# My Thesis

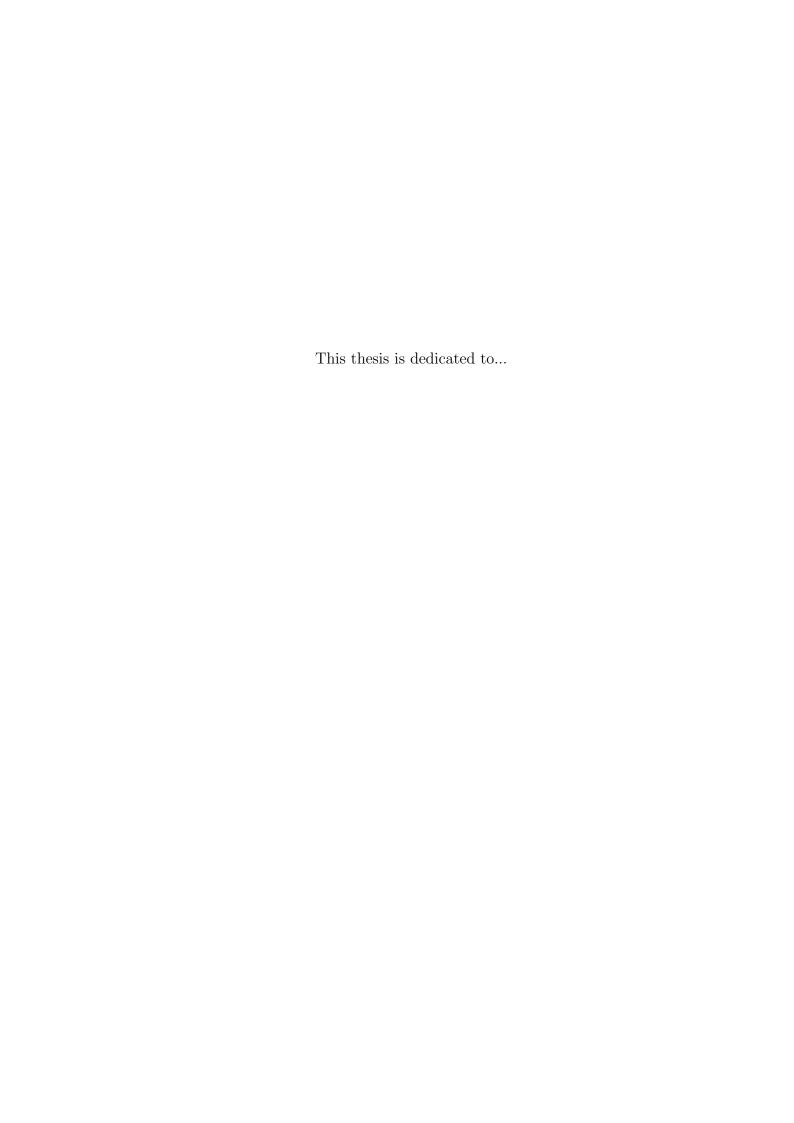


Joe Science

Department of Physics and Astronomy

University of Leeds

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## Acknowledgements

My thanks to...

### Abstract

My abstract in here...  $\,$ 

### Abbreviations

 $k_B$  Boltzmann's constant

 $k_BT$  Thermal energy

... ...

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## Chapter 1

### Introduction

#### 1.1 Introduction

Defining and distinguishing phases of matter has been a continuing effort by physicists for many years.

### 1.2 Free Fermion Systems

#### 1.2.1 Dirac Fermions

The basic constituent of fermionic systems is the Dirac fermion. In the language of quantum field theory, a Dirac fermion can be represented by a second quantised field operator a, called the annihilation operator, and its conjugate partner  $a_i^{\dagger}$ , called the creation operator. They obey the anticommutation relations  $\{a, a^{\dagger}\} = 1$  and  $\{a^{(\dagger)}, a^{(\dagger)}\} = 0$ . These operators act on a number state, also known as a Fock state, in the following way

$$a |0\rangle = 0$$
  $a^{\dagger} |0\rangle = |1\rangle$   $a |1\rangle = |0\rangle$   $a^{\dagger} |1\rangle = 0$  (1.1)

which follow from the commutation relations. A general state of a system consisting of a single Dirac fermion can be written in the form

$$|\psi\rangle = \alpha_0 |0\rangle + \alpha_1 |1\rangle , \qquad (1.2)$$

where  $\alpha_0, \alpha_1 \in \mathbb{C}$ .

