

Fra Kreyszig (10th), avsnitt 12.6

- 5 Randbetingelser: $u(0, t) = u(10, t) = 0 \quad \forall t$.
 Initalbetingelse: $u(x, 0) = f(x) = \sin(0.1\pi x)$.
 $f(x)$ er sin egen fourierrekke med $B_1 = 1$ og $B_n = 0 \quad n \in \mathbb{N} \setminus \{1\}$

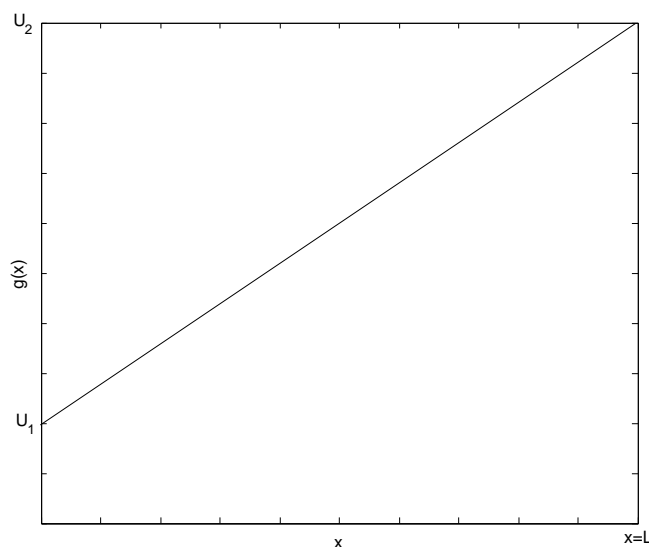
$$\begin{aligned} \Rightarrow u(x, t) &= \sin\left(\frac{1 \cdot \pi}{10}x\right) e^{-\left(\frac{c \cdot 1 \cdot \pi}{10}\right)^2 t} \\ &= \sin\left(\frac{\pi}{10}x\right) e^{-\frac{1.75\pi^2}{100}t} \end{aligned}$$

der $c^2 = \frac{\kappa}{\rho\sigma}$ med $\kappa = 1.04$, $\rho = 10.6$ og $\sigma = 0.056$.

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$$\begin{aligned} u_t &= c^2 u_{xx} \\ u(0, t) &= U_1 \quad t \geq 0 \\ u(L, t) &= U_2 \quad t \geq 0 \\ u(x, 0) &= f(x) \end{aligned}$$

Gjetter på formen $g(x) = U_1 + \frac{U_2 - U_1}{L}x$.



$$\text{La } v(x, t) = u(x, t) - g(x)$$

$$\implies v_t = u_t = c^2 u_{xx} = c^2 v_{xx} \quad \text{siden } g^2(x) = 0$$

Rand- og initialbetingelser for v blir

$$v(0, t) = U_1 - g(0) = U_1 - U_1 = 0$$

$$v(L, t) = U_2 - g(L) = U_2 - U_2 = 0$$

$$v(x, 0) = f(x) - g(x)$$

Løsning:

$$v(x, t) = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{L}\right) e^{-\lambda_n^2 t} \quad \left(\lambda_n = \frac{cn\pi}{L}\right) \quad \text{Se seksjon 12.6, (9)-(10)}$$

$$B_n = \frac{2}{L} \int_0^L (f(x) - g(x)) \sin\left(\frac{n\pi x}{L}\right) dx$$

$$\implies u(x, t) = g(x) + v(x, t) = g(x) + \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{L}\right) e^{-\lambda_n^2 t}$$

Vi har at

$$\sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{L}\right) e^{-\lambda_n^2 t} \rightarrow 0 \text{ når } t \rightarrow \infty$$

$$\implies u(x, t) \rightarrow g(x) \text{ når } t \rightarrow \infty$$

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$$u_x(0, t) = 0$$

$$u_x(L, t) = 0$$

$$u(x, 0) = x$$

$$L = \pi$$

$$c = 1$$

$$\xrightarrow{Eks.4} u(x, t) = A_0 + \sum_{n=1}^{\infty} A_n \cos(nx) e^{-n^2 t}$$

$$\text{med } A_0 = \frac{1}{\pi} \int_0^{\pi} f(x) dx = \frac{1}{\pi} \int_0^{\pi} x dx = \frac{x^2}{2\pi} \Big|_0^{\pi} = \frac{\pi}{2}$$

$$A_n = \frac{2}{\pi} \int_0^{\pi} x \cos(nx) dx = \frac{2}{\pi} x \frac{\sin(nx)}{n} \Big|_0^{\pi} - \frac{2}{\pi} \int_0^{\pi} \frac{\sin(nx)}{n} dx$$

$$= \frac{2}{\pi} \frac{\cos(nx)}{n^2} \Big|_0^{\pi} = \frac{2}{\pi n^2} ((-1)^n - 1)$$

