

# Mathematical Biology Assignment 1

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1. A boat carries  $N$  similar rowers each of whom puts in the same  $P$  power to propelling the boat

- (a) Assuming each rower occupies the same  $V$  volume of the boat, show the wetted area of the boat is  $A \propto (NV)^{2/3}$ .  $A$  has units  $[L]^2$ ,  $NV$  has units  $[L]^3$  since  $NV$  is the volume occupied by  $N$  rowers

Want to find non-dimensional groupings which work, since the system will have form

$$f(\lambda_1, \lambda_2, \dots) = 0$$

Where  $\lambda_i$  are non-dimensional groupings of terms. In this case we can generate the non-dimensional grouping

$$\frac{A^3}{(NV)^2}$$

And hence

$$A^3 \propto (NV)^2 \implies A \propto (NV)^{2/3}$$

- (b) Assuming that  $F_{drag}$  depends on the wetted area of the boat,  $A$ , its speed,  $U$ , and the density of the water,  $\rho$ , show that  $F_{drag}$  is proportional to  $\rho U^2 A$  and that the rate of energy dissipation due to drag must be proportional to  $\rho U^3 A$ .

Force has units  $[M][L][T]^{-2}$ , density  $[M][L]^{-3}$ , speed  $[L][T]^{-1}$ , area  $[L]^2$ . If the LHS is the quantity, and the RHS is the units (I will write  $F$  instead of  $F_{drag}$ )

$$\begin{aligned} F &= \frac{[M][L]}{[T]^2} \\ \frac{F}{\rho} &= \frac{[M][L][L]^3}{[M][T]^2} \\ \frac{F}{\rho} &= \frac{[L]^4}{[T]^2} \\ \frac{F}{\rho U^2} &= \frac{[L]^4[T]^2}{[L]^2[T]^2} \\ \frac{F}{\rho U^2} &= [L]^2 \\ \frac{F}{\rho U^2 A} &= \frac{[L]^2}{[L]^2} = [1] \end{aligned}$$

Hence

$$f\left(\frac{F}{\rho U^2 A}\right) = const \implies F \propto \rho U^2 A$$

And the rate of energy dissipation, has same units as power:  $dE$  has units  $[M][L]^2[T]^{-3}$  (energy/time)

$$\begin{aligned} dE &= [M][L]^2[T]^{-3} \\ \frac{dE}{\rho U^2 A} &= \frac{[L]}{[T]} \\ \frac{dE}{\rho U^3 A} &= [1] \end{aligned}$$

And as before we get

$$f\left(\frac{dE}{\rho U^3 A}\right) = \text{const} \implies dE \propto \rho U^3 A$$

(c) Hence show  $U \propto N^{1/9} P^{1/3} \rho^{-1/3} V^{-2/9}$

From part (a),  $A \propto (NV)^{2/3}$  and the power provided,  $NP$  with units  $[M][L]^2[T]^{-3}$

$$\begin{aligned} dE &\propto \rho U^3 A \\ U &\propto \left(\frac{dE}{\rho A}\right)^{1/3} \\ U &\propto dE^{1/3} \rho^{-1/3} (NV)^{-2/9} \end{aligned}$$

But  $dE \propto NP$

$$\begin{aligned} U &\propto dE^{1/3} \rho^{-1/3} (NV)^{-2/9} \\ U &\propto (NP)^{1/3} \rho^{-1/3} (NV)^{-2/9} \\ U &\propto N^{1/9} P^{1/3} \rho^{-1/3} V^{-2/9} \end{aligned}$$

(d) If we assume  $P, V$  are both propto body mass, is size an advantage to a rower? A rower will ideally generate the most speed, so it is an advantage if  $U$  is bigger.

If  $P \propto V \propto M$  then we can sub it into the  $U$  equation

$$\begin{aligned} U &\propto N^{1/9} P^{1/3} \rho^{-1/3} V^{-2/9} \\ U &\propto N^{1/9} M^{1/3} \rho^{-1/3} M^{-2/9} \\ U &\propto N^{1/9} M^{1/9} \rho^{-1/3} \end{aligned}$$

Since the power of  $M$  is positive, yes it would be an advantage.

