

Assignment - 2

Q.1.

$$\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

Given: propagation along X-axis $\Rightarrow \vec{k} \parallel \hat{x}$
 polarization along Y-axis $\Rightarrow \vec{E}_0 \parallel \hat{y}$

$$\therefore \vec{E} = E_0 e^{i(kx - \omega t)} \hat{y}$$

Now, $\vec{B} \perp \vec{E}$ & $\vec{B} \perp \vec{k}$

$$\therefore \vec{B}_0 \parallel \hat{z}$$

$$\therefore \vec{B} = B_0 e^{i(kx - \omega t)} \hat{z}$$

To check that they satisfy the wave equation

$$\frac{\partial^2 \vec{A}}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2}$$

$$\begin{aligned} \frac{\partial^2 \vec{E}}{\partial x^2} &= E_0 \hat{y} (ik) \frac{\partial}{\partial x} e^{i(kx - \omega t)} = -E_0 k^2 e^{i(kx - \omega t)} \hat{y} \\ &= -k^2 \vec{E} \hat{y} \end{aligned}$$

$$\frac{\partial^2 \vec{E}}{\partial t^2} = -\omega^2 \vec{E} \hat{y}$$

Now, for an EM wave in vacuum
 $k = \frac{\omega}{c}$

$$\text{Thus, } \frac{\partial^2 \vec{E}}{\partial x^2} = -\frac{\omega^2}{c^2} \vec{E} = -\frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\text{Similarly, } \frac{\partial^2 \vec{B}}{\partial x^2} = -k^2 \vec{B} \hat{z}; \quad \frac{\partial^2 \vec{B}}{\partial t^2} = -\omega^2 \vec{B} \hat{z}$$

a)

To prove: $\frac{\partial E_y}{\partial x} = \frac{\partial B_z}{\partial t}$

Now, $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$ (Maxwell's relation) (1)

holds since both \vec{E} & \vec{B} satisfy the wave equations that are derived from Maxwell's relations

$$\vec{\nabla} \times \vec{E} = (\partial_x E_y - \partial_y E_x) \hat{z} + (\partial_z E_x - \partial_x E_z) \hat{y} + (\partial_y E_z - \partial_z E_y) \hat{x}$$

Now, $\vec{E} = E_x \hat{x} + E_y \hat{y} + E_z \hat{z}$
 $= E_0 e^{i(kx - \omega t)} \hat{y}$

$$\Rightarrow E_x = E_z = 0$$

$$\therefore \vec{\nabla} \times \vec{E} = \partial_x E_y \hat{z} - \partial_z E_y \hat{x}$$

But $\frac{\partial E_y}{\partial z} = \frac{\partial}{\partial z} E_0 e^{i(kx - \omega t)} = 0$

$$\therefore \vec{\nabla} \times \vec{E} = \frac{\partial E_y}{\partial x} \hat{z} \quad \text{--- (2)}$$

Now, $\frac{\partial \vec{B}}{\partial t} = \frac{\partial B_x}{\partial t} \hat{x} \quad (\because B_x, B_y = 0)$ (3)

\therefore from (1), (2) & (3) we have

$$\frac{\partial E_y}{\partial x} \hat{z} = -\frac{\partial B_x}{\partial t} \hat{x}$$

$$\Rightarrow \boxed{\frac{\partial E_y}{\partial x} = -\frac{\partial B_x}{\partial t}}$$

Now,

$$\frac{\partial E_y}{\partial x} = E_0 (ik) e^{i(kx - \omega t)} \hat{z}$$

$$\& \quad \frac{\partial B_x}{\partial t} = B_0 (-i\omega) e^{i(kx - \omega t)} \hat{x}$$

$$\Rightarrow E_0 (ik) = -B_0 (i\omega)$$

$$\Rightarrow \frac{E_0}{B_0} = \frac{\omega}{k} = c$$

$$\therefore E_0 = c B_0$$

$$\Rightarrow |\vec{E}| = c |\vec{B}|$$

b)

$$Q = CV \quad \& \quad I = \frac{\Phi}{L}$$

Now,

energy stored in a capacitor : $U_E = \frac{Q^2}{2C}$

energy stored in an inductor : $U_B = \frac{1}{2} LI^2$

$$\& \quad |\vec{E}| = \frac{dV}{dx} ; \quad C = \frac{A \epsilon_0}{x} \quad (\text{in vacuum})$$