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$$\begin{aligned}
u_t &= c^2 u_{xx} + H \\
L &= \pi \\
u(0, t) &= u(\pi, t) = 0 \quad t \geq 0 \\
u(x, 0) &= f(x)
\end{aligned}$$

La $u(x, t) = v(x, t) - Hx \frac{x-\pi}{2c^2}$. Dette gir

$$\begin{aligned}
u_t(x, t) &= v_t(x, t) \\
u_{xx}(x, t) &= v_{xx}(x, t) - \frac{H}{c^2} \\
\implies v_t(x, t) &= u_t(x, t) = c^2 u_{xx}(x, t) + H = c^2 v_{xx} - H + H = c^2 v_{xx}(x, t)
\end{aligned}$$

Vi har dermed

$$\begin{aligned}
v_t(x, t) &= c^2 v_{xx}(x, t) \\
v(0, t) &= u(0, t) = 0 \\
v(\pi, t) &= u(\pi, t) = 0 \\
v(x, 0) &= u(x, 0) + Hx \frac{x-\pi}{2c^2} = f(x) + Hx \frac{x-\pi}{2c^2}
\end{aligned}$$

$$\stackrel{(9)-(10)}{\implies} v(x, t) = \sum_{n=1}^{\infty} B_n \sin(nx) e^{-\lambda n^2 t}$$

med

$$\lambda_n = cn \quad \text{og} \quad B_n = \frac{2}{\pi} \int_0^{\pi} \left(f(x) + Hx \frac{x-\pi}{2c^2} \right) \sin(nx) dx$$

Dermed er

$$u(x, t) = \sum_{n=1}^{\infty} B_n \sin(nx) e^{-\lambda n^2 t} - Hx \frac{x-\pi}{2c^2}$$

med

$$\lambda_n = cn \quad \text{og} \quad B_n = \frac{2}{\pi} \int_0^{\pi} \left(f(x) + Hx \frac{x-\pi}{2c^2} \right) \sin(nx) dx$$

21 Oppgitte betingelser:

$$u(0, y) = u(a, y) = u(x, 0) = 0, \quad \text{og} \quad u(x, a) = 25, \quad \text{med } a = 24$$

Steady-state temperatur vil si at temperaturen ikke lenger endrer seg med tiden: $u_t = 0$. Varmeledningsligningen i to dimensjoner blir dermed

$$\begin{aligned}
\nabla^2 u &= 0 \\
u_{xx} + u_{yy} &= 0, \quad (1)
\end{aligned}$$

som er Laplaces ligning. Bruker separasjon av variable:

$$u(x, y) = F(x)G(y)$$

Innsatt i (1):

$$\begin{aligned}\frac{d^2 F}{dx^2} G + F \frac{d^2 G}{dy^2} &= 0 \\ \frac{1}{F} \frac{d^2 F}{dx^2} &= -\frac{1}{G} \frac{d^2 G}{dy^2} = k \quad (\text{en konstant})\end{aligned}$$

Vi har dermed to ordinære differensialligninger:

$$F'' - kF = 0, \quad G'' + kG = 0$$

Med $k = \mu^2 > 0$:

$$F(x) = C_1 e^{\mu x} + C_2 e^{-\mu x}$$

Initialbetingelsene $u(0, y) = u(a, y) = 0$ gir $C_1 = C_2 = 0$.

Med $k = 0$:

$$F(x) = C_3 x + C_4$$

Initialbetingelsene $u(0, y) = u(a, y) = 0$ gir $C_3 = C_4 = 0$.

Med $k = -\mu^2 < 0$:

$$F(x) = C_5 \cos(\mu x) + C_6 \sin(\mu x)$$

Initialbetingelsen $u(0, y) = 0$ gir:

$$\begin{aligned}C_5 \cos 0 + C_6 \sin 0 &= 0 \\ C_5 &= 0\end{aligned}$$

Initialbetingelsen $u(a, y) = 0$ gir:

$$\begin{aligned}C_6 \sin(\mu a) &= 0 \\ \mu a &= n\pi \\ \mu &= \frac{n\pi}{a} \quad n = 0, \pm 1, \pm 2, \dots\end{aligned}$$

$$F(x) = C_6 \sin\left(\frac{n\pi x}{a}\right)$$

Med $k = -\mu^2$ blir ligningen for $G(y)$:

$$G'' - \mu^2 G = 0$$

Med løsning

$$G(y) = C_7 e^{\mu y} + C_8 e^{-\mu y}$$

Initialbetingelsen $u(x, 0) = 0$ gir $C_8 = -C_7$:

$$\begin{aligned} G(y) &= C_7 (e^{\mu y} - e^{-\mu y}) \\ &= 2C_7 \sinh(\mu y) \\ &= 2C_7 \sinh\left(\frac{n\pi y}{a}\right) \end{aligned}$$