## 2. Investigate the Coriolis effect

Navier-Stokes gives

$$\rho \left( \frac{\partial \mathbf{u}}{\partial t} + 2\boldsymbol{\Omega} \times \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = -\nabla p + \mu \nabla^2 \mathbf{u} + \rho(\mathbf{g} - \boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{x}))$$

Assuming the origin of the coordinate system is the centre of the Earth.  $\boldsymbol{\Omega}$  is  $2\pi$  per 24 hours (or  $7.3 \times 10^{-5} \text{ s}^{-1}$ ) in the direction of the Earth's axis of rotation. The radius of

the earth is  $\approx 6,400km$ . Assume the water body is a bathtub with lengthscale  $\sim 1m$  and water flows  $\sim 1ms^{-1}$ . Take density  $= 1000kg\ m^{-3}$  and viscosity  $8.9 \times 10^{-4}Pa\ s$ . Non dimensionalise and determine if the LHS (Coriolis acceleration) is significant (and hence if swirl direction will change depending on which hemisphere you are in)

Where  $2\Omega \times \mathbf{u}$  is the Coriolis gravity term.

Let  $x = L\tilde{x}$ ,  $u = U\tilde{u}$ ,  $t = \frac{L}{U}\tilde{t}$ ,  $p = P\tilde{p}$  and  $\Omega = \frac{U}{L}\tilde{\Omega}$

$$\begin{aligned} \rho \left( \frac{\partial \mathbf{u}}{\partial t} + 2\Omega \times \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} \right) &= -\nabla p + \mu \nabla^2 \mathbf{u} + \rho(\mathbf{g} - \Omega \times (\Omega \times \mathbf{x})) \\ \rho \left( \frac{U^2}{L} \frac{\partial \tilde{\mathbf{u}}}{\partial \tilde{t}} + \frac{U^2}{L} 2\tilde{\Omega} \times \tilde{\mathbf{u}} + \frac{U^2}{L} \tilde{\mathbf{u}} \cdot \nabla \tilde{\mathbf{u}} \right) &= -\frac{P}{L} \nabla \tilde{p} + \mu \frac{U}{L^2} \nabla^2 \tilde{\mathbf{u}} + \rho \left( \frac{U^2}{L} \tilde{\mathbf{g}} - \frac{U^2}{L^2} L\tilde{\Omega} \times (\tilde{\Omega} \times \tilde{\mathbf{x}}) \right) \\ \frac{\partial \tilde{\mathbf{u}}}{\partial \tilde{t}} + 2\tilde{\Omega} \times \tilde{\mathbf{u}} + \tilde{\mathbf{u}} \cdot \nabla \tilde{\mathbf{u}} &= -\frac{P}{U^2 \rho} \nabla \tilde{p} + \frac{\mu}{U \rho L} \nabla^2 \tilde{\mathbf{u}} + \tilde{\mathbf{g}} - \tilde{\Omega} \times (\tilde{\Omega} \times \tilde{\mathbf{x}}) \end{aligned}$$

$$\frac{\partial \tilde{\mathbf{u}}}{\partial \tilde{t}} + 2\tilde{\Omega} \times \tilde{\mathbf{u}} + \tilde{\mathbf{u}} \cdot \nabla \tilde{\mathbf{u}} = -\frac{P}{U^2 \rho} \nabla \tilde{p} + \frac{1}{Re} \nabla^2 \tilde{\mathbf{u}} + \tilde{\mathbf{g}} - \tilde{\Omega} \times (\tilde{\Omega} \times \tilde{\mathbf{x}})$$

$$Re \left( \frac{\partial \tilde{\mathbf{u}}}{\partial \tilde{t}} + 2\tilde{\Omega} \times \tilde{\mathbf{u}} + \tilde{\mathbf{u}} \cdot \nabla \tilde{\mathbf{u}} \right) = -\frac{P\mu L}{U} \nabla \tilde{p} + \nabla^2 \tilde{\mathbf{u}} + Re(\tilde{\mathbf{g}} - \tilde{\Omega} \times (\tilde{\Omega} \times \tilde{\mathbf{x}}))$$

