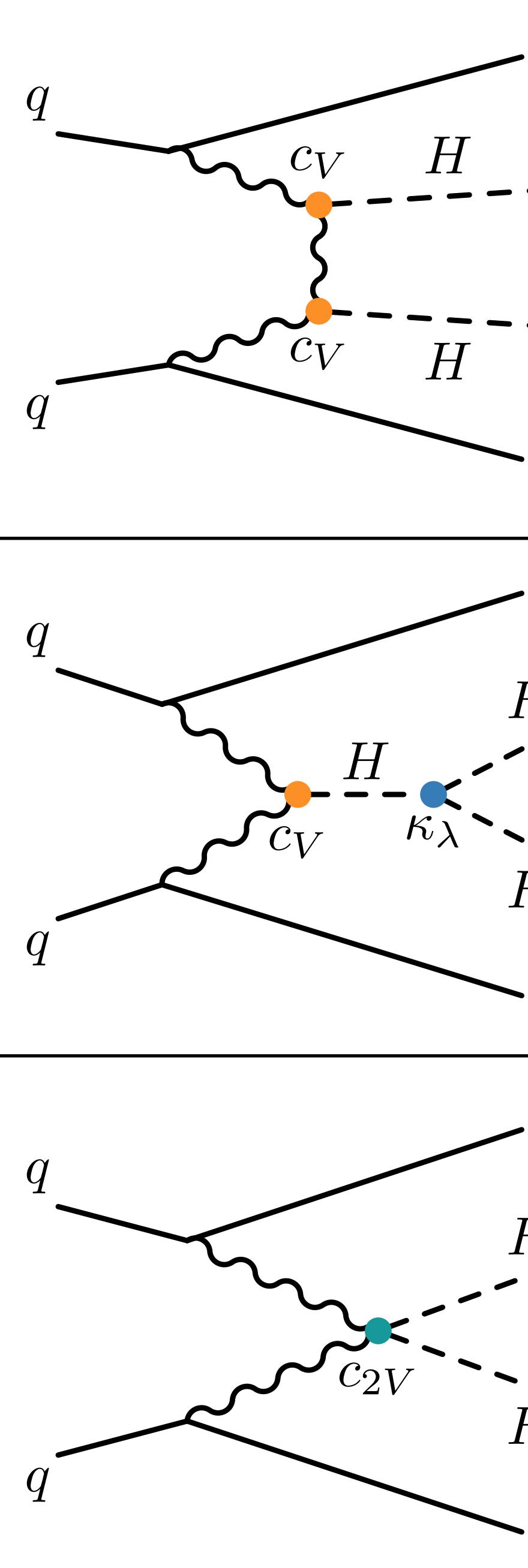
**However** - distinct topology that makes for easy signal discrimination

- Production yields 2 extra jets from quarks radiating vector bosons
- Extra jets have unique properties
  - Very forward in the detector
  - High  $m_{jj}$ , compared to  $H \rightarrow bb$  jets



# BSM Vector Boson Fusion Production

What if coupling strengths deviate from SM predictions?



$C_V$  (V VH) - Constrained by mono-Higgs analyses

[ATLAS-CONF-2018-031](#)

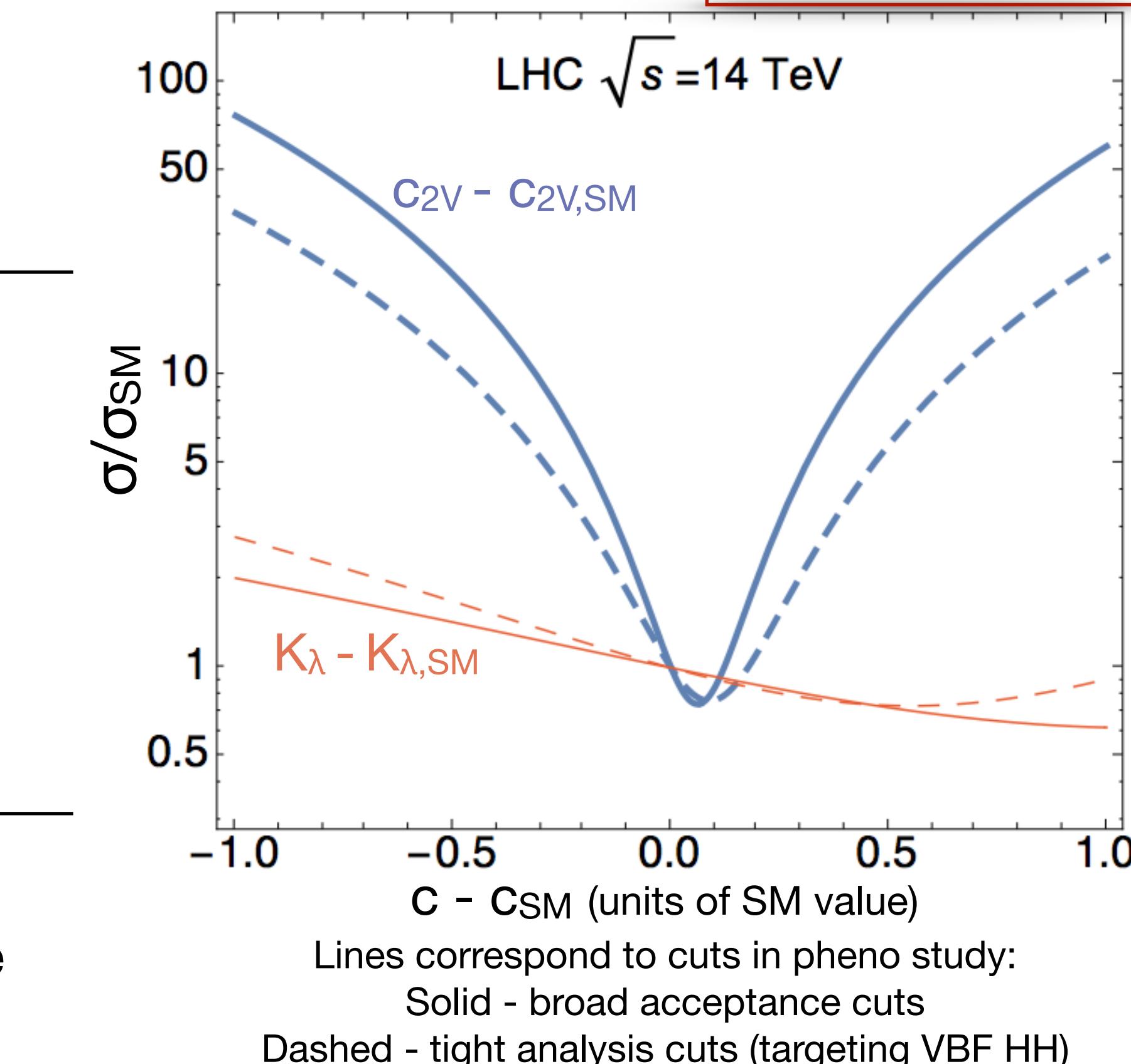
$$\sigma_{C_V}/\sigma_{C_V,SM} = 1.21 \pm^{0.22}_{0.21}$$

$\kappa_\lambda$  (HHH) - VBF HH production is also sensitive to deviations in  $\kappa_\lambda$ , and can contribute to limits

$C_{2V}$  (VVHH) - VBF HH provides unique handle at tree level to this coupling!

Greatly enhances HH cross-section, even at small deviations from predictions

Bishara, Contino, Rojo  
Eur. Phys. J. C 77 (2017) 481

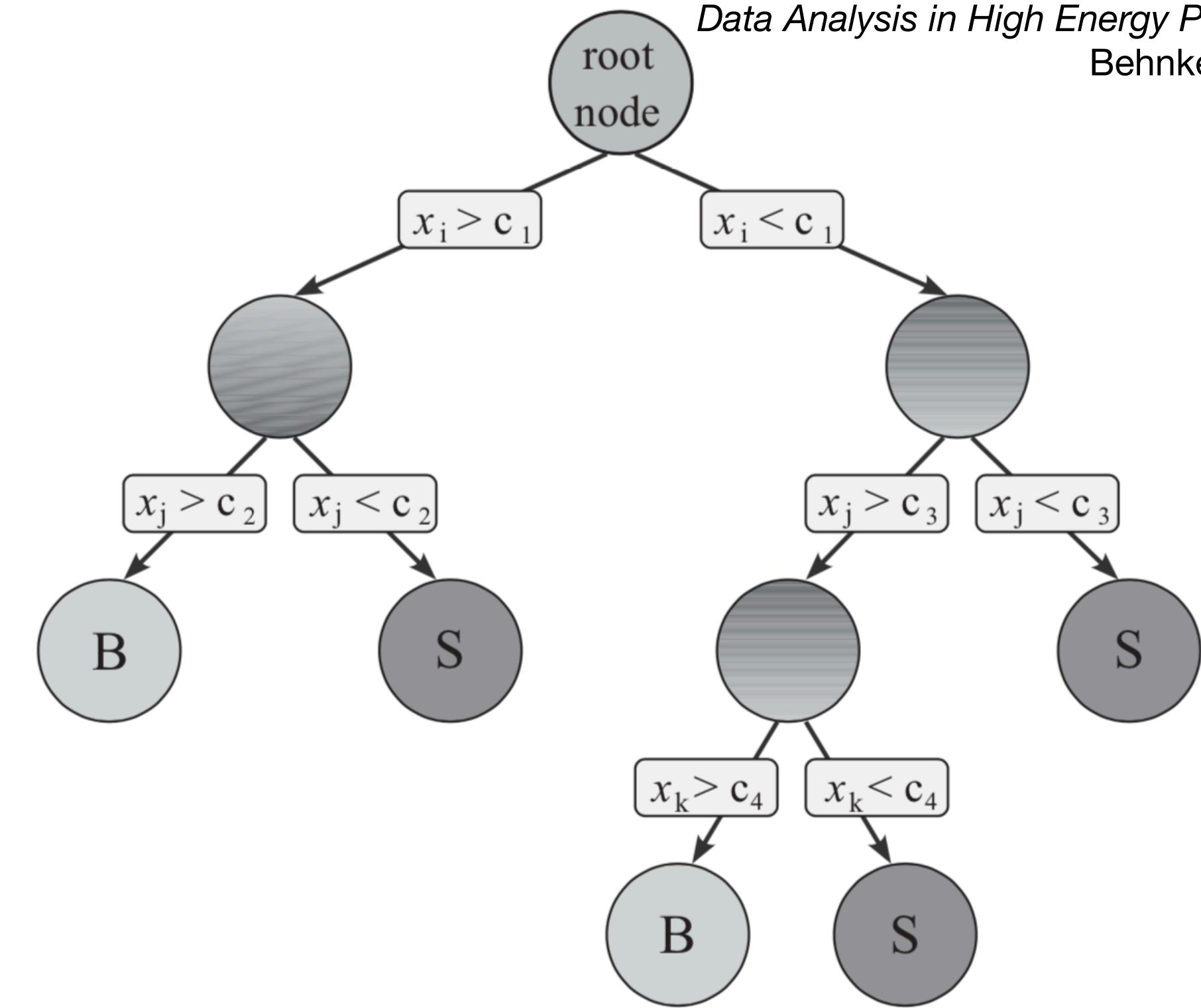


# Machine Learning Interlude #1

## Boosted Decision Trees

Decision Tree:

- Familiar “tree” based learning - split data using input variables, maximize a splitting criteria (e.g. signal/background separation)
- Continue splitting until a stopping criteria is met (e.g. depth, samples in a node)



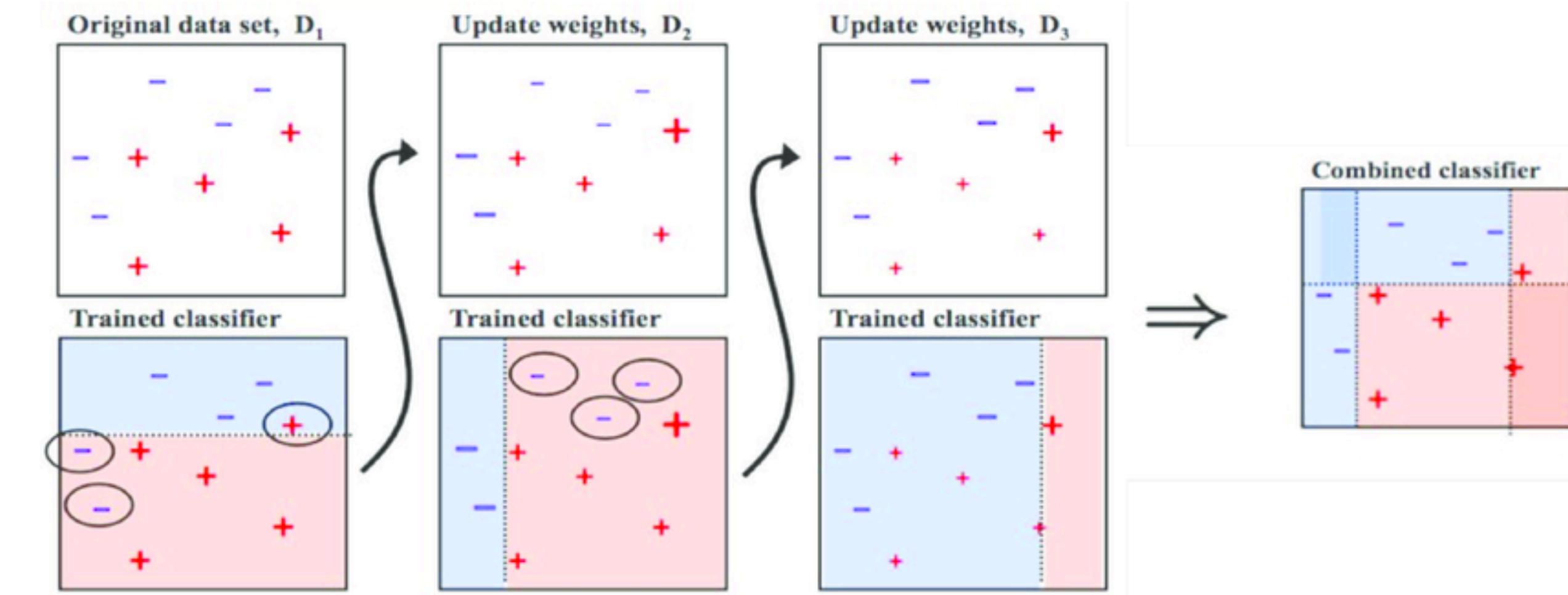
Boosting:

- After building tree, find misclassifications, upweight, and derive a new tree, do this for a defined number of trees
- Yields complex output space!
- “Ensemble method” - combining many weak learners to make one strong learner

Parameters in ML adapt by learning from observations (data and/or MC) - “learning” stage is called *training*

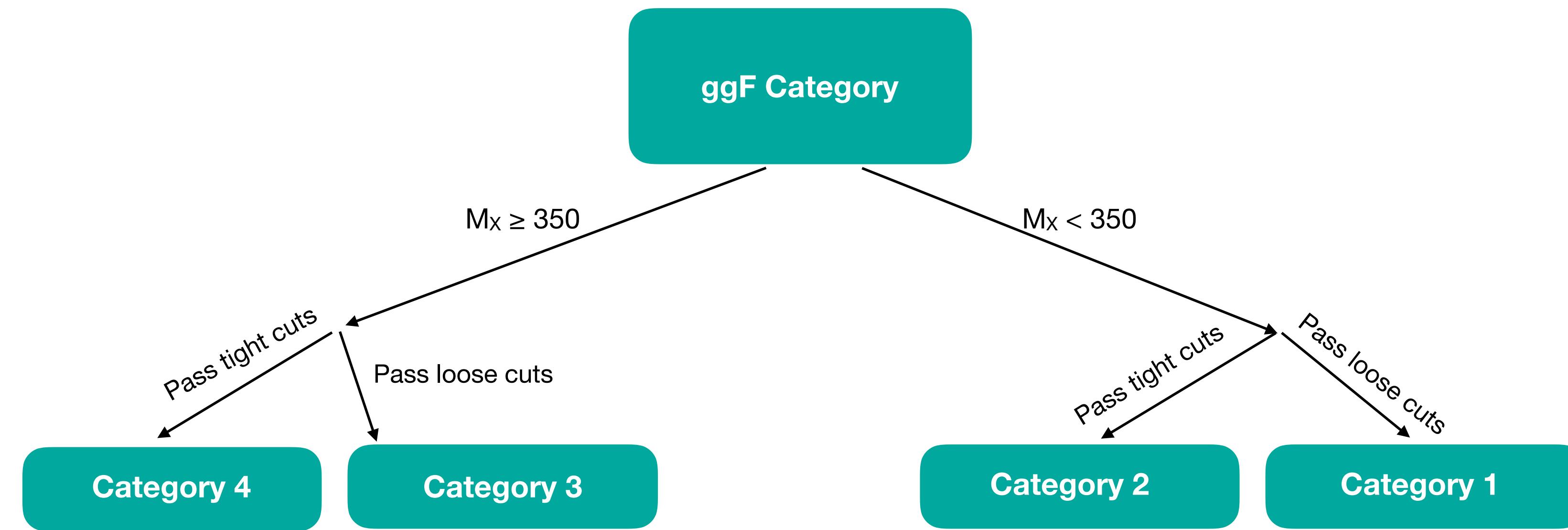
### Notes

- Commonly used in physics, easy to integrate into ATLAS software
- Good interpretability



# Full Run 2 Analysis Selection

Next iteration of analysis will construct 4 signal regions using a BDT trained on ggF signal  
- Targets differences in di-Higgs kinematics from backgrounds



Tight and loose cuts defined using BDT scores

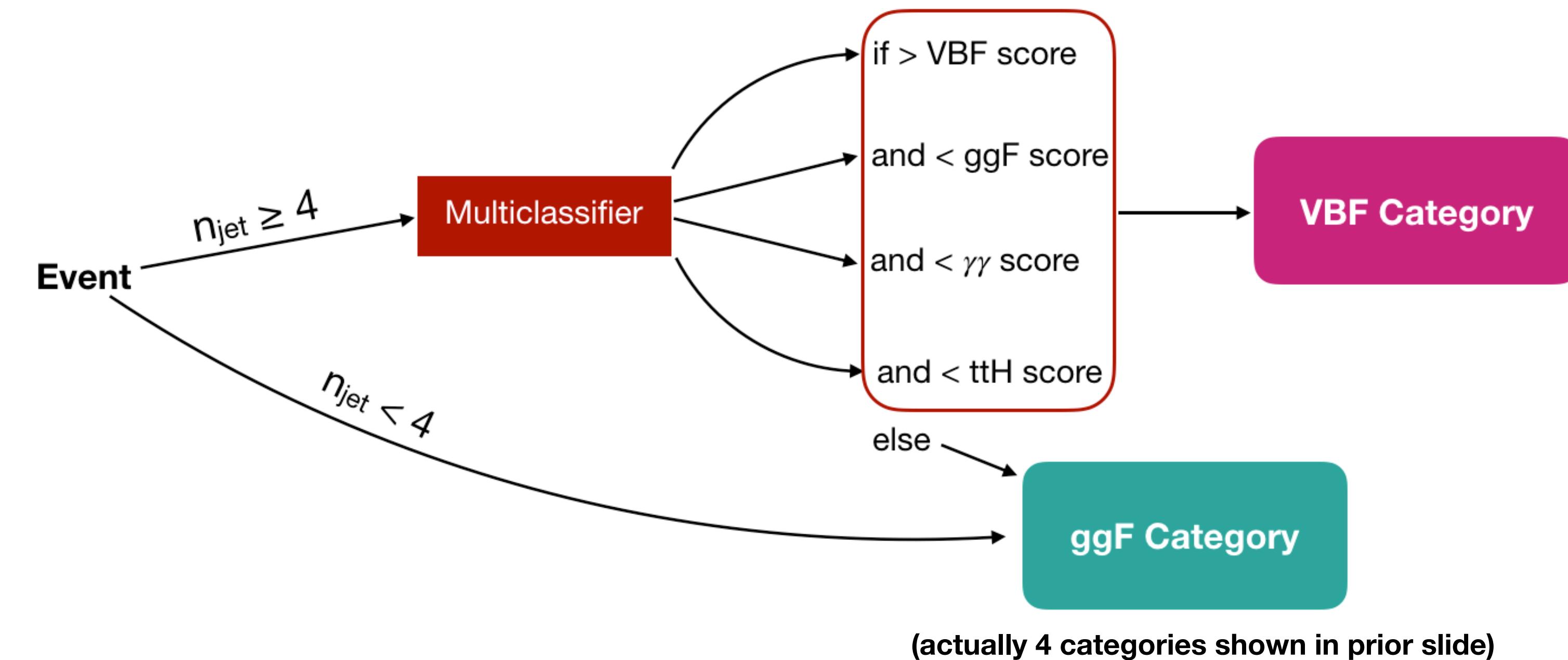
How to fit VBF production into this?



# Implementing into Analysis

Create a **VBF-enriched signal region** - use unique VBF topology to isolate this production from background and ggF HH production

- First considered just  $\gamma\gamma$ -continuum background, but found large contamination from ttH so developed a **multiclassifier** BDT, independent probabilities for VBF HH, ggF HH,  $\gamma\gamma$ , and ttH



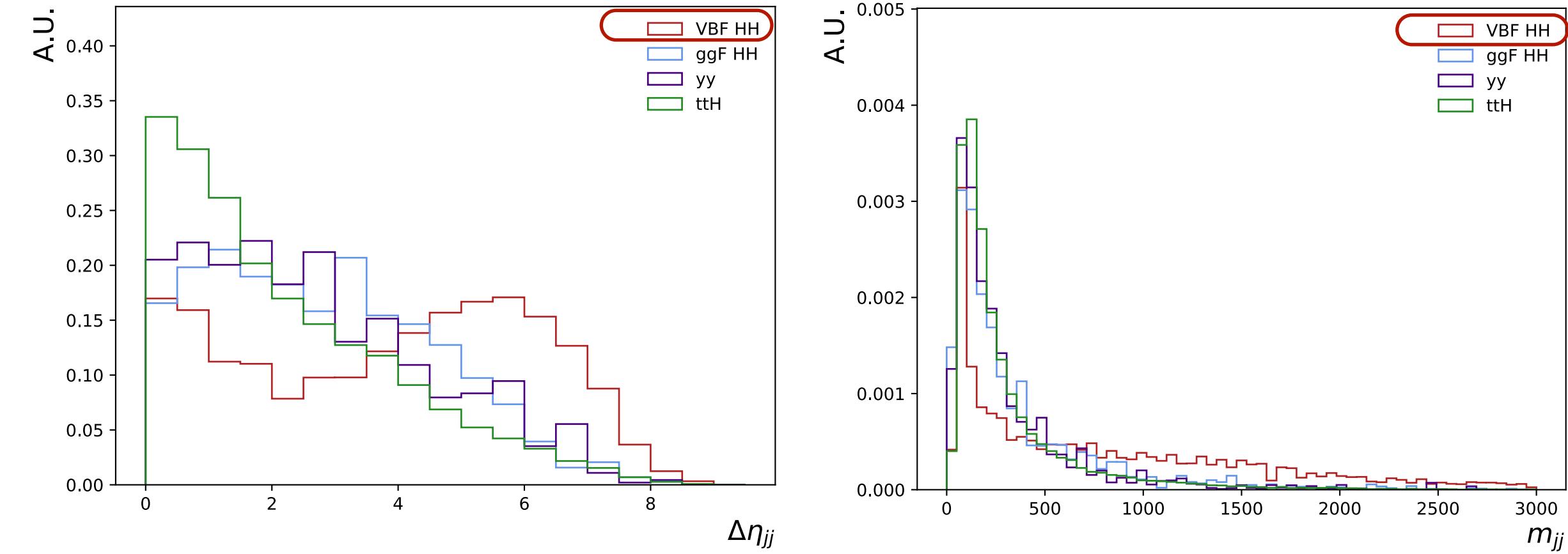
In total **26 variables** investigated as BDT inputs, through a variable pruning, **reduced to 10** with equivalent performance  
- Pruning motivated by cross correlations, correlation to final output, and final exhaustive search of all possible subsets



# BDT Inputs (not all shown here)

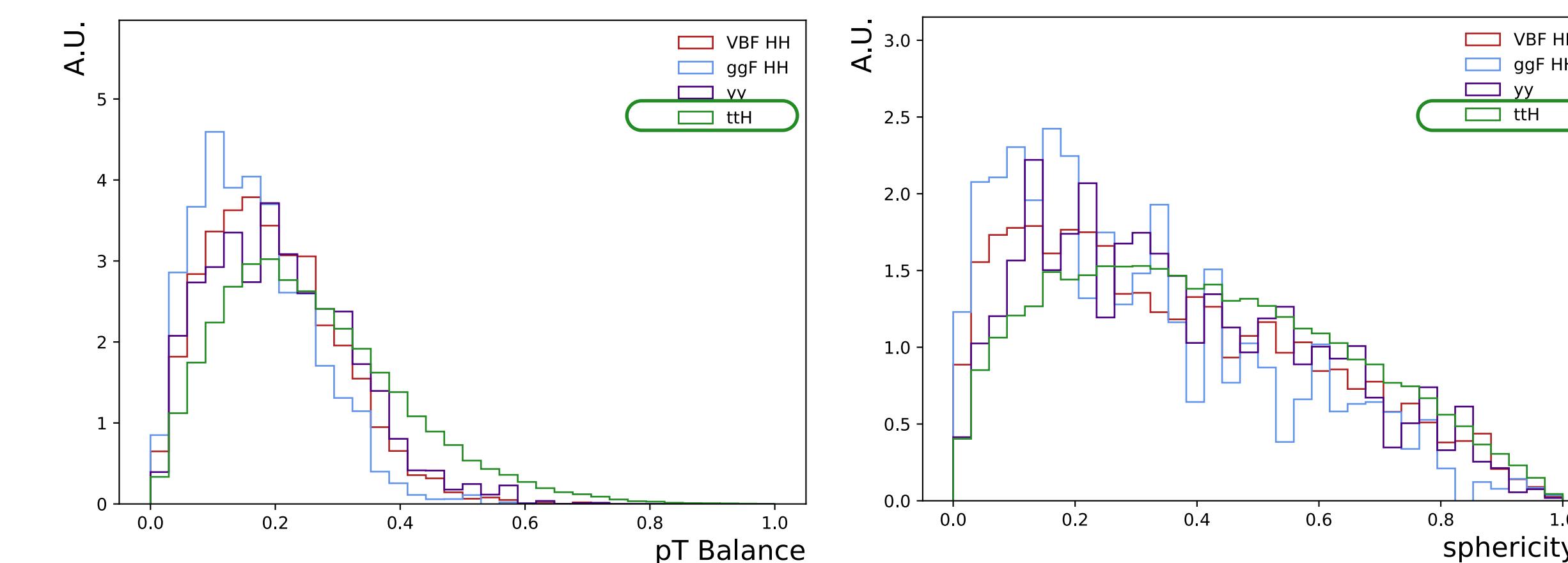
## VBF-targeting

Take advantage of unique VBF topology to separate out production mode



## ttH-Rejecting

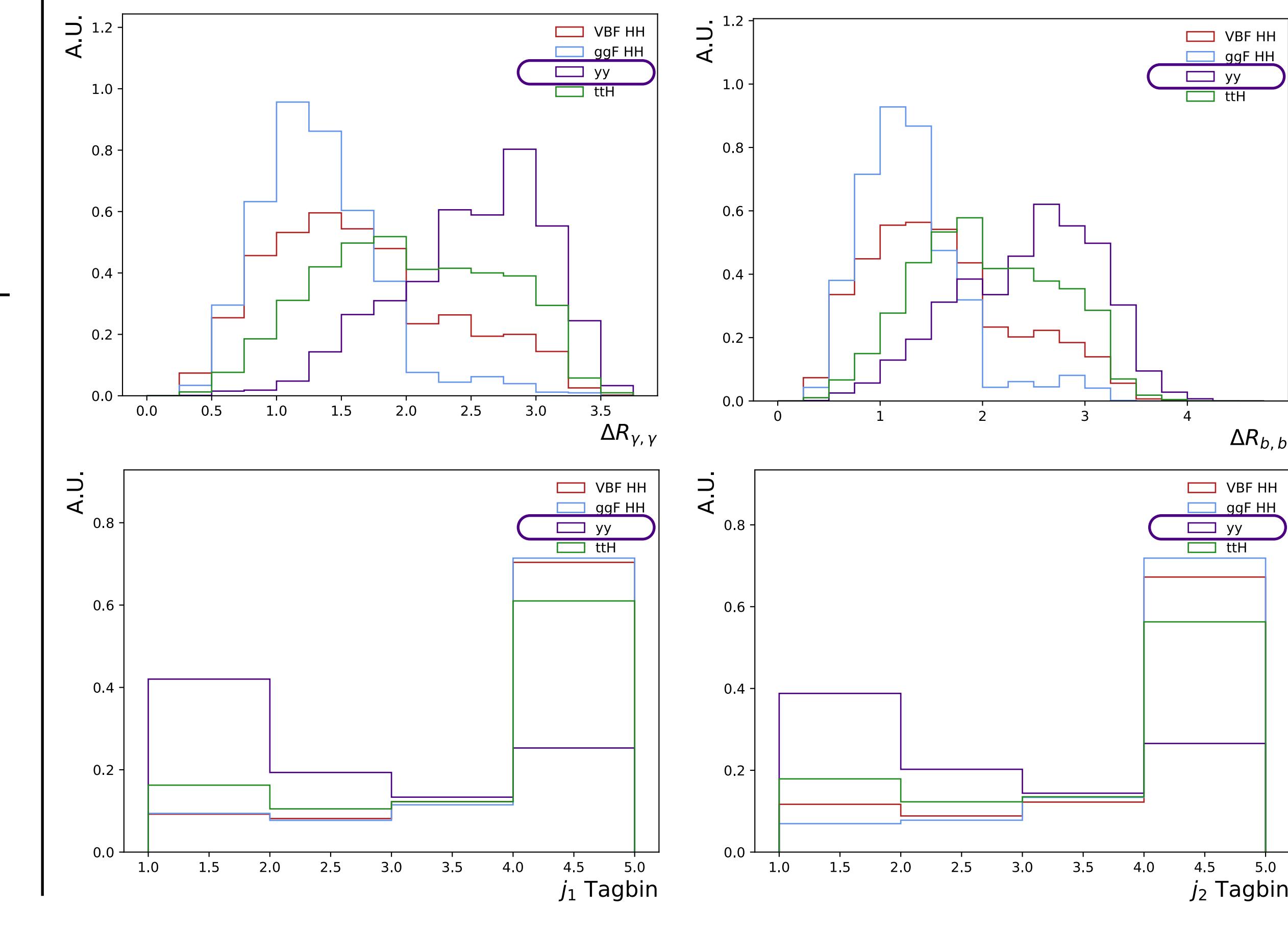
"Event shape" variables - describe the flow and distribution of energy in an event



