## System Simulation on 1-D Mask

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#### Outline

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Fourier Transform
Sampling Theorem
Discrete Fourier Transform

Simulation

Simulation Systems

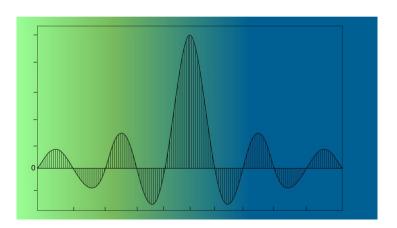
Simulation Results - Single Slit

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- Future Work
- A Reference

#### Introduction

The Nature of Curiosity Deepens our Horizons



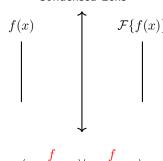
#### Fourier Transform (1/1)

1-D Fourier Transform pair is shown below

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} g(k) e^{jkx} dk \longleftrightarrow g(k) = \int_{-\infty}^{+\infty} f(x) e^{-jkx} dx$$

Below shows a system that has a relationship with Fourier Transform pairs.

#### Condensed Lens



With this arrangement, we will have

$$\mathsf{Output} = \mathcal{F}\{f(x)\} = \mathcal{F}\{\mathsf{Input}\}$$

and the notation  ${\mathcal F}$  represents

$$\mathcal{F}: f(x) \longrightarrow g(k)$$

## Sampling Theorem (1/5)

Assume we have a continuous time signal  $\boldsymbol{x}(t)$  and

$$\mathcal{F}: x(t) \longleftrightarrow X(j\omega)$$

Typically, we will use an impulse train s(t) to discretize the signals that

$$s(t) = \sum_{n=-\infty}^{\infty} \delta(t-nT_s)$$
 where  $T_s$  is the sampling period

Hence, we can obtain x[n] from x(t) and s(t) that

$$x[n] = x(t)s(t) = x(t)\sum_{n=-\infty}^{+\infty} \delta(t - nT_s)$$

With sifting property, we rewrite it as

$$x[n] = x(nT_s)$$



## Sampling Theorem (2/5)

Since s(t) is periodic, we can find its FS coefficient

$$\mathcal{FS}: s(t) = \sum_{k=-\infty}^{\infty} S[k]e^{jk\omega_0 t} \longleftrightarrow S[k] = \frac{1}{T}$$

and shifting property that

$$\mathcal{F}: e^{j\omega_0 t} \longleftrightarrow 2\pi\delta(\omega - \omega_0)$$

Thus, we can get  $S(j\omega)$  that

$$S(j\omega) = \frac{2\pi}{T} \sum_{k=-\infty}^{\infty} \delta(\omega - k\omega_0)$$

Hence, for signal x(t)s(t), we could have the complete form that

$$\mathcal{F}: x(t)s(t) \longleftrightarrow \frac{1}{2\pi}X(j\omega) * S(jw)$$

# Sampling Theorem (3/5)

Consider a signal of x(t) having this  $X(j\omega)$  and impulse train s(t).

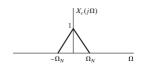


Figure:  $X(j\omega)$  of signal x(t)

It is bandlimited because

$$X(j\omega) = 0 \quad |\omega| \ge \omega_N$$

It is better to sample bandlimited signals.



Figure:  $S(j\omega)$  of signal s(t)

We can choose the sampling period  $T_s$ .

$$\omega_s = \frac{2\pi}{T_s}$$

# Sampling Theorem (4/5)

Consider  $x_s(t) = x(t)s(t)$  with different sampling period.

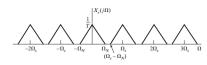


Figure:  $X_s(j\omega)$  of signal  $x_s(t)$ 

The sampling period is quite small because

$$\omega_s > \omega_N$$

which ensures a intact pattern of  $X(j\omega)$ .

The sampling period is quite large because

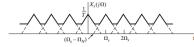


Figure:  $X_s(j\omega)$  of signal  $x_s(t)$ 

 $\omega_s < \omega_N$ 

which shows a mixed pattern of  $X(j\omega)$ .

The importance of the sampling period  $T_s$  is now shown above.

#### Sampling Theorem (5/5)

Generally, a mixed pattern for a bandlimited signal is called aliasing.