Where  $Re = \frac{U\rho L}{\mu}$  If  $Re \rightarrow 0$  Then the Coriolis effect is negligible, i.e. we get (letting  $P = U^2 \rho$ )

$$-\nabla \tilde{p} + \nabla^2 \tilde{\mathbf{u}} = 0$$

With no curl terms. If  $Re \rightarrow \infty$ , then the Coriolis effect is not negligible:

$$\frac{\partial \tilde{\mathbf{u}}}{\partial \tilde{t}} + 2\tilde{\Omega} \times \tilde{\mathbf{u}} + \tilde{\mathbf{u}} \cdot \nabla \tilde{\mathbf{u}} = -\nabla \tilde{p} + \tilde{\mathbf{g}} - \tilde{\Omega} \times (\tilde{\Omega} \times \tilde{\mathbf{x}})$$

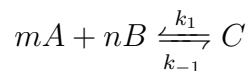
Where  $P = \frac{U}{\mu L}$

Using the given numbers:  $U \sim 1$  and  $L \sim 1$ ,  $\rho \sim 1000$  and  $\mu \sim 8.9 \times 10^{-4}$ .

$$Re = \frac{1000}{8.9 \times 10^{-4}} = 8.9 \times 10^7$$

Which is quite large. Hence the Coriolis effect is not negligible.

3. Consider the chemical equation



(a) Given that the concentration  $c$  of  $C$  is

$$\frac{dc}{dt} = k_1 a^m b^n - k_{-1} c$$

Write the equations for  $a, b$  (the concentrations of  $A, B$ ) The DEs for  $a, b$  are:

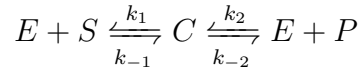
$$\begin{aligned}\frac{da}{dt} &= -k_1 a^m b^n + k_{-1} c \\ \frac{db}{dt} &= -k_1 a^m b^n + k_{-1} c\end{aligned}$$

Subject to initial conditions with form  $a(0) = a_i, b(0) = b_i$  for some  $a_i, b_i$ .

(b) Using conservation, eliminate  $a$ , and  $b$  from the equation for  $c$

$$\begin{aligned}\frac{d(a+c)}{dt} &= 0 \implies a+c = \text{const} = a_i \\ \frac{d(b+c)}{dt} &= 0 \implies b+c = \text{const} = b_i \\ \frac{dc}{dt} &= k_1 a^m b^n - k_{-1} c \\ &= k_1 (c - a_i)^m (c - b_i)^n - k_{-1} c\end{aligned}$$

4. Consider Michaelis-Menten, but relax irreversibility, i.e.



(a) Write the set of equations for concentrations  $s, e, c$ , and  $p$ . With ICs

$$s(0) = s_i > 0, \quad e(0) = e_i = \epsilon s_i > 0, \quad c(0) = p(0) = 0$$

$$\begin{aligned}\frac{de}{dt} &= -k_1 es + k_{-1} c + k_2 ep - k_{-2} ep \\ \frac{ds}{dt} &= -k_1 es + k_{-1} c \\ \frac{dc}{dt} &= k_1 es - k_{-1} es - k_2 c + k_{-2} ep \\ \frac{dp}{dt} &= k_2 c - k_{-2} ep\end{aligned}$$

With

$$s(0) = s_i > 0, \quad e(0) = e_i = \epsilon s_i > 0, \quad c(0) = p(0) = 0$$

(b) Obtain a conservation law and eliminate  $e$  from the system

$$\begin{aligned}\frac{d(e+c)}{dt} &= 0 \implies e(t) + c(t) = \text{const} = e_i = \epsilon s_i \\ &\implies e = \epsilon s_i - c\end{aligned}$$

Hence the system of equations reduces to

$$\begin{aligned}\frac{ds}{dt} &= -k_1 (\epsilon s_i - c)s + k_{-1} c \\ \frac{dc}{dt} &= k_1 (\epsilon s_i - c)s - k_{-1} (\epsilon s_i - c)s - k_2 c + k_{-2} (\epsilon s_i - c)p \\ \frac{dp}{dt} &= k_2 c - k_{-2} (\epsilon s_i - c)p\end{aligned}$$

- (c) Nondimensionalise as in lectures with  $\epsilon = e_i/s_i \ll 1$ . Show the scaled system has 3 dimensionless parameters (give in terms of the original parameters)

Let

$$s = s_i \tilde{s}, \quad c = c_* \tilde{c}, \quad p = p_* \tilde{p}, \quad t = t_* \tilde{t}$$