For hver eneste n får vi dermed en løsning:

$$u_n(x, y) = C_9 \sin\left(\frac{n\pi x}{a}\right) \sinh\left(\frac{n\pi y}{a}\right)$$

Nå er fortsatt $n = 0, \pm 1, \pm 2$, men siden u_n er odde mtp n (som betyr at u_{-n} er proporsjonal med u_n) holder det å summere $u(x, y)$ for bare positive n . Løsningen for $n = 0$ er $u = 0$, som ikke er særlig interessant.

Generell løsning:

$$u(x, y) = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{a}\right) \sinh\left(\frac{n\pi y}{a}\right)$$

Bruker nå siste betingelse: $u(x, a) = 25$

$$25 = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{a}\right) \sinh(n\pi)$$

$$\begin{aligned} B_n \sinh(n\pi) &= \frac{2}{a} \int_0^a 25 \sin\left(\frac{n\pi x}{a}\right) dx \\ \frac{B_n \sinh(n\pi)a}{50} &= \left[\frac{-a}{n\pi} \cos\left(\frac{n\pi x}{a}\right) \right]_0^a \\ \frac{B_n \sinh(n\pi)}{50} &= \frac{1}{n\pi} (1 - (-1)^n) \\ B_n &= \frac{50}{n\pi \sinh(n\pi)} (1 - (-1)^n) \end{aligned}$$

Som gir løsningen:

$$u(x, y) = \sum_{n=1}^{\infty} \frac{50}{n\pi \sinh(n\pi)} (1 - (-1)^n) \sinh\left(\frac{n\pi y}{24}\right) \sin\left(\frac{n\pi x}{24}\right)$$

23 Oppgitte betingelser:

$$u(0, y) = 0, \quad u(a, y) = f(y), \quad u_y(x, 0) = u_y(x, a) = 0 \quad \text{med } a = 24$$

Samme metode som i oppgave 12.6.23:

$$u_{xx} + u_{yy} = 0 \quad (\text{Laplace's ligning})$$

Seperasjon av variable:

$$u(x, y) = F(x)G(y)$$

$$\Rightarrow \quad \frac{1}{F} \frac{d^2 F}{dx^2} = -\frac{1}{G} \frac{d^2 G}{dy^2} = k \quad (\text{en konstant})$$

$$F'' - kF = 0, \quad G'' + kG = 0$$

Med $k = \mu^2 > 0$:

$$\begin{aligned} G(y) &= C_1 \cos(\mu y) + C_2 \sin(\mu y) \\ G'(y) &= \mu C_2 \cos(\mu y) - \mu C_1 \sin(\mu y) \end{aligned}$$

Initialbetingelsen $u_y(x, a) = 0$ gir:

$$\begin{aligned} -\mu C_1 \sin(\mu a) &= 0 \\ \mu a &= n\pi \\ \mu &= \frac{n\pi}{a}, \quad n = 0, \pm 1, \pm 2, \dots \end{aligned}$$

$$G(y) = C_1 \cos\left(\frac{n\pi y}{a}\right)$$

Med $k = \mu^2 > 0$ blir løsningen for $F(x)$:

$$\begin{aligned} F'' - \mu^2 F &= 0 \\ \Rightarrow F(x) &= C_3 e^{\mu x} + C_4 e^{-\mu x} \end{aligned}$$

Initialbetingelsen $u(0, y) = 0$ gir $C_4 = -C_3$.

$$F(x) = C_3 (e^{\mu x} - e^{-\mu x}) = 2C_3 \sinh(\mu x)$$

$$\Rightarrow u_n(x, y) = A_n \sinh\left(\frac{n\pi x}{a}\right) \cos\left(\frac{n\pi y}{a}\right), \quad (1), \quad n = 1, 2, 3, \dots$$

(Av samme grunn som i oppgave 12.6.21 får vi ikke noen flere løsninger av å inkludere negative n , mens $n = 0$ bare gir $u = 0$).