BDT developed using XGBoost

## Higgs Kinematics (reject $\gamma\gamma$ )

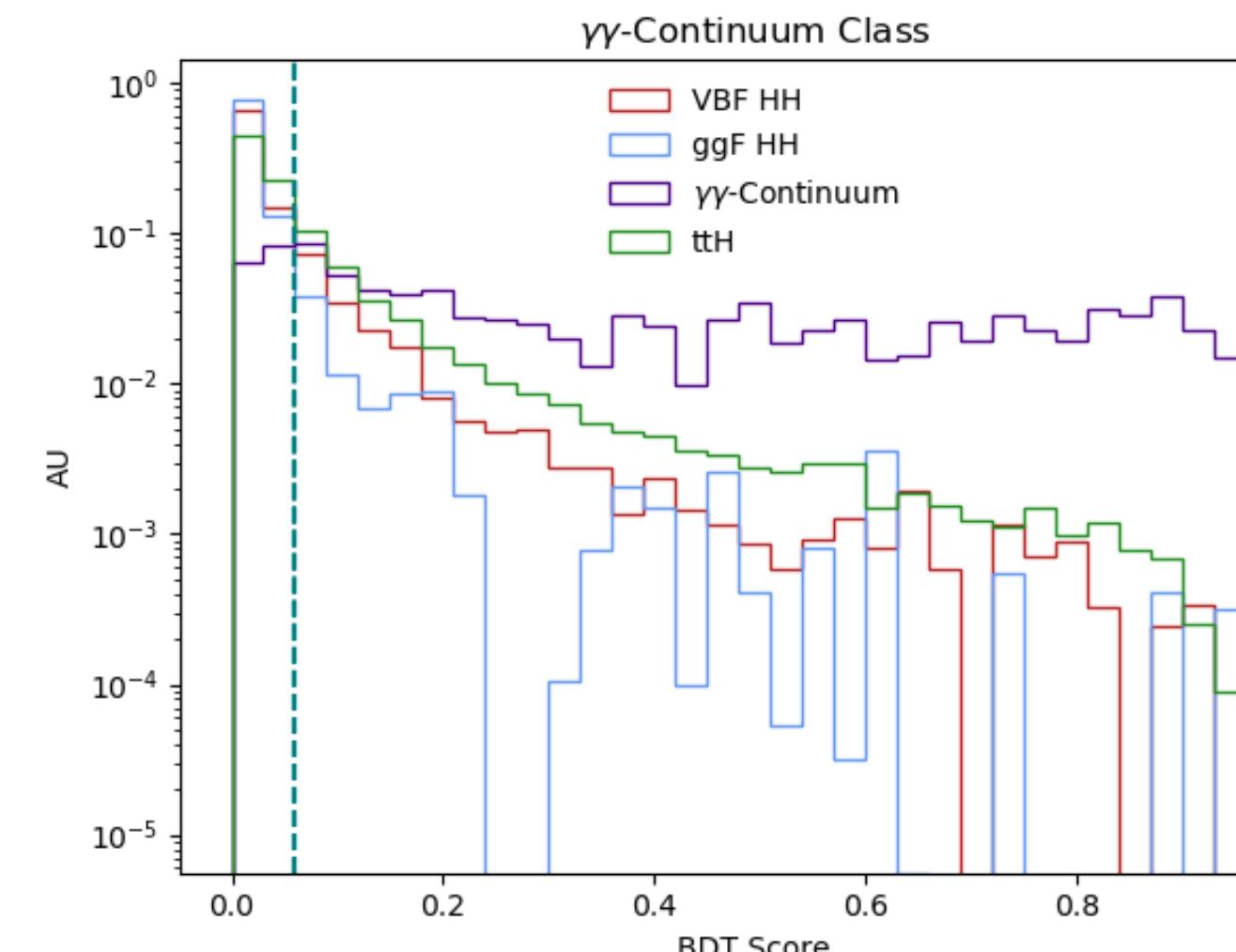
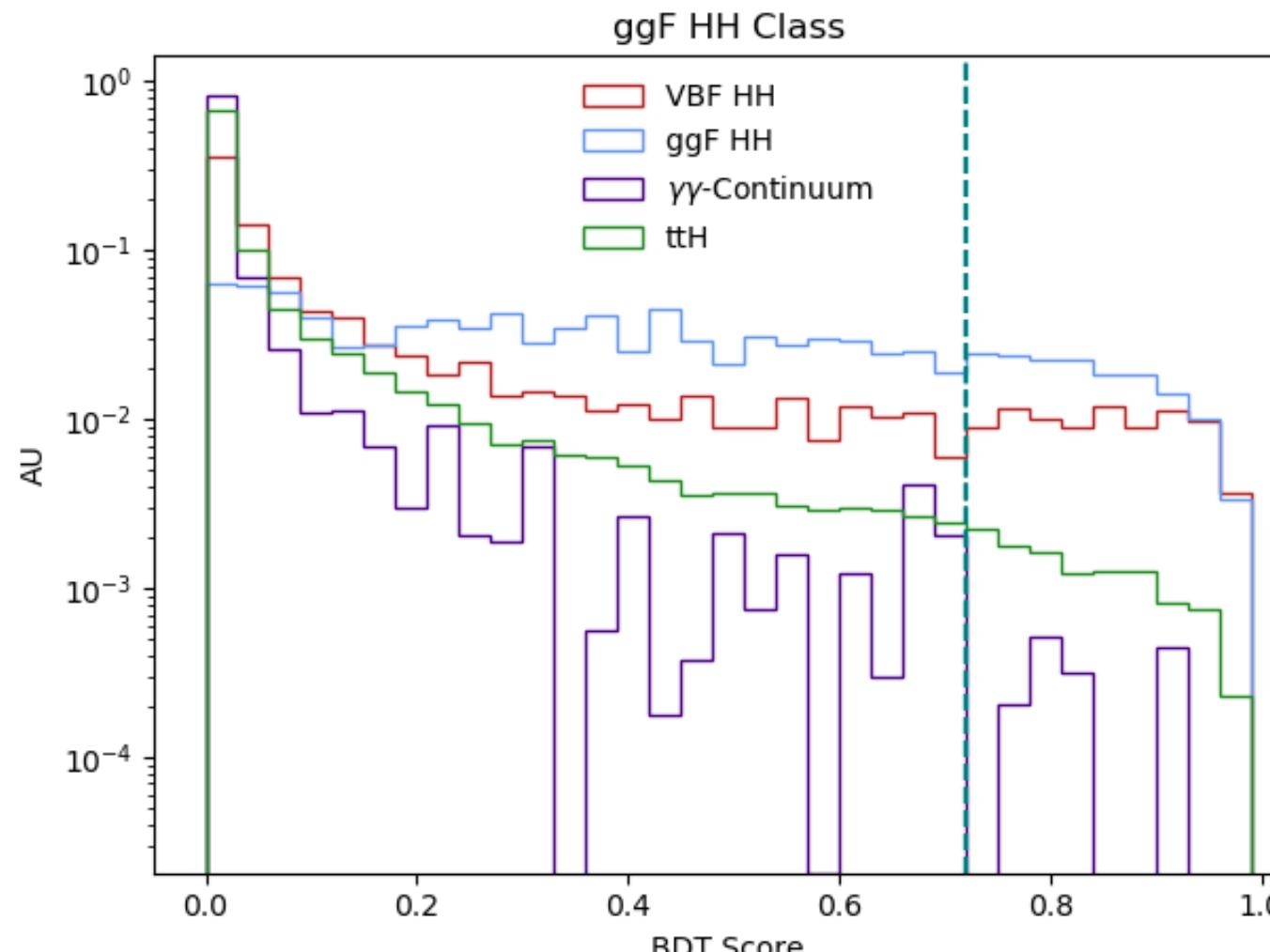
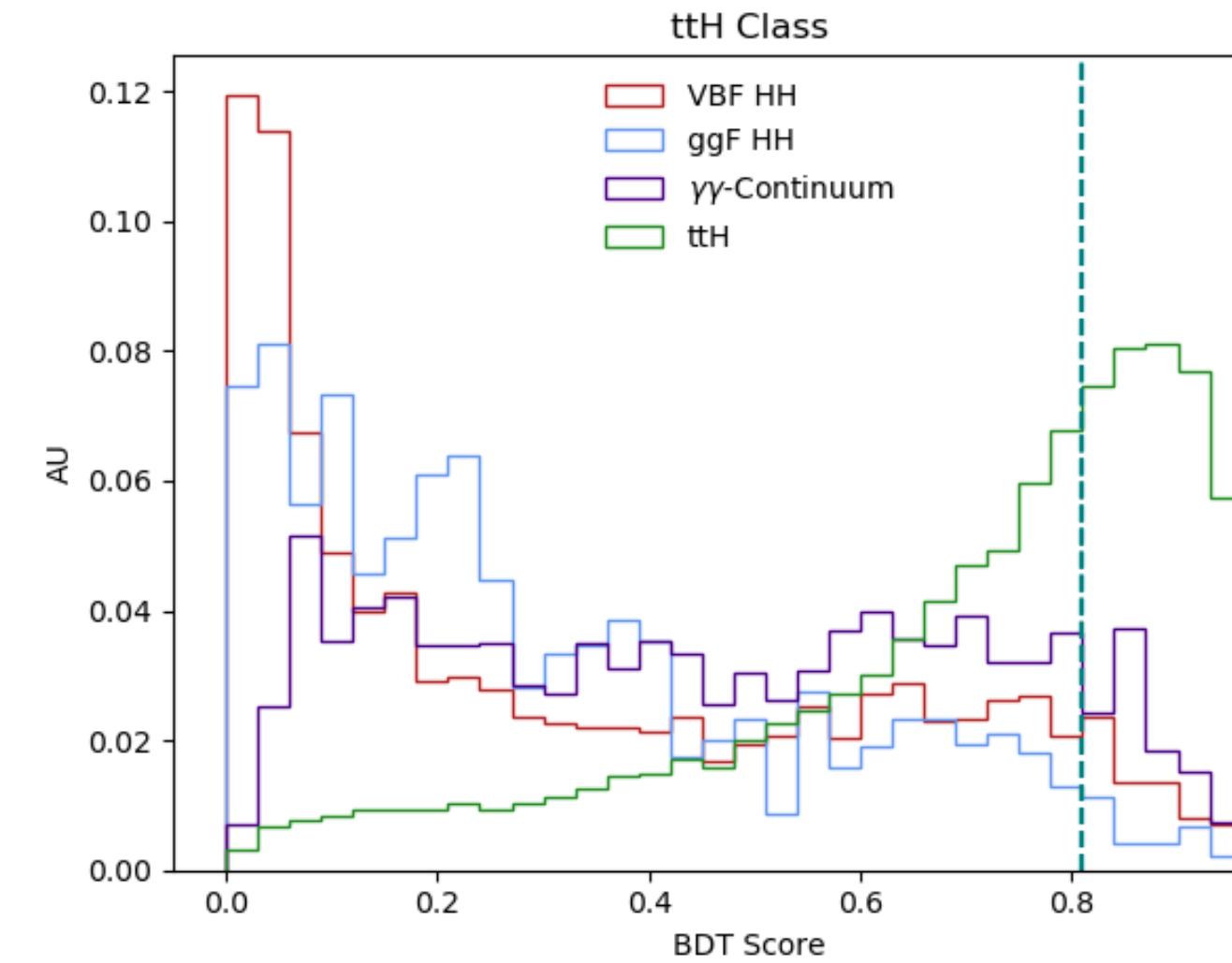
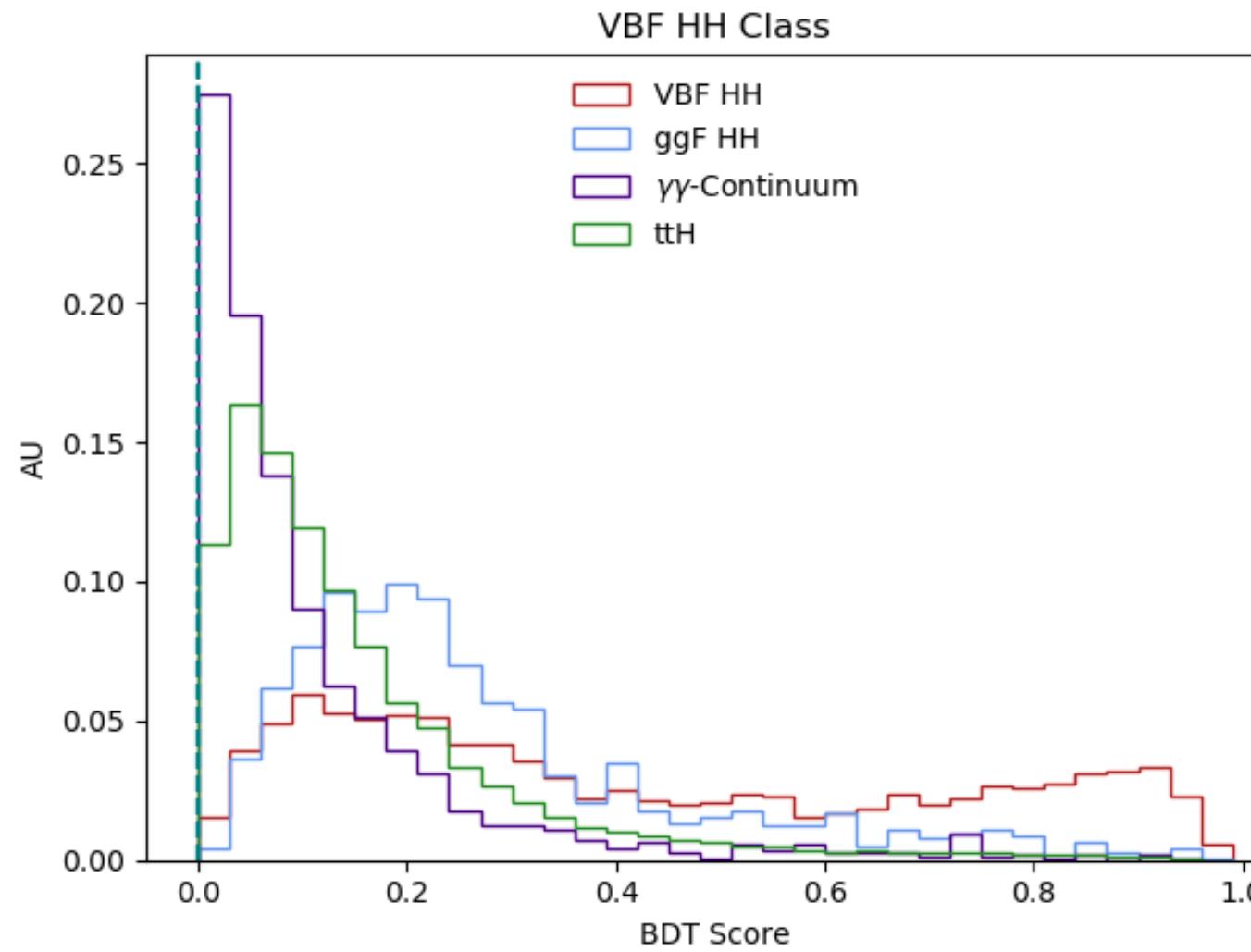
Use differences in Higgs decay kinematics and b-tagging to differentiate from dominant background



# BDT Optimization

Probability cuts to define category selected via a 4-d scan over the 4 BDT scores.

- On each step, calculate the combined Asimov significance, and select highest value



Asimov significance given by

$$Z = \sqrt{2[(s + b)\ln(1 + s/b) - s]}$$

At each step, calculate value for each category, add in quadrature (sensitivity to full HH cross section)

$$Z_{\text{total}} = \sqrt{Z_{ggF_{cat1}}^2 + Z_{ggF_{cat2}}^2 + Z_{ggF_{cat3}}^2 + Z_{ggF_{cat4}}^2 + Z_{VBF}^2}$$

selecting most optimum point gives

$$Z = 0.459$$

which is a 9.7% improvement over not defining such a category ( $Z=0.415$ )



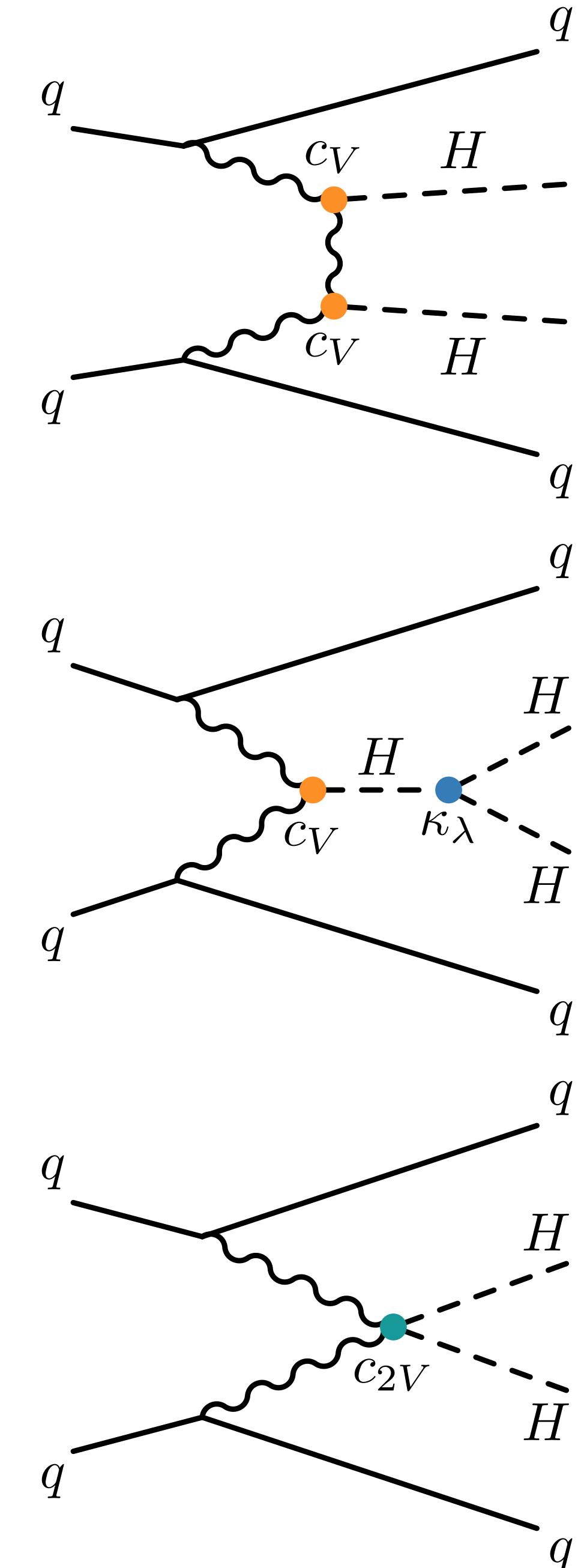
# VBF HH Conclusion

**Studying VBF HH production provides interesting insights into di-Higgs production**

- VBF HH production is characterized by a unique topology, good for signal/background discrimination
- Studying VBF HH production provides sensitivity to new couplings, specifically  $c_{2V}$  (VHH vertex), and deviations from SM prediction can cause large enhancements

**The upcoming  $\text{HH} \rightarrow \gamma\gamma\text{bb}$  analysis will incorporate a VBF-enriched signal region**

- A multiclassifier BDT is used to separate VBF HH production from ggF HH as well as  $\gamma\gamma$ -continuum and ttH backgrounds
- After optimizing cuts on 4 probability scores, an improvement of 9.7% in Asimov significance can be found by adding this category



# Photon Identification

$\gamma\gamma bb$  is sensitive to the square of the photon identification efficiency -  
motivates searching for improvements in photon ID



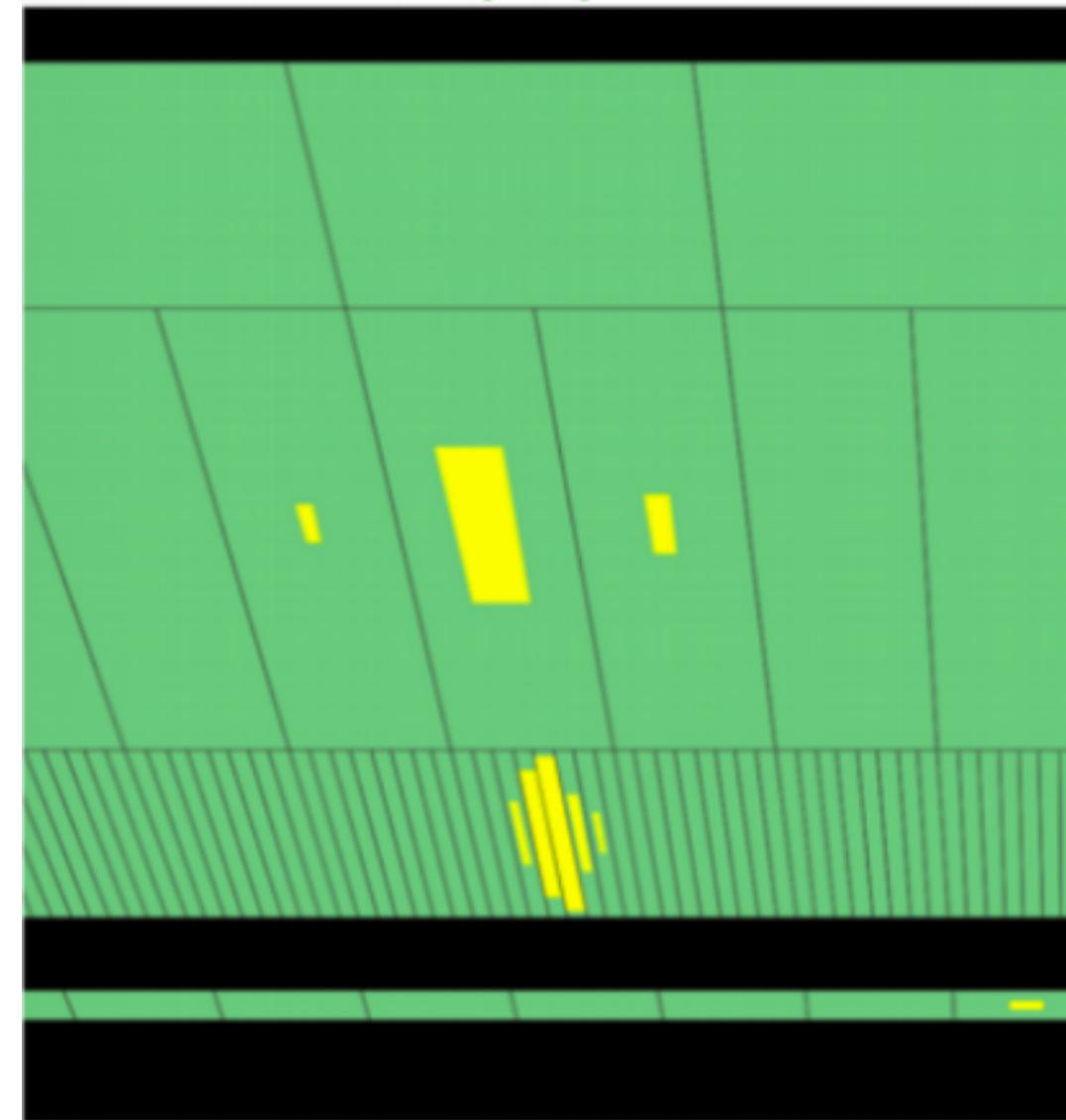
# Photon Identification in ATLAS

**Goal:** Classify energy depositions in the EM calorimeter

- Identify “prompt” photons - come from colliding protons, process of interest
- Reject photons resulting from background processes (e.g. hadronic decays in jets)

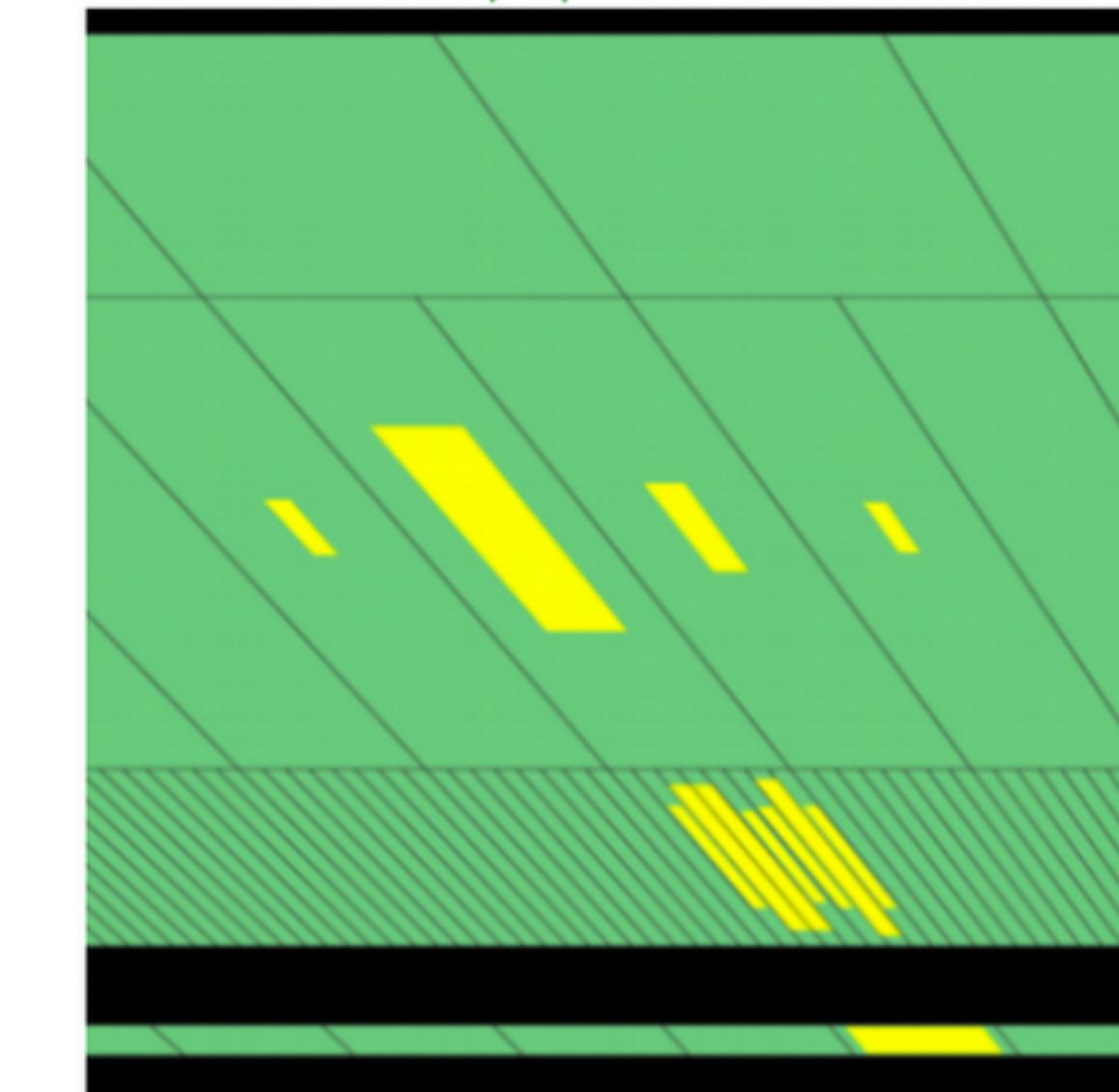
Want to  
discriminate

Prompt photon ( $\gamma$ )