(Assuming  $s_i$  for  $s$  is sensible as it is given as the IC)  $s$  equation:

$$\begin{aligned} \frac{s_i}{t_*} \frac{d\tilde{s}}{d\tilde{t}} &= -k_1(\epsilon s_i - c_* \tilde{c}) \tilde{s} + k_{-1} c_* \tilde{c} \\ \frac{d\tilde{s}}{d\tilde{t}} &= -t_* k_1 (\epsilon s_i - c_* \tilde{c}) \tilde{s} + \frac{k_{-1} t_* c_*}{s_i} \tilde{c} \end{aligned}$$

$c$  equation

$$\begin{aligned} \frac{c_*}{t_*} \frac{d\tilde{c}}{d\tilde{t}} &= k_1 s_i (\epsilon s_i - c_* \tilde{c}) \tilde{s} - k_{-1} s_i (\epsilon s_i - c_* \tilde{c}) \tilde{s} - k_2 c_* \tilde{c} + k_{-2} (\epsilon s_i - c_* \tilde{c}) p_* \tilde{p} \\ \frac{d\tilde{c}}{d\tilde{t}} &= k_1 t_* s_i \left( \frac{\epsilon s_i}{c_*} - \tilde{c} \right) \tilde{s} - k_{-1} t_* \left( \frac{\epsilon s_i}{c_*} - \tilde{c} \right) \tilde{s} - k_2 t_* \tilde{c} + k_{-2} t_* p_* \left( \frac{\epsilon s_i}{c_*} - \tilde{c} \right) \tilde{p} \end{aligned}$$

$p$  equation

$$\begin{aligned} \frac{p_*}{t_*} \frac{d\tilde{p}}{d\tilde{t}} &= k_2 c_* \tilde{c} - k_{-2} (\epsilon s_i - c_* \tilde{c}) p_* \tilde{p} \\ \frac{d\tilde{p}}{d\tilde{t}} &= k_2 \frac{t_* c_*}{p_*} \tilde{c} - k_{-2} (\epsilon s_i - c_* \tilde{c}) t_* \tilde{p} \end{aligned}$$

Let  $c_* = \epsilon s_i$ . The system becomes:

$$\begin{aligned} \frac{d\tilde{s}}{d\tilde{t}} &= -t_* k_1 \epsilon s_i (1 - \tilde{c}) \tilde{s} + k_{-1} t_* \epsilon \tilde{c} \\ \frac{d\tilde{c}}{d\tilde{t}} &= k_1 t_* s_i (1 - \tilde{c}) \tilde{s} - k_{-1} t_* (1 - \tilde{c}) \tilde{s} - k_2 t_* \tilde{c} + k_{-2} t_* p_* (1 - \tilde{c}) \tilde{p} \\ \frac{d\tilde{p}}{d\tilde{t}} &= k_2 \frac{t_* c_*}{p_*} \tilde{c} - k_{-2} \epsilon s_i t_* (1 - \tilde{c}) \tilde{p} \end{aligned}$$

Incomplete

- (d) Neglect  $\mathcal{O}(\epsilon)$  and smaller terms, find the leading order expression for the dimensionless complex concentration, and show that the dimensionless reaction velocity takes form

$$\frac{d\tilde{p}}{d\tilde{t}} = \frac{A_3 \tilde{s} - A_1 A_2 \tilde{p}}{\tilde{s} + A_2 \tilde{p} + A_1 + A_3}$$

We have the conservation

$$\frac{d(p + s + c)}{dt} = 0$$

Where  $A_i$  are constants

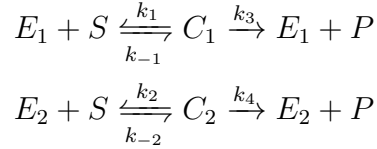
Incomplete

- (e) Show that to leading order, the steady state of product and substrate concentrations (known as the Haldane relationship) is

$$\frac{\tilde{p}}{\tilde{s}} = \frac{k_1 k_2}{k_{-1} k_{-2}}$$

Incomplete

5. Consider



- (a) Write down the equations for the concentrations of  $S, E_1, E_2, C_1, C_2, P$ . Show that there are two conserved quantities and hence reduce the system to 3 equations only containing  $S, C_1, C_2$ .

