

Differential Equations in Geophysical Fluid Dynamics

XX. Wave in rotation: Kelvin wave

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Recap

In the long wave (big) limit...

$$\frac{\partial \bar{u}}{\partial t} - f\bar{v} = -g\frac{\partial \eta}{\partial x} \quad (1a)$$

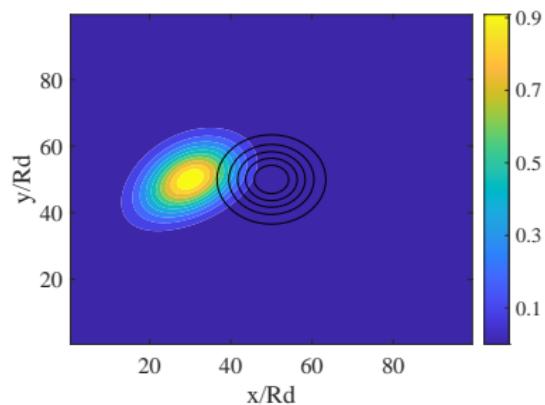
$$\frac{\partial \bar{v}}{\partial t} + f\bar{u} = -g\frac{\partial \eta}{\partial y} \quad (1b)$$

$$\frac{\partial \eta}{\partial t} + h \left(\frac{\partial \bar{u}}{\partial x} + \frac{\partial \bar{v}}{\partial y} \right) = 0. \quad (1c)$$



$$\frac{\partial \eta}{\partial t} - \beta \frac{gh}{f^2} \frac{\partial \eta}{\partial x} = 0 \quad (2)$$

↑
Phase velocity!
(nondispersive wave)



Recap

Weak consideration of inertia via perturbation (e.g., $\bar{u} = \bar{u}_0 + \bar{u}_1$):

$$\frac{\partial(\bar{u}_0 + \bar{u}_1)}{\partial t} - (\bar{f}_0 + \beta y)(\bar{v}_0 + \bar{v}_1) = -g \frac{\partial \eta}{\partial x} \quad (3)$$

Big

Small

O(1) balance (big of big):

$$-f_0 \bar{v}_0 = -g \frac{\partial \eta}{\partial x} \quad (4)$$

Substituting (4) and (5) into continuity yields

O(Ro) balance (small of big):

$$\frac{\partial \bar{u}_0}{\partial t} - f_0 \bar{v}_1 + \beta y \bar{v}_0 = 0 \quad (5)$$

$$\begin{aligned} \frac{\partial \eta}{\partial t} - \beta \frac{gh}{f_0^2} \frac{\partial \eta}{\partial x} \\ - \frac{gh}{f_0^2} \frac{\partial}{\partial t} \left(\frac{\partial^2 \eta}{\partial x^2} + \frac{\partial^2 \eta}{\partial y^2} \right) = 0 \end{aligned} \quad (6)$$

Topographic Rossby wave

Let us consider f -plane with a slope $h(y) = -\alpha y + h_0$:

$$-f_0 \bar{v} = -g \frac{\partial \eta}{\partial x} \quad (7a)$$

so-called topographic Rossby wave.

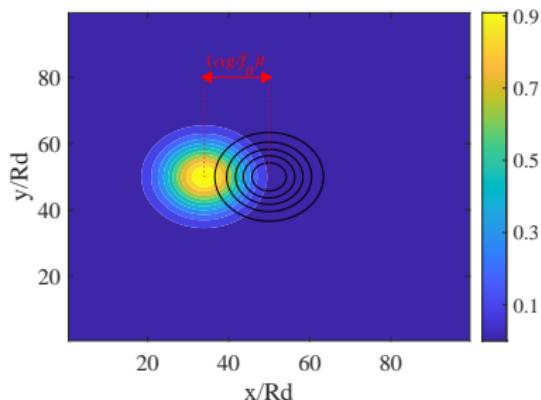
$$f_0 \bar{u} = -g \frac{\partial \eta}{\partial y} \quad (7b)$$

$$\frac{\partial \eta}{\partial t} + h \frac{\partial \bar{u}}{\partial x} + \frac{\partial(\bar{v})}{\partial y} = 0 \quad (7c)$$



$$\frac{\partial \eta}{\partial t} - \frac{\alpha g}{f_0} \frac{\partial \eta}{\partial x} = 0 \quad (8)$$

↑
Phase velocity
(nondispersive wave)



Topographic Rossby wave perturbed by inertia

$$\frac{\partial \bar{u}}{\partial t} - f_0 \bar{v} = -g \frac{\partial \eta}{\partial x} \quad (9a)$$

$$\frac{\partial \bar{v}}{\partial t} + f_0 \bar{u} = -g \frac{\partial \eta}{\partial y} \quad h = -\alpha y + h_0 \quad (9b)$$

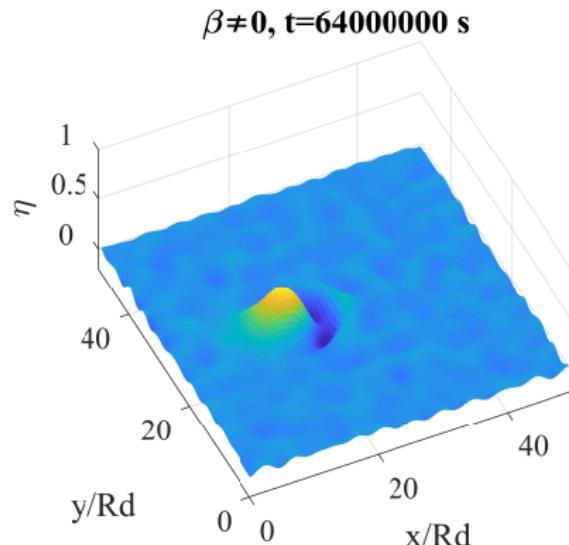
$$\frac{\partial \eta}{\partial t} + h \frac{\partial \bar{u}}{\partial x} + \frac{\partial(h \bar{v})}{\partial y} = 0 \quad (9c)$$