$\pi^0 \rightarrow \gamma\gamma$

from



# Current Methodology

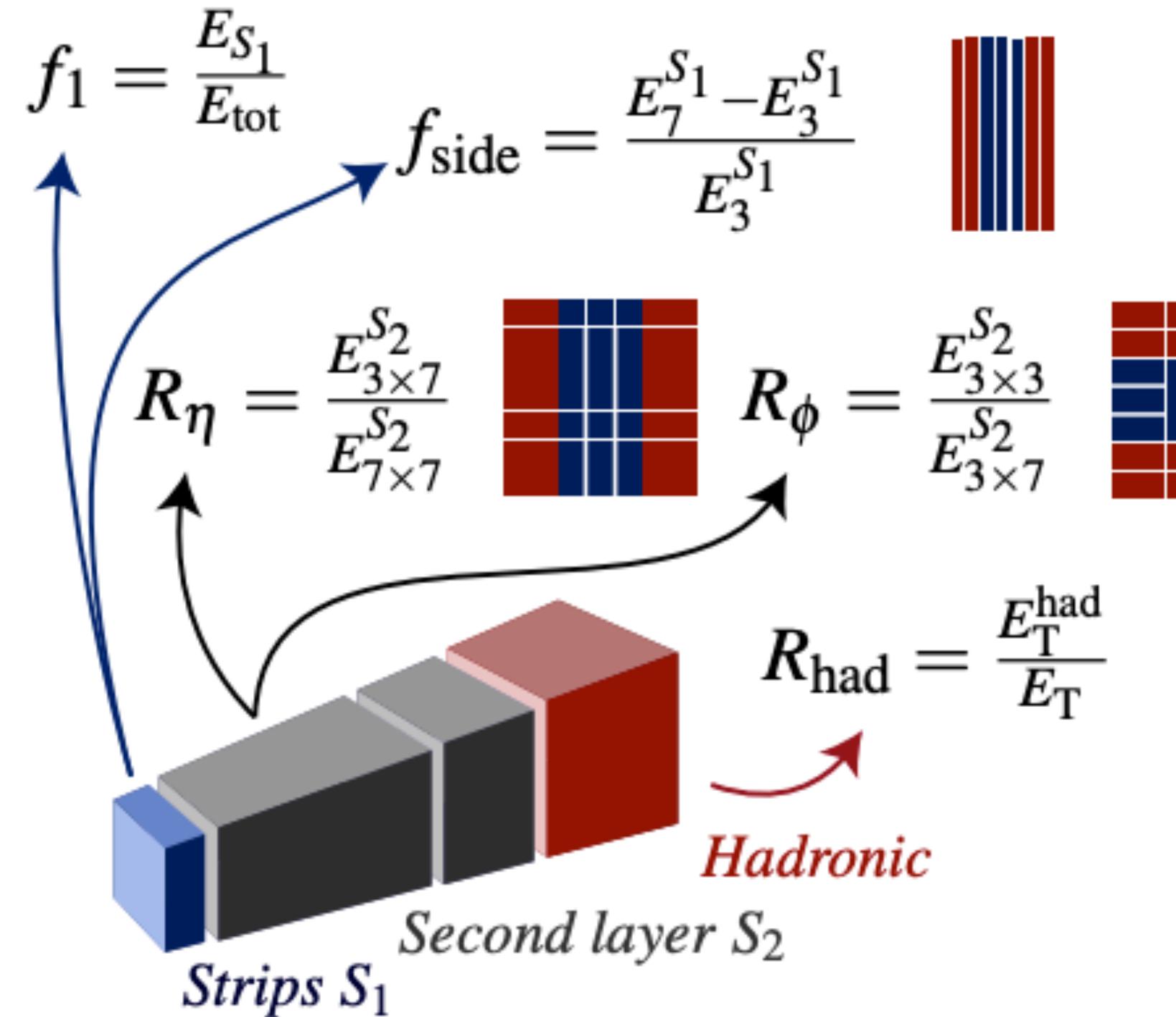
**Showering** - Description of particle interaction with detector material

- For photons, this is successive conversion to electron-positron pairs and bremsstrahlung photon emission

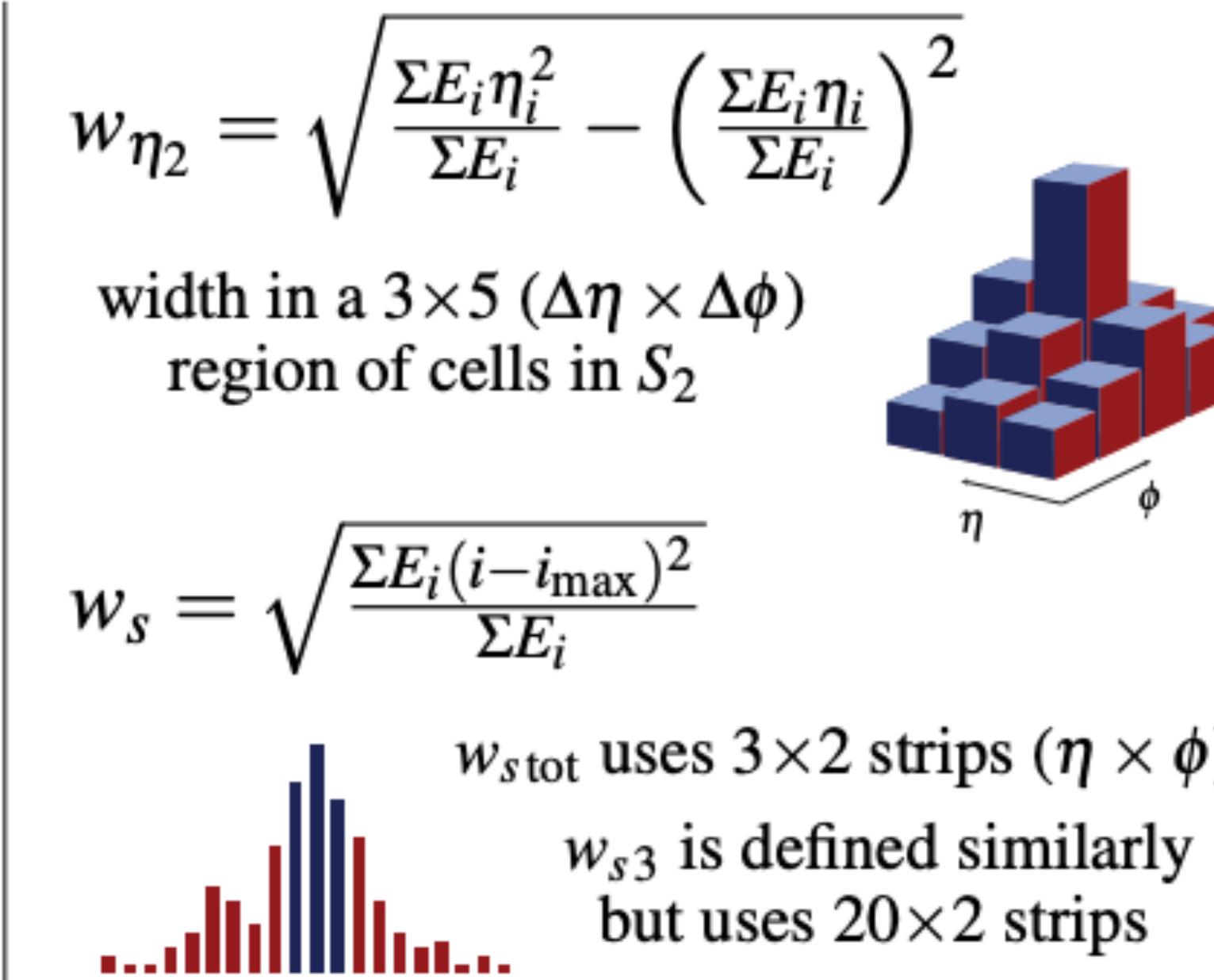
**Cutting** on shower shape variables

- 9 variables that describe longitudinal and lateral shower development in the calorimeter
- Showering dependent on detector segmentation - **bin in  $\eta$ ,  $p_T$ , and conversion status**

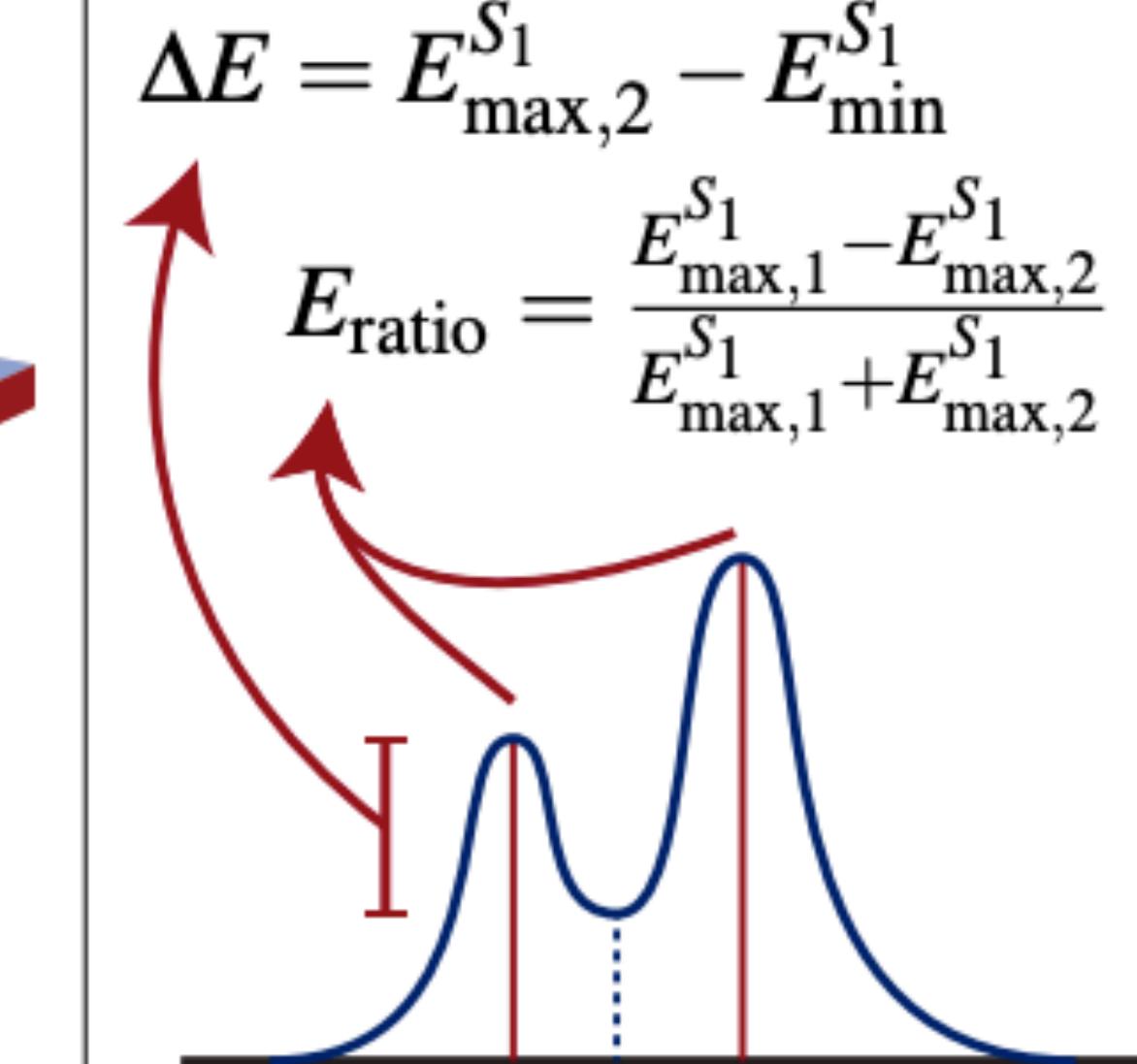
## Shower Shapes



## Widths



## Energy Ratios



**How can we improve this?**

New Variables

Better algorithm to separate signal/background

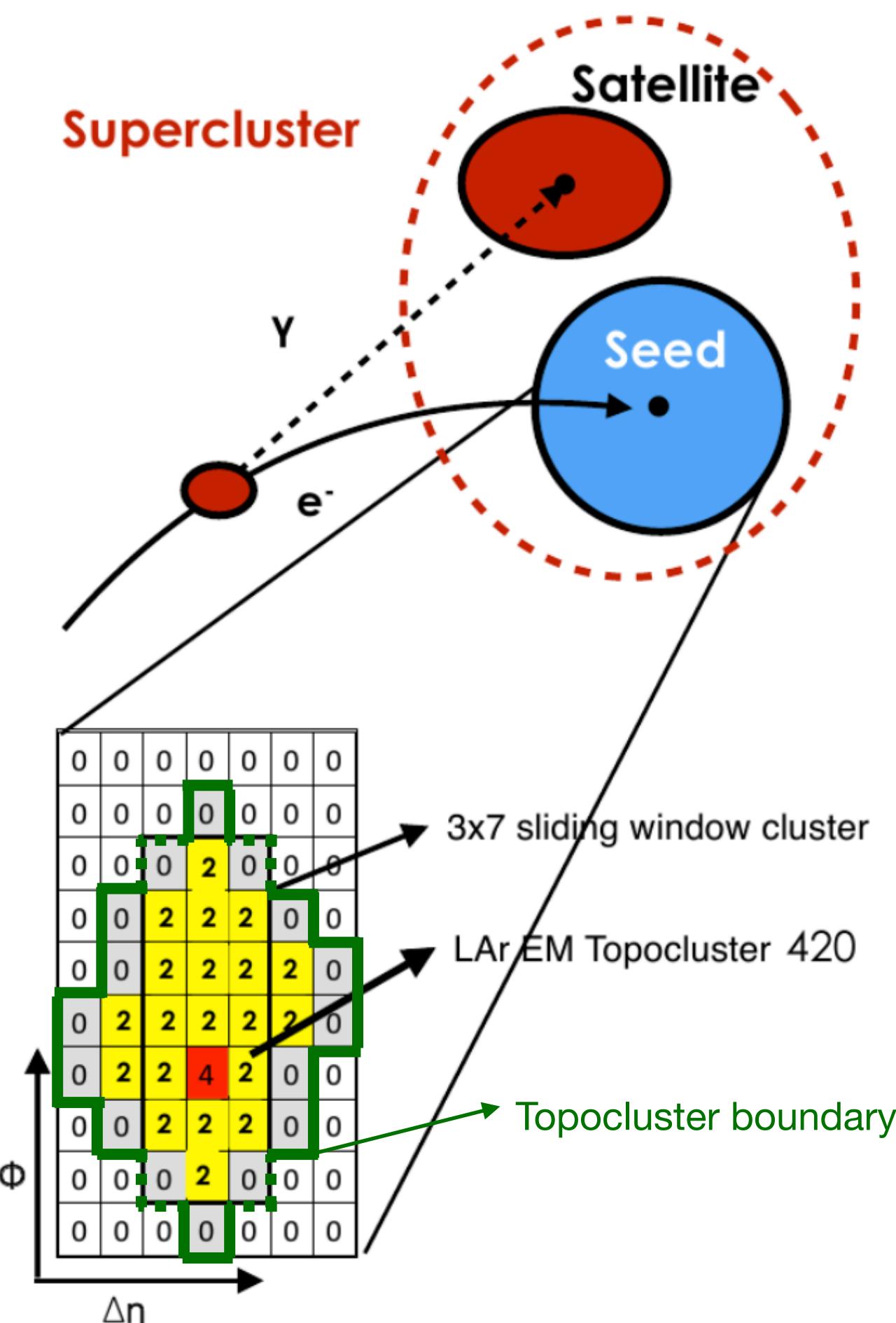
# Topological Clusters for Photon ID

## Topological Clusters (topo-clusters)

- Photons in ATLAS are defined from objects known as “superclusters”
  - These are groupings of topological clusters, made through the logic shown
  - Studied various topo-cluster observables (moments) as inputs to photon ID algorithm to see if they can bring additional discriminating power

## Samples Used

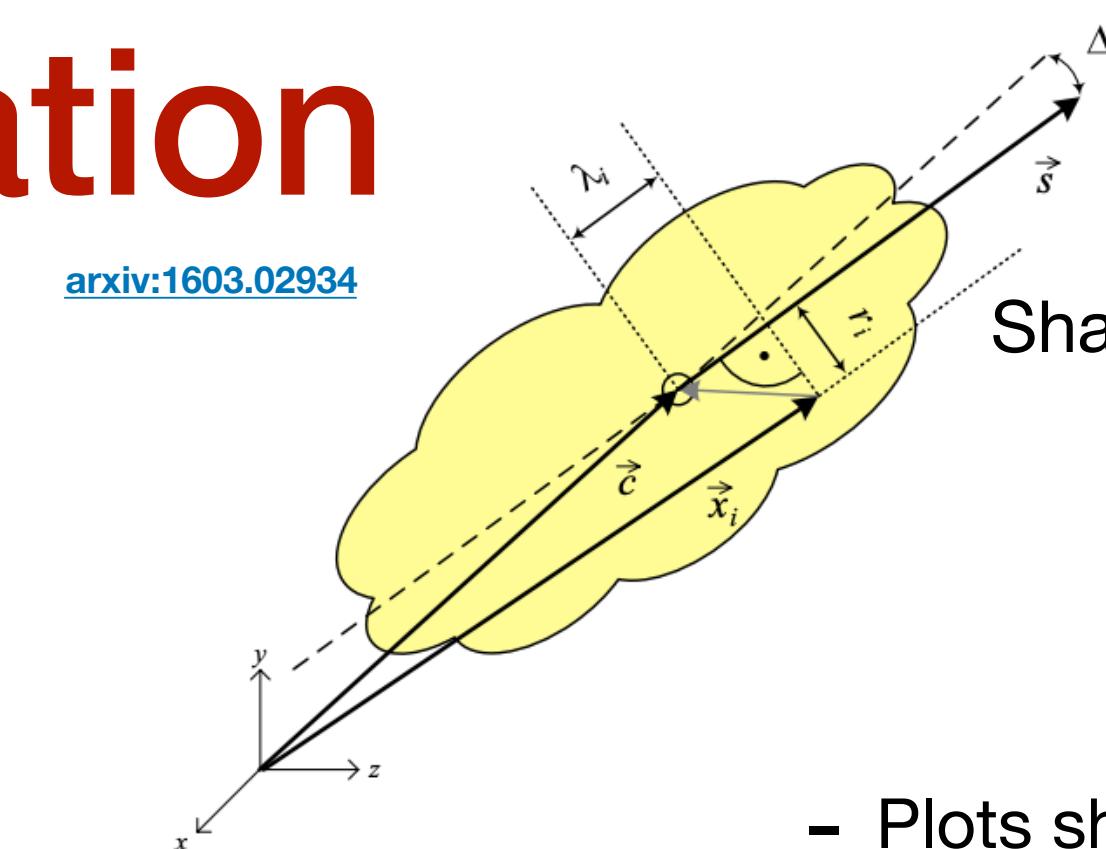
- $\gamma + \text{jets}$  signal, jet-fakes background
- Monte Carlo truth information used to require signal sample be a prompt photon, vetoed for background sample



## Formed by:

Find “seed” - cell significance  $> 4\sigma$   
Scan neighbors. If  $> 2\sigma$ , add cell and neighbors to cluster

# Topo-Cluster Shape Information



$$r_i = |(\vec{x} - \vec{c}) \times \vec{s}|$$

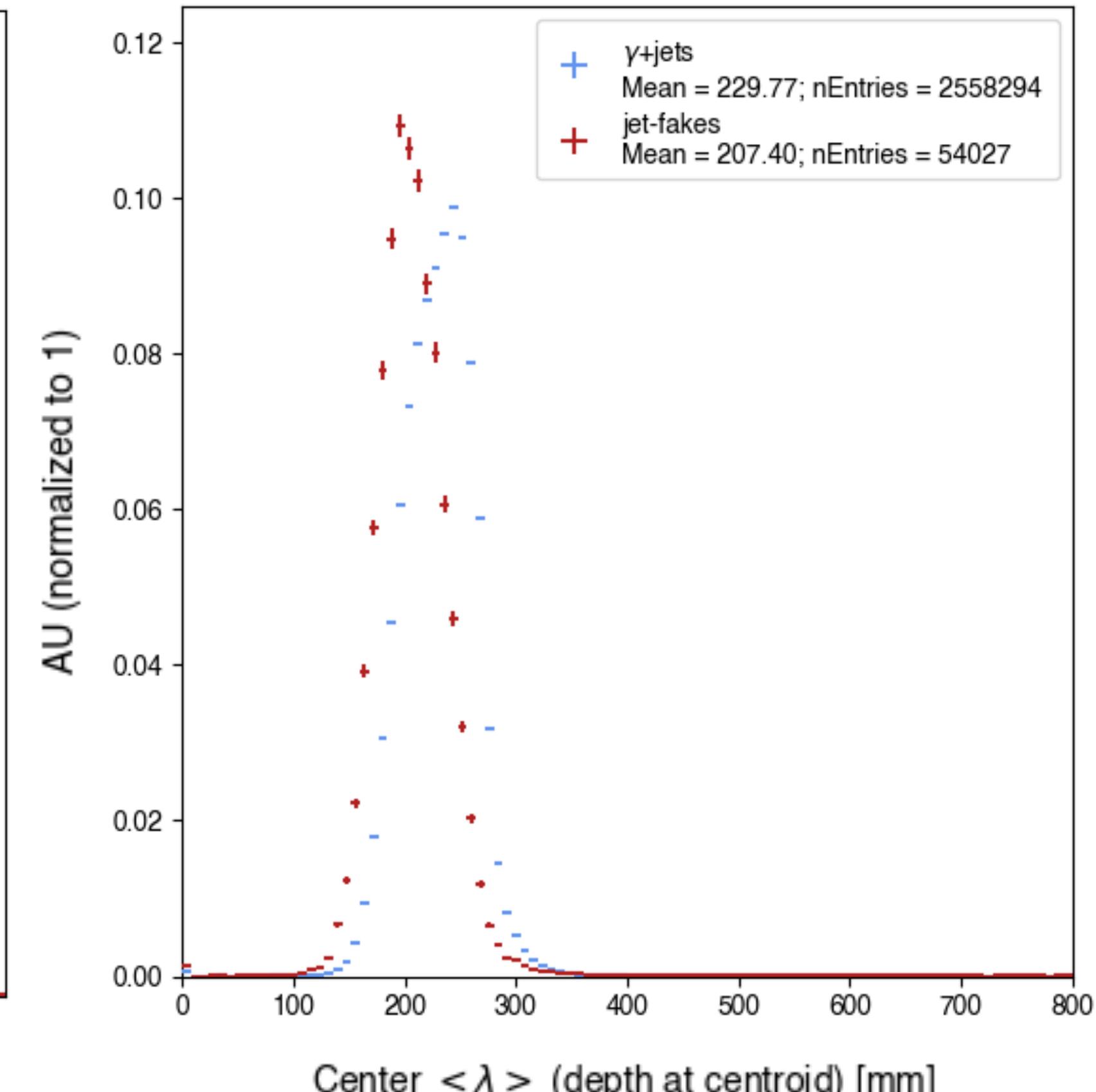
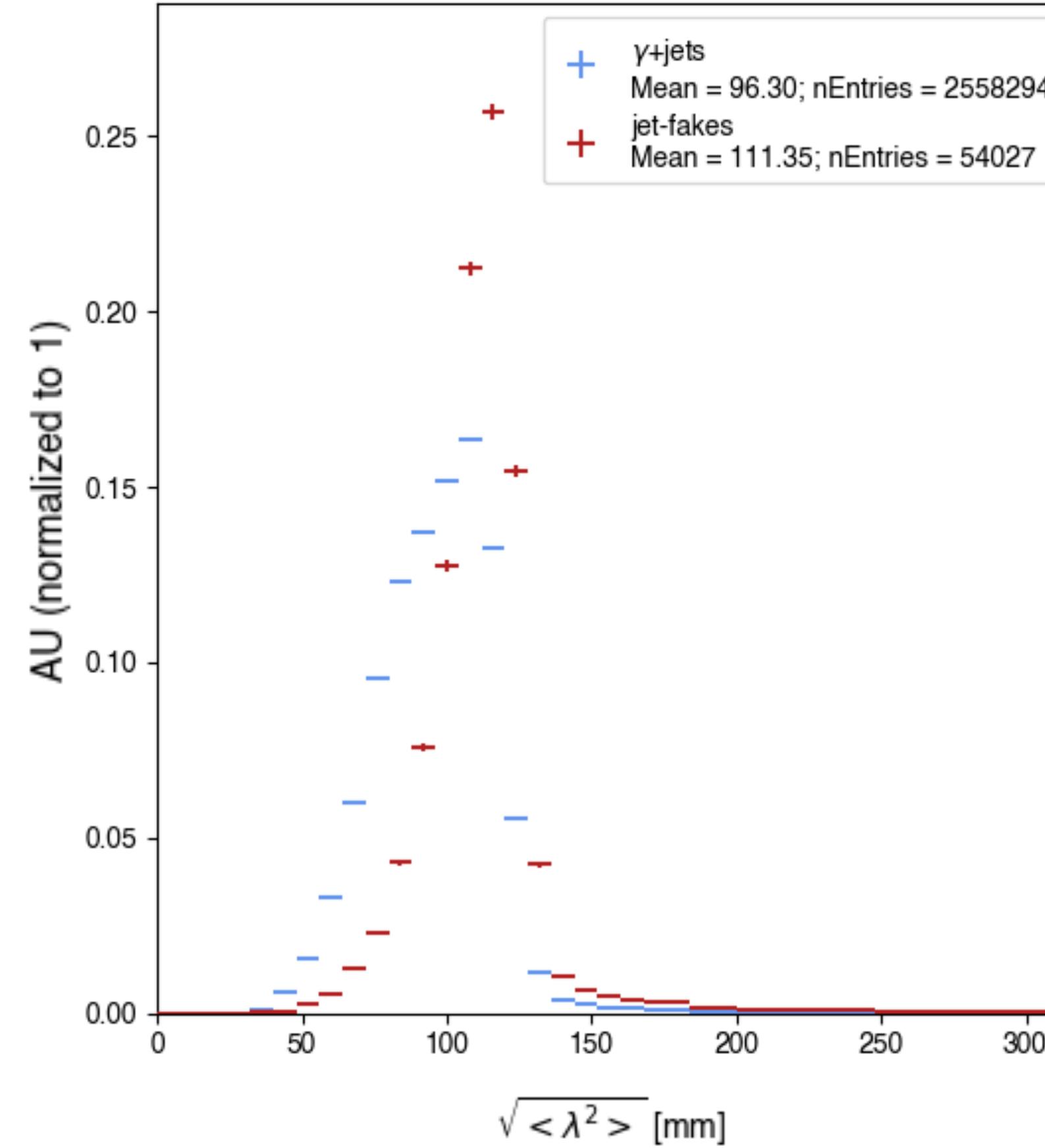
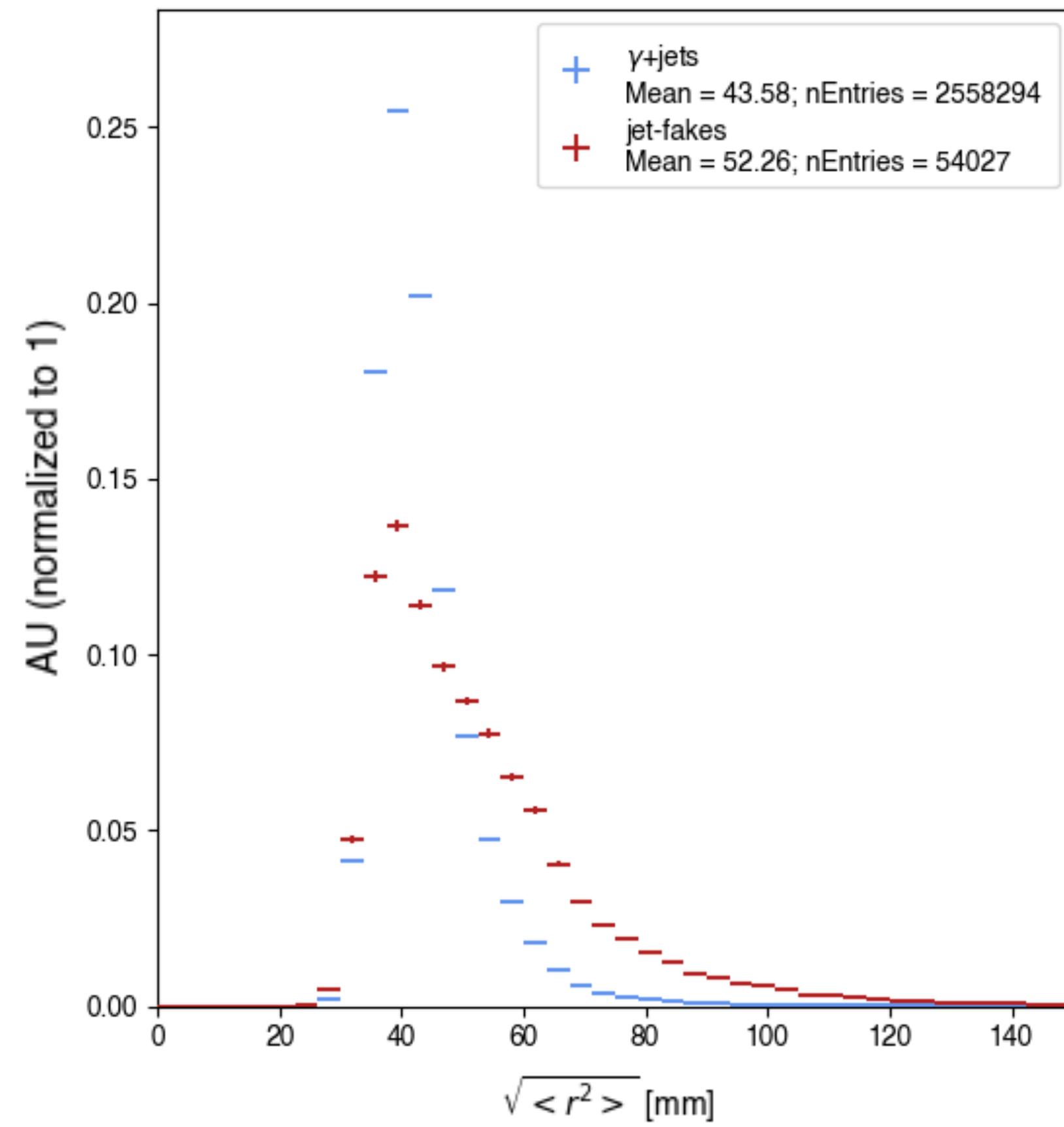
$$\lambda_i = (\vec{x}_i - \vec{c}) \cdot \vec{s}$$