#### Nyquist-Shannon Sampling Theorem

Let  $x_c(t)$  be a bandlimited signals with

$$X_c(j\omega) \approx 0 \quad \text{for } |\omega| \ge \omega_N$$

Then  $x_c(t)$  is uniquely determined by its samples  $x[n] = x_c(nT_s)$  where

$$\omega_s = \frac{2\pi}{T_s} \ge 2\omega_N$$

Also, it is also referred to as

 $\omega_N \longrightarrow \mathsf{Nyquist}$  frequency

 $2\omega_N \longrightarrow \mathsf{Nyquist}$  rate

To prevent aliasing, it is required to sample at Nyquist rate where

$$T_s < \frac{\pi}{\omega_N} \Longrightarrow f_s > \frac{\omega_N}{\pi}$$

## Discrete Fourier Transform (1/3)

During simulation, we still can't get a continuous frequency response. Recall that

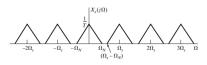


Figure:  $X_s(j\omega)$  of signal  $x_s(t)$ 

Observation :

- Input : discrete
- Output : continuous

Problem occurs with continuous signals.

Is there any method to solve this problem? DFT.

## Discrete Fourier Transform (2/3)

Discrete Fourier Transform is used for finite discrete signals.

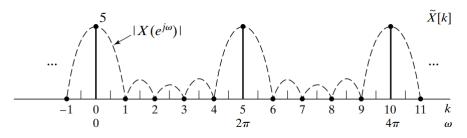


Figure: Example for relationship between  $X(e^{j\omega})$  and X[k]

The picture shown above easily depicts X[k] just samples from  $X(e^{j\omega})$ .

# Discrete Fourier Transform (3/3)

What are the benefits using DFT ?

- It provides a discrete structure for computers to store.
- It makes the number of sampling from input and output the same.

It will be helpful with Nyquist-Shannon Sampling Theorem.

- Find the appropriate sampling frequency at Nyquist rate.
- Divide this valid region to the number of input.

Hence, it is easy to get the value of every discrete point in frequency space.

frequency of 
$$X[k] = \frac{k}{N} \cdot \omega_s$$

because the valid region based on  $\omega_s$ .



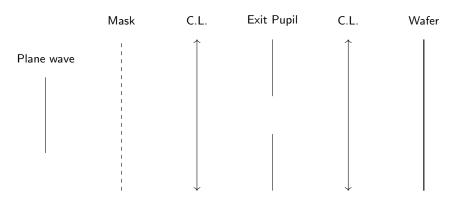
#### Simulation

#### Putting Everything into Practice Prevents Cheap Talks

```
requests.get(url)
                                    .oad from the website
       # checking response.status_code (if you get SR2, to
       if response.status_code != 200:
              print(f"Status: {response.status_code} - Try renaming the control
        else:
              print(f"Status: {response.status_code)\n")
        # using BeautifulSoup to parse the response edges
6
        SOUP = BeautifulSoup(response.content, "tral_pareer")
        images = soup, find_all("img", attractatt: "max ame")
18
19
20
        # downloading images
21
22
```

## Simulation Systems (1/4)

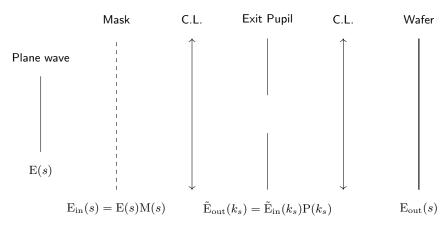
Below shows the structure of our simulation system.



Note that C.L. means condensed lens.

# Simulation Systems (2/4)

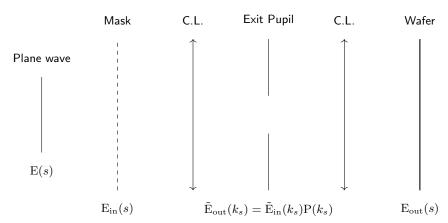
We assume the focal length of both C.L. are f, which Fraunhofer diffraction is FT.



The states are shown above respectively.

# Simulation Systems (3/4)

Let's take a closer look at all the states.



All relationships between  $E_{\{\cdot\}}$  and  $\tilde{E}_{\{\cdot\}}$  are FT pairs.

# Simulation Systems (4/4)

Simulations are under these conditions.

