The conformal Killing spinor initial data equations

Edgar Gasperín

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Spinorial notation in a nutshell

Translation spinors and world tensors $a \to AA'$. Given T_{ab} the spinorial counterpart is given by

$$T_{\mathbf{A}\mathbf{A}'\mathbf{B}\mathbf{B}'} = T_{\mathbf{a}\mathbf{b}}\sigma^{\mathbf{a}}{}_{\mathbf{A}\mathbf{A}'}\sigma^{\mathbf{b}}{}_{\mathbf{B}\mathbf{B}'}$$

where $\sigma^a{}_{AA'}$ are the Infeld-van der Waerden symbols (Pauli matrices & Identity)

$$(\alpha_{\mathbf{0}},\alpha_{\mathbf{1}},\alpha_{\mathbf{2}},\alpha_{\mathbf{3}}) \rightarrow \frac{1}{\sqrt{2}} \begin{bmatrix} \alpha_{\mathbf{0}} + \alpha_{\mathbf{3}} & \alpha_{\mathbf{1}} - \mathrm{i}\alpha_{\mathbf{2}} \\ \alpha_{\mathbf{1}} + \mathrm{i}\alpha_{\mathbf{2}} & \alpha_{\mathbf{0}} - \alpha_{\mathbf{3}} \end{bmatrix}$$

(frame-) metric:
$$g_{AA'BB'} = \epsilon_{AB}\epsilon_{A'B'}$$

Raise and lower indices:
$$\xi_B = \xi^A \epsilon_{AB}$$

Curvature spinors

Riemann → Weyl, tracefree Ricci, Ricci scalar.

$$\Psi_{ABCD}, \qquad \Phi_{ABA'B'}, \qquad \Lambda$$

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 $\mathsf{Riemann} \to \mathsf{Weyl}, \quad \mathsf{tracefree} \ \mathsf{Ricci}, \quad \mathsf{Ricci} \ \mathsf{scalar}.$

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Some Killing objects

Killing vectors

$$\tilde{\nabla}_{(a}\tilde{\xi}_{b)} = 0.$$

Killing tensors

$$\tilde{K}_{ab} = \tilde{K}_{(ab)}, \quad \tilde{\nabla}_{(a}\tilde{K}_{bc)} = 0,$$

Killing-Yano tensors

$$\begin{split} \tilde{Y}_{ab} &= \tilde{Y}_{[ab]}, \quad \tilde{\nabla}_{(a} \tilde{Y}_{b)c} = 0, \\ \tilde{K}_{ab} &= \tilde{Y}_a{}^c \tilde{Y}_{cb}. \end{split}$$

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$$\tilde{\kappa}_{AB} = \tilde{\kappa}_{(AB)}, \quad \tilde{\nabla}_{A'(A}\tilde{\kappa}_{BC)} = 0$$

$$\tilde{Y}_{AA'BB'} = \mathrm{i}(\tilde{\kappa}_{AB}\tilde{\epsilon}_{A'B'} - \bar{\tilde{\kappa}}_{A'B'}\tilde{\epsilon}_{AB})$$
Uling spinors are more primitive objects

An example: Kerr spacetime

$$\tilde{\xi}^a = (\partial_t)^a, \ \tilde{\eta}^a = (\partial_{\varphi})^a,$$

$$\exists \tilde{K}_{ab}, \qquad \tilde{u}^a \tilde{\nabla}_a \tilde{u}^b = 0.$$

$$\mu = \tilde{u}^a \tilde{u}_a, \quad e = \tilde{\xi}^a \tilde{u}_a,$$

$$\ell = \tilde{\eta}^a \tilde{u}_a, \quad C = \tilde{K}_{ab} \tilde{u}^a \tilde{u}^b$$

- Geodesic motion completely integrable!
- Hidden symmetry

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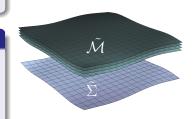
Rigidity thru symmetries. Symmetries thru intial data

Spacetime symmetries

- Black hole uniqueness problems
- Spacetime characterisations

Killing Initial Data KIDs (Chrusciel & Beig 97)

• Given initial data $(\tilde{\Sigma}, \tilde{h}_{ij}, \tilde{\chi}_{ij})$ for the vacuum Einstein Field Equations (EFE) if the Killing vector initial data equations (KID) are satisfied on $\tilde{\Sigma}$ then the spacetime developement admits a Killing vector ξ^a .