Shape can be defined using these quantities:

$\sqrt{\langle r^2 \rangle}$  = **Semi-major axis in width**

$\sqrt{\langle \lambda^2 \rangle}$  = **Semi-major axis in depth**

- Plots shown are in  $|\eta| < 0.6$  region, pT inclusive



Differences exist in topo-cluster shape after current tight ID - additional information that can be used to further separate signal and background

\*current Tight ID applied

# Improvement by Adding Topo-clusters

Run existing cut-based optimization, adding the topo-cluster variables shown in prior slide

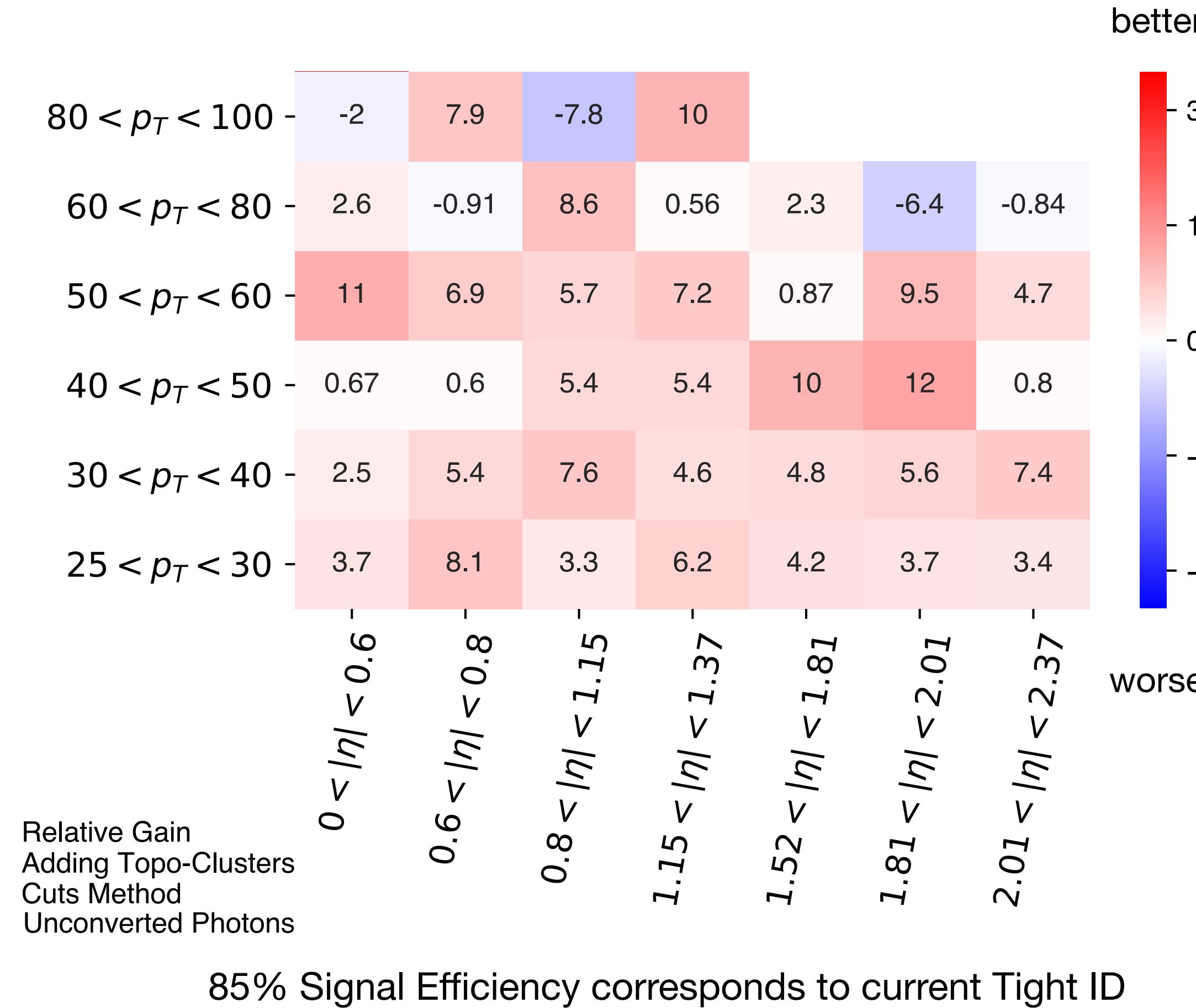


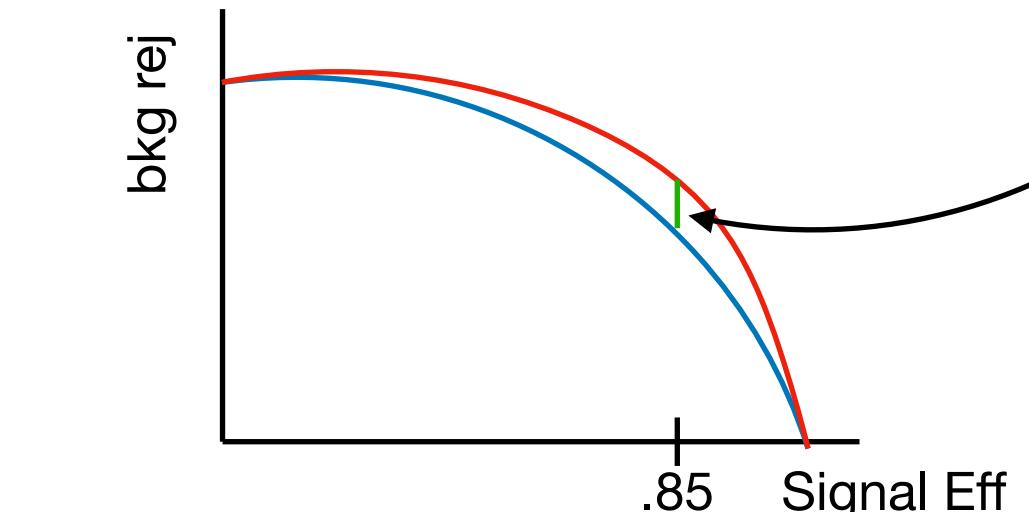
Figure of merit - improvement in background rejection at same signal efficiency as the current “tight” identification (85%)

Values shown:

$$Z_{rel} = \frac{\% Rej_{bkg,topo} - \% Rej_{bkg,SS-only}}{\% Rej_{bkg,SS-only}}$$

Equivalently,

$$Z_{rel} = \frac{\epsilon_{bkg,SS-only} - \epsilon_{bkg,topo}}{1 - \epsilon_{bkg,SS-only}}$$



Adding topo-cluster information can add **up to an additional ~10% background rejection** for the same signal efficiency using the current cut based method



# Machine Learning Approaches

Current algorithm is a 9-dimensional cut optimization - could we consider a more sophisticated algorithm?

**Machine learning** - Two models considered:

## Boosted Decision Tree

- Explained earlier! Ensemble method following tree logic
- Developed using TMVA, library for MVA in ROOT

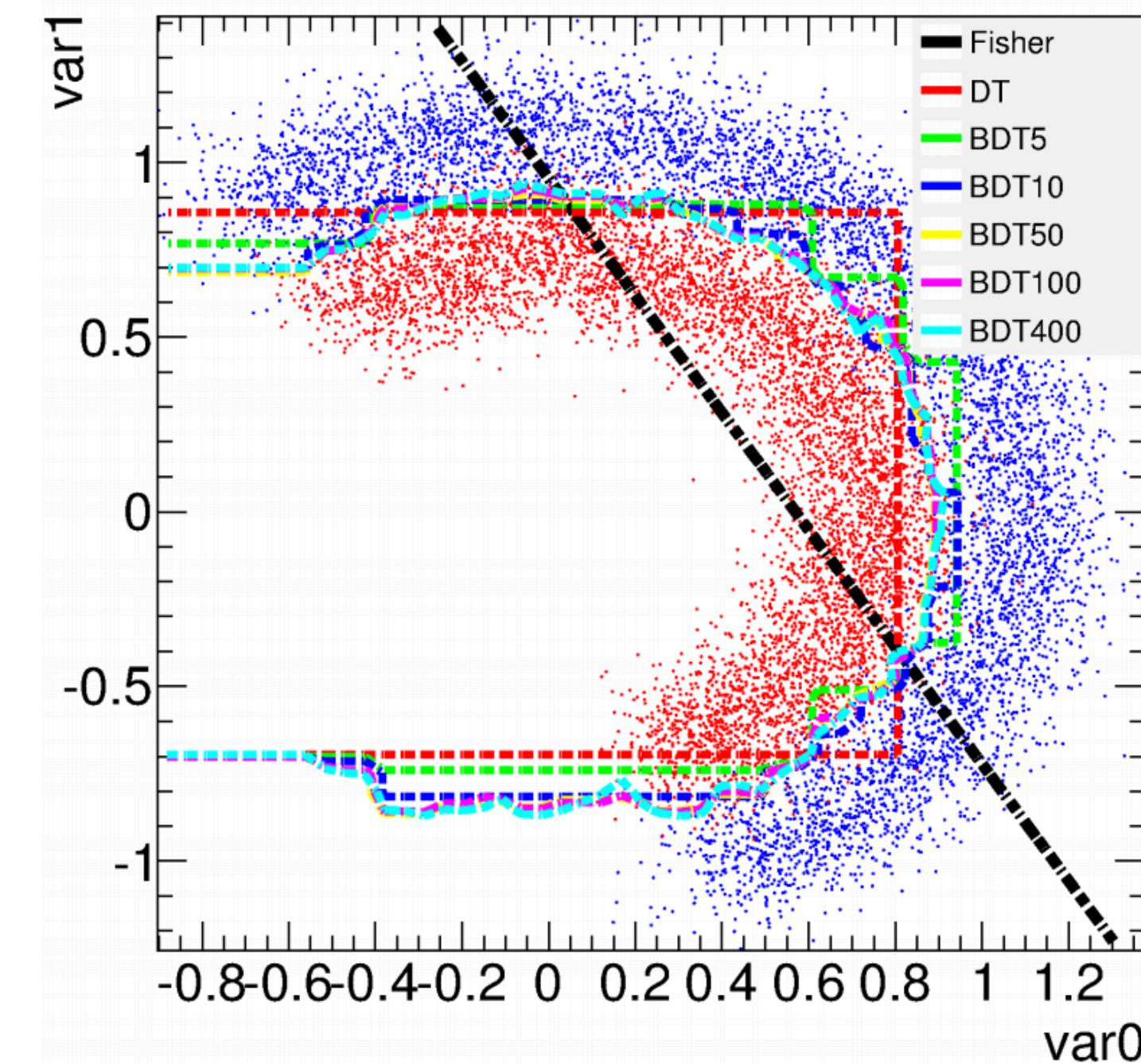
## Neural Network

- More flexible, can model more arbitrary functions
- Developed in Keras using Tensorflow backend

MVA Advantages	MVA Disadvantages
----------------	-------------------

- More flexible solution space compared to cuts
- Can exploit non-linearities

- Loss of clearly defined cuts, transparency



Of course, this is a toy example - cuts could do just as well with a variable transformation!



# Boosted Decision Tree

Evaluated improvement replacing cut-based method used presently with a BDT

While training, found statistics to be problematic in high  $p_T$ , high  $\eta$  bins

- Navigated by training *inclusively* in  $p_T$ . Variables uncorrelated to  $p_T$ , and provided better performance in all bins

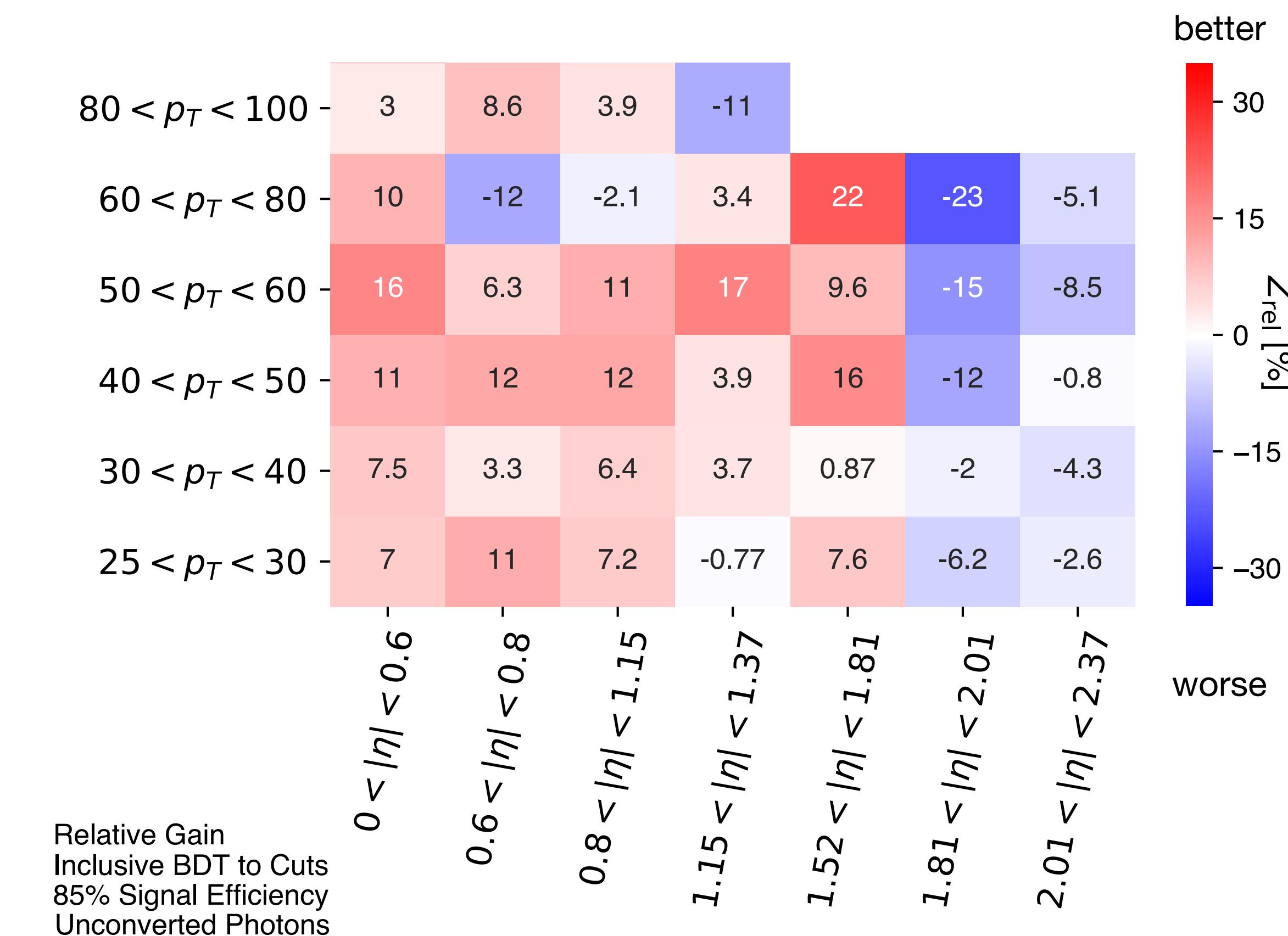
Same FOM used:

$$Z_{rel} = \frac{\% Rej_{BDT} - \% Rej_{cuts}}{\% Rej_{cuts}}$$

Only current shower-shape variables as input

Bins in which  $Z$  is worse with the BDT can be interpreted as the cut-method being better

- Limited statistics



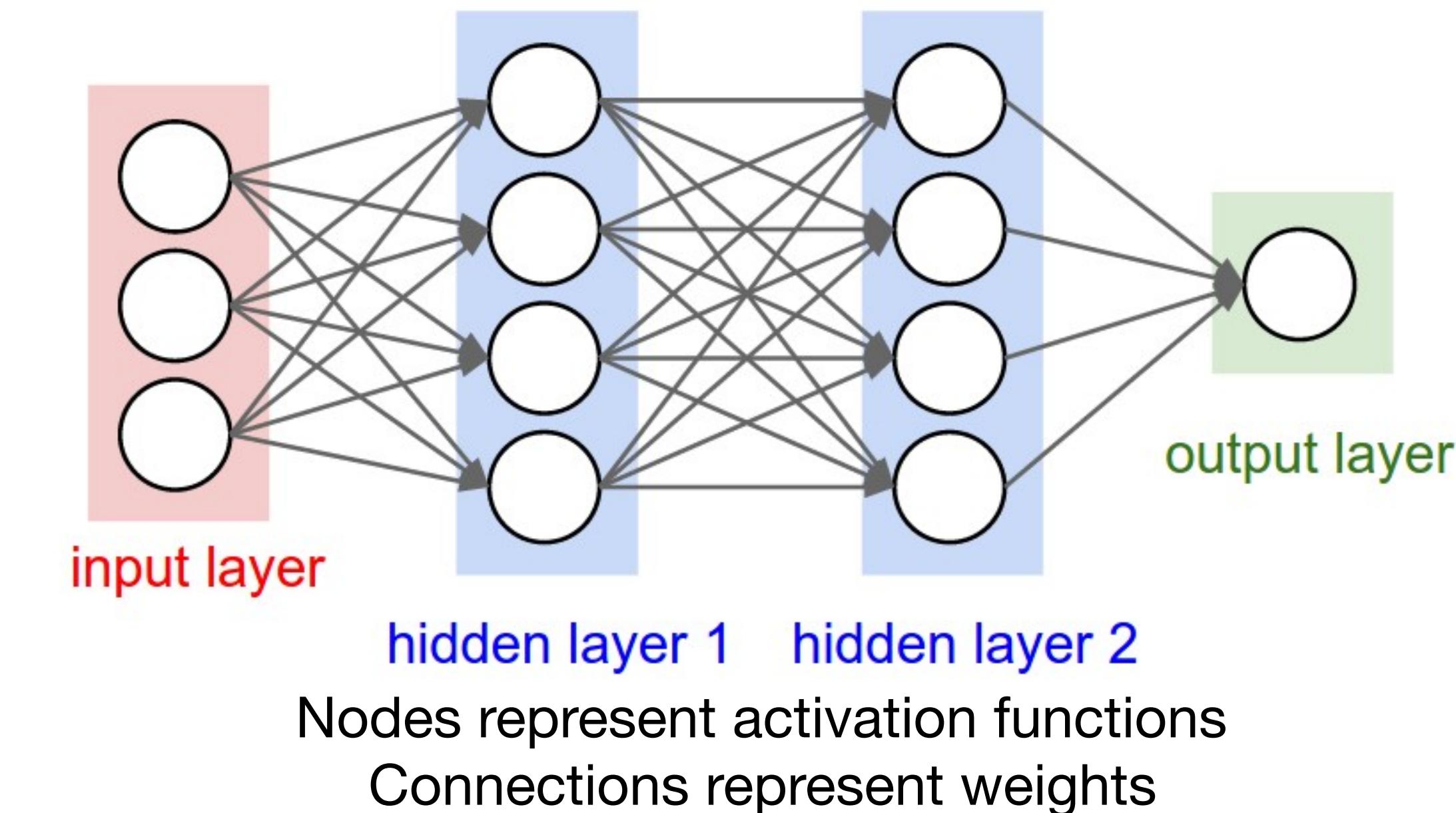
# Machine Learning Interlude #2

## Neural Networks

Successive layers of nonlinear basis functions (*activation* functions), connected by weights

- Training stage learns the values of weights and the parameters
- In standard NNs (used in this work), data proceeds strictly from input to output (*feed forward*)
- Name comes from neuroscience analogy, where activation functions are analogous to neurons
- Over a sufficient number of activation functions, any possible function can be approximated

[Andrej Karpathy](#)



### Notes

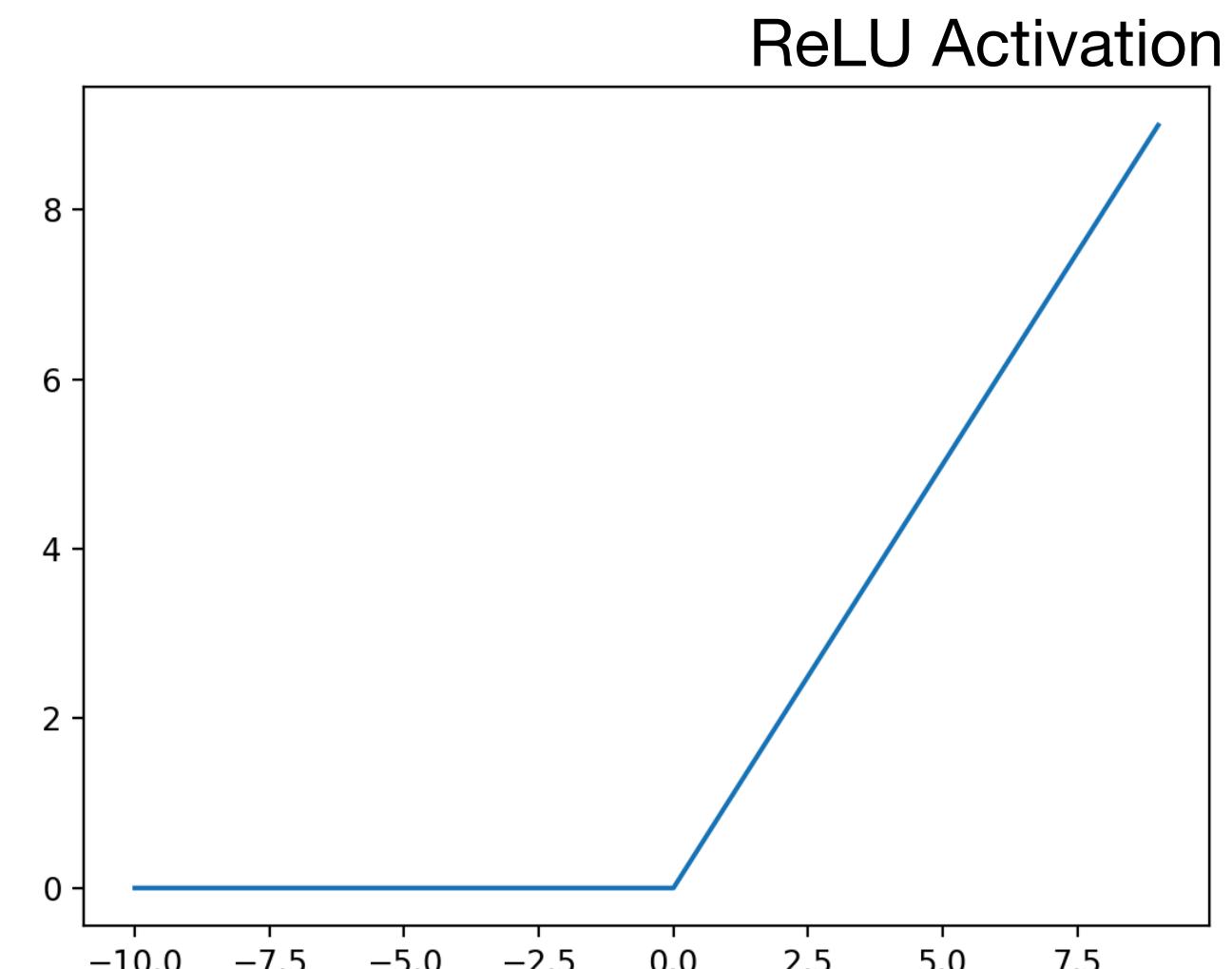
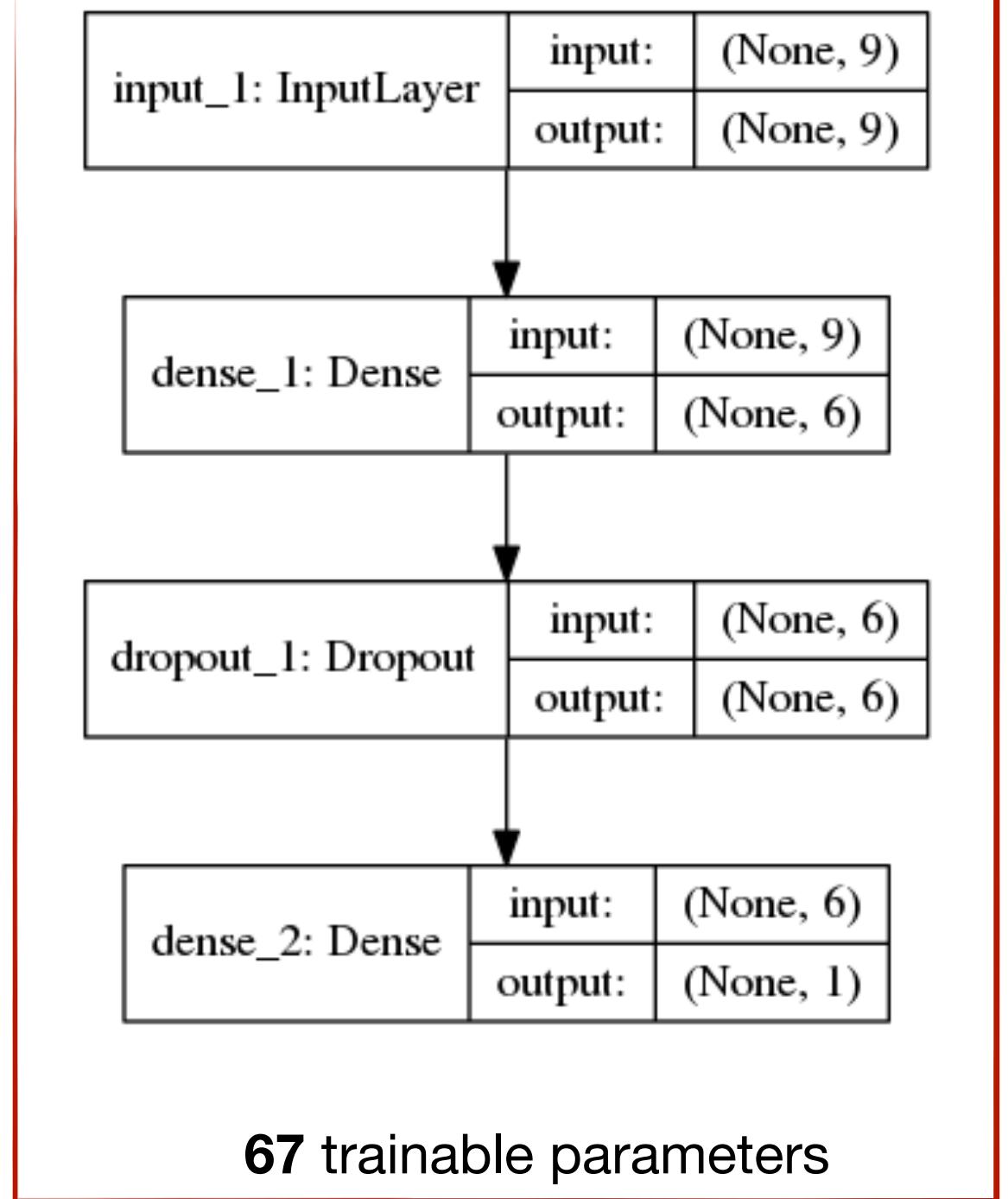
- Incredibly flexible algorithm!
- Requires *significant* training data