By assuming $Ro \ll 1$ and $\alpha L \ll h_0$, based on $O(1)$ and $O(Ro)$ balances, equation for η is given by

$$\frac{\partial \eta}{\partial t} - \frac{\alpha g}{f_0} \frac{\partial \eta}{\partial x} - \frac{gh_0}{f_0^2} \left(\frac{\partial^2 \eta}{\partial x^2} + \frac{\partial^2 \eta}{\partial y^2} \right) = 0 \quad (10)$$

HW: Obtain (10) from (9) using the approximations and perturbation theory.

Topographic Rossby wave perturbed by inertia



https://jang-geun.github.io/vis_topo_rossby_small.gif

Governing equation

Near the western boundary where $\bar{u} \approx 0$, governing equations can be simplified to

$$-f_0\bar{v} = -g \frac{\partial \eta}{\partial x} \quad (11a)$$

$$\frac{\partial \bar{v}}{\partial t} = -g \frac{\partial \eta}{\partial y} \quad (11b)$$

$$\frac{\partial \eta}{\partial t} + h \frac{\partial \bar{v}}{\partial y} = 0. \quad (11c)$$

Note that (11b) and (11c) can be rewritten as

$$\frac{\partial^2 \eta}{\partial t^2} = gh \frac{\partial^2 \eta}{\partial y^2} \quad (12)$$

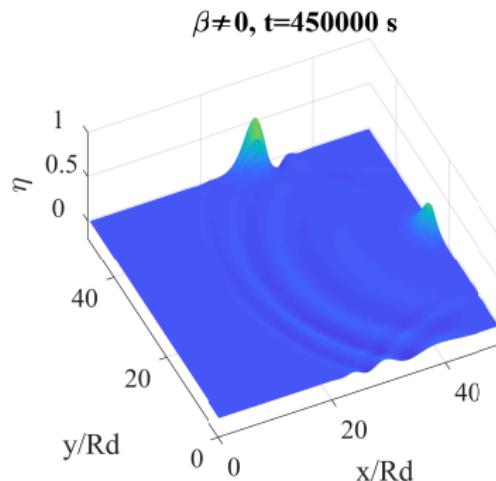
that is identical to the wave equation! But it is constrained by geostrophic balance (11a), that can be rewritten as

$$f_0 \frac{\partial \eta}{\partial y} = -\frac{\partial^2 \eta}{\partial t \partial x} \quad (13)$$

Kelvin wave

Solution to the problem, for a single mode near western boundary, is given by

$$\eta = \eta_0 e^{-x/Rd} e^{il(y+\sqrt{gh}t)} \quad (14)$$



https://jang-geun.github.io/vis_kelvin.gif

Equatorial Kelvin wave

Near the equator where $\bar{v} \approx 0$

$$\frac{\partial \bar{u}}{\partial t} = -g \frac{\partial \eta}{\partial x} \quad (15a)$$

$$\beta y \bar{u} = -g \frac{\partial \eta}{\partial y} \quad (15b)$$

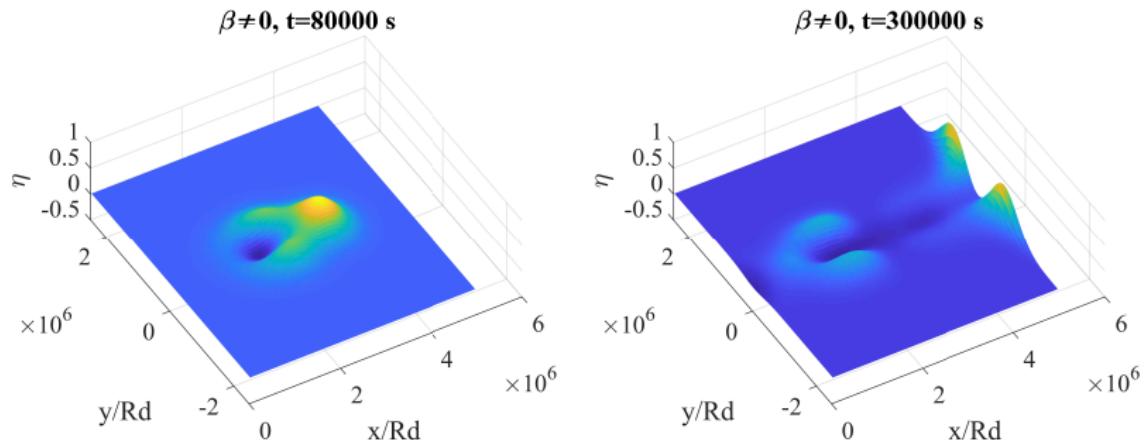
$$\frac{\partial \eta}{\partial t} + h \frac{\partial \bar{u}}{\partial x} = 0 \quad (15c)$$

that yields

$$\frac{\partial^2 \eta}{\partial x^2} = gh \frac{\partial^2 \eta}{\partial x^2} \quad (16a)$$

$$\beta y \frac{\partial \eta}{\partial x} = \frac{\partial^2 \eta}{\partial t \partial y} \quad (16b)$$

Conclusion



https://jang-geun.github.io/vis_equatorial_kelvin.gif