$$\begin{aligned} \frac{de_1}{dt} &= -k_1 e_1 s + (k_{-1} + k_3) c_1 \\ \frac{de_2}{dt} &= -k_2 e_2 s + (k_{-2} + k_4) c_2 \\ \frac{dc_1}{dt} &= k_1 e_1 s - (k_{-1} + k_3) c_1 \\ \frac{dc_2}{dt} &= k_2 e_2 s - (k_{-2} + k_4) c_2 \\ \frac{ds}{dt} &= -k_1 e_1 s - k_2 e_2 s + k_{-1} c_1 + k_{-2} c_2 \\ \frac{dp}{dt} &= k_3 c_1 + k_4 c_2 \end{aligned}$$

Clearly, by adding the equations,

$$\frac{d(e_1 + c_1)}{dt} = 0$$

$$\frac{d(e_2 + c_2)}{dt} = 0$$

And the  $p$  equation is redundant since

$$\frac{d(c_1 + c_2 + s + p)}{dt} = 0$$

And hence the system can reduce to

$$\begin{aligned} \frac{dc_1}{dt} &= k_1 e_1 s - (k_{-1} + k_3) c_1 \\ \frac{dc_2}{dt} &= k_2 e_2 s - (k_{-2} + k_4) c_2 \\ \frac{ds}{dt} &= -k_1 e_1 s - k_2 e_2 s + k_{-1} c_1 + k_{-2} c_2 \end{aligned}$$

(b) Assume

$$s(0) = s_i > 0, \quad e_1(0) = e_2(0) = e_i = \epsilon s_i > 0$$

Non-dimensionalise to get a system of form

$$\begin{aligned} \frac{ds}{dt} &= -s(1 + \alpha) + c_1(\mu_1 + s) + \alpha c_2(\mu_2 + s) \\ \epsilon \frac{dc_1}{dt} &= s(1 - c_1) - \lambda_1 c_1 \\ \epsilon \frac{dc_2}{dt} &= \alpha [s(1 - c_2) - \lambda_2 c_2] \end{aligned}$$

Defining the parameters in terms of the original dimensionless parameters

By using  $e_1 = e_i - c_1$  and  $e_2 = e_i - c_2$  Let

$$s = s_* \tilde{s}, \quad c_1 = c_1^* \tilde{c}_1, \quad c_2 = c_2^* \tilde{c}_2, \quad t = t^* \tilde{t}$$

$$\begin{aligned} \frac{ds}{dt} &= -k_1 e_1 s - k_2 e_2 s + k_{-1} c_1 + k_{-2} c_2 \\ &= -k_1 (e_i - c_1) s - k_2 (e_i - c_2) s + k_{-1} c_1 + k_{-2} c_2 \\ &= -k_1 e_i s + k_1 c_1 s - k_2 e_i s + k_2 c_2 s + k_{-1} c_1 + k_{-2} c_2 \\ &= -s(e_i k_1 + e_i k_2) + c_1(k_{-1} + k_1 s) + c_2(k_{-2} + k_2 s) \\ &= -s e_i k_1 (1 + \frac{k_2}{k_1}) + k_1 c_1 (\frac{k_{-1}}{k_1} + s) + k_2 c_2 (\frac{k_{-2}}{k_2} + s) \\ \frac{s_*}{t^*} \frac{d\tilde{s}}{d\tilde{t}} &= -s_* \tilde{s} e_i k_1 (1 + \frac{k_2}{k_1}) + k_1 c_1^* \tilde{c}_1 (\frac{k_{-1}}{k_1} + s_* \tilde{s}) + k_2 c_2^* \tilde{c}_2 (\frac{k_{-2}}{k_2} + s_* \tilde{s}) \\ \frac{d\tilde{s}}{d\tilde{t}} &= -t^* \tilde{s} e_i k_1 (1 + \frac{k_2}{k_1}) + t^* k_1 c_1^* \tilde{c}_1 (\frac{k_{-1}}{k_1 s_*} + \tilde{s}) + t^* k_2 c_2^* \tilde{c}_2 (\frac{k_{-2}}{k_2 s_*} + \tilde{s}) \end{aligned}$$