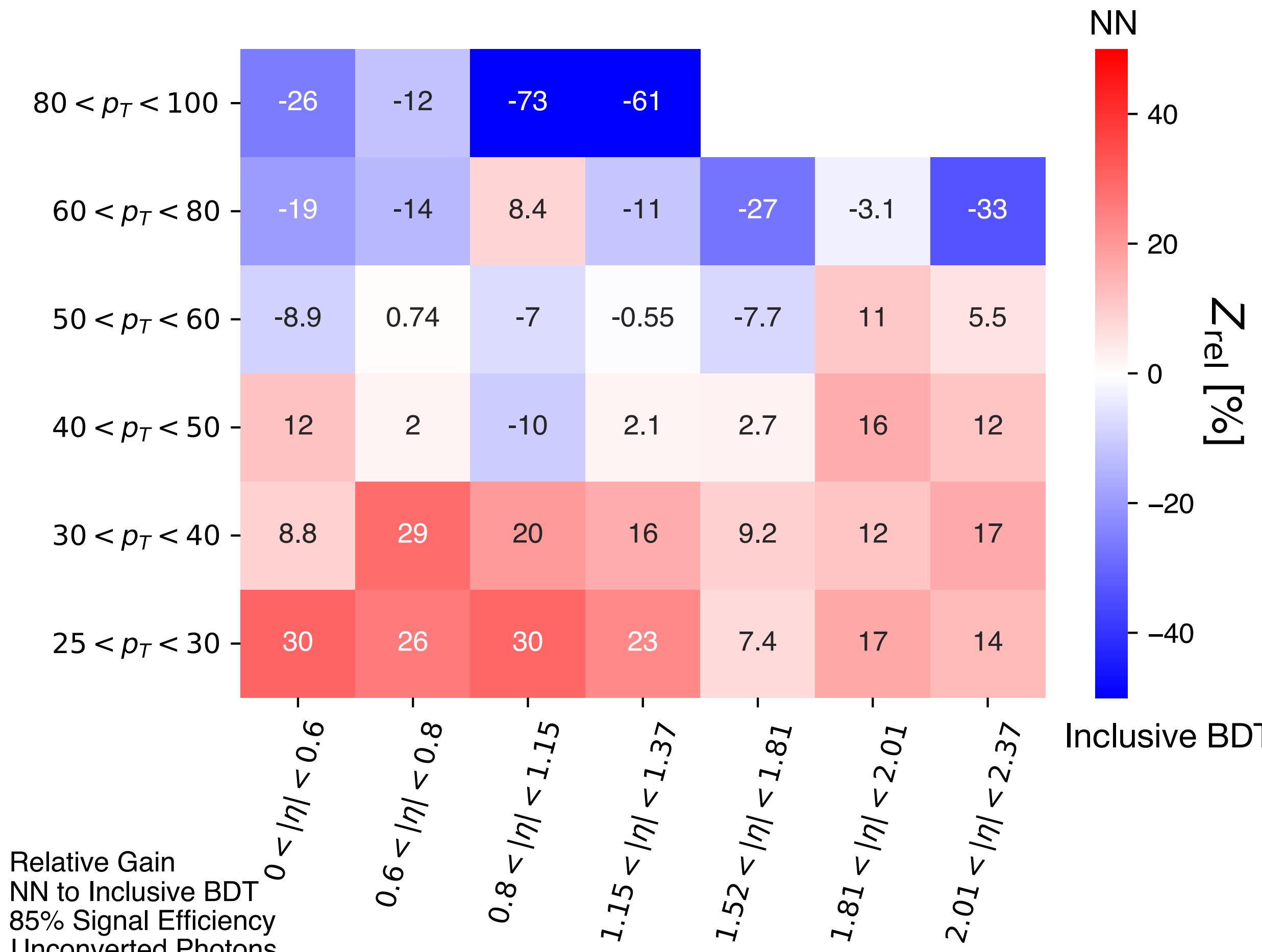
# NN Training Notes

## Topology and notes

- Constructed in Keras, topology shown, selected by a course optimization by hand
  - Very simple topology, few parameters compared to modern networks
  - Deeper/wider networks tested but found no additional improvement
- ReLU activation used for internal nodes, sigmoid on output
- Binary Cross-Entropy loss
- Epochs dictated by early stopping - stops after no improvement for 20 epochs



# Neural Network Results



$$\frac{Rej_{bkg,NN} - Rej_{bkg,BDT}}{Rej_{bkg,BDT}}$$

Again, can think of negative values as BDT performing better, positive as NN performing better

NN trained binwise does better than BDT in the lowest pT regions

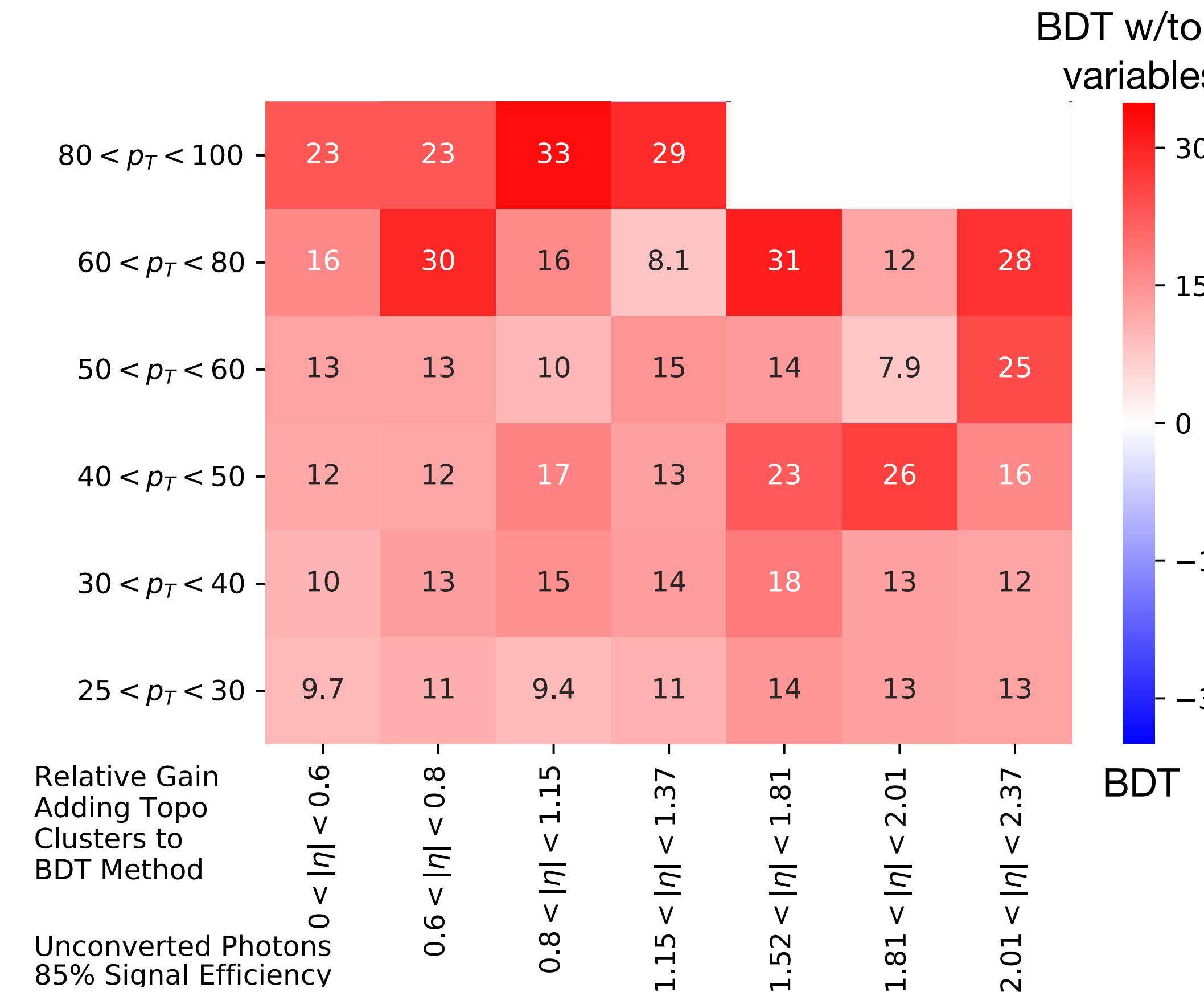
Inclusively trained NN was evaluated, but significant gains only seen at  $p_T > 80$  GeV, and much worse performance seen below



# Putting Improvements Together

Want to also consider how the algorithms can improve by including the aforementioned topo-cluster variables

- Shown for test case using a BDT

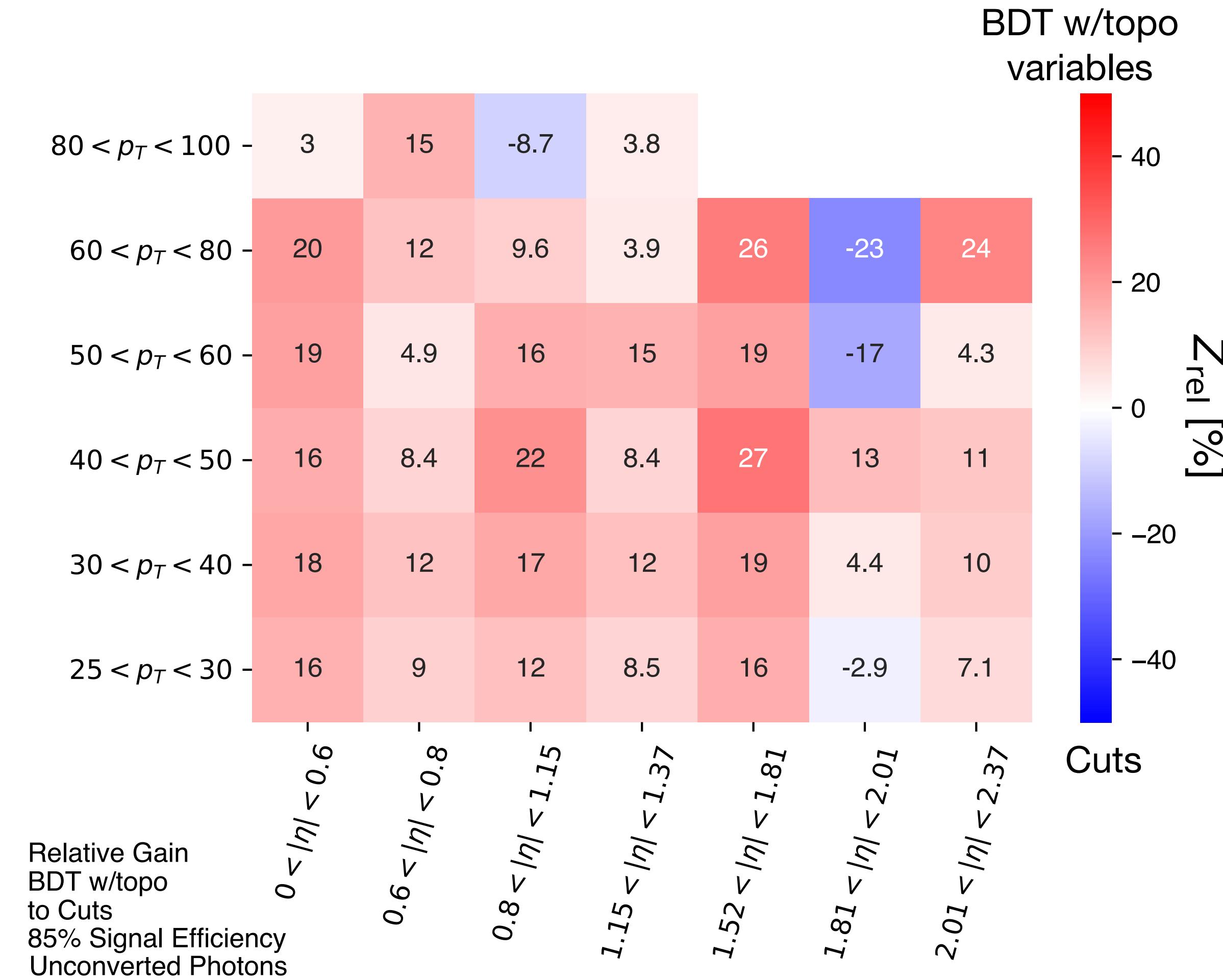


As much as a 30% improvement in the BDT by including the topo-cluster variables



# Putting Improvements Together

Last, as a final metric of improvement, compare the BDT including the topo-cluster variables to the present method



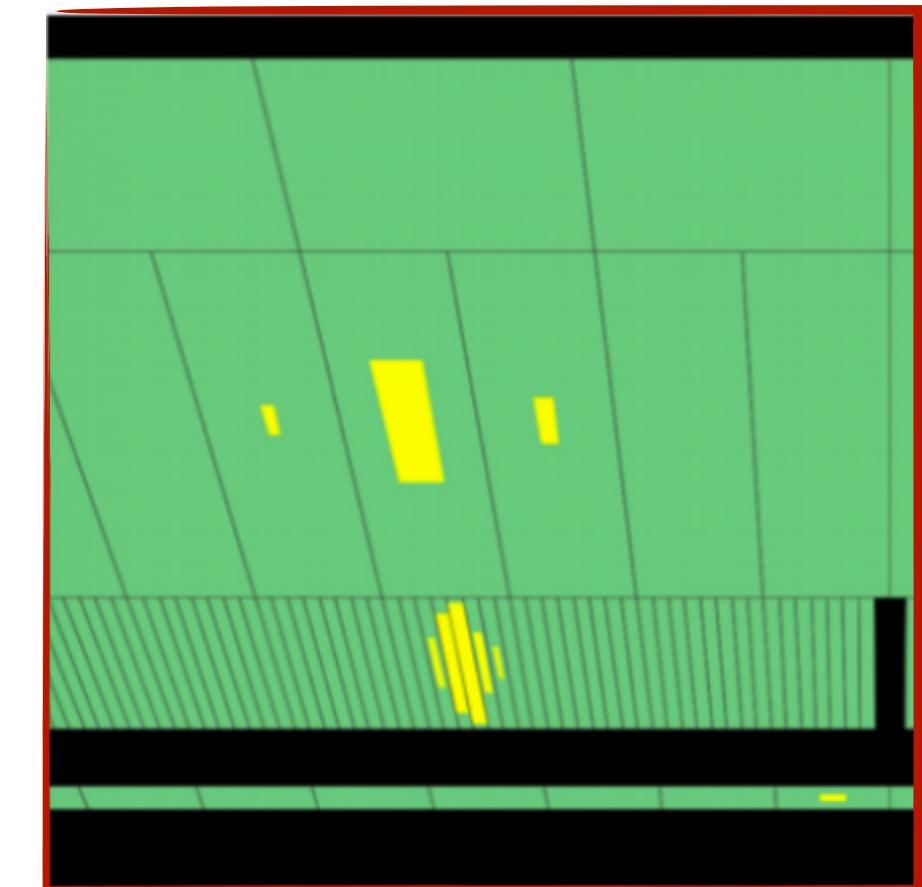
**As much as a 25% gain through all improvements compared to current methods**  
Most bins show an improvement of > 10%



# Photon Identification Conclusion

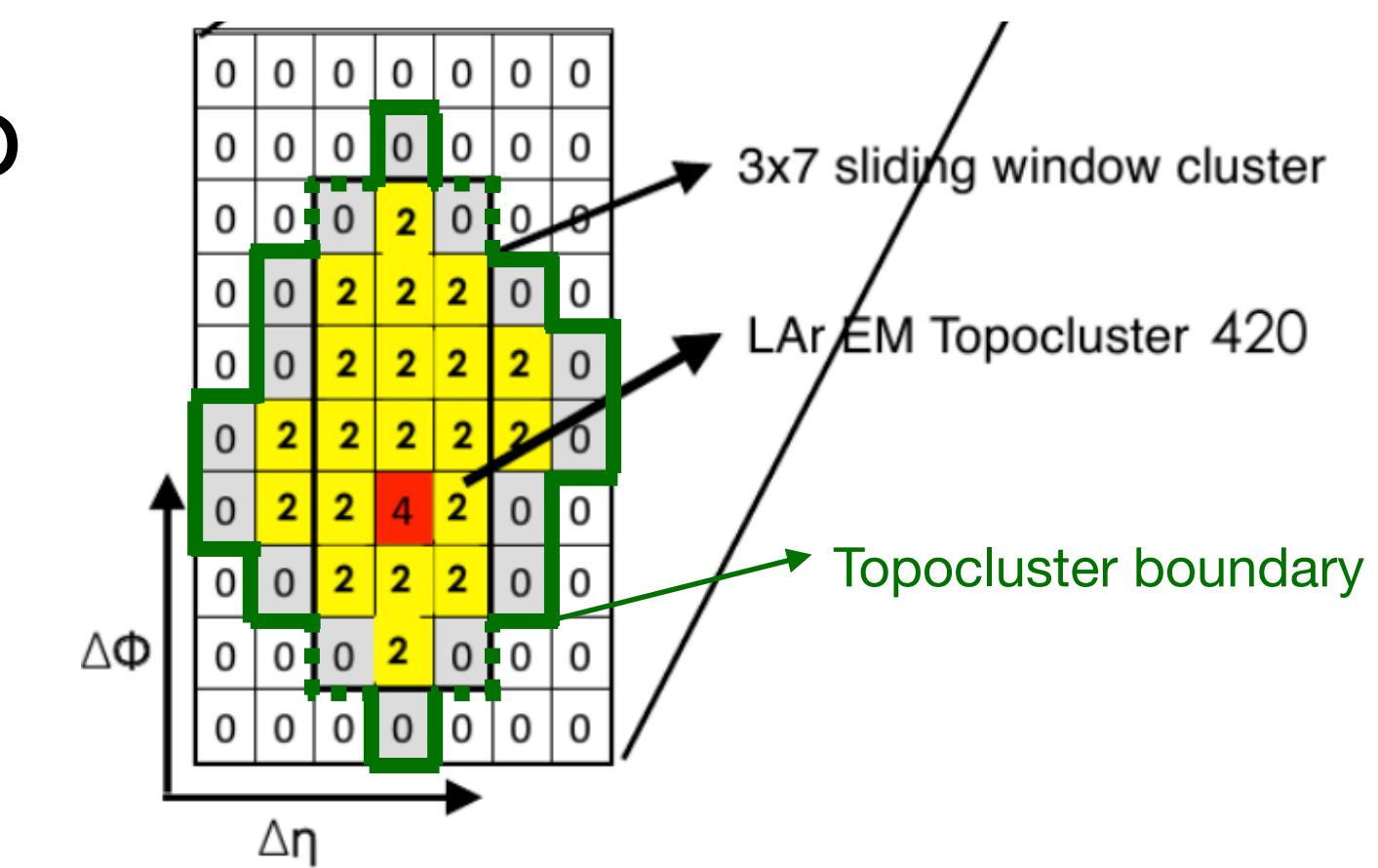
The  $\text{HH} \rightarrow \gamma\gamma\text{bb}$  analysis is sensitive to the square of the photon identification efficiency, motivates searching for improvements

Methods to optimize photon identification studied, using improvement in background at current tight WP as a FOM



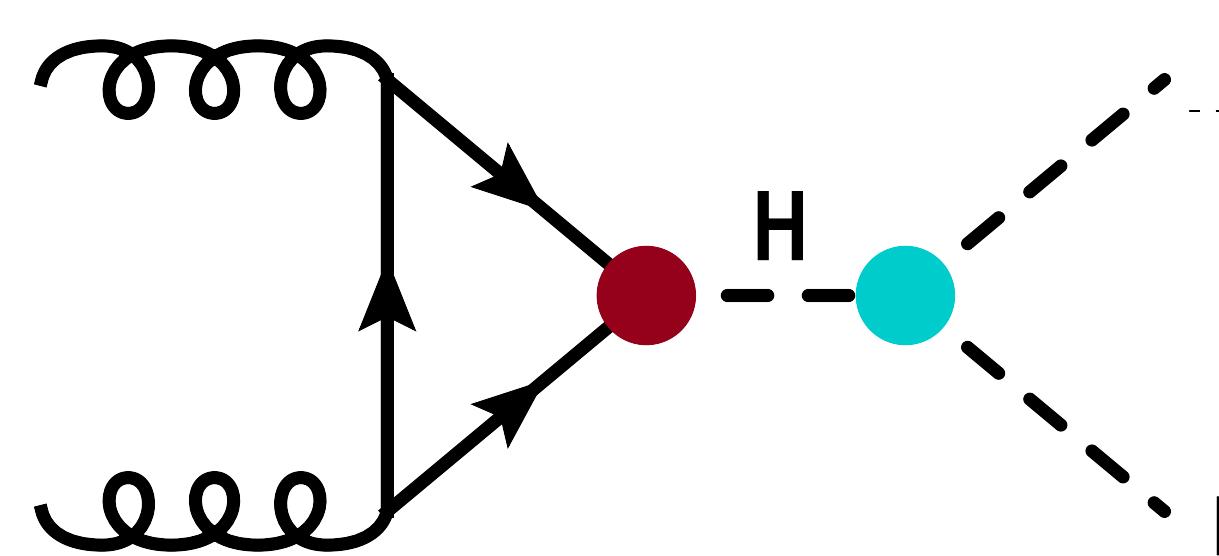
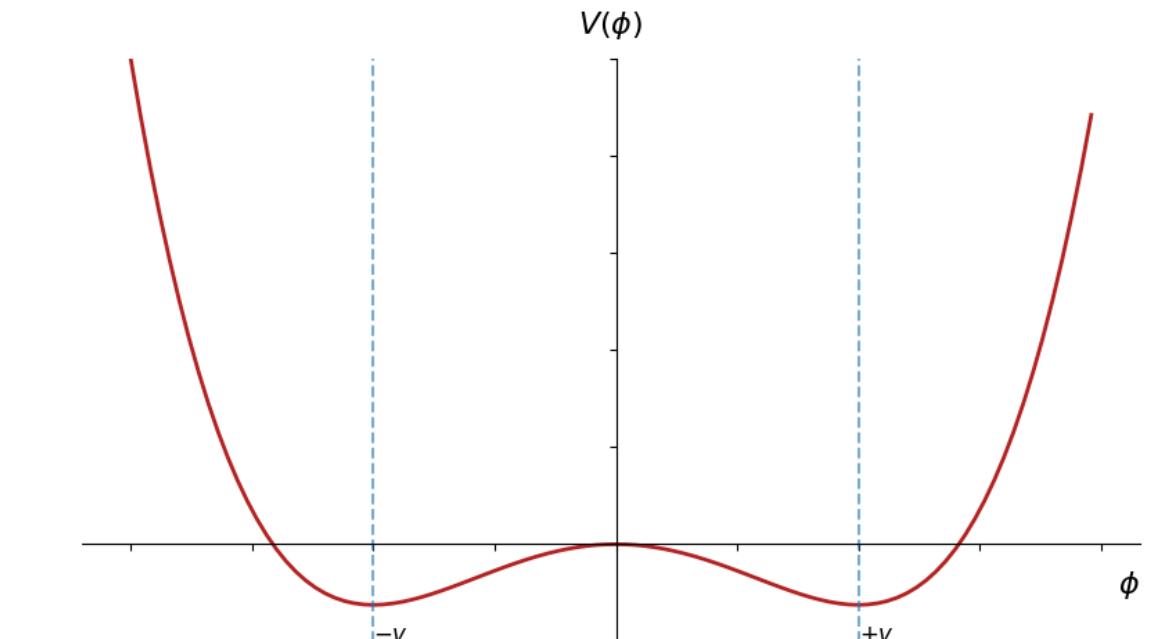
1. Including new variables as inputs based on the topo-cluster variables, notably the shape observables
2. Switching photon identification from a univariate cuts method to a multivariate algorithm, such as a BDT or NN

Through implementing these methods, an improvement of as much as 25% background rejection for the same signal efficiency can be found, greater than 10% improvement in most  $\eta$ - $p_T$  bins



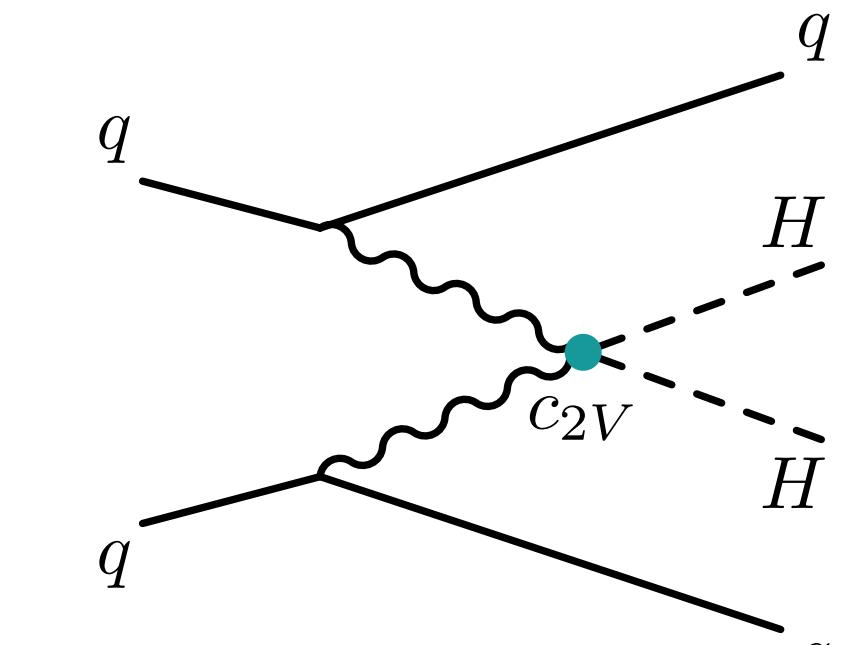
# Conclusion

The discovery of the Higgs boson has been a marker of success for the LHC physics program, and further understanding Electroweak Symmetry breaking is a leading priority of the future

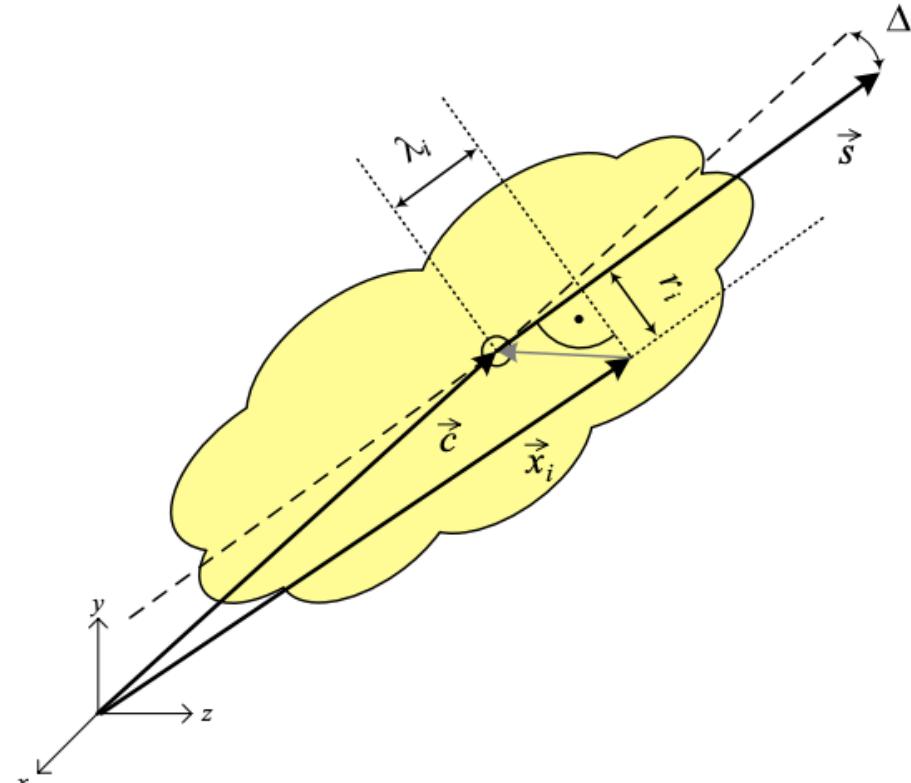


Studying di-Higgs production gives a handle on the Higgs self-coupling and the structure of the Higgs potential, ATLAS searches have provided the best current limits on the self-coupling

$HH \rightarrow \gamma\gamma bb$  is a key channel contributing to this limit, and is working toward results using the Full Run 2 data. Showed a projected 9.8% improvement by including a VBF-enriched category defined via a multiclass BDT



Improvement to future photon identification in ATLAS presented, through inclusion of new discriminating variables, as well as through implementing machine learning methods



# Questions?



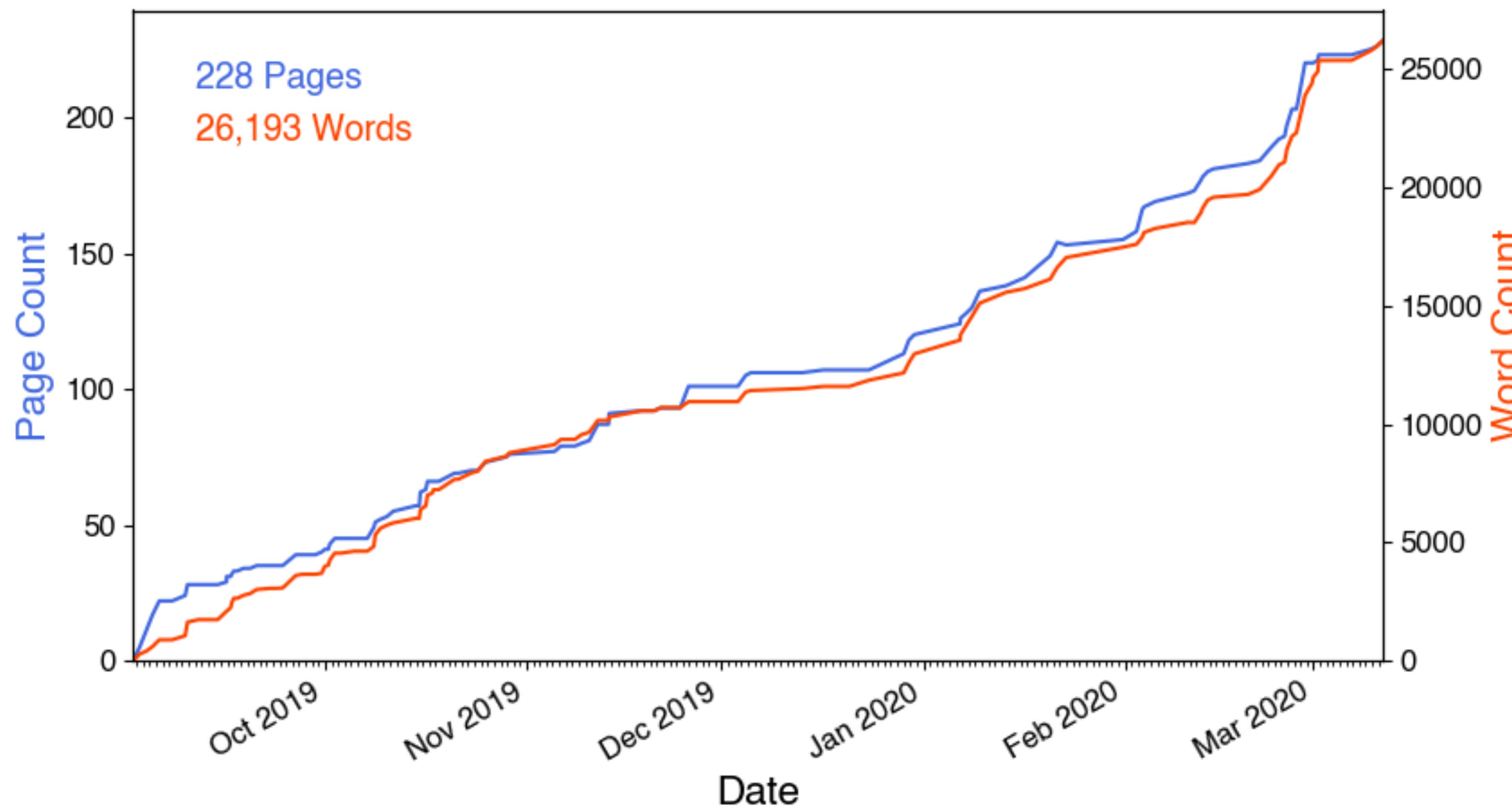
[xkcd](http://xkcd.com)



THE BEST THESIS DEFENSE IS A GOOD THESIS OFFENSE.

