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IYGB-MATHEMATICAL METHODS 3 - PAPER A - QUESTION 1

- $f(z)$ HAS A SIMPLE POLE AT $z=0$, WHICH IS VERY EASY TO FIND FROM ITS LAURENT EXPANSION

$$f(z) = \frac{\sin z}{z^2} = \frac{1}{z^2} \left[z - \frac{z^3}{3!} + \frac{z^5}{5!} + O(z^7) \right]$$

$$= \frac{1}{z} - \frac{1}{6}z + \frac{1}{120}z^3 + O(z^5)$$

∴ RESIDUE IS 1



- ALTERNATIVE IS TO USE THE STANDARD METHOD FOR A SIMPLE POLE AT $z=a$

$$\lim_{z \rightarrow a} \left[\frac{d}{dz} \left[(z-a)^2 f(z) \right] \right] = \lim_{z \rightarrow 0} \left[\frac{d}{dz} \left[z \cdot \frac{\sin z}{z^2} \right] \right]$$

$$= \lim_{z \rightarrow 0} \left[\frac{d}{dz} \left(\frac{\sin z}{z} \right) \right]$$

BY L'HOSPITAL RULE

$$= \lim_{z \rightarrow 0} \left(\frac{\cos z}{1} \right)$$

$$= 1$$

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AS REQUIRED

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① STARTING WITH A SUBSTITUTION

$$\Rightarrow t = \ln\left(\frac{1}{x}\right)$$

$$\Rightarrow t = -\ln x$$

$$\Rightarrow -t = \ln x$$

$$\Rightarrow x = e^{-t}$$

$$\Rightarrow dx = -e^{-t} dt$$

LIMITS

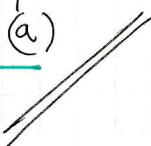
$$x=0 \rightarrow t = \infty$$

$$x=1 \rightarrow t = 0$$

② THE INTEGRAL TRANSFORMS TO

$$\begin{aligned} \int_0^1 \left[\ln\left(\frac{1}{x}\right) \right]^{a-1} dx &= \int_{\infty}^0 t^{a-1} (-e^{-t}) dt \\ &= \int_0^{\infty} t^{a-1} e^{-t} dt \end{aligned}$$

WHAT IS THE EXACT DEFINITION OF $\Gamma(a)$



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① WE PROCEED AS FOLLOWS

$$\frac{1}{z} = \frac{1}{1+(z-1)} = \frac{1}{z-1} \left(\frac{1}{1+\frac{1}{z-1}} \right)$$

② EXPAND USING $\frac{1}{1+x} = 1-x+x^2+x^3-\dots$ $|x| < 1$

$$\Rightarrow \frac{1}{z} = \frac{1}{z-1} \left[1 - \frac{1}{(z-1)} + \frac{1}{(z-1)^2} - \frac{1}{(z-1)^3} + \frac{1}{(z-1)^4} - \dots \right]$$

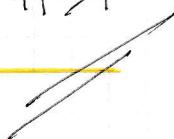
$$\Rightarrow \frac{1}{z} = \frac{1}{z-1} - \frac{1}{(z-1)^2} + \frac{1}{(z-1)^3} - \frac{1}{(z-1)^4} + \dots$$

VALID FOR $\left| \frac{1}{z-1} \right| < 1$

$$|z-1| > 1$$

$$\therefore \frac{1}{z} = \sum_{r=1}^{\infty} \frac{(-1)^{r+1}}{(z-1)^r}$$

FOR ANNUAL $|z-1| > 1$



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- START BY A SUGGESTING SUBSTITUTION

$$\Rightarrow u = 8x^3$$

$$\Rightarrow x^3 = \frac{1}{8}u$$

$$\Rightarrow x = \frac{1}{2}u^{\frac{1}{3}}$$

$$\Rightarrow dx = \frac{1}{6}u^{-\frac{2}{3}}du$$

& UNITS ARE UNCHANGED

- THE INTEGRAL TRANSFORMS TO

$$\int_0^\infty \frac{x^3}{(1+8x^3)^2} dx = \int_0^\infty \frac{\frac{1}{8}u}{(1+u)^2} \left(\frac{1}{6}u^{-\frac{2}{3}} du \right) = \frac{1}{48} \int_0^\infty \frac{u^{\frac{1}{3}}}{(1+u)^2} du$$