KID: Killing vector initial data

• Let $(\tilde{\mathcal{M}}, \tilde{\boldsymbol{g}})$ solution to $\tilde{R}_{ab} = 0$.

$$\tilde{S}_{ab} := \tilde{\nabla}_a \tilde{\xi}_b + \tilde{\nabla}_b \tilde{\xi}_a$$

- ullet $ilde{\xi}_a$ is KV iff $ilde{S}_{ab}=0$ (zero-quantity)
- Does $\tilde{S}_{ab} = 0$ propagate?

Identities to eqs

• Assume vacuum EFE

$$\tilde{R}_{ab} = 0$$

• Assume **candidate** eqn

$$\tilde{\Box} \tilde{\xi}_a = 0 \text{ on } (\tilde{\mathcal{M}}, \tilde{\boldsymbol{g}})$$

Propagation equation

$$\tilde{\Box}\tilde{S}_{ab} = 2\tilde{R}^c{}_{ab}{}^d\tilde{S}_{cd}$$

 $\tilde{S}_{ab} = 0 \& \partial_t \tilde{S}_{ab} = 0 \text{ on } \tilde{\Sigma}$ $\Longrightarrow \tilde{S}_{ab} = 0 \text{ on } \mathcal{W} \subseteq \mathcal{D}^+(\tilde{\Sigma})$

Sketch of the KIDs derivation & proof

• Identity:

$$\tilde{\nabla}^a \tilde{S}_{ab} - \frac{1}{2} \nabla_b S_a{}^a = \tilde{\Box} \tilde{\xi}_b + \tilde{R}_{cb}$$

- ullet Definition: A **candidate** KV $ilde{\zeta}^a$ is a vector satisfying $ilde{\Box} ilde{\zeta}_a=- ilde{R}_{ab} ilde{\zeta}^b$
- Strategy: find ID $(\zeta_a, \partial_t \zeta_a)$ on $\tilde{\Sigma}$ ensuring a candidate gives a true KV.
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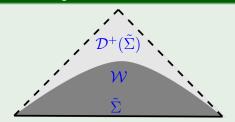
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KID: Killing vector initial data



Trivial data: $\tilde{S}_{ab}=0$ & $\partial_t \tilde{S}_{ab}=0$ on $\tilde{\Sigma}$ \Longrightarrow $(\tilde{\xi}_a,\partial_t \tilde{\xi}_a)$ on $\tilde{\Sigma}$ aka KIDs.

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KSID zero-quantities definitions

$$\tilde{H}_{A'ABC} := 3\tilde{\nabla}_{A'(A}\tilde{\kappa}_{BC)},
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 $\tilde{\Phi}=0$ & $\tilde{H}=0$ \Longrightarrow $\tilde{\xi}$ is a Killing vector!

(SID (Valiente Kroon & Garcia-Parrado 08)

- Candidate KS eq for $\tilde{\kappa}$, Candidate KV eq for $\tilde{\xi}$.
- Assume $\Phi = 0$.
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The Einstein equations under conformal transformations

• The Einstein field equations in vacuum

$$\tilde{R}_{ab} = \lambda \tilde{g}_{ab}$$

are not conformally invariant!.

A calculation shows that

$$g_{ab} = \Omega^2 \tilde{g}_{ab}, \quad \Longrightarrow$$

$$R_{ab} = \tilde{R}_{ab} - 2\Omega^{-1}\nabla_a\nabla_b\Omega - g_{ab}(\Omega^{-1}\nabla^c\nabla_c\Omega - 3\Omega^{-2}\nabla_c\Omega\nabla^c\Omega),$$

where R_{ab} , R and ∇_a are associated to g_{ab} .