# Backup



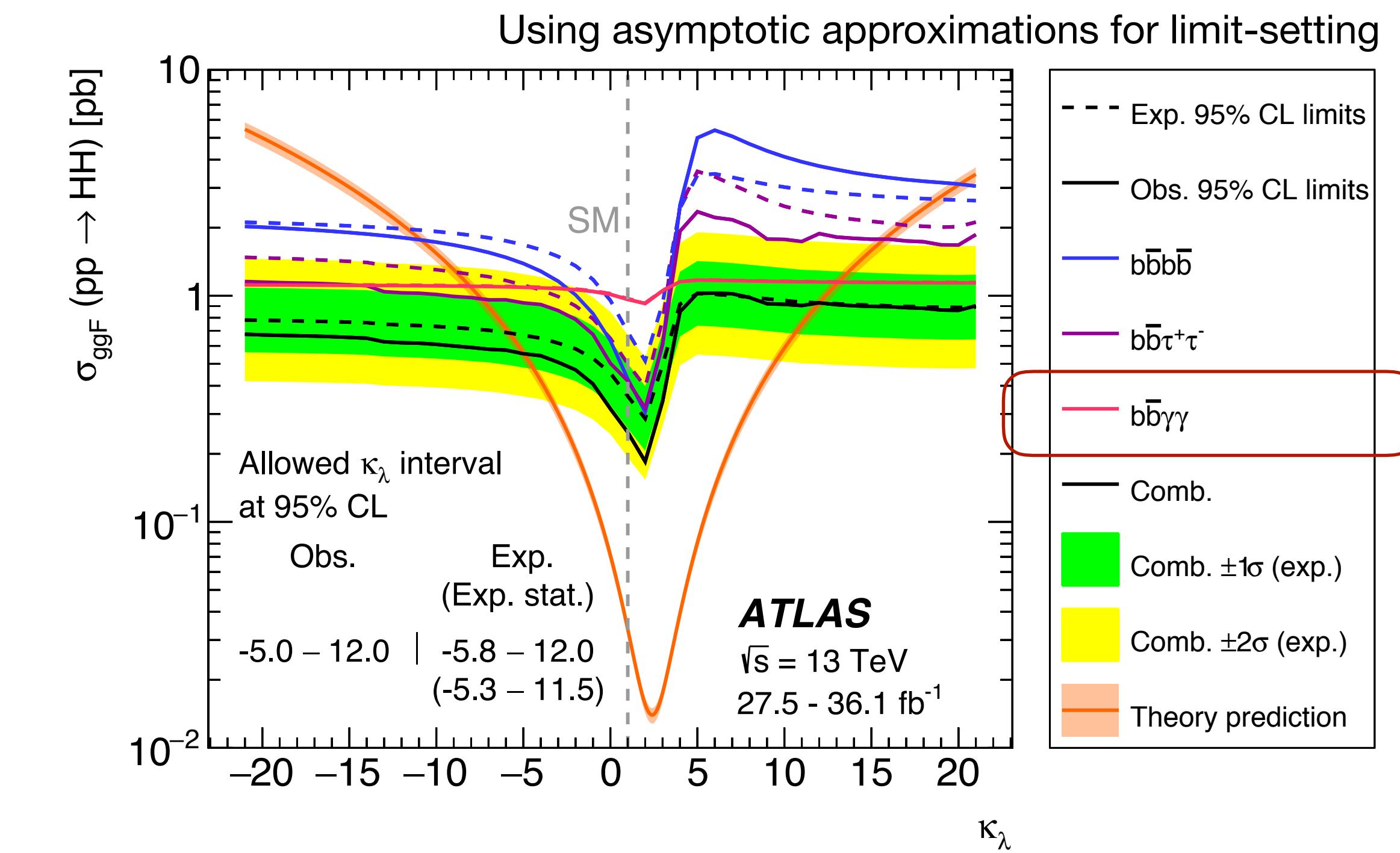
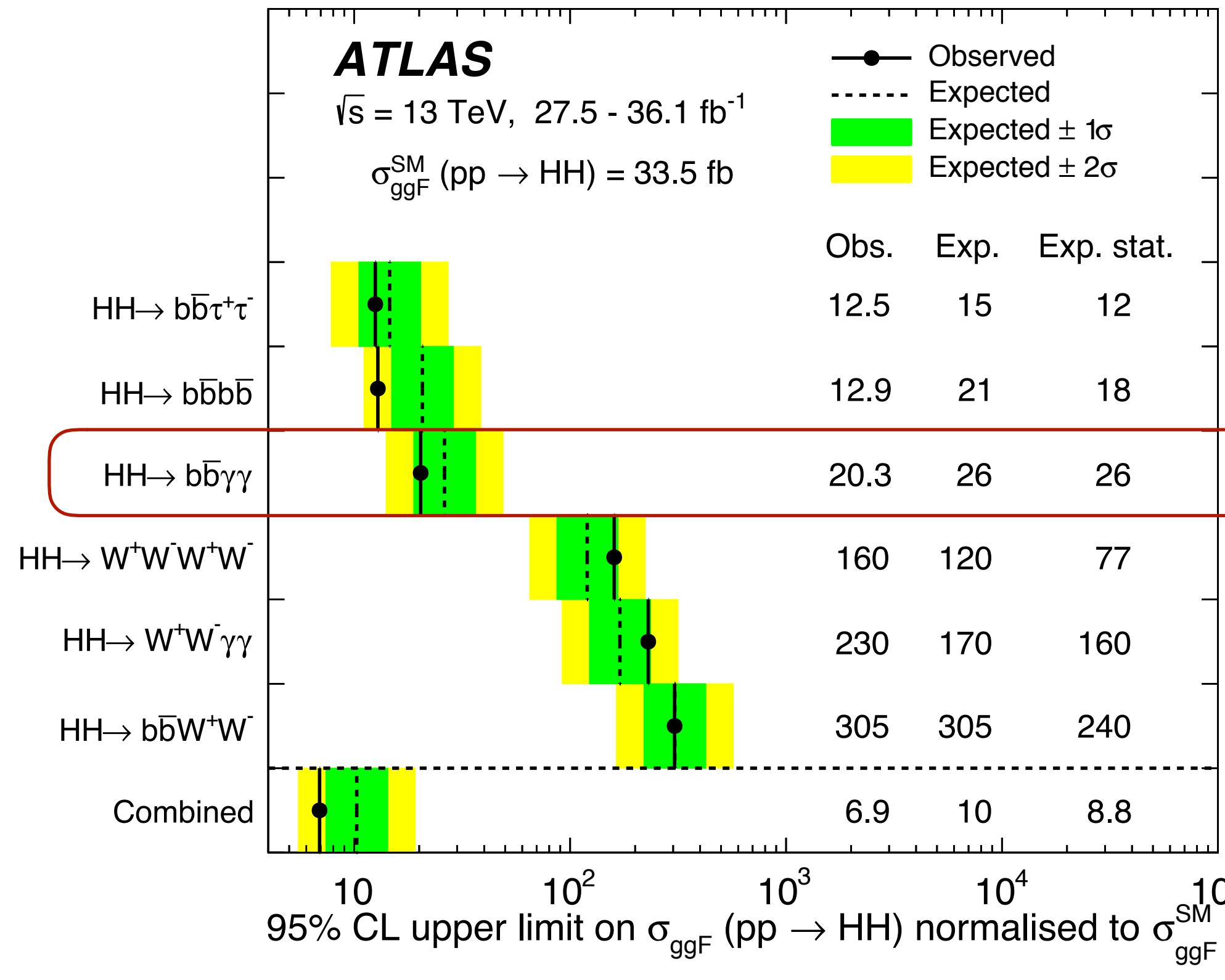


Tyler James Burch

# Combined Limit

$\gamma\gamma bb$  combined with other leading channels to set best ever di-Higgs limits

1906.02025



## Trilinear Coupling

Expected:  $-5.8 < K_\lambda < 12.0$

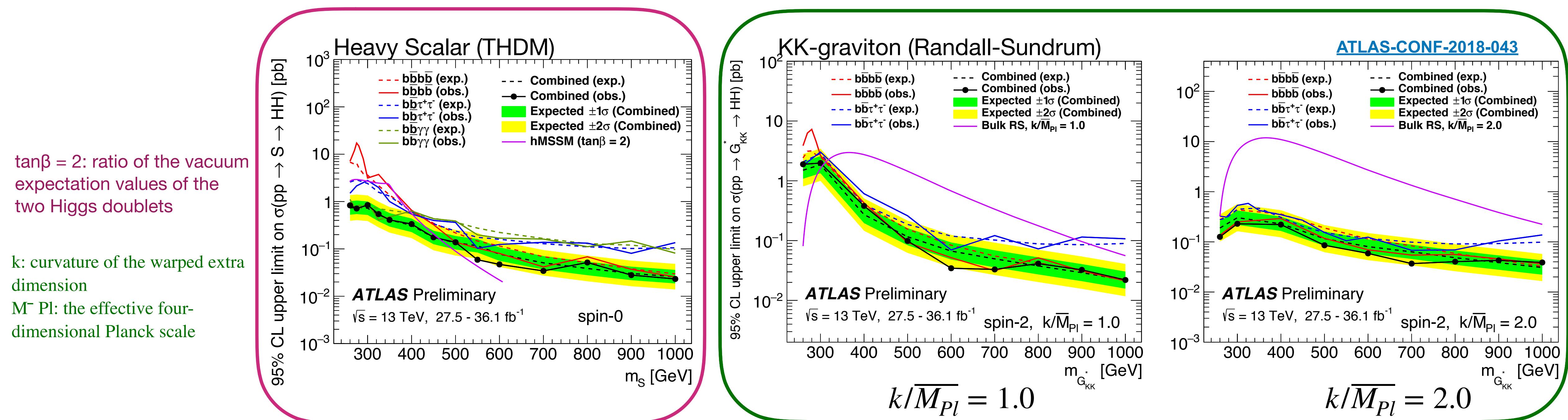
Observed:  $-5.0 < K_\lambda < 12.0$



# Resonant HH Combination Limits

Two-Higgs Doublet Model  
 $m_X < 462 \text{ GeV} @ 95\% \text{ CL in hMSSM}$

Randall-Sundrum Model  
 $k/M_{Pl} = 1.0$  constraints:  $307 < m_G < 1362 \text{ GeV}$   
 $k/M_{Pl} = 2.0$  constraints:  $m_G < 1744 \text{ GeV}$



# VBF Current Constraints

Current constraints set by  $HH \rightarrow bbbb$  publication

$$-1.02 < c_{2V} < 2.71$$

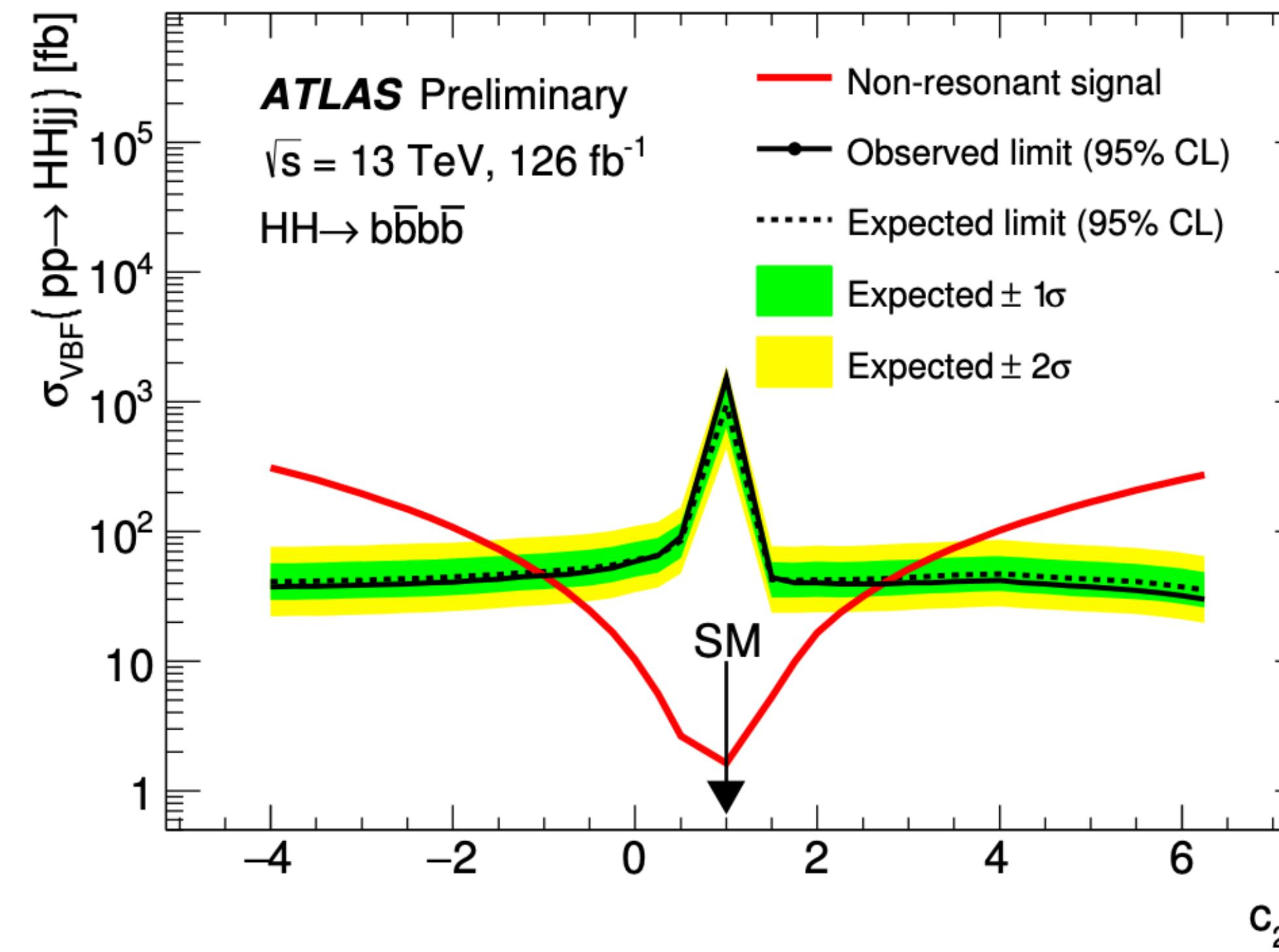


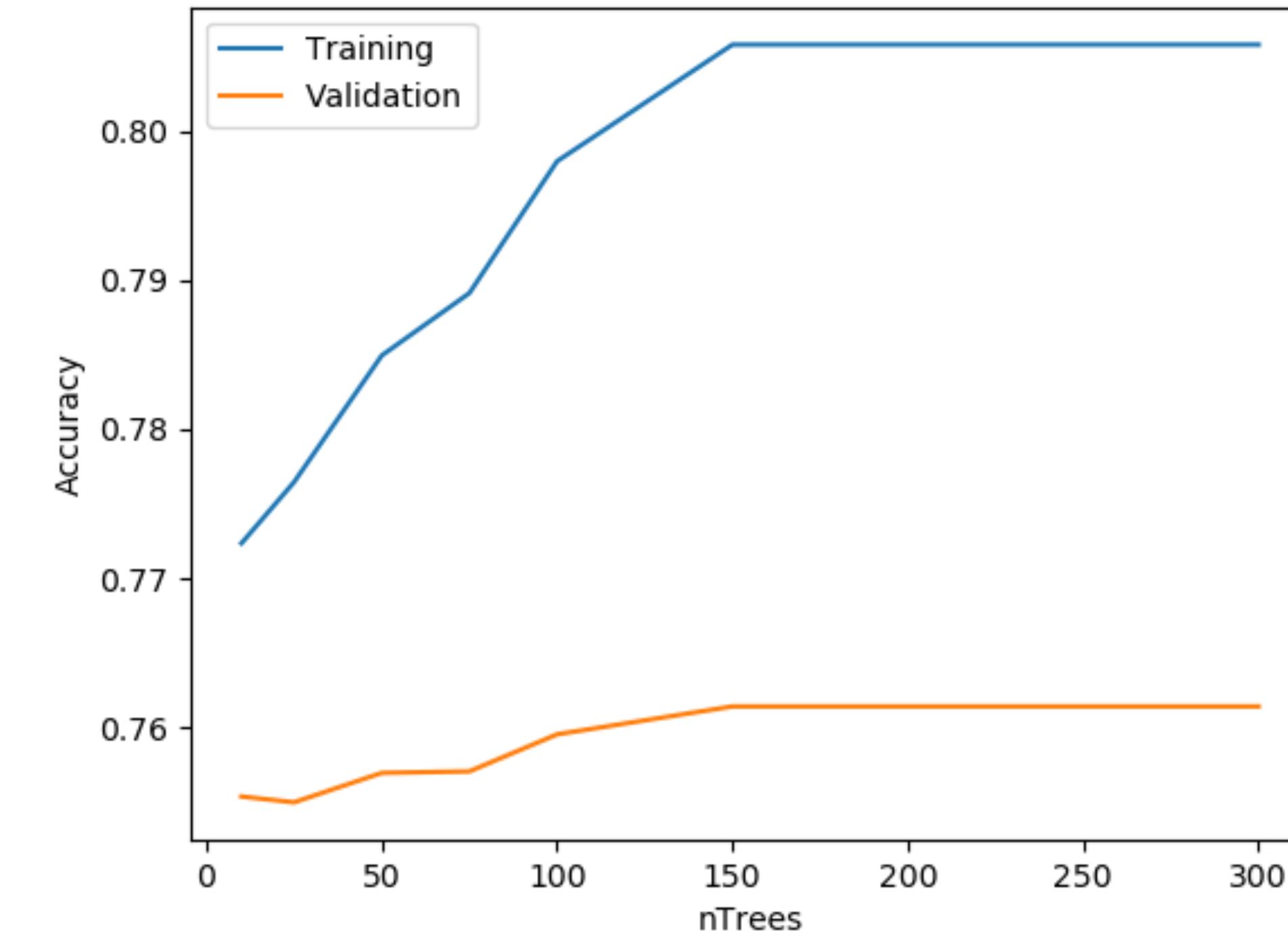
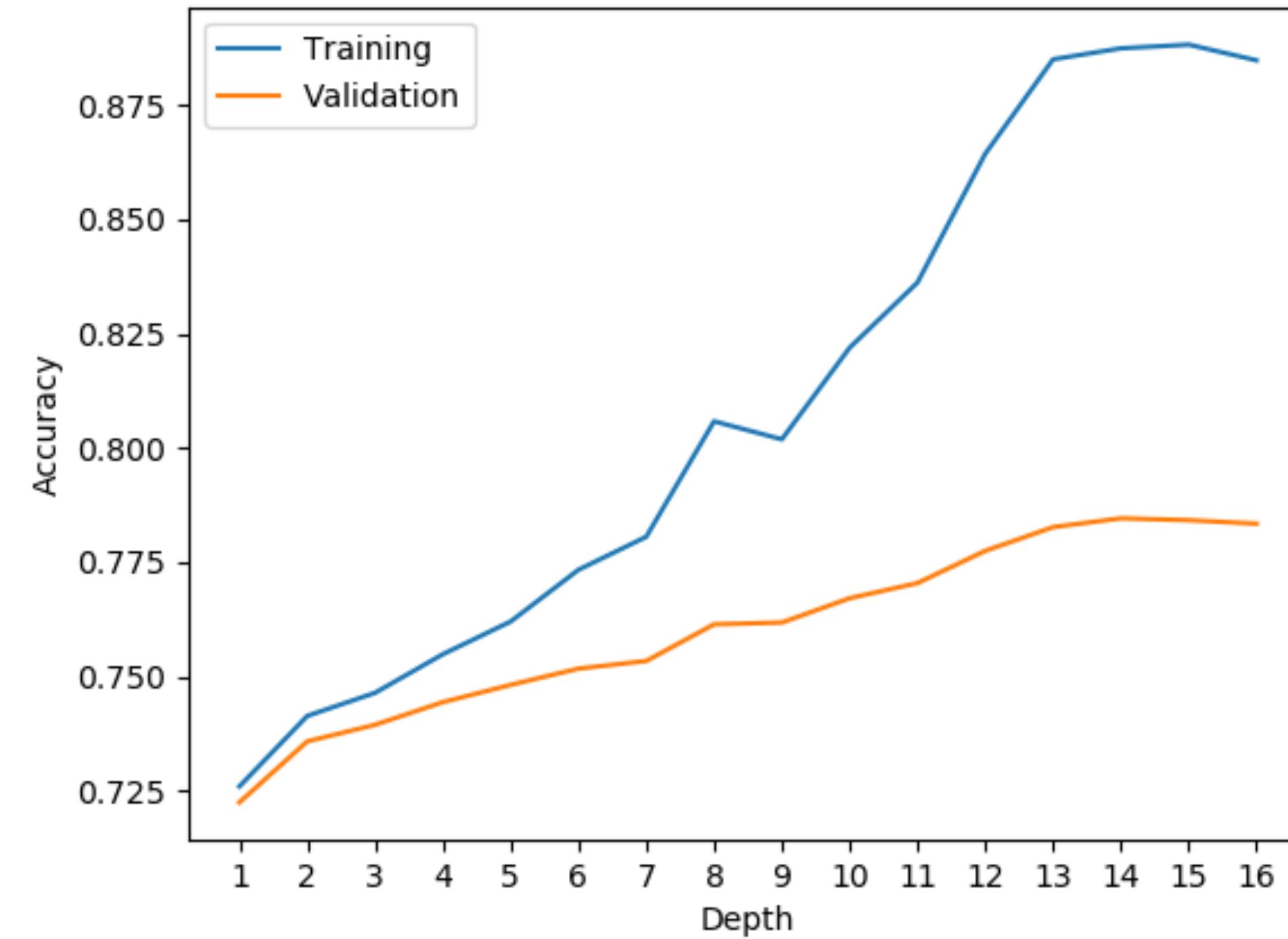
Figure 6: Observed and expected 95% CL upper limits on the production cross-section for non-resonant  $HH$  production via VBF as a function of the di-vector-boson di-Higgs-boson coupling  $c_{2V}$ . The red line corresponds to variations of the leading-order cross-section with  $c_{2V}$ .



# Input List:

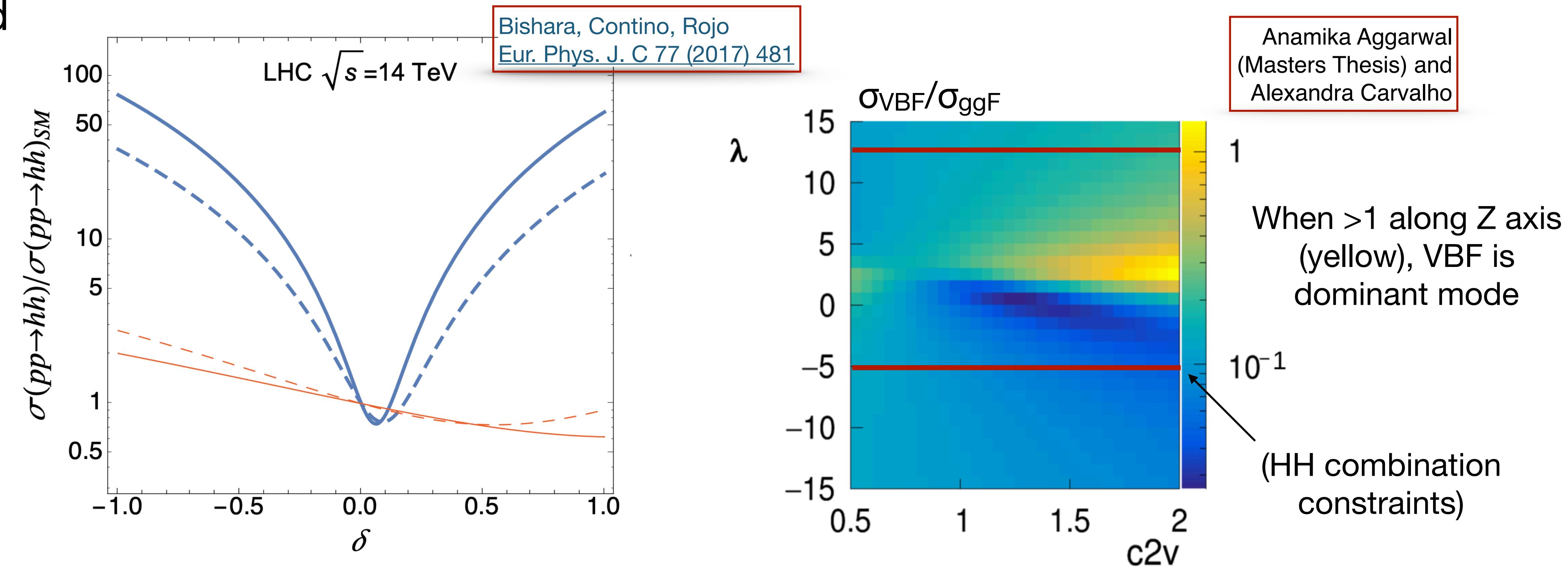
- $\Delta\eta_{jj}$ ,  $m_{jj}$ ,  $m_{yybb}$ ,  $dR_{\gamma\gamma}$ ,  $dR_{bb}$ ,  $j_1$  tagbin,  $j_2$  tagbin, sphericityT, planar flow, pT balance

## Hyperparameters for VBF BDT



# A Note on Enhancements

Coupling variations, specifically to  $c_{2v}$ , could cause scenario in which VBF HH production *only* is enhanced



4b analysis constrained to  $-1.0 < c_{2v} < 2.82$ , so still regimes in which large enhancements to VBF production are allowed, including some in which VBF would theoretically be dominant



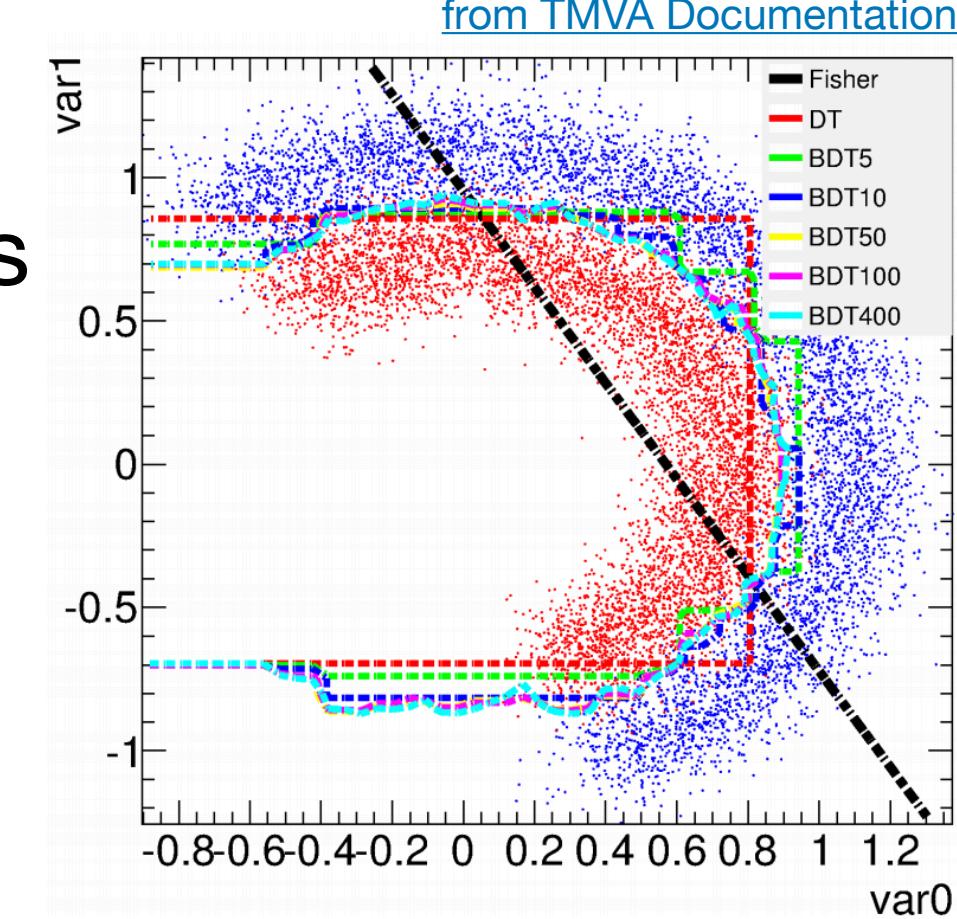
# MVA Techniques for Photon ID

Current methods of photon ID use rectangular cuts optimization over 9 shower shape variables