- NOW USING THE ALTERNATIVE FORM OF A BETA FUNCTION

$$B(m, n) = \int_0^\infty \frac{u^{m-1}}{(1+u)^{m+n}} du$$

- THE INTEGRAL BECOMES

$$\dots = \frac{1}{48} \int_0^\infty \frac{u^{\frac{4}{3}-1}}{(1+u)^{\frac{4}{3}+\frac{2}{3}}} du = \frac{1}{48} B\left(\frac{4}{3}, \frac{2}{3}\right) = \frac{1}{48} \frac{\Gamma\left(\frac{4}{3}\right)\Gamma\left(\frac{2}{3}\right)}{\Gamma\left(\frac{4}{3} + \frac{2}{3}\right)}$$

$$= \frac{1}{48} \frac{\Gamma\left(\frac{4}{3}\right)\Gamma\left(\frac{2}{3}\right)}{\Gamma(2)} = \frac{1}{48} \frac{\frac{1}{3}\Gamma\left(\frac{1}{3}\right)\Gamma\left(\frac{2}{3}\right)}{1!} = \frac{1}{144} \Gamma\left(\frac{1}{3}\right)\Gamma\left(\frac{2}{3}\right)$$

- FINALLY USING

$$\Gamma(x)\Gamma(1-x) \equiv \frac{\pi}{\sin \pi x}$$

$$= \frac{1}{144} \Gamma\left(\frac{1}{3}\right)\Gamma\left(1-\frac{1}{3}\right) = \frac{1}{144} \frac{\pi}{\sin \frac{\pi}{3}} = \frac{1}{144} \frac{\pi}{\frac{\sqrt{3}}{2}} = \frac{\pi}{72\sqrt{3}}$$

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a)

START BY DIFFERENTIATING THE GENERATING FUNCTION FOR $J_y(x)$, WITH RESPECT TO x

$$\Rightarrow e^{\frac{1}{2}x(t-\frac{1}{t})} = \sum_{n=-\infty}^{\infty} [t^n J_n(x)]$$

$$\Rightarrow \frac{1}{2}(t-\frac{1}{t})e^{\frac{1}{2}x(t-\frac{1}{t})} = \sum_{n=-\infty}^{\infty} [t^n J'_n(x)]$$

$$\Rightarrow \frac{1}{2}te^{\frac{1}{2}x(t-\frac{1}{t})} - \frac{1}{2t}e^{\frac{1}{2}x(t-\frac{1}{t})} = \sum_{n=-\infty}^{\infty} [t^n J'_n(x)]$$

$$\Rightarrow \frac{1}{2}t \sum_{n=-\infty}^{\infty} [t^n J_n(x)] - \frac{1}{2t} \sum_{n=-\infty}^{\infty} [t^n J_n(x)] = \sum_{n=-\infty}^{\infty} [t^n J'_n(x)]$$

$$\Rightarrow \frac{1}{2} \sum_{n=-\infty}^{\infty} [t^{n+1} J_n(x)] - \frac{1}{2} \sum_{n=-\infty}^{\infty} [t^n J'_n(x)] = \sum_{n=-\infty}^{\infty} [t^n J'_n(x)]$$

EQUATING POWERS OF t , SAY $[t^n]$, GIVES

$$\Rightarrow \frac{1}{2} t^n J_{n-1}(x) - \frac{1}{2} t^n J'_{n+1}(x) = t^n J'_n(x)$$

$$\Rightarrow \frac{1}{2} J_{n-1}(x) - \frac{1}{2} J_{n+1}(x) = J'_n(x)$$

// ~~AS REQUIRED~~

b)

NEXT DIFFERENTIATE WITH RESPECT TO t

$$\Rightarrow e^{\frac{1}{2}x(t-\frac{1}{t})} = \sum_{n=-\infty}^{\infty} [t^n J_n(x)]$$

$$\Rightarrow \frac{1}{2}x(1+\frac{1}{t^2})e^{\frac{1}{2}x(t-\frac{1}{t})} = \sum_{n=-\infty}^{\infty} [nt^{n-1} J_n(x)]$$

$$\Rightarrow \frac{1}{2}x e^{\frac{1}{2}x(t-\frac{1}{t})} + \frac{x}{2t^2} e^{\frac{1}{2}x(t-\frac{1}{t})} = \sum_{n=-\infty}^{\infty} [nt^{n-1} J_n(x)]$$

