

1. Free motion

$$S(q) = \int_{t'}^{t''} dz \frac{1}{2} m \dot{q}^2(z)$$

$$\hat{H} = \frac{\hat{p}^2}{2m}, \quad \hat{V} = 0$$

Integrating by parts:

$$\int_{t'}^{t''} dz \dot{q}(z)^2 = q(z) \dot{q}(z) \Big|_{t'}^{t''} - \int_{t'}^{t''} dz q(z) \frac{d^2}{dz^2} q(z)$$

↳ Gaussian integral w. $A = -m \frac{d^2}{dz^2}$

$$q(t) = q_c(t) + r(t), \text{ with } q_c(t') = q', \quad q_c(t'') = q'', \quad r(t') = r(t'') = 0$$

$$S(q_c + r) = S(q_c) + S(r) + m \int_{t'}^{t''} dz \dot{q}_c(z) \dot{r}(z).$$

$$\text{and } \int_{t'}^{t''} dz \dot{q}_c(z) \dot{r}(z) = \cancel{\dot{q}_c(z) r(z)} \Big|_{t'}^{t''} - \int_{t'}^{t''} dz \ddot{q}_c(z) r(z)$$

$m \ddot{q}_c(z) = 0 \Rightarrow$ linear term in $r(t)$ vanishes.

$$\text{Then: } q_c(z) = q' + \frac{z-t'}{t''-t'} (q''-q')$$

$$\Rightarrow S(q_c) = \frac{m}{2} \frac{(q''-q')^2}{t''-t'}$$

$$\Rightarrow \langle q'' | e^{-i \hat{H} (t''-t')} | q \rangle = \mathcal{N} \exp \left\{ + i \frac{m}{2} \frac{(q''-q')^2}{t''-t'} \right\}$$

The normalization factor is given by:

$$\mathcal{N} = \int \mathcal{D}r \exp \left\{ i \frac{m}{2} \int_{t'}^{t''} d\tau \dot{r}(\tau)^2 \right\}.$$

↳ independent of q', q'' (i.e. independent of the b.c.).

2. Harmonic oscillator

$$\hat{H} = \frac{\hat{P}^2}{2} + \frac{1}{2} \omega^2 \hat{Q}^2 \quad m=1$$

$$S(q) = \int_{t'}^{t''} dt \left[\frac{1}{2} \dot{q}(t)^2 - \frac{1}{2} \omega^2 q(t)^2 \right]$$

$$\langle q'' t'' | q' t' \rangle = \int_{q' q''} \mathcal{D}q \exp \left\{ i S(q) \right\}$$

Matrix $A \rightarrow$ diff. op. $\left(-\frac{d^2}{dt^2} - \omega^2 \right)$

In order to separate the dependence on $q', q'' \rightarrow$ change of var.

$$q(t) = q_c(t) + r(t)$$

$$\begin{aligned} S(q_c + r) &= S(q_c) + S(r) + \int_{t'}^{t''} d\tau \left[\dot{q}_c(\tau) \dot{r}(\tau) - \omega^2 q_c(\tau) r(\tau) \right] \\ &= S(q_c) + S(r) + \int_{t'}^{t''} d\tau \left[-\ddot{q}_c(\tau) - \omega^2 q_c(\tau) \right] r(\tau) \end{aligned}$$

Classical sol.n

$$\ddot{q}_c(z) = -\omega^2 q_c(z)$$

$$\Rightarrow S(q) = S(q_c) + S(r)$$

$$\text{Hence: } \langle q''t'' | q't' \rangle = \mathcal{N}(T, \omega) \exp \{ i S(q_c) \}.$$

$$\mathcal{N}(\omega, T) = \int \mathcal{D}r \exp \left\{ i \int_{t'}^{t''} dt \left[\frac{1}{2} \dot{r}(t)^2 - \frac{1}{2} \omega^2 r(t)^2 \right] \right\},$$

$$\text{and } r(t') = r(t'') = 0.$$

\Rightarrow divergences in the normalization factor $\mathcal{N}(T, \omega)$.

3. Solution of Schrödinger eq. [IGNORE THIS!]

$$\hat{H} = \frac{\hat{p}^2}{2m}$$

$$|p\rangle = \int dq e^{iqp} |q\rangle$$

$$\langle p | p' \rangle = (2\pi) \delta(p - p')$$

Time evolution operator $U(t, t')$

$$i \frac{\partial}{\partial t} U(t, t') = \hat{H} U(t, t')$$

$$\Rightarrow i \frac{\partial}{\partial t} \langle p | U(t, t') | q \rangle = \langle p | \hat{H} U(t, t') | q \rangle$$

$$= \frac{p^2}{2m} \langle p | U(t, t') | q \rangle$$

$$\text{b.c.: } \langle p | U(t, t') | q \rangle = \langle p | q \rangle = e^{-i p q}$$

$$\Rightarrow \langle p | U(t, t') | q \rangle = \exp \left[-i q \cdot p - i (t - t') \frac{p^2}{2m} \right]$$

$$\langle q | U(t, t') | q' \rangle = \int \frac{dp}{2\pi} \langle q | p \rangle \langle p | U(t, t') | q' \rangle$$

$$= \int \frac{dp}{2\pi} \exp \left[-i (q - q') \cdot p - i (t - t') \frac{p^2}{2m} \right]$$

$$\propto \exp \left[i \frac{m (q - q')^2}{2 (t - t')} \right].$$

3. Weyl ordering

In order to compute $[A^2 B]$, consider $(\alpha A + \beta B)^3$ and expand.