Letting  $t^* = \frac{1}{e_i k_1}$  gives

$$\frac{d\tilde{s}}{d\tilde{t}} = -\tilde{s}(1 + \frac{k_2}{k_1}) + \frac{1}{e_i} c_1^* \tilde{c}_1 (\frac{k_{-1}}{k_1 s_*} + \tilde{s}) + \frac{1}{e_i k_1} k_2 c_2^* \tilde{c}_2 (\frac{k_{-2}}{k_2 s_*} + \tilde{s})$$

Letting  $c_1^* = e_i$ , and  $\frac{k_2}{k_1} =: \alpha$

$$\frac{d\tilde{s}}{d\tilde{t}} = -\tilde{s}(1 + \alpha) + \tilde{c}_1 (\frac{k_{-1}}{k_1 s_*} + \tilde{s}) + \frac{1}{e_i} \alpha c_2^* \tilde{c}_2 (\frac{k_{-2}}{k_2 s_*} + \tilde{s})$$

Letting  $c_2^* = e_i$  gives

$$\frac{d\tilde{s}}{d\tilde{t}} = -\tilde{s}(1 + \alpha) + \tilde{c}_1 (\frac{k_{-1}}{k_1 s_*} + \tilde{s}) + \alpha \tilde{c}_2 (\frac{k_{-2}}{k_2 s_*} + \tilde{s})$$

Considering the  $c$  equations:

$$\begin{aligned}
 \frac{dc_1}{dt} &= k_1 e_1 s - (k_{-1} + k_3) c_1 \\
 &= k_1 (e_i - c_1) s - (k_{-1} + k_3) c_1 \\
 e_i^2 k_1 \frac{d\tilde{c}_1}{d\tilde{t}} &= k_1 (e_i - e_i \tilde{c}_1) s_* \tilde{s} - (k_{-1} + k_3) e_i \tilde{c}_1 \\
 e_i \frac{d\tilde{c}_1}{d\tilde{t}} &= (1 - \tilde{c}_1) s_* \tilde{s} - \frac{(k_{-1} + k_3)}{k_1} \tilde{c}_1 \\
 \frac{e_i}{s_*} \frac{d\tilde{c}_1}{d\tilde{t}} &= \tilde{s} (1 - \tilde{c}_1) - \frac{(k_{-1} + k_3)}{k_1 s_*} \tilde{c}_1
 \end{aligned}$$

The  $c_2$  equation

$$\begin{aligned}
 \frac{dc_2}{dt} &= k_2 e_2 s - (k_{-2} + k_4) c_2 \\
 &= k_2 (e_i - c_2) s - (k_{-2} + k_4) c_2 \\
 e_i^2 k_1 \frac{d\tilde{c}_2}{d\tilde{t}} &= k_2 (e_i - e_i \tilde{c}_2) s_* \tilde{s} - (k_{-2} + k_4) e_i \tilde{c}_2 \\
 e_i \frac{d\tilde{c}_2}{d\tilde{t}} &= \frac{k_2}{k_1} (1 - \tilde{c}_2) s_* \tilde{s} - \frac{(k_{-2} + k_4)}{k_1} \tilde{c}_2 \\
 \frac{e_i}{s_*} \frac{d\tilde{c}_2}{d\tilde{t}} &= \alpha \left( (1 - \tilde{c}_2) \tilde{s} - \frac{(k_{-2} + k_4)}{k_2 s_*} \tilde{c}_2 \right)
 \end{aligned}$$

Using  $s_* = s_i$  and  $e_i = \epsilon s_i$  gives, for the  $c$  equations:

$$\begin{aligned}
 \epsilon \frac{d\tilde{c}_1}{d\tilde{t}} &= \tilde{s} (1 - \tilde{c}_1) - \frac{(k_{-1} + k_3)}{k_1 s_i} \tilde{c}_1 \\
 \epsilon \frac{d\tilde{c}_2}{d\tilde{t}} &= \alpha \left( (1 - \tilde{c}_2) \tilde{s} - \frac{(k_{-2} + k_4)}{k_2 s_i} \tilde{c}_2 \right)
 \end{aligned}$$

And letting

$$\lambda_1 = \frac{k_{-1} + k_3}{k_1 s_i}, \quad \lambda_2 = \frac{k_{-2} + k_4}{k_2 s_i}$$