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$$\Rightarrow \frac{1}{2}x \sum_{n=-\infty}^{\infty} [t^n J_n(x)] + \frac{x}{2t^2} \sum_{n=-\infty}^{\infty} [t^n \bar{J}_n(x)] = \sum_{n=-\infty}^{\infty} [nt^{n-1} J_n(x)]$$

$$\Rightarrow \frac{1}{2}x \sum_{n=-\infty}^{\infty} [t^n J_n(x)] + \frac{1}{2}x \sum_{n=-\infty}^{\infty} [t^{n-2} \bar{J}_n(x)] = \sum_{n=-\infty}^{\infty} [nt^{n-1} J_n(x)]$$

EQUATING POWERS OF t , SAY $[t^{n-1}]$ GIVES

$$\Rightarrow \frac{1}{2}x [t^{n-1} J_{n-1}(x)] + \frac{1}{2}x [t^{n-1} \bar{J}_{n+1}(x)] = nt^{n-1} J_n(x)$$

$$\Rightarrow \frac{1}{2}x J_{n-1}(x) + \frac{1}{2}x \bar{J}_{n+1}(x) = n J_n(x)$$

$$\Rightarrow J_n(x) = \frac{x}{2n} [J_{n-1}(x) + \bar{J}_{n+1}(x)]$$

// AS REQUIRED

c) DIFFERENTIATING, USING THE PRODUCT RULE, WITH RESPECT TO x

$$\frac{d}{dx} [(x^n + \bar{x}^n) J_n(x)] = (nx^{n-1} - n\bar{x}^{n-1}) J_n(x) + (x^n + \bar{x}^n) J'_n(x)$$

USING PART (b) AND PART (a) IN THE RESPECTIVE TERMS GIVES

$$= (nx^{n-1} - n\bar{x}^{n-1}) \frac{x}{2n} [J_{n-1}(x) + \bar{J}_{n+1}(x)] + (x^n + \bar{x}^n) \left[\frac{1}{2} \bar{J}_{n-1}(x) - \frac{1}{2} J_{n+1}(x) \right]$$

$$= \frac{1}{2} (x^n - \bar{x}^n) [J_{n-1}(x) + \bar{J}_{n+1}(x)] + \frac{1}{2} (x^n + \bar{x}^n) [J_{n-1}(x) - \bar{J}_{n+1}(x)]$$

$$= \frac{1}{2} \left[\cancel{x^n J_{n-1}(x)} + \cancel{x^n \bar{J}_{n+1}(x)} - \cancel{\bar{x}^n J_{n-1}(x)} - \cancel{\bar{x}^n \bar{J}_{n+1}(x)} \right]$$

$$= x^n J_{n-1}(x) - \bar{x}^n \bar{J}_{n+1}(x)$$

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IYGB - MATHEMATICAL METHODS 3 - PAPER A - QUESTION 6

$$f(t) = \int^{-1} \left[\frac{1}{s^2} (2 - 2e^{-4s} - e^{-6s}) \right]$$

① INVERTING INTO "HAWKES"

$$f(t) = \int^{-1} \left[\frac{2}{s^2} - \frac{2e^{-4s}}{s^2} - \frac{e^{-6s}}{s^2} \right]$$

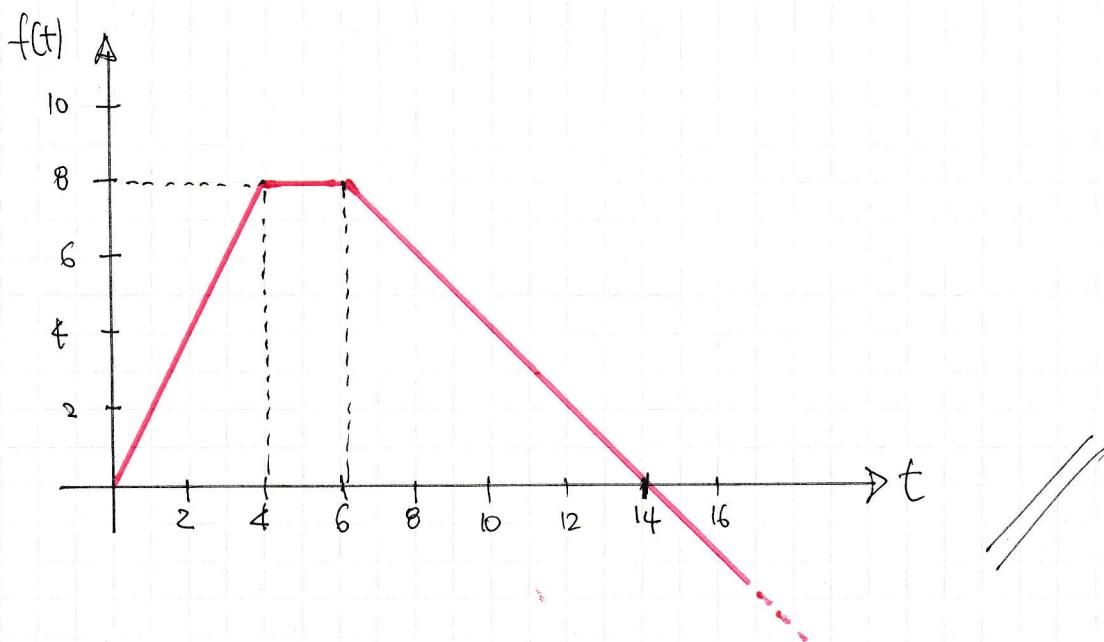
$$f(t) = 2t - 2(t-4)H(t-4) - (t-6)H(t-6)$$