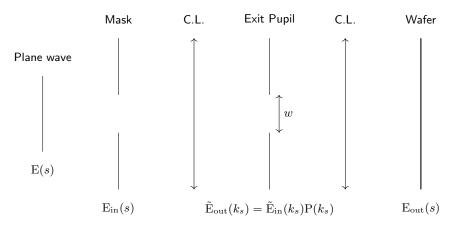
- Numerical Aperature N.A. = 0.86
- $\lambda = 248 \ nm$
- $\bullet \ \ \text{The wave number} \ k = \frac{2\pi}{\lambda}$
- The width of Exit pupil  $w = \frac{\mathrm{N.A.}}{\lambda}$
- Wafer Intensity Threshold = 0.5 (Normalized)

These conditions will be pointed out when we see the structure of the system.

Single Slit

### Simulation Results - Single Slit (1/11)

Take an easy look at the system using a single slit as mask.



If we consider a plane wave traveling at an angle of 0, we set  $E_{\mathrm{in}}(s)$  as a constant.

#### Simulation Results - Single Slit (2/11)

Since there is a rect is at "Mask", we expect a sinc before "Exit Pupil".

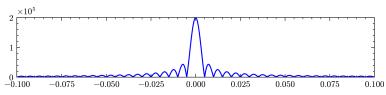


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

Also, at "Exit Pupil", the distribution will be affected by the size of it.

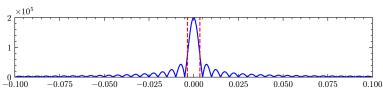


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{\mathrm{E_{in}}(s)\}$  and Cut-off Region

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#### Simulation Results - Single Slit (3/11)

Hence, the distribution at "Exit Pupil" is shown below.

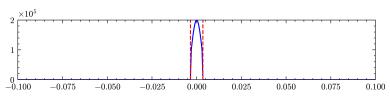


Figure: Absolute Value of Cut-off  $\mathcal{F}\{\mathrm{E_{in}}(s)\}$  after Exit Pupil

With the last condensed lens, the aerial image is just an inverse Fourier Transform.

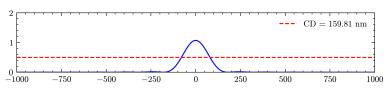
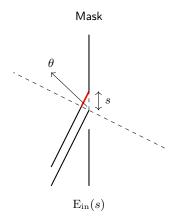


Figure: Magnitude Distribution on the Wafer

#### Simulation Results - Single Slit (4/11)

Now consider a plane wave traveling at an angle of  $\theta$  with respect to the x-axis.



The plot shows that

position s leads to path difference

For convenience, it is good to choose that

- ullet constant at the center because s=0
- others multiplied by a factor  $e^{jks\sin\theta}$ , ex. the red part

#### Simulation Results - Single Slit (5/11)

Let's see the results when traveling at angle  $20^{\circ}.$ 

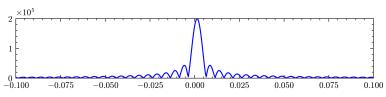


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

The pattern seems to have a shifting compared to one traveling at angle  $0^{\circ}$ .

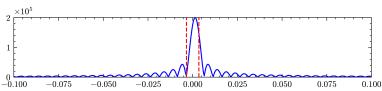


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{\mathrm{E_{in}}(s)\}$  and Cut-off Region

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## Simulation Results - Single Slit (6/11)

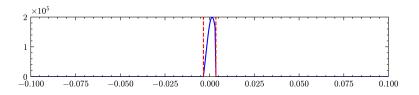


Figure: Absolute Value of Cut-off  $\mathcal{F}\{\mathrm{E_{in}}(s)\}$  after Exit Pupil

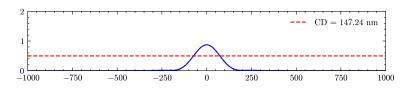


Figure: Magnitude Distribution on the Wafer

#### Simulation Results - Single Slit (7/11)

Let's see the results when traveling at angle  $40^{\circ}$ .