Given  $N \in \mathbb{N}^+$  Dirac fermions, they can be represented by a set of second quantised fermionic field operators  $\{a_i\}$  and their conjugate partners  $\{a_i^{\dagger}\}$ , where i=1,...,N. They obey the following commutation relations

$$\{a_i, a_i^{\dagger}\} = \delta_{ij} \qquad \{a_i^{(\dagger)}, a_i^{(\dagger)}\} \qquad (1.3)$$

where  $\delta_{ij}$  Kronercker delta function. These operators act on a tensor product of Fock states. A general state of the fermionic system,  $|\psi\rangle$  can be written as

$$|\psi\rangle = \sum_{n_i=0.1} \left(\alpha_{n_1,\dots,n_N} \bigotimes_{i=1}^N |n_i\rangle\right),$$
 (1.4)

where  $\alpha_{n_1,\dots,n_N} \in \mathbb{C}$  and

$$\bigotimes_{i=1}^{N} |n_i\rangle = \left(\bigotimes_{i=1}^{N} \left(a_i^{\dagger}\right)^{n_i}\right) \left(\bigotimes_{i=1}^{N} |0\rangle\right). \tag{1.5}$$

#### 1.2.2 Quadratic Hamiltonians

# Chapter 2

# Decomposition of the Chern Number

- 2.1 Introduction
- 2.1.1 Background

## Appendix A

# Code samples

#### A.1 Random Number Generator

The Bayes Durham Shuffle ensures that the psuedo random numbers used in the simulation are further shuffled, ensuring minimal correlation between subsequent random outputs (Press *et al.*, 1992).

```
#define IM1 2147483563
#define IM2 2147483399
#define AM (1.0/IM1)
#define IMM1 (IM1-1)
#define IA1 40014
#define IA2 40692
#define IQ1 53668
#define IQ2 52774
#define IR1 12211
#define IR2 3791
#define NTAB 32
#define NDIV (1+IMM1/NTAB)
#define EPS 1.2e-7
#define RNMX (1.0 - EPS)
double ran2(long *idum)
{
```

```
/* Minimum Standard Random Number Generator
/* Taken from Numerical recipies in C
/* Based on Park and Miller with Bays Durham Shuffle */
/* Coupled Schrage methods for extra periodicity
                                                      */
/* Always call with negative number to initialise
                                                      */
int j;
long k;
static long idum2=123456789;
static long iy=0;
static long iv[NTAB];
double temp;
if (*idum <=0)
{
  if (-(*idum) < 1)
    *idum = 1;
  }else
    *idum = -(*idum);
  }
  idum2=(*idum);
  for (j=NTAB+7; j>=0; j--)
   k = (*idum)/IQ1;
    *idum = IA1 *(*idum-k*IQ1) - IR1*k;
    if (*idum < 0)
    {
      *idum += IM1;
    }
    if (j < NTAB)
```

```
{
        iv[j] = *idum;
      }
    }
    iy = iv[0];
  }
  k = (*idum)/IQ1;
  *idum = IA1*(*idum-k*IQ1) - IR1*k;
  if (*idum < 0)
    *idum += IM1;
  }
  k = (idum2)/IQ2;
  idum2 = IA2*(idum2-k*IQ2) - IR2*k;
  if (idum2 < 0)
  {
    idum2 += IM2;
  }
  j = iy/NDIV;
  iy=iv[j] - idum2;
  iv[j] = *idum;
  if (iy < 1)
    iy += IMM1;
  if ((temp=AM*iy) > RNMX)
    return RNMX;
  }else
  {
    return temp;
  }
}
```

# References

NIELSEN, S., LOPEZ, C., SRINIVAS, G. & KLEIN, M. (2004). Coarse grain models and the computer simulation of soft materials. *J. Phys. Condens. Matter*, **16**, R481–R512.

PRESS, W. et al. (1992). Numerical recipes in C. Cambridge University Press Cambridge. 4