$$f(t) = 2tH(t) - 2(t-4)H(t-4) - (t-6)H(t-6)$$

② USE A TABLE TO EXTRACT THE FUNCTION

INTERVAL	$2tH(t)$	$-2(t-4)H(t-4)$	$-(t-6)H(t-6)$	$f(t)$
$0 \leq t \leq 4$	$2t$	0	0	$2t$
$4 < t \leq 6$	$2t$	$8-2t$	0	8
$t > 6$	$2t$	$8-2t$	$6-t$	$14-t$

③ FINALLY BEFORE SKETCHING CHECK FOR CONTINUITY AT $x=4, x=6$



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$$(x^2-1) \frac{d^2y}{dx^2} + 4x \frac{dy}{dx} + 2y = 0$$

- AS THE O.D.E IS ANALYTIC AT $x=0$ (WHICH WE DIVIDE BY x^2-1 , THE POLES AT ± 1), WE MAY SEEK FOR A SOLUTION OF THE FORM

$$y = \sum_{r=0}^{\infty} a_r x^r = a_0 + a_1 x + a_2 x^2 + a_3 x^3 + a_4 x^4 + \dots$$

- DIFFERENTIATE WITH RESPECT TO x

$$\frac{dy}{dx} = \sum_{r=1}^{\infty} a_r r x^{r-1}, \quad \frac{d^2y}{dx^2} = \sum_{r=2}^{\infty} a_r r(r-1) x^{r-2}$$

- SUBSTITUTE INTO THE O.D.E.

$$\Rightarrow (x^2-1) \sum_{r=2}^{\infty} a_r r(r-1) x^{r-2} + 4x \sum_{r=1}^{\infty} a_r r x^{r-1} + 2 \sum_{r=0}^{\infty} a_r x^r = 0$$

$$\Rightarrow \sum_{r=2}^{\infty} a_r r(r-1) x^r - \sum_{r=2}^{\infty} a_r r(r-1) x^{r-2} + \sum_{r=1}^{\infty} 4a_r r x^r + \sum_{r=0}^{\infty} 2a_r x^r = 0$$

↑ ↑ ↑ ↑
 LOWEST POWER IS x^2 LOWEST POWER IS x^0 LOWEST POWER IS x^1 LOWEST POWER OF x^0

- PULL OUT x^0 AND x^1 OUT OF THE SUMMATIONS

$$\Rightarrow \sum_{r=2}^{\infty} a_r r(r-1) x^r - a_2 \times 2 \times 1 \times x^0 + 4a_1 \times 1 \times x^1 + 2a_0 x^0 \} = 0$$

$$- a_3 \times 3 \times 2 \times x^1 + \sum_{r=2}^{\infty} 4a_r r x^r + 2a_1 x^1 \} = 0$$

$$- \sum_{r=4}^{\infty} a_r r(r-1) x^{r-2} + \sum_{r=2}^{\infty} 2a_r x^r \} = 0$$

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⑥ TIDYING UP

$$\Rightarrow (2a_0 - 2a_2) + (6a_1 - 6a_3)$$

$$+ \sum_{r=2}^{\infty} a_r r(r-1)x^r - \sum_{r=4}^{\infty} a_r r(r-1)x^{r-2} + \sum_{r=2}^{\infty} 4a_r r x^r + \sum_{r=2}^{\infty} 2a_r x^r = 0$$