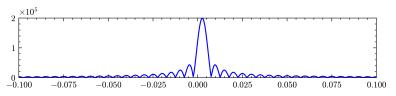


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

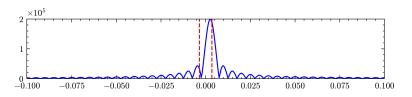


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{\mathrm{E_{in}}(s)\}$  and Cut-off Region

## Simulation Results - Single Slit (8/11)

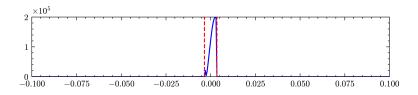


Figure: Absolute Value of Cut-off  $\mathcal{F}\{\mathrm{E_{in}}(s)\}$  after Exit Pupil

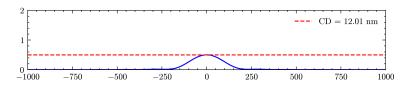


Figure: Magnitude Distribution on the Wafer

#### Simulation Results - Single Slit (9/11)

Let's see the results when traveling at angle  $60^{\circ}$ .

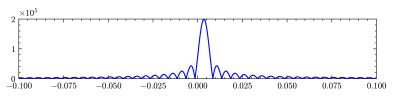


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

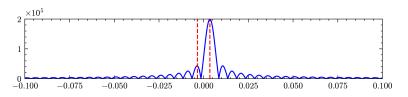


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$  and Cut-off Region

## Simulation Results - Single Slit (10/11)

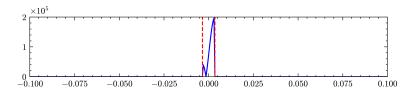


Figure: Absolute Value of Cut-off  $\mathcal{F}\{\mathrm{E_{in}}(s)\}$  after Exit Pupil

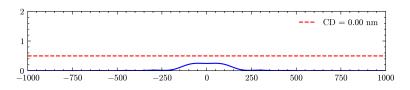


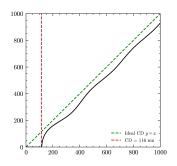
Figure: Magnitude Distribution on the Wafer

#### Simulation Results - Single Slit (11/11)

From these results, we can observe that when the pattern shifts too much

- Dominant intensity parts are lost.
- Measured C.D. are decreased.

Thus, let's change the width to see what happens when traveling at an angle 0.



We can see that

- Measured C.D. is getting closed to C.D.
- $\bullet$  No patterns under approximately 116~nm

Figure: C.D. v.s. Measured C.D.

Dense Line

#### Simulation Results - Dense Line (1/10)

Discuss the signal traveling at an angle of 0, which the distribution is periodic.

$$x(t) = \operatorname{rect}(t) * s_p(t) \longleftrightarrow X(j\omega) = \tilde{\operatorname{rect}}(j\omega)S(j\omega)$$

Thus,  $X(j\omega)$  can be regarded as sampling from a  $\operatorname{rect}(j\omega)$  which is a sinc.

By Nyquist Theorem, it is possible to find

$$X(j\omega) \approx 0 \Longrightarrow f_s = \frac{\omega_N}{\pi} \neq \frac{1}{p}$$

Also, we can expect a impulse graph when we start simulating.

### Simulation Results - Dense Line (2/10)

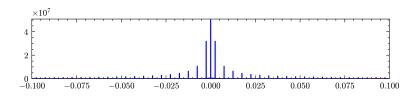


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

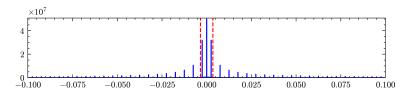


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

#### Simulation Results - Dense Line (3/10)

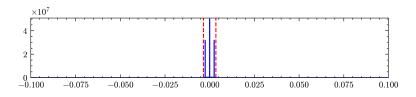


Figure: Absolute Value of Cut-off  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$  after Exit Pupil

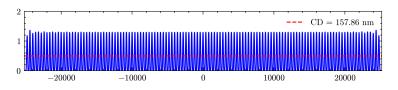


Figure: Magnitude Distribution on the Wafer

### Simulation Results - Dense Line (4/10)

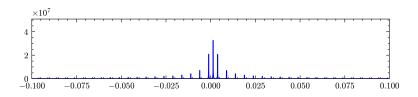


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

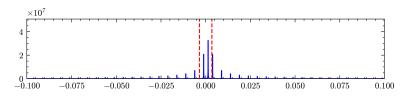


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

#### Simulation Results - Dense Line (5/10)

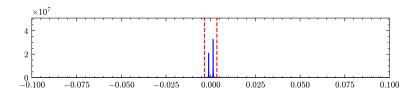


Figure: Absolute Value of Cut-off  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$  after Exit Pupil