$$\therefore \text{for a capacitor} \quad V = \frac{Q}{A \epsilon_0} \cdot x$$

$$\Rightarrow |\vec{E}| = \frac{Q}{A \epsilon_0}$$

$$\therefore U_E = \frac{A^2 \epsilon_0^2 |\vec{E}|^2}{2 \cdot A \epsilon_0 / x} \Rightarrow U_E = \frac{1}{2} \epsilon_0 |\vec{E}|^2 \cdot \left(\frac{A x}{\text{volume}} \right)$$

Similarly,

$$\Phi = BA, \quad L = \frac{\mu_0 N^2 A}{x} \quad (\text{put } N=1)$$

$$\therefore U_B = \frac{1}{2} \frac{(LI)^2}{L} = \frac{1}{2} \frac{B^2 A^2}{\mu_0 A} \cdot x$$

$$\therefore U_B = \frac{1}{2 \mu_0} B^2 \left(\frac{A \cdot x}{\text{volume}} \right)$$

Thus,

$$U_E = \frac{1}{2} \epsilon_0 |\vec{E}|^2 \quad \& \quad U_B = \frac{1}{2 \mu_0} |\vec{B}|^2$$

Now, for an EM wave

$$U_E = \frac{1}{2} \epsilon_0 |\vec{E}|^2 = \frac{1}{2} \epsilon_0 E_0^2$$

$$U_B = \frac{1}{2\mu_0} |\vec{B}|^2 = \frac{1}{2\mu_0} B_0^2$$

$$\& E_0 = c B_0 = \frac{1}{\sqrt{\epsilon_0 \mu_0}} B_0$$

$$\therefore U_E = \frac{1}{2} \cdot \epsilon_0 \cdot \frac{1}{\epsilon_0 \mu_0} B_0^2$$

$$\therefore U_E = \frac{1}{2\mu_0} B_0^2 = U_B$$

\therefore for an EM wave, the electric & magnetic field have same energy density.

(c) \vec{S} is the flux of energy flow.

Now, energy is flowing at the speed of light (radiation)

\therefore Energy flux over a length = $U_E \times \frac{c}{x \rightarrow \text{length}}$

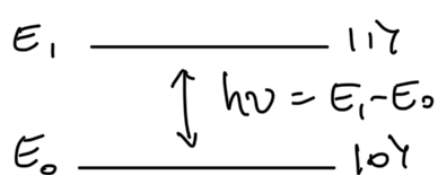
$$\therefore \vec{S} = U_E \times \frac{c}{x} \cdot \hat{x} \quad (\text{flux vector of EM radiation})$$

Intensity will then be $\frac{|\vec{S}|}{\text{Area}}$

$$\therefore I = \frac{U_E \times c}{x \times A} = \frac{1}{2} \epsilon_0 E_0^2 \times \frac{(A \cdot x) \times c}{x \times A}$$

$$\therefore \boxed{I = \frac{1}{2} \epsilon_0 c E_0^2}$$

Q. 2.



$N_0 \rightarrow$ populatⁿ of ground state

$N_1 \rightarrow$ populatⁿ of excited state

$$\frac{dN_1}{dt} = -A N_1 + B_{01} \rho N_0 - B_{10} \rho N_1$$

Einstein postulated that in equilibrium the two level system behaves like a black body

$$\text{i.e. } \rho_{eq} = \frac{8\pi h}{c^3} \cdot \nu^3 \cdot \frac{1}{e^{\beta h\nu} - 1} \quad \text{--- (a)}$$

$$\& \left(\frac{dN_1}{dt} \right)_{eq} = 0 \Rightarrow \rho_{eq} = \frac{A N_1}{B_{01} N_0 - B_{10} N_1}$$

$$\therefore \rho_{eq} = \frac{A}{B_{10}} \cdot \frac{1}{\frac{B_{01}}{B_{10}} \cdot \frac{N_0}{N_1} - 1}$$

