

1) Quantum assignment 2. pwn274, Team Luvie

$$1) H' = kr^2(1 - 3\cos^2\theta)$$

$$H^0 = -\frac{\hbar^2}{2m}\nabla^2 - \frac{e^2}{4\pi\epsilon_0 r}$$

$$E' = \langle 100 | H' | 100 \rangle = \int_0^\infty dr \int_0^\pi d\theta \int_0^{2\pi} d\phi (r^2 \sin\theta) kr^2(1 - \cos^2\theta) \frac{1}{a^3} e^{-2r/a} \frac{1}{4\pi}$$

$$= \frac{k}{\pi a^3} \int_0^\infty dr \int_0^\pi d\theta \int_0^{2\pi} d\phi \left(r^4 e^{-2r/a} \sin\theta + \frac{r^4 e^{-2r/a}}{3} \frac{d}{d\theta} \cos^3\theta \right)$$

$$= \frac{k2}{a^3} \int_0^\infty dr \left(-2 - \frac{2}{3} \right) r^4 e^{-2r/a} = \frac{k2}{a^3} \frac{8}{3} \cdot \frac{1}{2} \left(\frac{a}{2} \right)^5 = \frac{ka^2 \cdot 2 \cdot 2 \cdot 2 \cdot 4 \cdot 3 \cdot 2}{3 \cdot 2 \cdot 2 \cdot 2 \cdot 2 \cdot 2} = \underline{\underline{-2ka^2}}$$

2)

E_2^0 is the eigenvalue of $|200\rangle, |21-1\rangle, |210\rangle$ and $|211\rangle$

ϕ -part will be $\int_0^{2\pi} d\phi e^{-im\phi} e^{im'\phi} = \int_0^{2\pi} d\phi e^{i(m'-m)\phi}$

this will for all $m \neq m'$ be an integral over Δm periods of $e^{i\phi}$ which, by symmetry, always be 0.

Only for $m = m'$ we will have $\int_0^{2\pi} d\phi e^{i0\phi} = \int_0^{2\pi} d\phi = 2\pi$

3) The perturbation-matrix is diagonal in the basis of

$|21m\rangle$, as H' effectively is $-Y_2^0(\theta, \phi)$, meaning

$\langle n1m | H' | n00 \rangle = \langle n1m | n20 \rangle = 0$ as Y_l^m are orthogonal

The matrix is then:

$21m$	$11-1$	110	111	100
$11-1$	$\langle 11-1 11-1 \rangle$	0	0	0
110	0	$\langle 10 10 \rangle$	0	0
111	0	0	$\langle 11 11 \rangle$	0
100	0	0	0	$\langle 00 00 \rangle$

As I now have the perturbation-matrix in diagonal shape, I can see the first-order perturbation brings

$$E_2' = \langle 00 | H' | 00 \rangle, \langle 1-1 | H' | 1-1 \rangle, \langle 10 | H' | 10 \rangle, \langle 11 | H' | 11 \rangle$$

I have no room on this page, but having used wolfram-Alpha to solve the integrals I get

$$\langle 00 | H' | 00 \rangle = \langle 10 | H' | 10 \rangle = 0 \text{ and } \langle 11 | H' | 11 \rangle = \frac{45}{16} \langle 1-1 | H' | 1-1 \rangle = \frac{45}{16}$$

This means the total perturbation is either:

$$\{0, 0, \frac{45}{16}, -\frac{45}{16}\}$$

2.1 We have the formula from lab:

$$\Delta E = \gamma B_0 \hbar$$

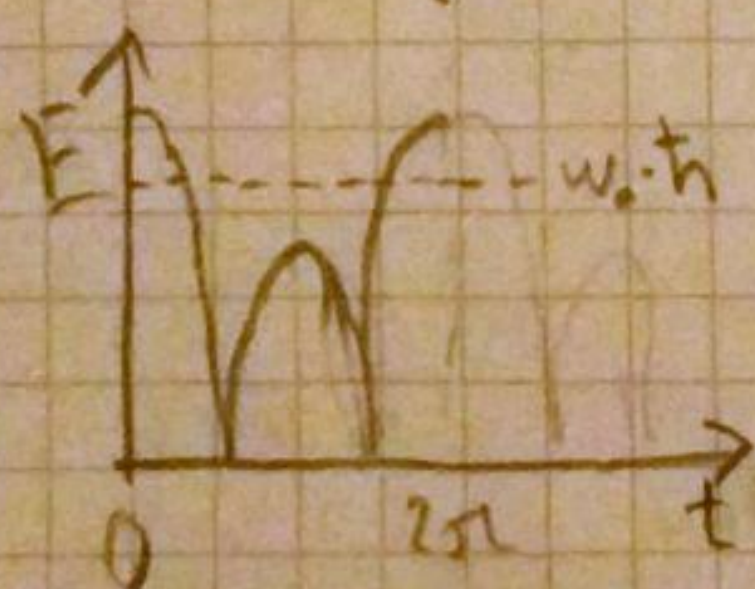
In our experiment B_0 was given by:

$$B_0 = \left(\frac{4}{3}\right)^{\frac{3}{2}} \frac{N}{R} N_0 I$$

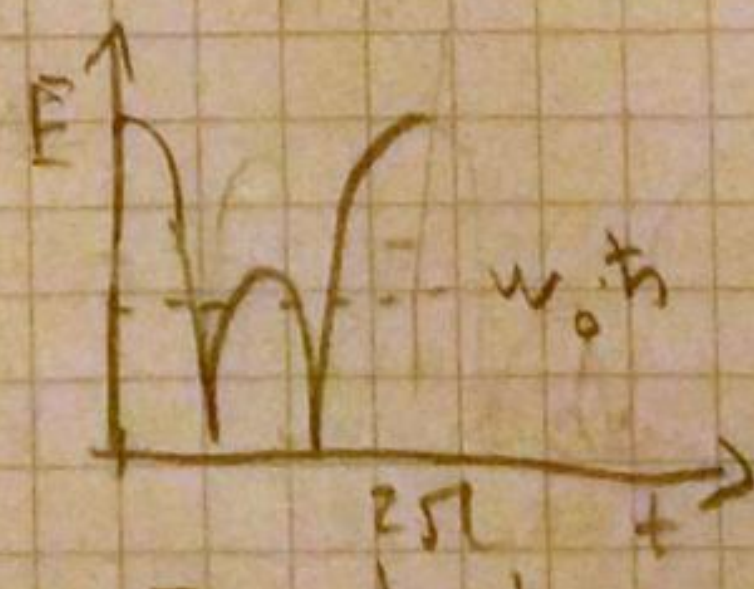
We designed the experiment with $I = (I_{pc} + I_{ac})$:

$$I = \frac{V_{pc}}{R} + \frac{V_{ac}^0}{R} \cos(\omega t)$$

This gives ($\omega=1$):



For low V_{ac}



For high V_{ac}

It is clear, that E hits the right ΔE four times per period when looking at a high enough V_{ac}

3. Kicked Rotor

- 1) The eigenfunctions are Y_l^m as they are simultaneous eigenfunctions for \hat{L}^2 and \hat{L}_z .

The eigenenergies are therefore $\frac{\hbar^2}{2I} l(l+1) + a\hbar m$ with the ground-state at $Y_0^0 \Rightarrow \frac{\hbar^2}{2I} - a\hbar$

2) $C_a^0 = 1 \quad C_b^0 = 0$

$C_a = 0 \Rightarrow C_a^1 = 1$

$\omega_0 = \frac{E_1 - E_0}{\hbar} = \frac{a^2 \hbar}{2I}$

$$C_b^1 = -\frac{i}{\hbar} \langle Y_1^0 | H | Y_0^0 \rangle e^{i\omega_0 t} = -\frac{i}{\hbar} \langle Y_1^0 | b \hat{L}_z | Y_0^0 \rangle \sum_{n=1}^N \delta(t - ndt) e^{i\omega_0 t}$$

$$= -\frac{ib}{2\hbar} \langle Y_1^0 | \hat{L}_+ + \hat{L}_- | Y_0^0 \rangle \sum_{n=1}^N \delta(t - ndt) e^{i\omega_0 t}$$

$$= -\frac{ib}{2\hbar} \frac{1}{\sqrt{2}} \langle Y_1^0 | Y_0^0 \rangle \sum_{n=1}^N \delta(t - ndt) e^{i\omega_0 t}$$

$$= -i \frac{\sqrt{2}}{2} b \sum_{n=1}^N \delta(t - ndt) e^{i\omega_0 t}$$

$$C_b^1 = -i \frac{\sqrt{2}}{2} b \int_0^t \sum_{n=1}^N \delta(t' - ndt) e^{i\omega_0 t'} dt' = -i \frac{\sqrt{2}}{2} b \sum_{n=1}^N e^{i\omega_0 ndt} = -i \frac{\sqrt{2}}{2} b \sum_{n=1}^N e^{iandt}$$

The probability is then the square of the absolute value:

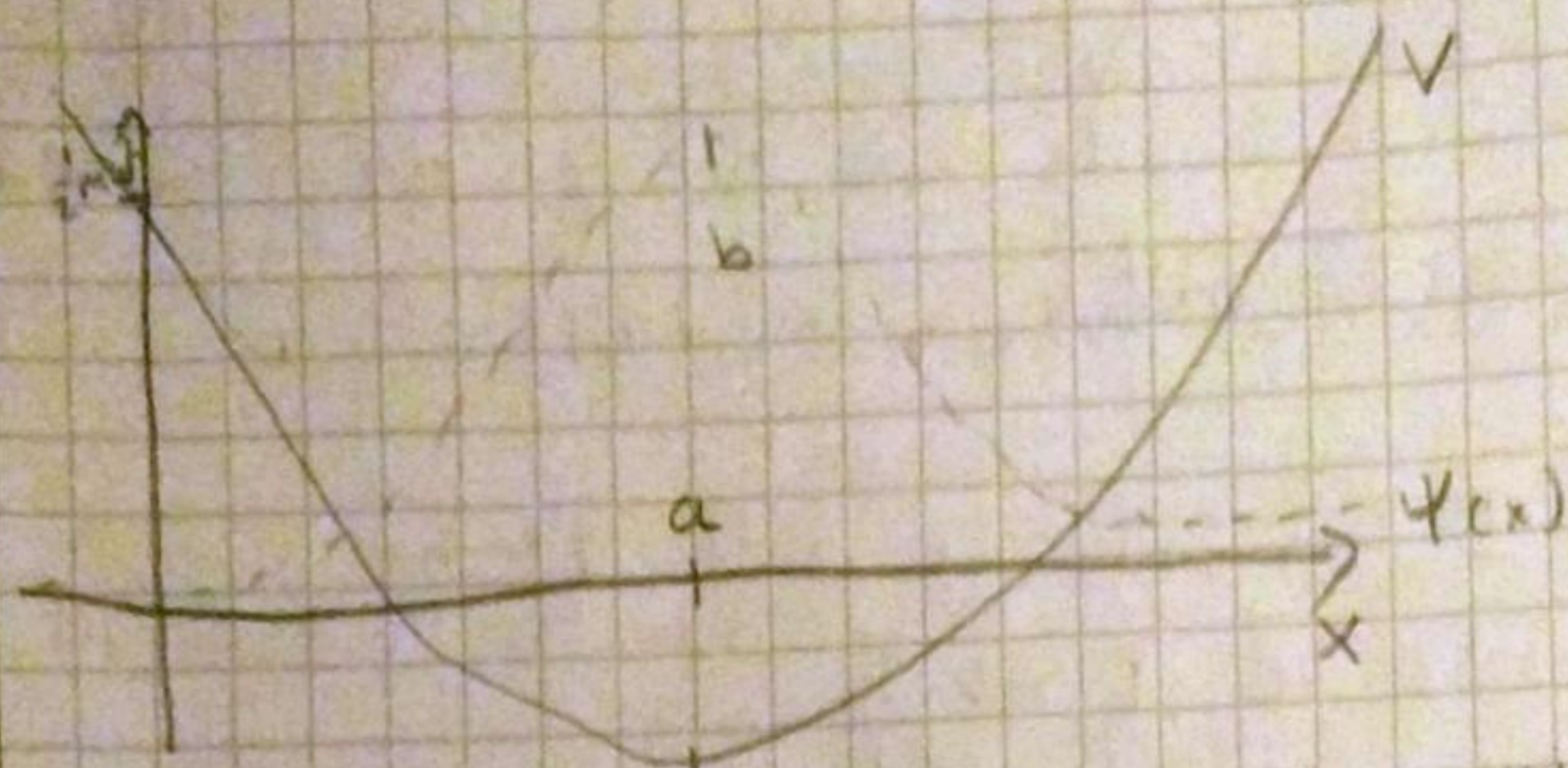
$$P = |C_b^1|^2 = 2b^2 \left| \sum_{n=1}^N e^{iandt} \right|^2 = 2b^2 \sum_{n,m=1}^N e^{i(n-m)adt}$$

- 3) This can be seen as adding N unit-vectors, the length of this is of course maximized when they all point in the same direction.

At $dt = \frac{2\pi}{a}$ all $(n-m)adt$ will do this exactly:

$$P = 2b^2 \sum_{n,m=1}^N e^{i2\pi(n-m)} = 2b^2 \sum_{n,m=1}^N 1 = 2Nb^2$$

$$4) H = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m \omega^2 (x^2 - 2ax)$$



1) My guess is $b=a$ minimizes $\langle \psi | H | \psi \rangle$

2) As the perturbation looks like $H' = -m\omega^2 x$, and x can be written as a raising and a lowering operator it is trivial to see, the first order perturbation is zero. To find the second order correction I write:

$$E_0^{(2)} = \sum_{m \neq 0} \frac{|\langle \psi_m^0 | -m\omega^2 \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}_+ + \hat{a}_-) | \psi_0^0 \rangle|^2}{E_0 - E_m}$$

As it is only the raising operator that will hit ψ_0^0 I get:

$$E_0^{(2)} = \frac{|-m\omega^2 \sqrt{\frac{\hbar}{2m\omega}}|^2}{\frac{1}{2}\hbar\omega - \frac{3}{2}\hbar\omega} = \frac{m^2 \omega^4 \frac{\hbar}{2m\omega}}{-\hbar\omega} = -\frac{m\omega^2 a^2}{2}$$

3) From the hint I can see the potential is simply a harmonic oscillator lowered by a constant. this constant is $\frac{m\omega^2 a^2}{2}$, just as we found in ②.