$$(\alpha A + \beta B)(\alpha A + \beta B) = \alpha^2 A^2 + \alpha \beta (AB + BA) + \beta^2 B^2$$

$$\Rightarrow (\alpha A + \beta B)^3 = \dots + \alpha^2 \beta (A^2 B + ABA + BA^2) + \dots$$

$$[A^2 B] = (A^2 B + ABA + BA^2) \left(\frac{2}{3!} \right) \rightarrow 1/3$$

4. Schrödinger eq. again

$$\langle q'', t'' | q', t' \rangle = U(q', q'', T)$$

$$T = t'' - t'$$

$$\langle q'' t'' | q' t' \rangle = \lim_{n \rightarrow \infty} \frac{1}{(2\pi i \epsilon)^{1/2}} \int \prod_{k=1}^{n-1} \frac{dx_k}{(2\pi i \epsilon)^{1/2}} \times \exp \left\{ i \epsilon \sum_{m=1}^n \left[\frac{1}{2} \left(\frac{x_m - x_{m-1}}{\epsilon} \right)^2 - V \left(\frac{x_m + x_{m-1}}{2} \right) \right] \right\}.$$

$$= \int \frac{dx_{n-1}}{(2\pi i \epsilon)^{1/2}} \exp \left\{ i \epsilon \left[\frac{1}{2} \left(\frac{x_n - x_{n-1}}{\epsilon} \right)^2 - V \left(\frac{x_n + x_{n-1}}{2} \right) \right] \right\} \times$$

$$\times \frac{1}{(2\pi i \epsilon)^{1/2}} \int \prod_{k=1}^{n-2} \frac{dx_k}{(2\pi i \epsilon)^{1/2}} \exp \left\{ i \epsilon \sum_{j=1}^{n-1} [\dots] \right\}.$$

$$= \int \frac{dx_{n-1}}{(2\pi i \epsilon)^{1/2}} \exp \left\{ i \epsilon \left[\frac{1}{2} \left(\frac{x_n - x_{n-1}}{\epsilon} \right)^2 - V \left(\frac{x_n + x_{n-1}}{2} \right) \right] \right\} U(q'; x_{n-1}, T-\epsilon)$$

Expand the integrand around $x_{n-1} = x_n = q''$

$$U(q', x_{n-1}, T-\epsilon) = U(q', q'', T-\epsilon) + \epsilon \frac{\partial}{\partial x} U(q', x, T-\epsilon) \Big|_{x=q''} +$$

$$+ \frac{1}{2} \frac{\partial^2}{\partial x^2} U(q', x, T-\epsilon) \Big|_{x=q''} \epsilon^2 + \dots$$

$$z = x_{n-1} - q'' \quad , \quad dz = dx_{n-1}$$

$$\langle q'' t'' | q' t' \rangle = e^{-i \epsilon V(q'')} \int \frac{dz}{(2\pi i \epsilon)^{1/2}} \exp \left\{ -\frac{1}{2 \epsilon i} z^2 \right\} \left\{ U(q', q'', T-\epsilon) + \right.$$

$$\left. + \epsilon \frac{\partial}{\partial x} U(q', x, T-\epsilon) \Big|_{x=q''} + \frac{1}{2} \frac{\partial^2}{\partial x^2} U(q', x, T-\epsilon) \Big|_{x=q''} \epsilon^2 + \dots \right\}$$

Gaussian integral: $\int dx e^{-ax^2} = \sqrt{\frac{\pi}{a}}$

$$\int dx x^2 e^{-ax^2} = \frac{1}{2a} \sqrt{\frac{\pi}{a}}$$

$$\Rightarrow \langle q'' t'' | q' t' \rangle = e^{-i \epsilon V(q'')} \frac{1}{(2\pi i \epsilon)^{1/2}} \left\{ (2\pi i \epsilon)^{1/2} + (2\pi i \epsilon)^{1/2} \frac{i \epsilon}{2} \frac{\partial^2}{\partial x^2} \right\} \times$$

$$\times U(q', x, T-\epsilon) \Big|$$

$$|x = q''$$

Expand in ϵ :

$$i \frac{\partial}{\partial \tau} U(q', q'', \tau) = \left[-\frac{\partial^2}{\partial x^2} + V(x) \right] U(q', x, \tau) \Big|_{x=q''}$$