Finally gives (dropping tildes)

$$\begin{aligned}
 \epsilon \frac{dc_1}{dt} &= s(1 - c_1) - \lambda_1 c_1 \\
 \epsilon \frac{dc_2}{dt} &= \alpha [s(1 - c_2) - \lambda_2 c_2]
 \end{aligned}$$

Returning to the  $s$  equation

$$\frac{d\tilde{s}}{d\tilde{t}} = -\tilde{s}(1 + \alpha) + \tilde{c}_1 \left( \frac{k_{-1}}{k_1 s_i} + \tilde{s} \right) + \alpha \tilde{c}_2 \left( \frac{k_{-2}}{k_2 s_i} + \tilde{s} \right)$$

And letting

$$\mu_1 = \frac{k_{-1}}{k_1 s_i}, \quad \mu_2 = \frac{k_{-2}}{k_2 s_i}$$



Gives (dropping the tildes)

$$\frac{ds}{dt} = -s(1 + \alpha) + c_1(\mu_1 + s) + \alpha c_2(\mu_2 + s)$$

I.e. we have the system

$$\begin{aligned}\frac{ds}{dt} &= -s(1 + \alpha) + c_1(\mu_1 + s) + \alpha c_2(\mu_2 + s) \\ \epsilon \frac{dc_1}{dt} &= s(1 - c_1) - \lambda_1 c_1 \\ \epsilon \frac{dc_2}{dt} &= \alpha [s(1 - c_2) - \lambda_2 c_2]\end{aligned}$$

Where

$$\alpha = \frac{k_2}{k_1}, \quad \mu_1 = \frac{k_{-1}}{k_1 s_i}, \quad \mu_2 = \frac{k_{-2}}{k_2 s_i}, \quad \lambda_1 = \frac{k_{-1} + k_3}{k_1 s_i}, \quad \lambda_2 = \frac{k_{-2} + k_4}{k_2 s_i}$$

(c) Find leading order solutions for  $c_1, c_2$  and hence  $s$  (technology is allowed)

To leading order, take perturbation series:

$$c_1 = c_{10} + c_{11}\epsilon + \dots$$

$$c_2 = c_{20} + c_{21}\epsilon + \dots$$

$$s = s_0 + s_1\epsilon + \dots$$

Where  $\epsilon \ll 1$ . The leading order system is:

$$\begin{aligned}\frac{ds_0}{dt} &= -s_0(1 + \alpha) + c_{10}(\mu_1 + s_0) + \alpha c_{20}(\mu_2 + s_0) \\ 0 &= s_0(1 - c_{10}) - \lambda_1 c_{10} \\ 0 &= \alpha [s_0(1 - c_{20}) - \lambda_2 c_{20}]\end{aligned}$$

Solve the last two:

$$\begin{aligned}0 &= s_0(1 - c_{10}) - \lambda_1 c_{10} \\ s_0 c_{10} + \lambda_1 c_{10} &= s_0 \\ c_{10} &= \frac{s_0}{s_0 + \lambda_1}\end{aligned}$$

$$\begin{aligned}0 &= s_0(1 - c_{20}) - \lambda_2 c_{20} \\ c_{20} &= \frac{s_0}{s_0 + \lambda_2}\end{aligned}$$

And use in the first equation:

$$\begin{aligned}\frac{ds_0}{dt} &= -s_0(1 + \alpha) + c_{10}(\mu_1 + s_0) + \alpha c_{20}(\mu_2 + s_0) \\ &= -s_0(1 + \alpha) + \frac{s_0}{s_0 + \lambda_1}(\mu_1 + s_0) + \alpha \frac{s_0}{s_0 + \lambda_2}(\mu_2 + s_0)\end{aligned}$$

$$\int s_0(1 + \alpha) - \frac{s_0}{s_0 + \lambda_1}(\mu_1 + s_0) - \alpha \frac{s_0}{s_0 + \lambda_2}(\mu_2 + s_0) ds = \int dt$$

$$\begin{aligned} \frac{1}{2}s_0^2(1 + \alpha) + s_0(\lambda_1 - \mu_1) + s_0\alpha(\lambda_2 - \mu_2) - \alpha\lambda_2 \log(\lambda_2 + s_0)(\lambda_2 - \mu_2) \\ - \frac{1}{2}\alpha s_0^2 + \lambda_1 \log(\lambda_1 + s_0)(-\lambda_1 + \mu_1) - \frac{1}{2}s_0^2 = k \end{aligned}$$

$$\begin{aligned} s_0(\lambda_1 - \mu_1) + s_0\alpha(\lambda_2 - \mu_2) - \alpha\lambda_2 \log(\lambda_2 + s_0)(\lambda_2 - \mu_2) \\ + \lambda_1 \log(\lambda_1 + s_0)(-\lambda_1 + \mu_1) = k \end{aligned}$$