⑦ ADJUST THE SUMMATIONS SO THEY ALL START FROM $r=2$, BY MAPPING
" $r \mapsto r+2$ " IN THE 2ND SUMMATION (IGNORE THE LOOSE THREHS)

$$\Rightarrow \sum_{r=2}^{\infty} a_r r(r-1)x^r - \sum_{r=2}^{\infty} a_{r+2}(r+2)(r+1)x^r + \sum_{r=2}^{\infty} 4a_r r x^r + \sum_{r=2}^{\infty} 2a_r x^r = 0$$

$$\Rightarrow \sum_{r=2}^{\infty} \left[[a_r [r(r-1) + 4r + 2] - a_{r+2}(r+2)(r+1)] x^r \right] = 0$$

⑧ EQUATING POWERS IN x IN THE SUMMATION

$$\Rightarrow a_{r+2}(r+2)(r+1) = a_r [r(r-1) + 4r + 2]$$

$$\Rightarrow a_{r+2}(r+2)(r+1) = a_r(r^2 + 3r + 2)$$

$$\Rightarrow a_{r+2} = \frac{(r+1)(r+2)}{(r+2)(r+1)} a_r$$

$$\Rightarrow \underline{a_{r+2}} = \underline{a_r}$$

⑨ THIS IS A TRIVIAL RECURRANCE YIELDING

$$a_0 = a_2 = a_4 = a_6 = \dots$$

$$a_1 = a_3 = a_5 = a_7 = \dots$$

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① WE FINALLY HAVE

$$\Rightarrow y = a_0 + a_1 x + a_2 x^2 + a_3 x^3 + a_4 x^4 + \dots$$

$$\Rightarrow y = a_0 + a_1 x + a_0 x^2 + a_1 x^3 + a_0 x^4 + a_1 x^5 + a_0 x^6 + \dots$$

$$\Rightarrow y = a_0 [1 + x^2 + x^4 + x^6 + \dots] + a_1 [x + x^3 + x^5 + x^7 + \dots]$$

$$\Rightarrow y = A(1 + x^2 + x^4 + x^6 + \dots) + Bx(1 + x^2 + x^4 + x^6 + \dots)$$

$$\Rightarrow y = \frac{A}{1-x^2} + \frac{Bx}{1-x^2}$$

$$\Rightarrow y = \frac{A+Bx}{1-x^2}$$

↙ ↘

IVGB - MATHEMATICAL METHODS 3 - PAPER A - QUESTION 8

- USING THE CONTOUR $|z|=1$ OR $z = e^{i\theta}$, $0 \leq \theta < 2\pi$

Thus $z = e^{i\theta}$

$$dz = ie^{i\theta} d\theta$$

- HENCE THE INTEGRAL BECOMES

$$\int_0^{2\pi} \cos^6 \theta \sin^6 \theta d\theta$$

$$= \int_0^{2\pi} (\cos \theta \sin \theta)^6 d\theta = \int_0^{2\pi} \left(\frac{1}{2} \sin 2\theta\right)^6 d\theta$$

$$= \frac{1}{64} \int_0^{2\pi} \sin^6 2\theta d\theta = \frac{1}{64} \int_0^{2\pi} \left[\frac{1}{2i} (e^{i\theta} - e^{-i\theta}) \right]^6 d\theta$$

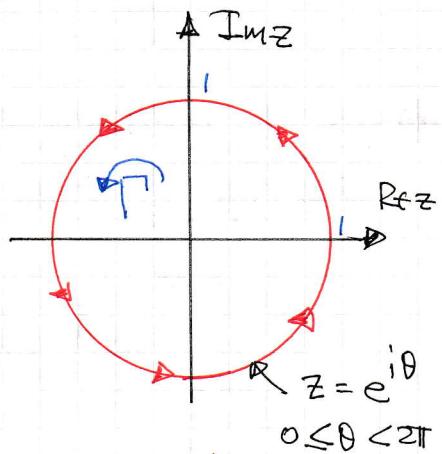
$$= \frac{1}{64} \left(\frac{1}{2i} \right)^6 \int_0^{2\pi} \left[(e^{i\theta})^2 - (e^{-i\theta})^2 \right]^6 d\theta$$

$$= \frac{1}{64} \times \frac{1}{-64} \int_{\Gamma} \left(z^2 - \frac{1}{z^2} \right)^6 \left(\frac{dz}{iz} \right) \quad \leftarrow \boxed{d\theta = \frac{dz}{ie^{i\theta}}}$$

$$= \frac{1}{64^2} \int_{\Gamma} \frac{1}{z} \left(z^2 - \frac{1}{z^2} \right)^6 dz$$

- EXPAND BINOMIALLY

$$= \frac{i}{2^{12}} \int_{\Gamma} \frac{1}{z} \left[z^{12} - 6z^8 + 15z^4 - 20 + \frac{15}{z^4} - \frac{6}{z^8} + \frac{1}{z^{12}} \right] dz$$