- Formally singular whenever $\Omega = 0!$.
- Friedrich's regularisation trick ... and use the curvature as a variable

$$\nabla_a \nabla_b \Omega = -\frac{\Omega}{2} R_{ab} + \dots$$

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The conformal Einstein field equations (H. Friedrich 81)

Equations:

$$\begin{split} \nabla_a \nabla_b \Omega &= -\Omega L_{ab} + s g_{ab}, \\ \nabla_a s &= -L_{ac} \nabla^c \Omega, \\ \nabla_c L_{db} - \nabla_d L_{cb} &= d^a{}_{bcd} \nabla_a \Omega, \\ \nabla_a d^a{}_{bcd} &= 0, \end{split} \qquad \begin{aligned} &\text{Fields:} \quad \left(\Omega, s, L_{ab}, d^a{}_{bcd}\right) \\ L_{ab} &:= \frac{1}{2} R_{ab} - \frac{1}{12} R g_{ab}, \quad \text{(Schouten tensor)} \\ s &:= \frac{1}{4} \nabla_a \nabla^a \Omega + \frac{1}{24} R \Omega, \quad \text{(Friedrich scalar)} \\ d^a{}_{bcd} &:= \Omega^{-1} C^a{}_{bcd}. \quad \text{(Rescaled Weyl tensor)} \end{aligned}$$

Key equation: Rescaled Weyl spinor

Applications

- Potentially turn global problems into local ones.
- $\mathscr{I}\left(\Omega=0\right)$ legit hypersurface to prescribe ID: Asympt. initial value problem.

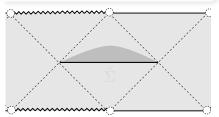
Combining these ideas?

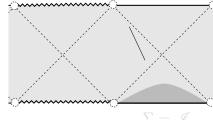
Unphysical Killing vector equation

Given a Killing vector $\tilde{\xi}^a$ on $(\tilde{\mathcal{M}}, \tilde{g})$ and (\mathcal{M}, g) with $g = \Omega^2 \tilde{g}$ then $X_a = \Omega^2 \tilde{\xi}_a$ satisfies the Unphysical Killing vector equations

$$\nabla_a X_b + \nabla_b X_a = \frac{1}{2} \nabla^c X_c g_{ab}, \quad X^a \nabla_a \Omega = \frac{1}{4} \nabla_a X^a.$$

- Field equations: CEFE.
- KID's on spacelike conformal





$$\Sigma = \mathscr{I}$$

Combining these ideas?

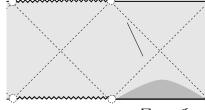
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Given a Killing vector $\tilde{\xi}^a$ on $(\tilde{\mathcal{M}}, \tilde{\boldsymbol{g}})$ and $(\mathcal{M}, \boldsymbol{g})$ with $\boldsymbol{g} = \Omega^2 \tilde{\boldsymbol{g}}$ then $X_a = \Omega^2 \tilde{\xi}_a$ satisfies the **Unphysical Killing vector equations**

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CKID: The conformal Killing vector initial data equations (Paetz 14)

- Field equations: CEFE.
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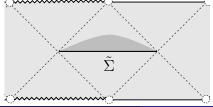


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Apps (Mars Paetz Senovilla 16, 17)

 Characterisations of Kerr-de-Sitter-like spacetimes

←□ → ←□ → ← = → ← = →



Killing spinors and conformal transformations

• KS are conf. inv. $\nabla_{A'(A}\tilde{\kappa}_{BC)}=0 \implies$ $\nabla_{A'(A}\kappa_{BC)} = 0$ with $\kappa_{AB} = \Omega^2 \tilde{\kappa}_{AB}$.