- Limited solution space in classification

Consider other methods for more sophisticated solution space - MVA techniques

- Two considered, boosted decision tree and neural network

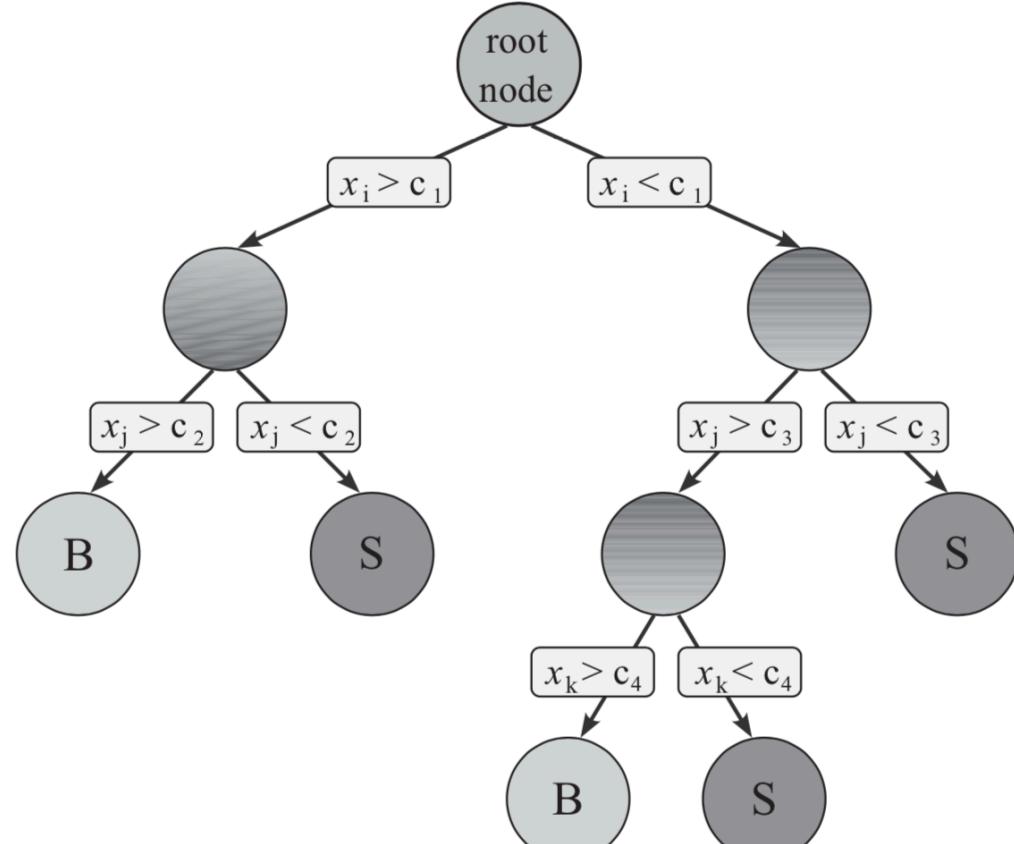


## Boosted decision Tree

Single decision tree: familiar “flow-chart” style method of classification

Boosting - combine many weak learners to make one strong learner

Use **gradient boosting** - subsequent trees upweight misclassified events based on residuals of previous misclassifications

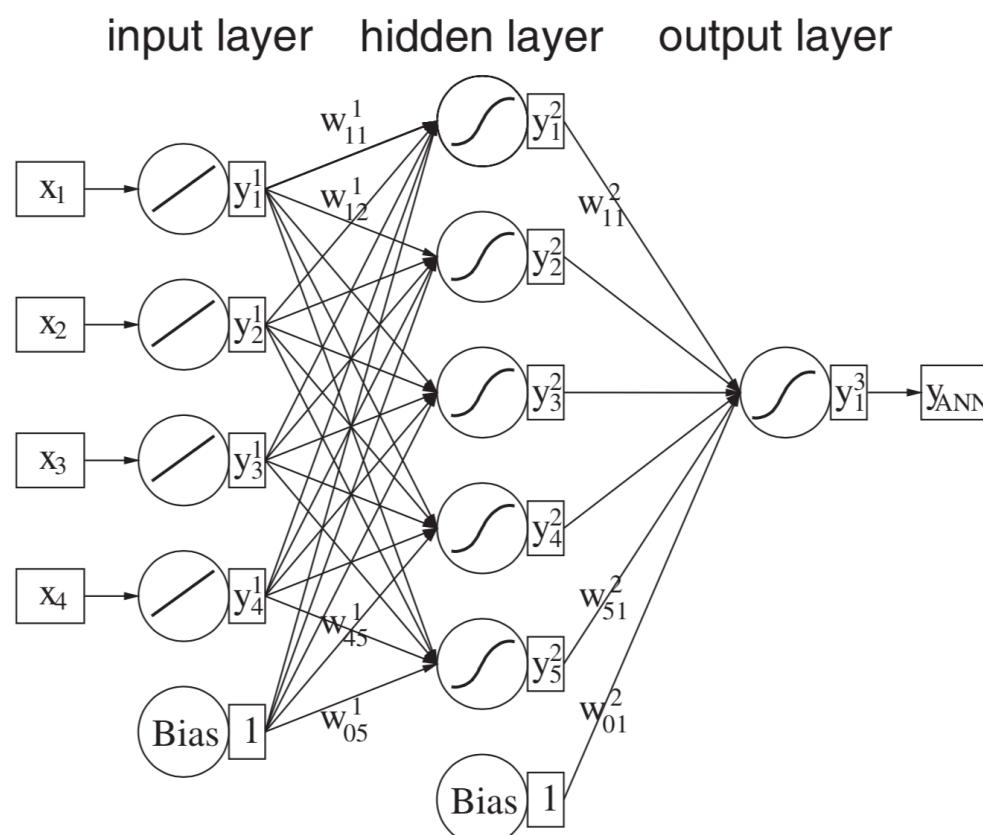


## Neural Network

Inputs evaluated over connected nonlinear basis (“activation”) functions, and ultimately issue a classification probability

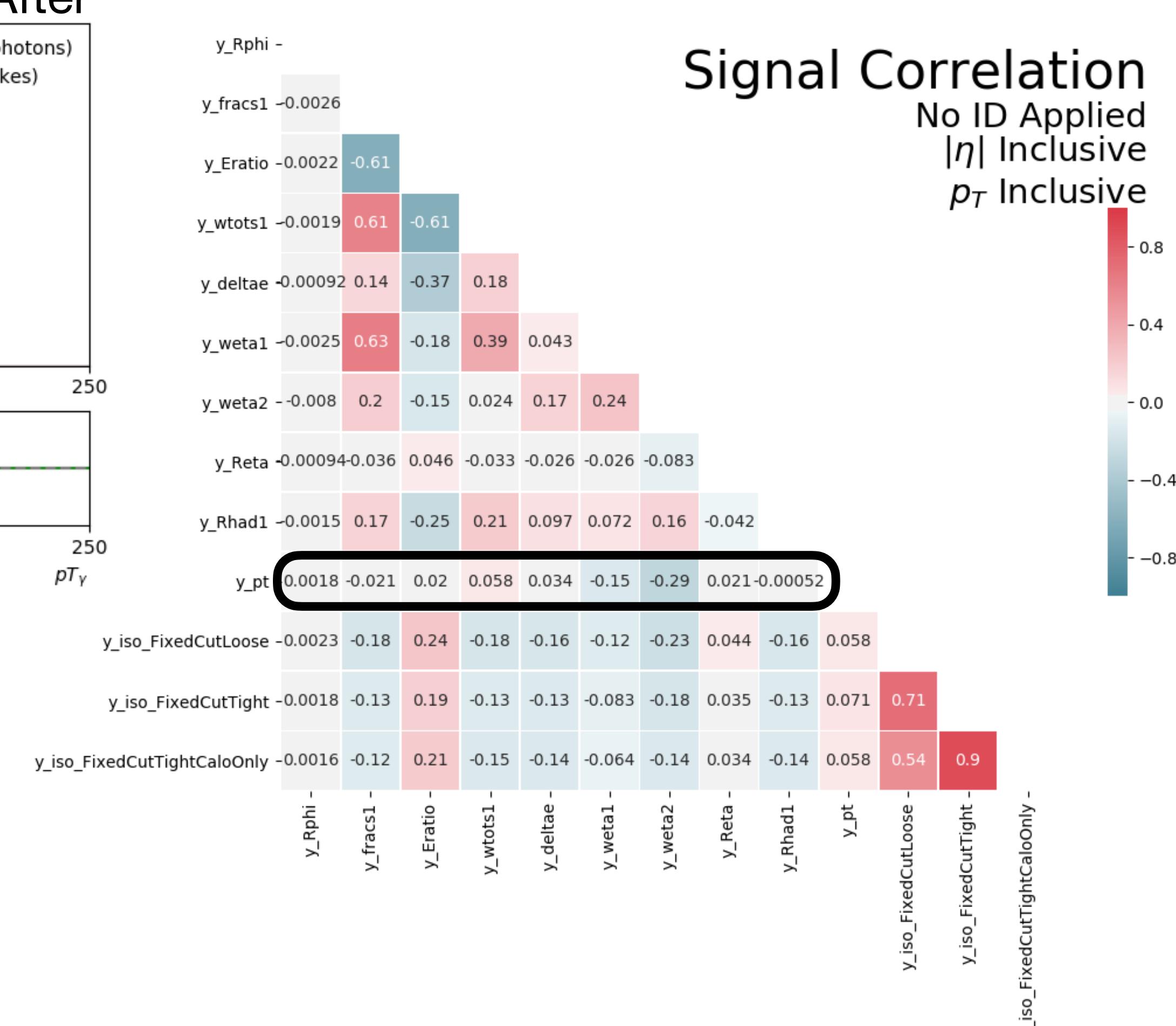
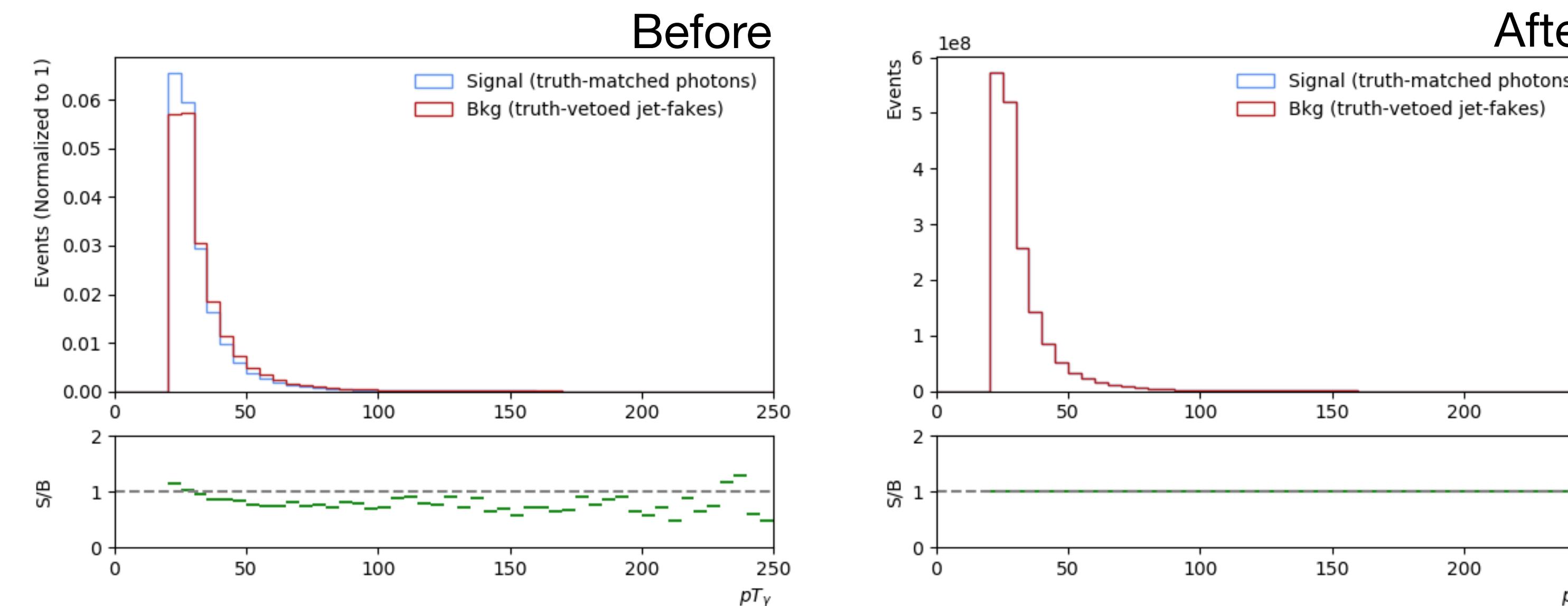
Any possible function can be approximated over sufficient number of activation functions<sup>1</sup>

Weights and biases learned during the training procedure



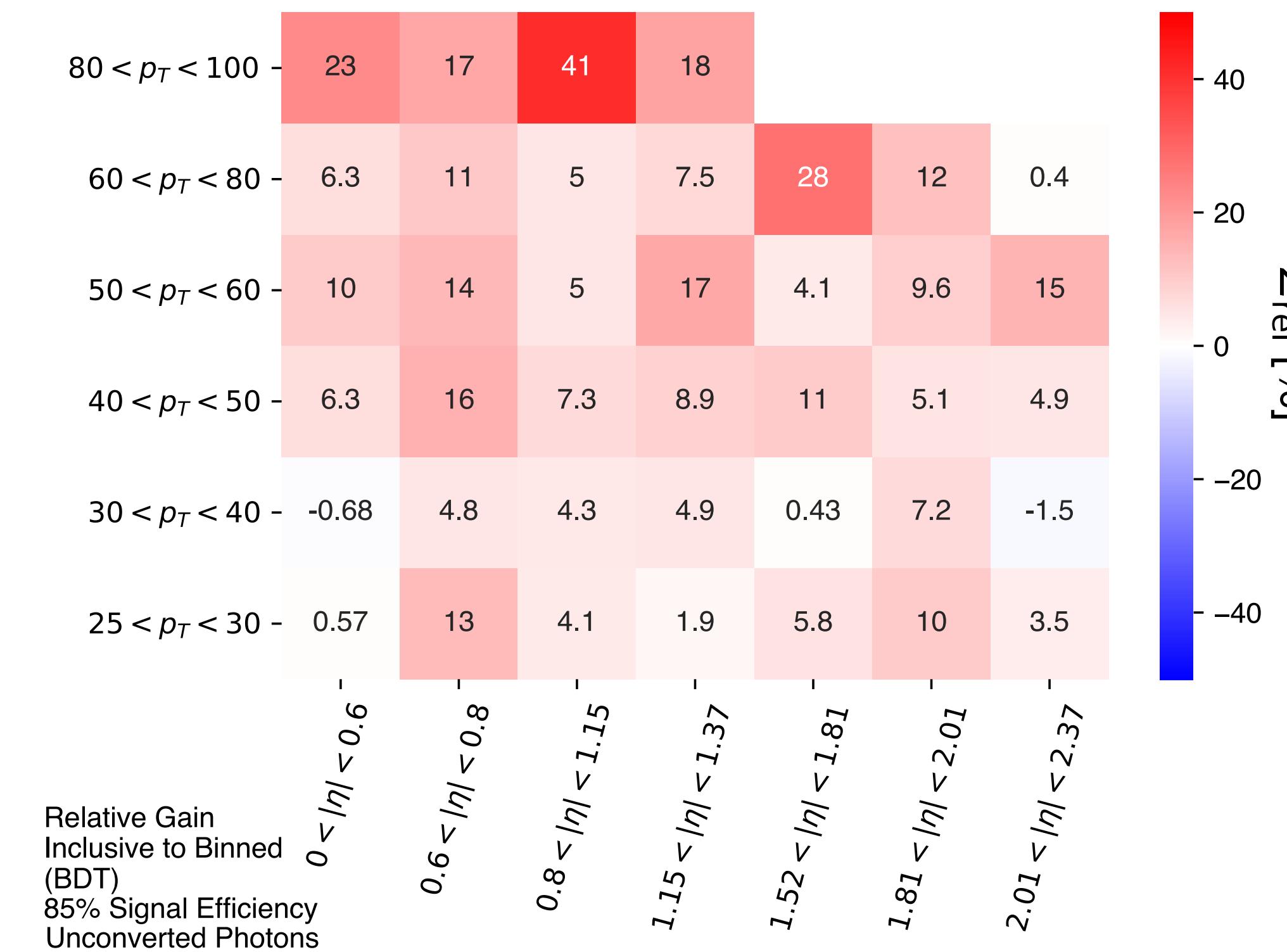
# Training Notes

- Training statistics
  - Specifically at high pT, training statistics become a major problem - approach taken was to train inclusively in pT and apply bin-wise
  - Concern about pT dependence, want to make sure it's not just learning pT differences between signal/ background sample - rescale background pT spectrum to look like signal



- Found rescaling didn't change training at all, didn't understand why for some time
- Ultimately discovered that shower shape has generally low correlation to pT, so correction not needed

# BDT Inclusively Trained vs binwise



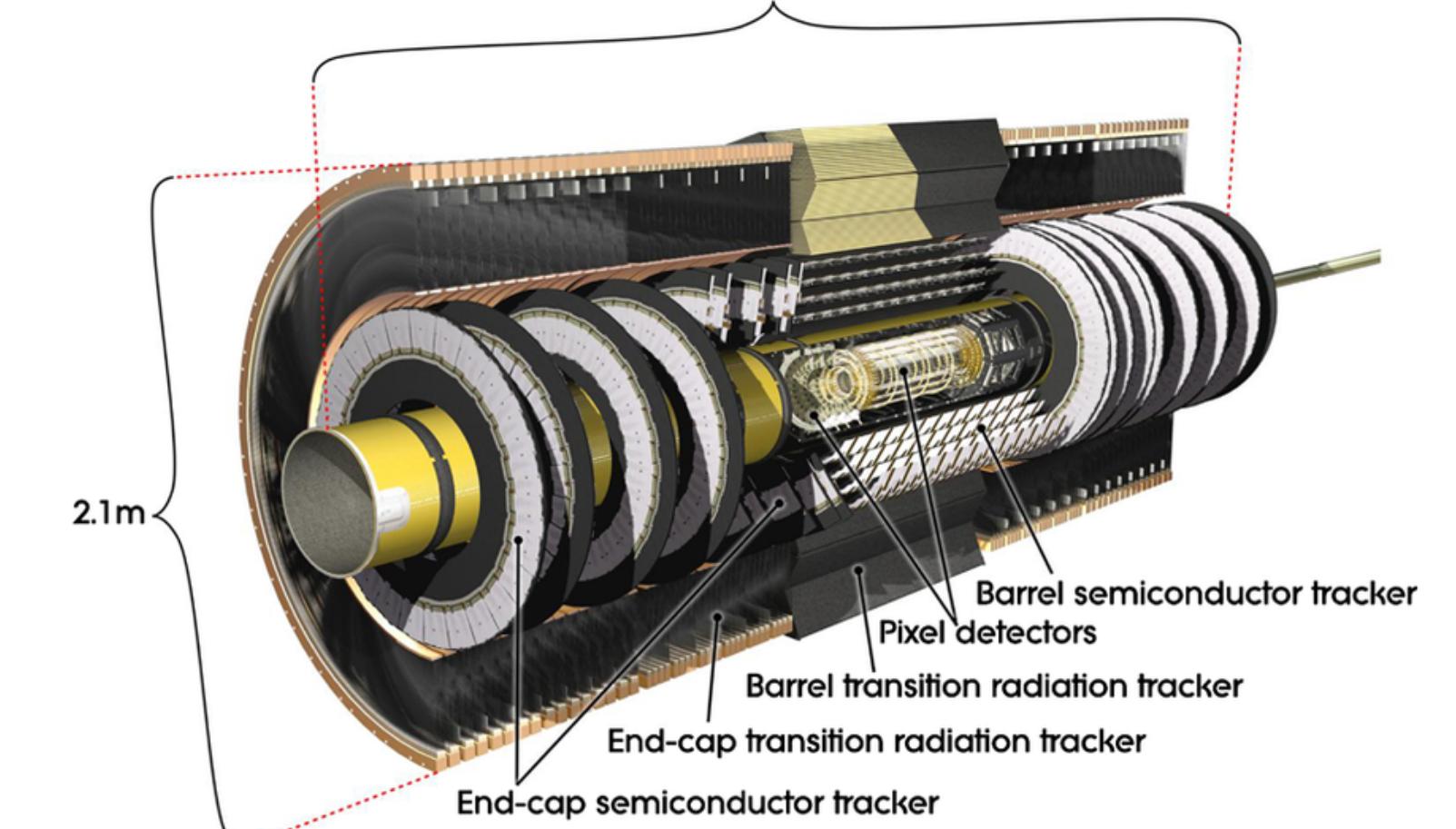
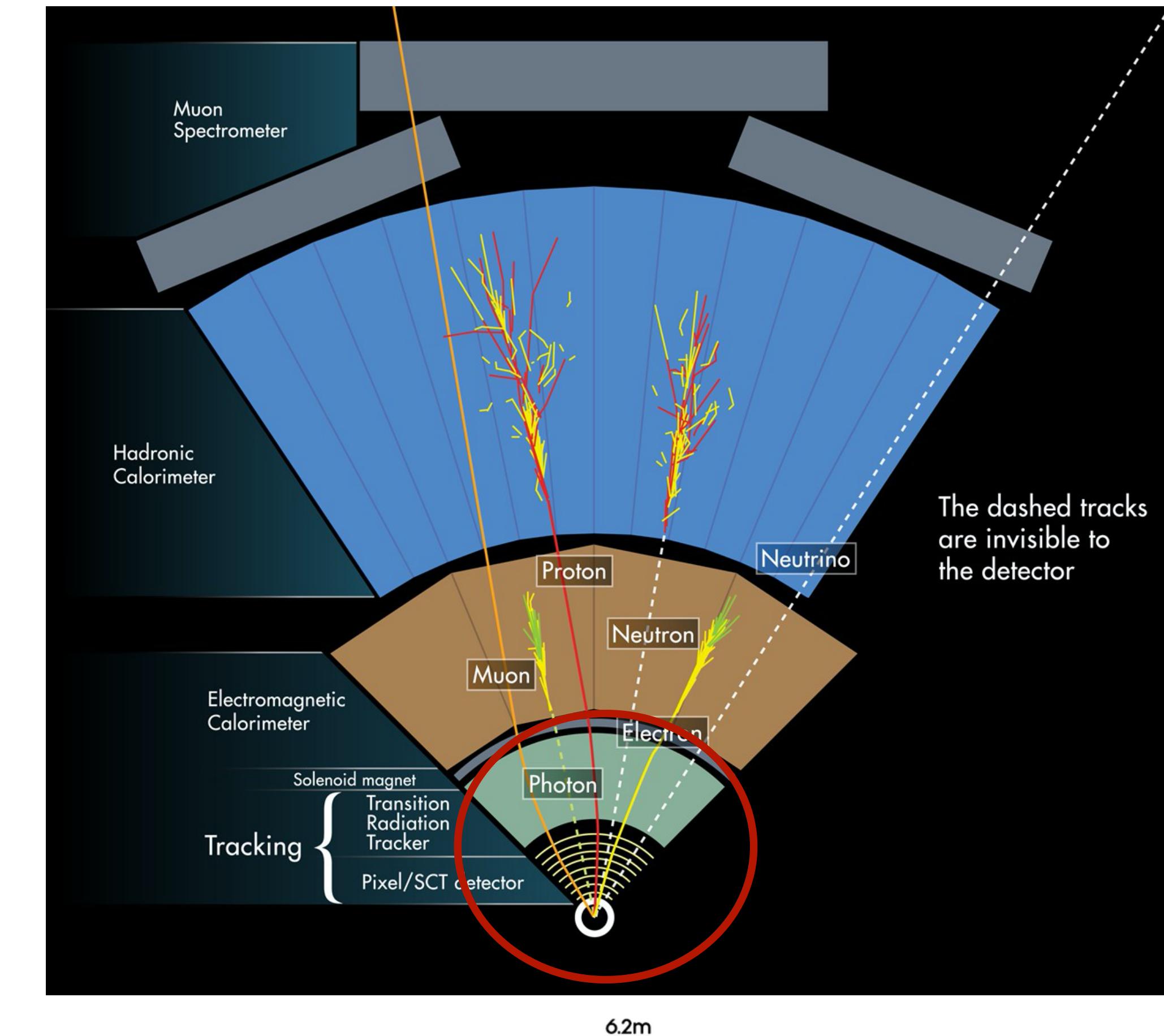
# Particles in...