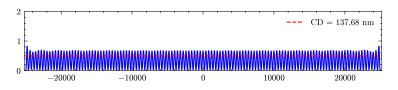


Figure: Magnitude Distribution on the Wafer

### Simulation Results - Dense Line (6/10)

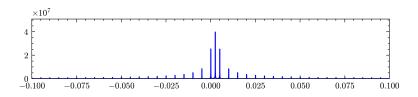


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

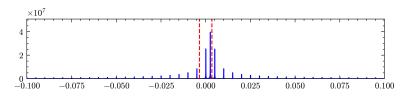


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

#### Simulation Results - Dense Line (7/10)

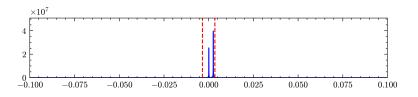


Figure: Absolute Value of Cut-off  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$  after Exit Pupil

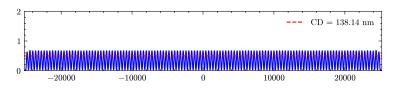


Figure: Magnitude Distribution on the Wafer

## Simulation Results - Dense Line (8/10)

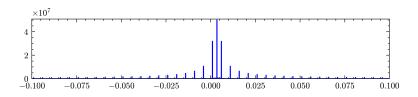


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

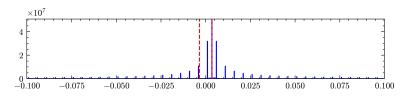


Figure: Absolute Value of Fourier Transform  $\mathcal{F}\{E_{\mathrm{in}}(s)\}$ 

#### Simulation Results - Dense Line (9/10)

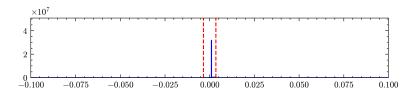


Figure: Absolute Value of Cut-off  $\mathcal{F}\{\mathrm{E_{in}}(s)\}$  after Exit Pupil

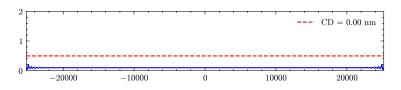


Figure: Magnitude Distribution on the Wafer

## Simulation Results - Dense Line (10/10)

From these results, we can observe that

- The Fourier Transform patterns truly comprise impulses.
- At least 2 impulses to have pattern on Wafer.

Hence, it is important to let C.D. have enough width based on the observations.

#### Rayleigh Criterion

The Rayleigh Criterion equation is

$$C.D. = k_1 \frac{\lambda}{N.A.}$$

where

C.D. is critical dimension

 $\lambda$  is the wavelength of light

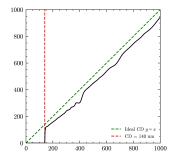
N.A is the numerical aperture

Physical limit of  $k_1$  is 0.25.

#### Simulation Results - Dense Line (11/)

From these results, we can observe that when the pattern shifts too much

- Dominant intensity parts are lost.
- Measured C.D. are decreased.



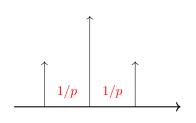
Also, if we just change the width (C.D.)

- distance between impulses gets closer
- pattern shows when features are enough

Figure: C.D. v.s. Measured C.D.

## Simulation Results - Dense Line (12/)

As the distance between each order is 1/p, the sampling frequency of  $s_p(t)$ .



For the dominant part, we have

$$\frac{1}{p} = \frac{1}{2 \text{ C.D.}} = \frac{\text{N.A.}}{\lambda}$$

Thus, it can be found that

$$k_1 = 0.5$$

#### Future Work

Discovering new ideas is a boost for the future.



#### Future Work

As we complete 1-D simulation, it helps me get the structure of the system.

- How a point source affects the aerial image.
- What single slits patterns and dense line patterns are.

However, it is an obstacle for me to simulate from source to image.

- Refer to books to know how a lens affects in a system.
- It is required to know 2-D Sampling.

Last, computing such a big matrix is a big problem.

- Find the minimum sampling rate to compute less.
- Learn how to use parallel computing.

#### Reference

- F. Pedrotti, Introduction to Optics, 3rd Edition
  - Chapter 21 Fourier Optics
- Oppenheim, Discrete-Time Signal Processing, 3rd Edition
  - Chapter 4 Sampling of Continuous-Time Signals
  - Chapter 8 The Discrete Fourier Transform