Now, from Boltzmann Statistics we know the fraction of each energy state

$$\frac{N_0}{N_0 + N_1} = \frac{g_0 e^{-\beta E_0}}{Z} \quad ; \quad Z = \sum_{i=1}^2 g_i e^{-\beta E_i}$$

where g_i is the multiplicity/degeneracy of E_i

$$\therefore \frac{N_0}{N_1} = \frac{g_0 e^{-\beta E_0}}{g_1 e^{-\beta E_1}} = \frac{g_0}{g_1} e^{\beta(E_1 - E_0)}$$

$$\text{let } h\nu = E_1 - E_0$$

$$\therefore \frac{N_0}{N_1} = \frac{g_0}{g_1} e^{\beta h\nu}$$

$$\therefore \rho_{eq} = \frac{A}{B_{10}} \cdot \frac{1}{\left(\frac{g_0 B_{01}}{g_1 B_{10}} \right) e^{\beta h\nu} - 1} \quad \text{--- (b)}$$

\therefore Comparing eqn (a) & (b)

$$\frac{A}{B_{10}} = \frac{8\pi h}{c^3} \cdot \nu^3 \quad \& \quad \frac{g_0 B_{01}}{g_1 B_{10}} = 1$$

$$\Rightarrow B_{01} = \frac{g_1}{g_0} B_{10}$$

Q. 3.

$$\frac{dN_1}{dt} = -A N_1 + (B_{01} \rho N_0 - B_{10} \rho N_1)$$

$$\begin{aligned} \therefore \frac{dN_1}{dt} &= -A N_1 - B_{10} N_1 \rho + B_{01} \rho (N - N_1) \\ &= -A N_1 - (B_{10} + B_{01}) \rho N_1 + B_{01} \rho N \end{aligned}$$

$$\therefore \frac{dN_1}{dt} + [A + (B_{10} + B_{01}) \rho] N_1 = B_{01} \rho N$$

$$\text{Now, } f_1(t) = \frac{N_1(t)}{N}$$

$$\therefore \frac{d}{dt} f_1(t) + [A + (B_{10} + B_{01}) \rho] f_1(t) = B_{01} \rho$$

$$\therefore \text{I.F.} = e^{\int [A + (B_{10} + B_{01}) \rho] dt} = e^{[A + (B_{10} + B_{01}) \rho] t} = g(t)$$

$$B_{10} + B_{01} = B_{10} + \frac{g_1}{g_0} B_{10} = \left(\frac{g_1 + g_0}{g_0} \right) B_{10}$$

$$\begin{aligned} \therefore f_1(t) e^{(\cdot)} \Big|_0^t &= \int_0^t \exp \left\{ [A + \left(\frac{g_1 + g_0}{g_0} \right) B_{10} \rho] t \right\} B_{01} \rho dt \\ &= \frac{e^{(\cdot)} - 1}{A + \left(\frac{g_1 + g_0}{g_0} \right) B_{10} \rho} \times B_{01} \rho \end{aligned}$$

$$\therefore f_1(t) = f_1(0) + \frac{B_{01} \rho}{A + \left(\frac{g_1 + g_0}{g_0} \right) B_{10} \rho} \left\{ 1 - e^{-[A + \left(\frac{g_1 + g_0}{g_0} \right) B_{10} \rho] t} \right\}$$

Assume all molecules are in the ground state at $t=0$ & assume no spontaneous emission

$$\Rightarrow f_1(0) = 0 \quad \& \quad A = 0$$

$$\therefore f_1(t) = \frac{g_0}{g_1 + g_0} \times \frac{B_{01}}{B_{10}} \times \left[1 - \exp \left\{ -B_{10} \left(\frac{g_1 + g_0}{g_0} \right) t \right\} \right]$$

$$\therefore \frac{B_{01}}{B_{10}} = \frac{g_1}{g_0} ; \quad \frac{g_0}{g_1 + g_0} \cdot \frac{B_{01}}{B_{10}} = \frac{g_1}{g_1 + g_0}$$

$$\therefore f_1(t) = \frac{g_1}{g_1 + g_0} \cdot \left[1 - \exp \left\{ -B_{10} \left(\frac{g_1 + g_0}{g_0} \right) t \right\} \right]$$

This formula is valid under the assumptions that;

- a) start with zero population of excited state
- b) ignore spontaneous emission effects

<plot attached below>