**ATLAS is a many-layered detector**

Each layer targets specific information from collision decay products

Inner detector - Tracking information

- Describing the particle trajectory through the detector and magnetic fields
- Provides tracking information, vertex measurement, momentum and charge information
- Consists of 3 subsystems:
  - Pixel Detector
  - SemiConductor Tracker
  - Transition Radiation Tracker



# Particles in...

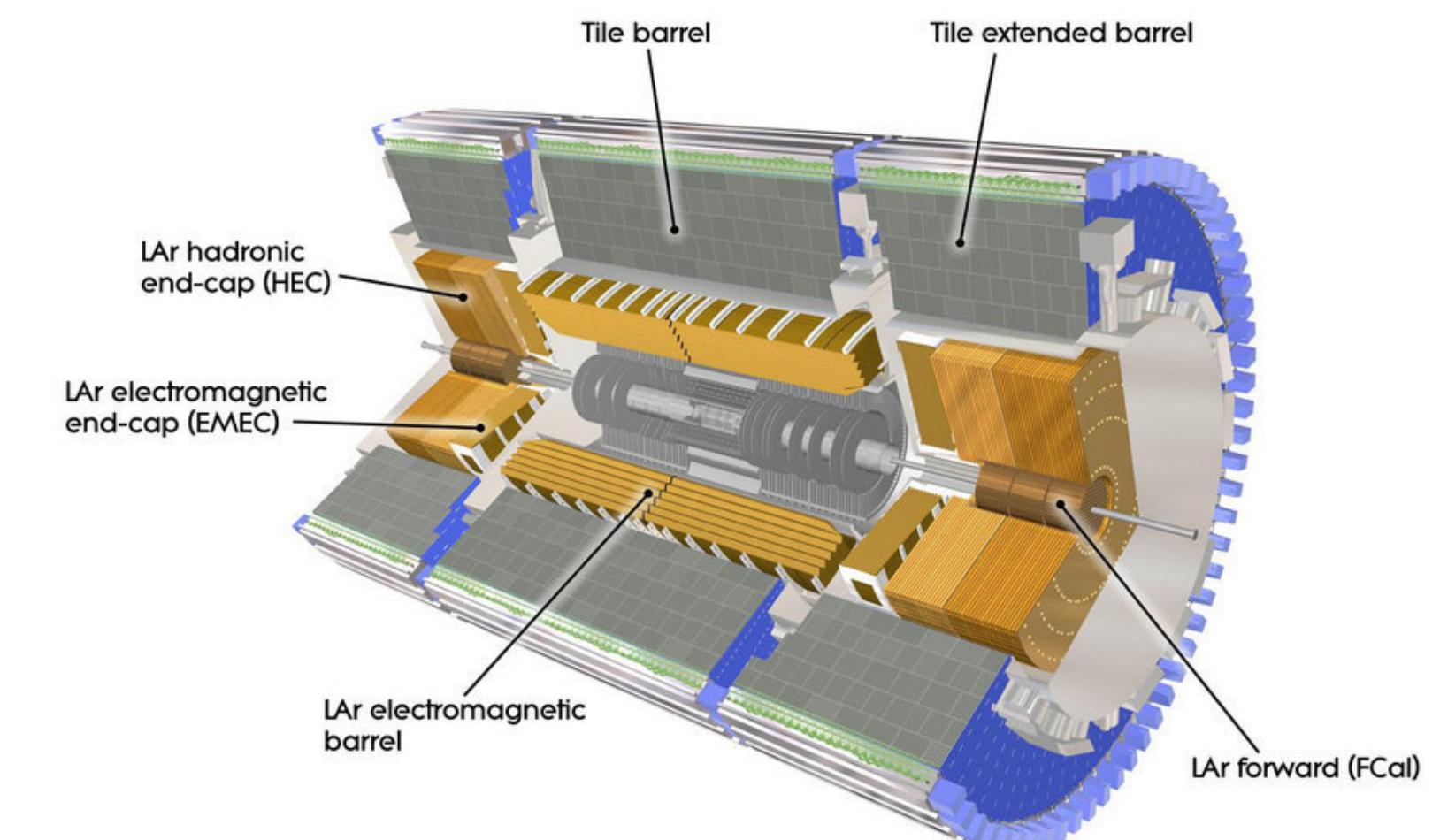
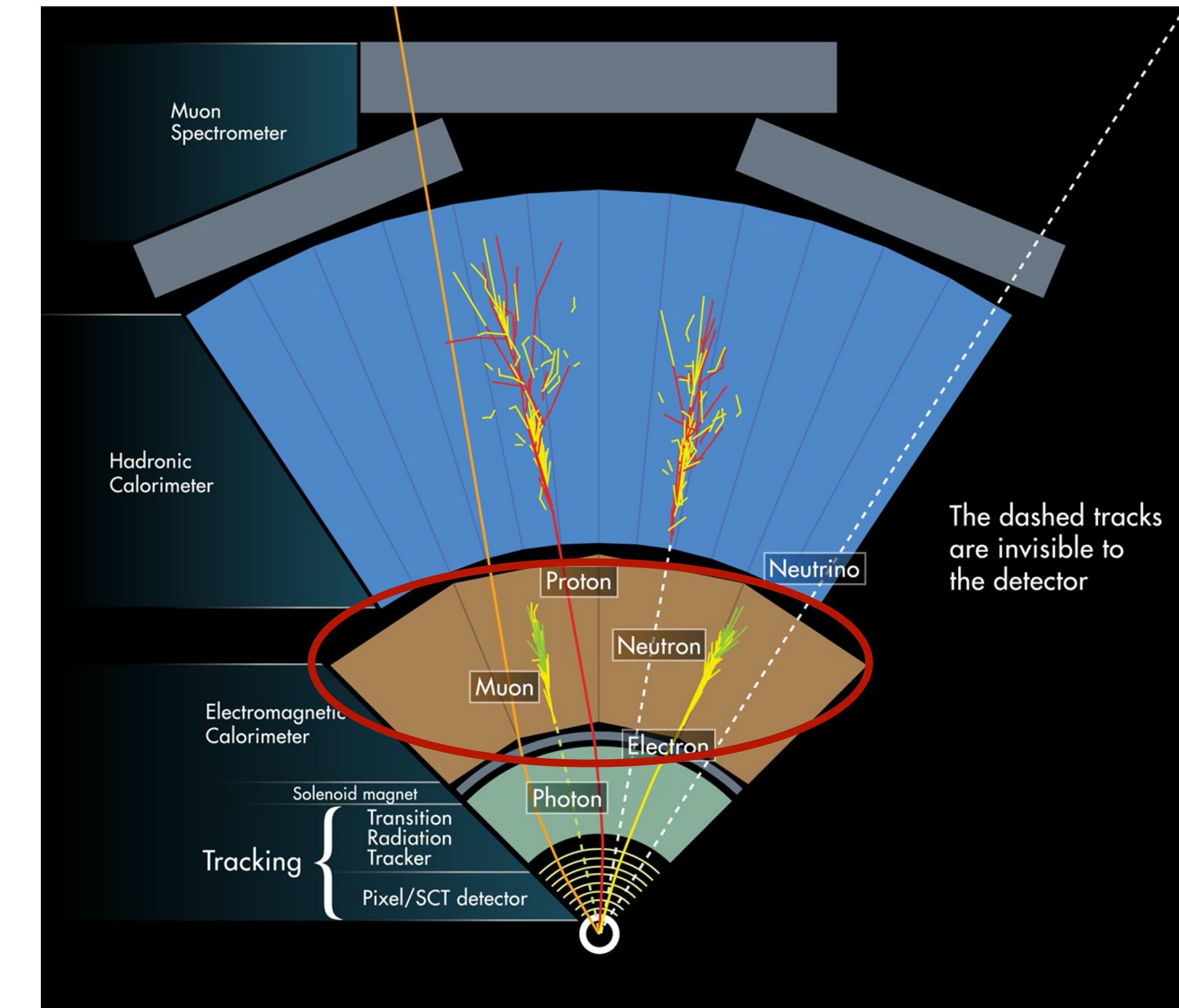
**ATLAS is a many-layered detector**

Each layer targets specific information from collision decay products

Inner detector - Tracking information

Electromagnetic Calorimeter - Electromagnetic (EM) signatures ( $\gamma, e$ )

- Liquid Argon calorimeter
- Accordion-style geometry - high granularity
- In material, electromagnetic showers occur, which are based on 2 processes
  - Bremsstrahlung photon emission
  - Electron-positron conversion



# Particles in...

**ATLAS is a many-layered detector**

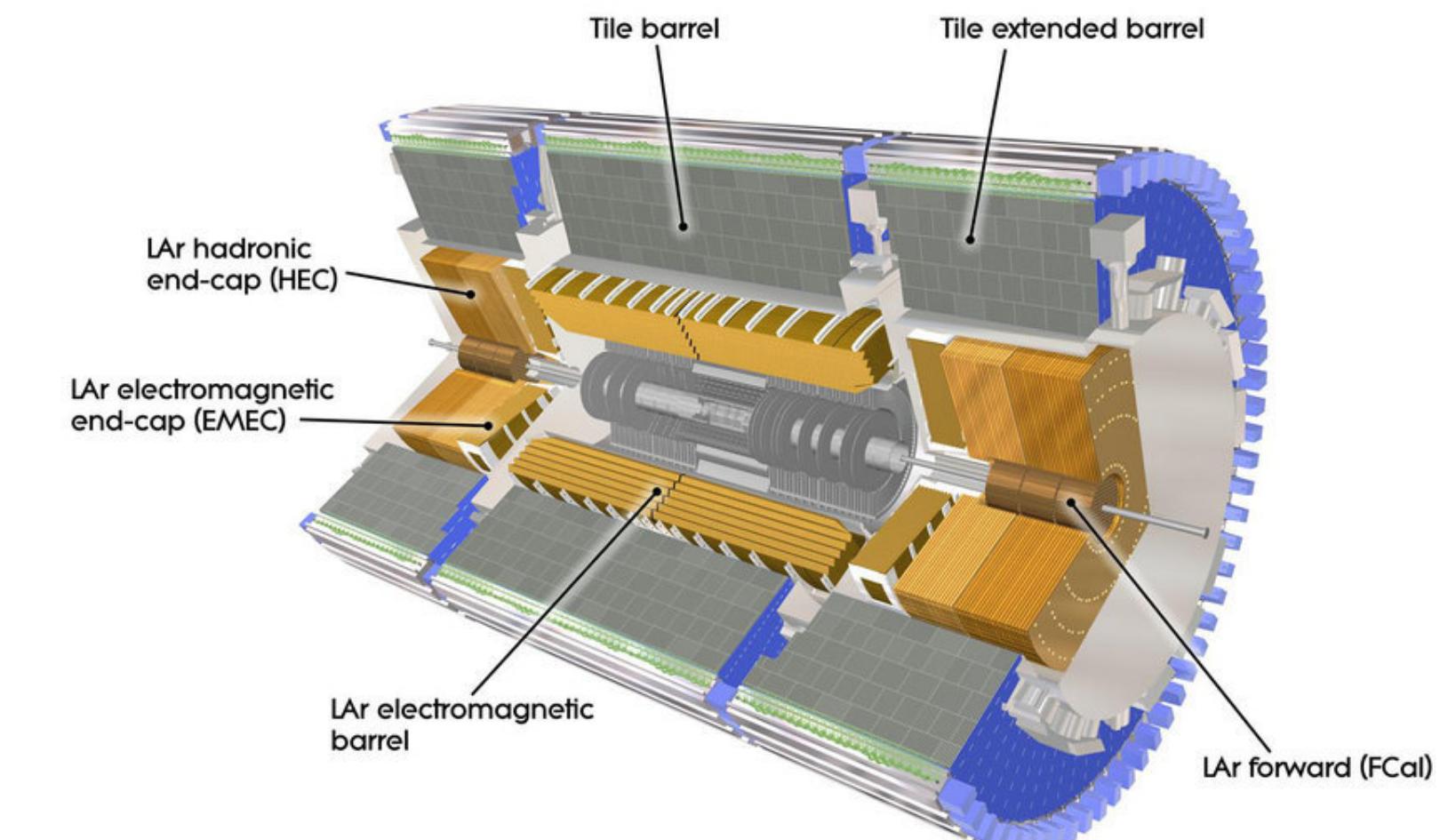
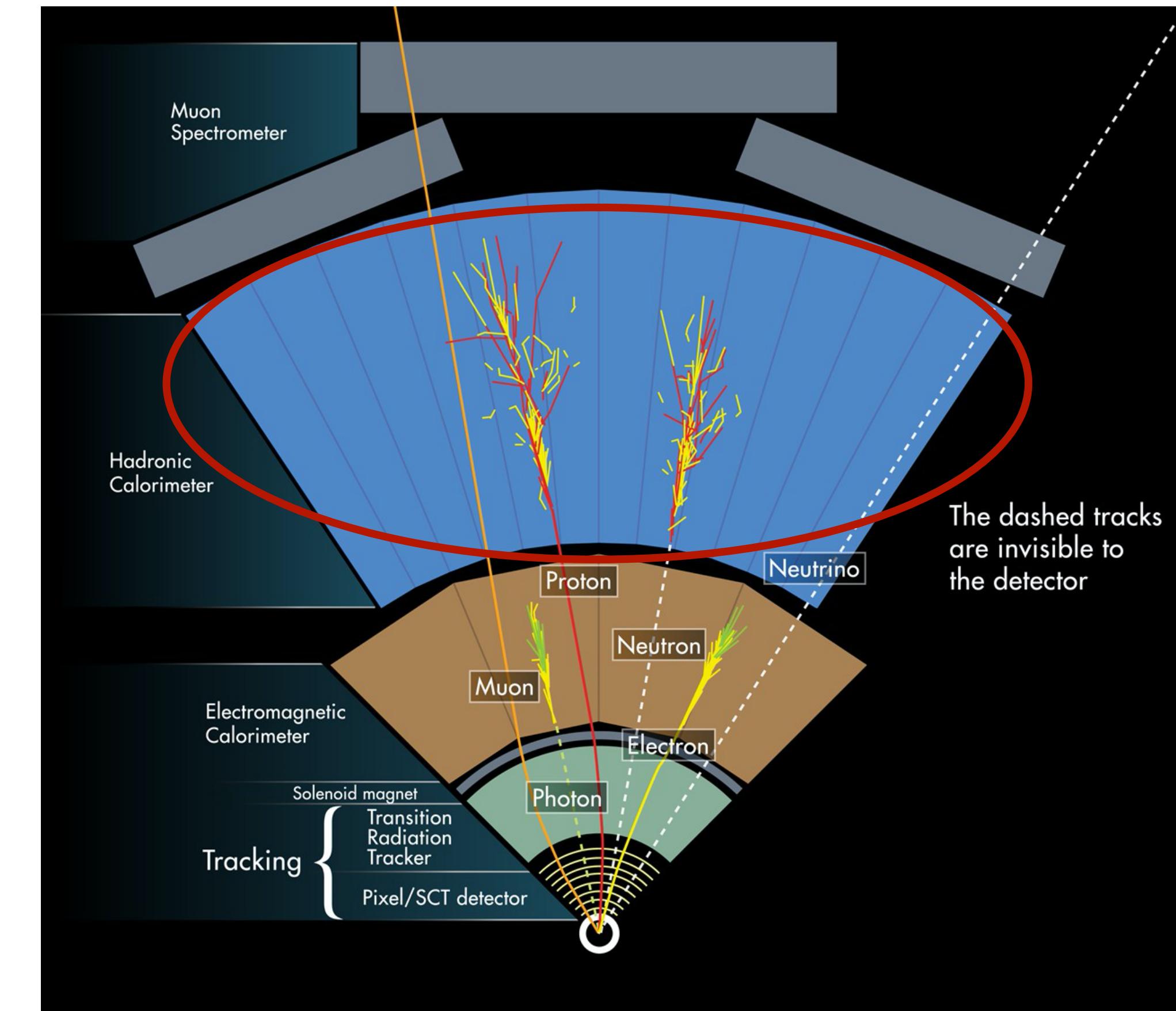
Each layer targets specific information from collision decay products

Inner detector - Tracking information

Electromagnetic Calorimeter - Electromagnetic (EM) signatures ( $\gamma, e$ )

Hadronic Calorimeter - Particles that interact via the strong force (quarks, gluons)

- Initiating strongly interacting particles may radiate additional quarks and gluons
- Must all be in colorless bound states due to *confinement*, so eventually turn into hadrons ( $\pi, K, p, n$ , etc.)
- Final signature is a collimated spray of hadrons “jets”
- As hadrons interact with matter they scatter and produce additional hadrons, cascade known as hadron shower
- Calorimeter to measure showers - steel absorber, scintillating tile active material



# Particles in...

**ATLAS is a many-layered detector**

Each layer targets specific information from collision decay products

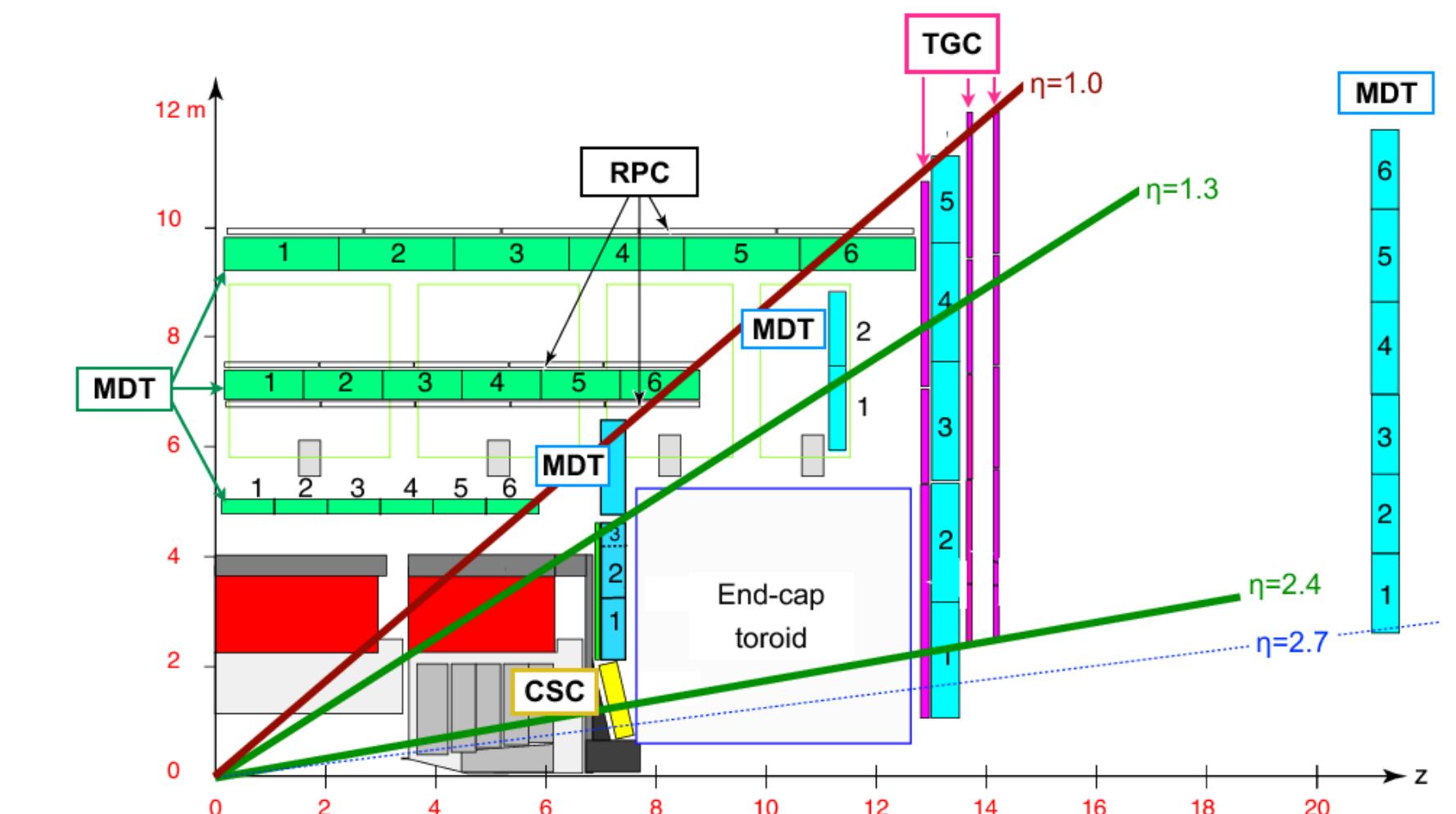
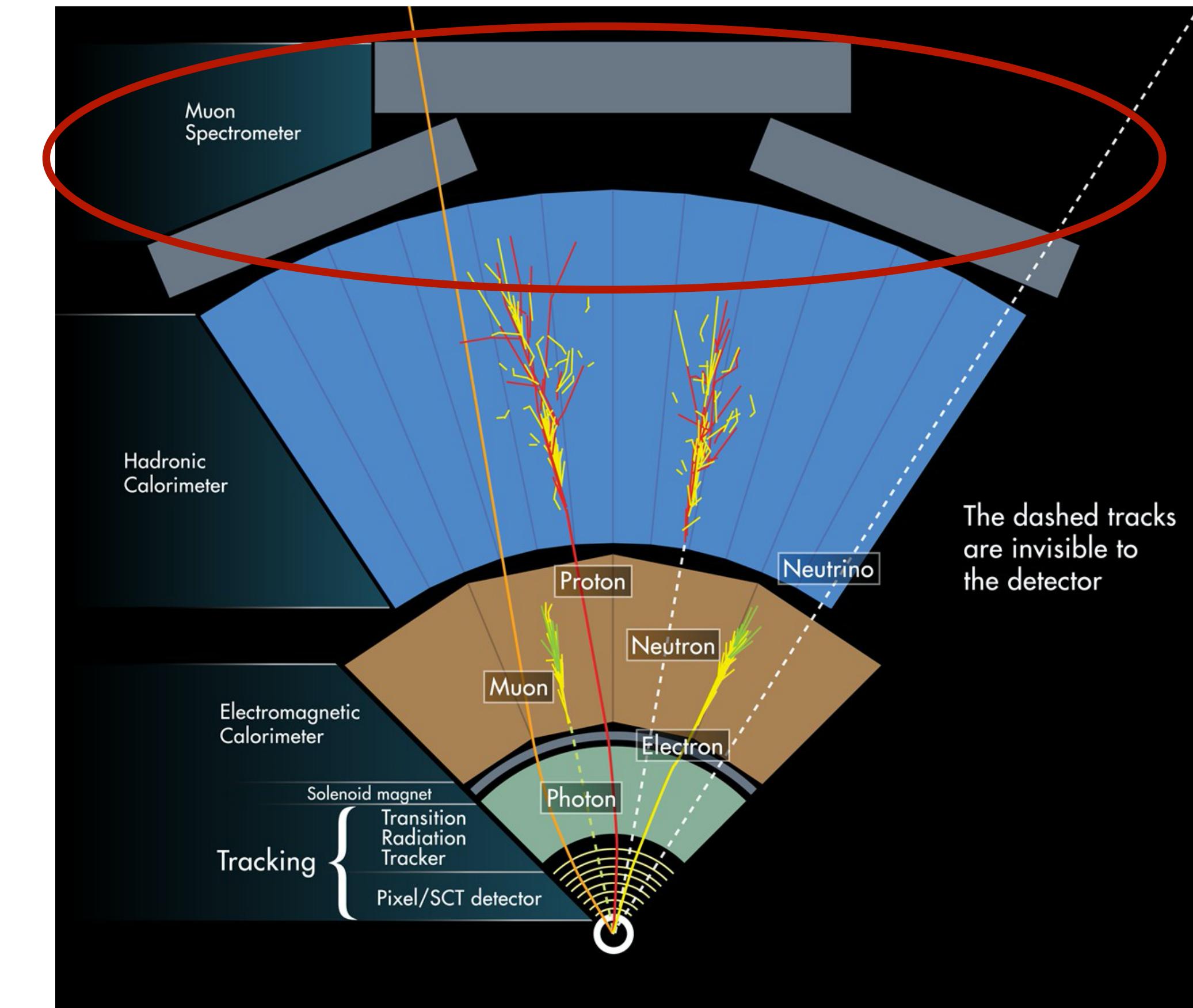
Inner detector - Tracking information

Electromagnetic Calorimeter - Electromagnetic (EM) signatures ( $\gamma, e$ )

Hadronic Calorimeter - Particles that interact via the strong force (quarks, gluons)

Muon Spectrometer - Muons

- Don't interact in upstream material, so dedicated subsystem for detecting
- Due to distance from IP, independent triggering, combine with ID information for final decision



# Particles in...

**ATLAS is a many-layered detector**

Each layer targets specific information from collision decay products

Inner detector - Tracking information

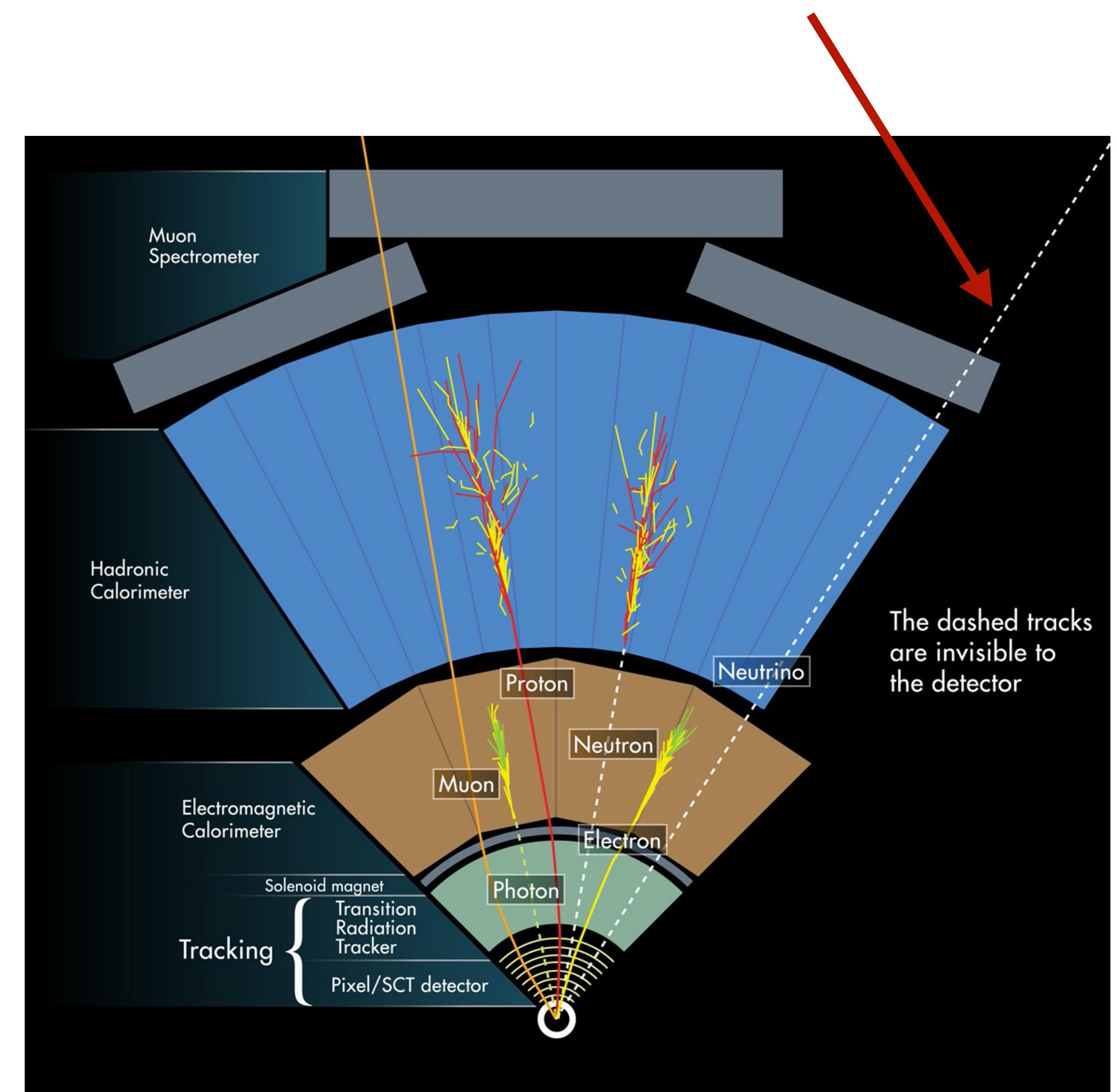
Electromagnetic Calorimeter - Electromagnetic (EM) signatures ( $\gamma, e$ )

Hadronic Calorimeter - Particles that interact via the strong force (quarks, gluons)

Muon Spectrometer - Muons

Missing Transverse energy (MET) - Neutrinos, BSM Signatures

- Particles that don't interact with any portion of the detector
- Conservation of momentum - all particles reconstructed should sum to 0, imbalances in this are taken as MET



# CLs

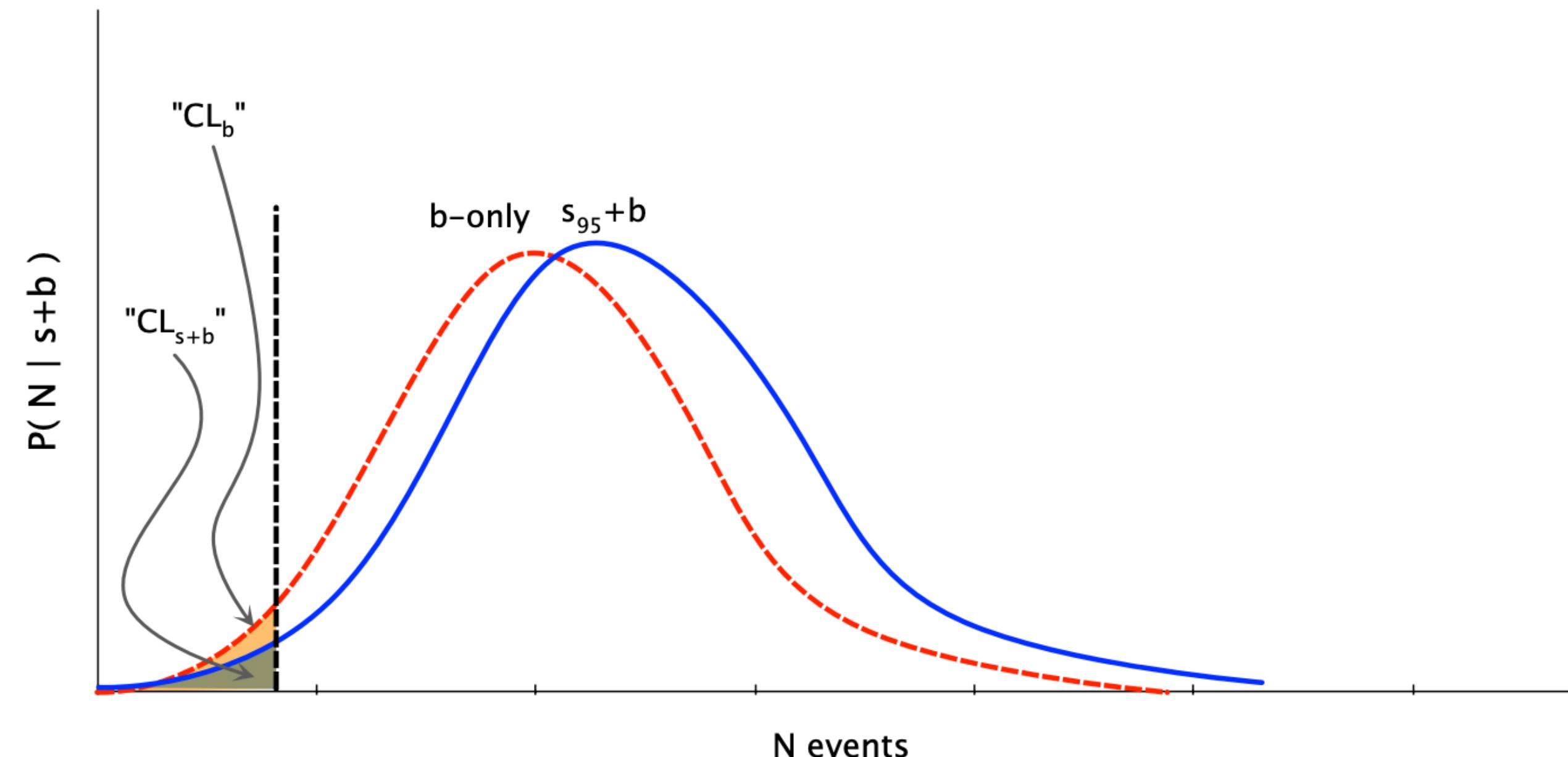
To address the sensitivity problem, CLs was introduced

<http://inspirehep.net/record/599622>

- common (misused) nomenclature:  $CL_s = CL_{s+b}/CL_b$
- idea: only exclude if  $CL_s < 5\%$  (if  $CL_b$  is small,  $CL_s$  gets bigger)

$CL_s$  is known to be “conservative” (over-cover): expected limit covers with 97.5%

- Note:  $CL_s$  is NOT a probability



Cranmer