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NYGB - MATHEMATICAL METHODS 3 - PAPER A - QUESTION 8

$$= \frac{i}{2^{12}} \int_{\Gamma} z^{11} - 6z^7 + 15z^3 - \frac{20}{z} + \frac{15}{z^5} - \frac{6}{z^9} + \frac{1}{z^{13}} dz$$

As the integrand is in real form the only contribution is
from the $\frac{1}{z}$ term if the residue is -20

$$\Rightarrow \int_0^{2\pi} \cos^6 \theta \sin^6 \theta d\theta = 2\pi i \times \sum \text{(residues inside } \Gamma)$$

$$\Rightarrow \int_0^{2\pi} \cos^6 \theta \sin^6 \theta d\theta = 2\pi i \times \frac{i}{2^{12}} \times (-20)$$

$$\Rightarrow \int_0^{2\pi} \cos^6 \theta \sin^6 \theta d\theta = \frac{5\pi \times 2^3}{2^{12}} = \frac{5\pi}{2^9}$$

$$\Rightarrow \int_0^{2\pi} \cos^6 \theta \sin^6 \theta d\theta = \frac{5\pi}{512} \quad \cancel{\text{}}$$

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IVGB - MATHEMATICAL METHODS 3 - PAPER A - QUESTION 9

- a) ● STARTING FROM THE GENERATING FUNCTION FOR LEGENDRE'S POLYNOMIALS

$$(1 - 2xt + t^2)^{-\frac{1}{2}} = \sum_{n=0}^{\infty} [t^n P_n(x)]$$

- LETTING $x=1$, IN THE ABOVE RELATIONSHIP

$$\Rightarrow (1 - 2t + t^2)^{-\frac{1}{2}} = \sum_{n=0}^{\infty} [t^n P_n(1)]$$

$$\Rightarrow [(1-t)^2]^{-\frac{1}{2}} = \sum_{n=0}^{\infty} [t^n P_n(1)]$$

$$\Rightarrow (1-t)^{-1} = \sum_{n=0}^{\infty} [t^n P_n(1)]$$

$$\Rightarrow \underline{1} + \underline{t} + \underline{t^2} + \underline{t^3} + \dots = \underline{P_0(1)} + \underline{tP_1(1)} + \underline{t^2P_2(1)} + \underline{t^3P_3(1)} + \dots$$

- HENCE THE RESULT FOLLOWS BY COMPARISON $P_n(1) = 1$



- b) ● STARTING WITH LEGENDRE'S EQUATION, WHOSE SOLUTION IS $y = P_n(x)$

$$\Rightarrow (1-x^2) \frac{d^2y}{dx^2} - 2x \frac{dy}{dx} + n(n+1)y = 0$$

$$\Rightarrow (1-x^2)y'' - 2xy' + n(n+1)y = 0$$

$$\Rightarrow (1-x^2)P_n''(x) - 2xP_n'(x) + n(n+1)P_n(x) = 0$$

- LET $x=1$ & NOTE FROM PART (a) , $P_n(1) = 1$

$$\Rightarrow -2P_n'(1) + n(n+1) = 0$$

$$\Rightarrow P_n'(1) = \frac{1}{2}n(n+1)$$

AS REQUIRED

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① STARTING BY THE DIFFERENTIAL EQUATION

$$\frac{d^2x}{dt^2} = a^2 x$$

WITH GENERAL SOLUTION

$$x = A \cosh at + B \sinh at$$

$$\dot{x} = A a \sinh at + B a \cosh at$$

② PICK INITIAL CONDITIONS FOR EACH CASE.