Where  $k$  would be obtained using  $s(0) = s_i$ . This is an implicit solution in  $s_0$ .  
Solved the integrals using **MATLAB**

```

1 syms s lambda1 lambda2 mu1 mu2 a
2 p1 = int(s*(1+a), s)
3 p2 =int(s*(mu1+s)/(s+lambda1), s)
4 p3 =int(a*s*(mu2+s)/(s+lambda2), s)

```

(d) Rescale time and find the inner solutions for  $s, c_1, c_2$  Rescale time so that  $T = \frac{t}{\epsilon}$

$$\frac{d}{dt} = \frac{\partial T}{\partial t} \frac{\partial}{\partial T} = \frac{1}{\epsilon} \frac{\partial}{\partial T}$$

Giving (to leading order) for the inner solutions:

$$\begin{aligned} \frac{1}{\epsilon} \frac{ds_0}{dT} &= -s_0(1 + \alpha) + c_{10}(\mu_1 + s_0) + \alpha c_{20}(\mu_2 + s_0) \\ \frac{\partial c_1}{\partial T} &= s_0(1 - c_{10}) - \lambda_1 c_{10} \\ \frac{\partial c_2}{\partial T} &= \alpha [s_0(1 - c_{20}) - \lambda_2 c_{20}] \end{aligned}$$

And hence

$$\begin{aligned} \frac{ds_0}{dT} &= 0 \\ \frac{\partial c_1}{\partial T} &= s_0(1 - c_{10}) - \lambda_1 c_{10} \\ \frac{\partial c_2}{\partial T} &= \alpha [s_0(1 - c_{20}) - \lambda_2 c_{20}] \end{aligned}$$

Trivially  $s_0(T) = s_i$

$$\begin{aligned} \frac{\partial c_1}{\partial T} &= s_0(1 - c_{10}) - \lambda_1 c_{10} \\ \frac{\partial c_1}{\partial T} &= s_i(1 - c_{10}) - \lambda_1 c_{10} \\ \frac{1}{s_i(1 - c_{10}) - \lambda_1 c_{10}} dc_1 &= \int dT \end{aligned}$$

Let  $u = s_i(1 - c_{10}) - \lambda_1 c_{10}$ ,  $du = -\lambda_1 - s_i dc_1$

$$\begin{aligned}
 \frac{1}{s_i(1 - c_{10}) - \lambda_1 c_{10}} dc_1 &= \int dT \\
 -\frac{1}{s_i + \lambda_1} \int \frac{1}{u} du &= T + a \\
 -\frac{1}{s_i + \lambda_1} \log(u) &= T + a \\
 -\frac{1}{s_i + \lambda_1} \log(s_i(1 - c_{10}) - \lambda_1 c_{10}) &= T + a \log(s_i(1 - c_{10}) - \lambda_1 c_{10}) = -T(s_i + \lambda_1) + b \\
 s_i(1 - c_{10}) - \lambda_1 c_{10} &= e^{-T(s_i + \lambda_1) + b} c_{10} = k_1 e^{-T(s_i + \lambda_1)} + \frac{s_i}{s_i + \lambda_1}
 \end{aligned}$$

For the inner  $c_2$  equation, the outcome is very similar:

$$\begin{aligned}
 \frac{\partial c_2}{\partial T} &= \alpha [s_i(1 - c_{20}) - \lambda_2 c_{20}] \log(s_i(1 - c_{20}) - \lambda_2 c_{20}) = -Tc(s_i + \lambda_2) + a \\
 c_{20} &= k_2 e^{-Tc(s_i + \lambda_2)} + \frac{s_i}{s_i + \lambda_2}
 \end{aligned}$$

The full solution could be obtained using a matching condition, but this would be difficult since the outer solution for  $s$  is implicit.