Med $k = 0$:

$$G(y) = C_5 y + C_6$$

Initialbetingelsene $u_y(x, 0) = u_y(x, a) = 0$ gir begge $C_5 = 0$, slik at $G(y) = C_6$ er en gyldig løsning. Da blir $F(x)$:

$$F(x) = C_7 x + C_8$$

Initialbetingelsen $u(0, y) = 0$ gir $C_8 = 0$. Dermed er en gyldig løsning

$$u(x, y) = F(x)G(y) = A_0 x, \quad (2) \quad (A_0 = C_6 C_7)$$

Med $k = -\mu^2$ gir initialbetingelsene at $u(x, y) = 0$.

Den mest generelle løsningen blir dermed en sum av alle løsningene fra (1) pluss løsningen fra (2):

$$u(x, y) = A_0 x + \sum_{n=1}^{\infty} A_n \sinh\left(\frac{n\pi x}{a}\right) \cos\left(\frac{n\pi y}{a}\right)$$

Bruker siste initialbetingelsen $u(a, y) = f(y)$ for å bestemme konstantene:

$$f(y) = A_0 a + \sum_{n=1}^{\infty} A_n \sinh(n\pi) \cos\left(\frac{n\pi y}{a}\right)$$

$$A_0 a = \frac{1}{a} \int_0^a f(y) dy$$

$$A_n \sinh(n\pi) = \frac{2}{a} \int_0^a f(y) \cos\left(\frac{n\pi y}{a}\right) dy$$

Setter dette inn i uttrykket for $u(x, y)$:

$$u(x, y) = \left[\frac{1}{24^2} \int_0^{24} f(y) dy \right] x + \sum_{n=1}^{\infty} \left[\frac{1}{12 \sinh(n\pi)} \int_0^{24} f(y) \cos\left(\frac{n\pi y}{24}\right) dy \right] \sinh\left(\frac{n\pi x}{24}\right) \cos\left(\frac{n\pi y}{24}\right)$$

Fra Kreyszig (10th), avsnitt 12.7

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$$f(x) = e^{-|x|}$$

$$u(x, t) = \int_0^{\infty} [A(p) \cos(px) + B(p) \sin(px)] e^{-c^2 p^2 t} dp$$

$$\implies u(x, 0) = \int_0^{\infty} [A(p) \cos(px) + B(p) \sin(px)] dp = f(x)$$

Vi ser dette er et Fourier-integral:

$$\begin{aligned} A(p) &= \frac{1}{\pi} \int_{-\infty}^{\infty} e^{-|v|} \cos(pv) dv \\ &= \frac{1}{\pi} \int_{-\infty}^0 e^v \cos(pv) dv + \frac{1}{\pi} \int_0^{\infty} e^{-v} \cos(pv) dv \\ &= \frac{1}{\pi} e^v \cos(pv) \Big|_{-\infty}^0 + \frac{1}{\pi} \int_{-\infty}^0 e^v \sin(pv) dv - \frac{1}{\pi} e^{-v} \cos(pv) \Big|_0^{\infty} - \frac{1}{\pi} \int_0^{\infty} e^{-v} \sin(pv) dv \\ &= \frac{1}{\pi} + \frac{1}{\pi} + \frac{1}{\pi} e^v \sin(pv) p \Big|_{-\infty}^0 - \frac{1}{\pi} \int_{-\infty}^0 e^v \cos(pv) p^2 dv + \frac{1}{\pi} e^{-v} \sin(pv) p \Big|_0^{\infty} \\ &\quad - \frac{1}{\pi} \int_0^{\infty} e^{-v} \cos(pv) p^2 dv \\ &= \frac{2}{\pi} - p^2 A(p) \end{aligned}$$

$$\implies A(p) = \frac{2}{\pi(1+p^2)}$$

$$B(p) = \frac{1}{\pi} \int_{-\infty}^{\infty} e^{-|v|} \sin(pv) dv = 0 \quad \text{siden integranden er odde}$$

$$\implies u(x, t) = \int_0^{\infty} \frac{2}{\pi(1+p^2)} \cos(px) e^{-c^2 p^2 t} dp$$

Alternativt via Fouriertransformasjon:

$$\hat{u}_t(w, t) = -c^2 w^2 \hat{u}(w, t) \implies \hat{u}(w, t) = C(w) e^{-c^2 w^2 t} \implies \hat{u}(w, t) = \hat{f}(w) e^{-c^2 w^2 t}$$

siden $\hat{u}(w, 0) = \hat{f}(w)$.

$$\implies u(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(p) \frac{1}{\sqrt{2c^2 t}} e^{-\frac{1}{4c^2 t} (x-p)^2} dp$$