$$t=0, x=1, \dot{x}=0$$

$$\Rightarrow x = \cosh at$$

$$\Rightarrow \dot{x} = a \sinh at$$

$$t=0, x=0, \dot{x}=a$$

$$\Rightarrow x = \sinh at$$

$$\Rightarrow \dot{x} = a \cosh at$$

③ TAKING THE LAPLACE TRANSFORM OF THE O. D. E

$$\Rightarrow \ddot{x} = a^2 x$$

$$\Rightarrow s^2 \bar{x} - s x_0 - \dot{x}_0 = a^2 \bar{x}$$

$$\Rightarrow (s^2 - a^2) \bar{x} = s x_0 + \dot{x}_0$$

$$\Rightarrow \bar{x} = \frac{s x_0 + \dot{x}_0}{s^2 - a^2}$$

$$\Rightarrow \bar{x} = \frac{s}{s^2 - a^2}$$

$$\Rightarrow \mathcal{L}[\cosh at] = \frac{s}{s^2 - a^2}$$

$$\Rightarrow \bar{x} = \frac{a}{s^2 - a^2}$$

$$\Rightarrow \mathcal{L}[\sinh at] = \frac{a}{s^2 - a^2}$$

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- ① START BY THE DEFINITION OF THE FOURIER TRANSFORM

$$\begin{aligned}\hat{f}(k) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-ikx} f(x) dx = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-ikx} (Ae^{-\alpha x^2}) dx \\ &= \frac{A}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-ikx} e^{-\alpha x^2} dx = \frac{2A}{\sqrt{2\pi}} \int_0^{\infty} e^{-\alpha x^2} \cos kx dx \\ &\quad (\text{EVEN INTEGRAND})\end{aligned}$$

- ② THIS INTEGRAL CAN BE DONE BY DIFFERENTIATION UNDER THE INTEGRAL SIGN, WHICH WILL THEN ALLOW INTEGRATION BY PARTS

$$\begin{aligned}\Rightarrow I &= \int_0^{\infty} e^{-\alpha x^2} \cos kx dx \\ \Rightarrow \frac{\partial I}{\partial k} &= \frac{\partial}{\partial k} \int_0^{\infty} e^{-\alpha x^2} \cos kx dx = \int_0^{\infty} e^{-\alpha x^2} \frac{\partial}{\partial k} (\cos kx) dx \\ \Rightarrow \frac{\partial I}{\partial k} &= \int_0^{\infty} (-2x \sin kx) e^{-\alpha x^2} dx = \int_0^{\infty} -x e^{-\alpha x^2} \sin kx dx\end{aligned}$$

- ③ FOLLOWING BY INTEGRATION BY PARTS

$\sin kx$	$k \cos kx$
$\frac{1}{2x} e^{-\alpha x^2}$	$-2e^{-\alpha x^2}$

$$\begin{aligned}\Rightarrow \frac{\partial I}{\partial k} &= \left[\frac{1}{2x} e^{-\alpha x^2} \sin kx \right]_0^\infty - \frac{k}{2x} \int_0^\infty e^{-\alpha x^2} \cos kx dx \\ \Rightarrow \frac{\partial I}{\partial k} &= -\frac{k}{2\alpha} I\end{aligned}$$

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① SOLVING THE O.D.E FOR I

$$\Rightarrow \frac{dI}{dk} = \frac{-k}{2\alpha} I$$

$$\Rightarrow \int \frac{dI}{I} = \int -\frac{k}{2\alpha} dk$$

$$\Rightarrow \ln I = -\frac{k^2}{4\alpha} + C$$

$$\Rightarrow I = Be^{-\frac{k^2}{4\alpha}} \quad (\text{B ARBITRARY})$$

② HENCE SO FAR WE HAVE

$$I = \int_0^\infty e^{-\alpha x^2} \cos kx dx = Be^{-\frac{k^2}{4\alpha}}$$

LET $k=0$

$$\int_0^\infty e^{-\alpha x^2} dx = B$$

— USING A SUBSTITUTION —

$$y^2 = \alpha x^2$$

$$y = \sqrt{\alpha} x$$

$$dy = \sqrt{\alpha} dx$$

UNITS UNCHANGED

$$B = \int_0^\infty e^{-y^2} \frac{dy}{\sqrt{\alpha}}$$

$$B = \frac{1}{\sqrt{\alpha}} \int_0^\infty e^{-y^2} dy$$

$$B = \frac{1}{\sqrt{\alpha}} \left(\frac{\sqrt{\pi}}{2} \right)$$

③ FINALLY WE HAVE

$$\hat{f}(k) = \frac{2A}{\sqrt{2\pi}} I = \frac{2A}{\sqrt{2\pi}} \left(\frac{1}{\sqrt{\alpha}} \frac{\sqrt{\pi}}{2} \right) e^{-\frac{k^2}{4\alpha}} = \frac{A}{\sqrt{2\alpha}} e^{-\frac{k^2}{4\alpha}}$$

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ALTERNATIVE APPROACH BY CONTOUR INTEGRATION

- ① STARTING BY THE DEFINITION OF THE FOURIER TRANSFORM

$$\hat{f}(k) = \frac{A}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\alpha x^2} e^{-ikx} dx = \frac{A}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\alpha[x^2 + \frac{ik}{2\alpha}x]} dx$$

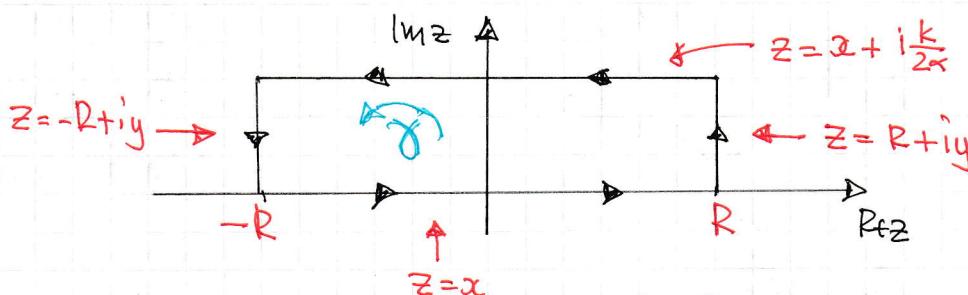
- ② COMPLETING THE SQUARE IN THE EXPONENT, IN x

$$x^2 + \frac{ik}{2\alpha}x = \left(x + \frac{ik}{2\alpha}\right)^2 - \left(\frac{ik}{2\alpha}\right)^2 = \left(x + \frac{ik}{2\alpha}\right)^2 + \frac{k^2}{4\alpha^2}$$

- ③ RETURNING TO THE TRANSFORM

$$\hat{f}(k) = \frac{A}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\alpha(x + \frac{ik}{2\alpha})^2} \times e^{-\frac{k^2}{4\alpha}} dx = \frac{Ae^{-\frac{k^2}{4\alpha}}}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\alpha(x + \frac{k}{2\alpha}i)^2} dx$$

- ④ NOW CONSIDER $f(z) = e^{-\alpha z^2}$ OVER THE CONTOUR SHOWN BELOW



NO POLES INSIDE γ

$$\Rightarrow \int_{\gamma} e^{-\alpha z^2} dz = 0$$

$$\Rightarrow \int_{-R}^R e^{-\alpha x^2} dx + \int_{y=0}^{y=\frac{k}{2\alpha}} e^{-\alpha(R+iy)^2} (idy) + \int_R^{-R} e^{-\alpha(x+\frac{ik}{2\alpha})^2} dx + \int_{y=\frac{k}{2\alpha}}^{y=0} e^{-\alpha(-R+iy)^2} (idy) = 0$$

$$\begin{array}{c} \oplus \\ z=x \\ dz=dx \end{array}$$

$$\begin{array}{c} \oplus \\ z=R+iy \\ dz=idy \end{array}$$

$$\begin{array}{c} \oplus \\ z=x+\frac{ik}{2\alpha} \\ dz=dx \end{array}$$

$$\begin{array}{c} \oplus \\ z=-R+iy \\ dz=idy \end{array}$$

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- Now as $R \rightarrow \infty$, the 2nd & 4th integrals vanish because of the term $-\alpha R^2$ in the exponent

$$\Rightarrow \int_{-\infty}^{\infty} e^{-\alpha x^2} dx + \int_{\infty}^{-\infty} e^{-\alpha(x+i\frac{k}{2\alpha})^2} dx = 0$$

$$\Rightarrow \int_{-\infty}^{\infty} e^{-\alpha(x+i\frac{k}{2\alpha})^2} dx = \int_{-\infty}^{\infty} e^{-\alpha x^2}$$

- Returning to the transform

$$\hat{f}(k) = \frac{Ae^{-\frac{k^2}{4\alpha}}}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\alpha(x+i\frac{k}{2\alpha})^2} dx = \frac{A}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\alpha x^2} dx$$

$$\hat{f}(k) = \frac{Ae^{-\frac{k^2}{4\alpha}}}{\sqrt{2\pi}} \times \frac{1}{\sqrt{\alpha}} \int_{-\infty}^{\infty} e^{-y^2} dy$$

$$\hat{f}(k) = \frac{Ae^{-\frac{k^2}{4\alpha}}}{\sqrt{2\pi}} \times \frac{1}{\sqrt{\alpha}} \times \sqrt{\pi}$$

Let $y = \sqrt{\alpha}x$

$dy = \sqrt{\alpha}dx$

$dx = \frac{1}{\sqrt{\alpha}}dy$

LIMITS UNCHANGED

$$\hat{f}(k) = \frac{A}{\sqrt{2\alpha}} e^{-\frac{k^2}{4\alpha}}$$

AS BEFORE