- $\xi_a \mapsto \xi_{AA'} := \nabla^B{}_{A'} \kappa_{AB}$ is **not** a (C)KV.

$$\begin{split} H_{A'ABC} &:= 3\nabla_{A'(A}\kappa_{BC)} \\ S_{AA'BB'} &:= \nabla^Q_{A'}H_{B'QAB} \\ &= \nabla_{AA'}\xi_{BB'} + \nabla_{BB'}\xi_{AA'} + 6\kappa_{(A}{}^Q\Phi_{B)QA'B'} \end{split}$$

- Φ satisfies diff conditions in (\mathcal{M}, g) .
- S not geom motivated.
- Relate Φ to Φ ? \leadsto Singular eqs Ω^{-1} -terms

$$\nabla_{(A}{}^{A'}H_{|A'|BCD)}$$
$$= 6\kappa_{(A}{}^{Q}\Psi_{BCD)Q} = 0$$

$$\phi_{ABCD} = \Omega^{-1} \Psi_{ABCD}$$

$$\nabla \times \boldsymbol{H} = 6\Omega \boldsymbol{E}$$

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Killing spinors and conformal transformations

 $\begin{array}{l} \bullet \ \ \mathsf{KS} \ \mathsf{are} \ \mathsf{conf.} \ \mathsf{inv.} \ \tilde{\nabla}_{A'(A} \tilde{\kappa}_{BC)} = 0 \\ \nabla_{A'(A} \kappa_{BC)} = 0 \ \mathsf{with} \ \kappa_{AB} = \Omega^2 \tilde{\kappa}_{AB}. \end{array}$

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Buchdah

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Buchdahl zero-quantity

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Edgar Gasperín (CENTRA-IST)

Conformal ("Unphysical") Killing Spinor ID

Unphysical zero-quantities

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KSID (G. & Williams 22

- ullet Candidate KS eq for κ
- Assume vacuum CEFE are satisfied
- Closed system of homogeneous wave eqs

$$\Box oldsymbol{H} = \mathcal{G}_1(oldsymbol{H}, oldsymbol{B})$$

$$\Box \boldsymbol{B} = \mathcal{G}_2(\boldsymbol{H}, \boldsymbol{B}, \boldsymbol{F})$$

$$\Box F = \mathcal{G}_3(H, B, F)$$

• Trivial ID on $\Sigma \Longrightarrow$ $\boldsymbol{H} = \boldsymbol{B} = \boldsymbol{F} = 0$ on $\mathcal{W} \subseteq \mathcal{D}^+(\Sigma)$

ullet $(\kappa,\partial_t\kappa)$ on Σ .

KS Candidate equation:

$$\Box \kappa_{AB} = -\frac{1}{6} R \; \kappa_{AB} + \Omega \phi_{ABCD} \kappa^{CD}$$

Initial data on ∑

$$\kappa_{AB} = \mathring{\kappa}_{AB} \nabla_{\boldsymbol{\tau}} \kappa_{AB} = \mathcal{D}_{Q(A} \mathring{\kappa}_{B)}^{Q}$$

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$$\mathcal{D}_{(AB} \mathring{\kappa}_{CD)} = 0, \quad \mathring{\kappa}_{(A}{}^{Q} \phi_{BCD)Q} = 0$$

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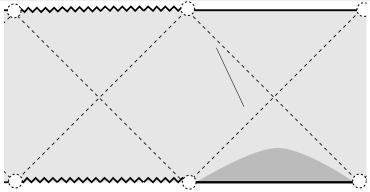
Thm CKSID (G. & Williams 22)

• Given **ID** for the vacuum **CEFE**, if the CKSID conditions

$$\mathcal{D}_{(AB}\kappa_{CD)} = 0, \qquad \kappa_{(A}{}^{Q}\phi_{BCD)Q} = 0$$

are satisfied on Σ , then there exists a Killing spinor in $\mathcal{W} \subseteq \mathcal{D}^+(\Sigma)$.

• In the asymptotic IVP set-up $\Sigma = \mathscr{I}$



Corollary: Weyl collineation

- ullet If $(\mathcal{M}, oldsymbol{g})$ is a conformally Einstein manifold
- $\xi_a \mapsto \xi_{AA'} := \nabla^B_{A'} \kappa_{AB}$ is a **curvature collineation** of the rescaled Weyl spinor:

Associated conformal Killing vector

$$X_{AA'} = \Omega \xi_{AA'} - 3\kappa_{AQ} \nabla_{A'}{}^{Q} \Omega$$

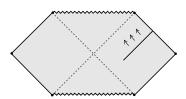
- X_a is a conformal KV of (\mathcal{M}, g)
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Future applications

- Characterisation Kerr-de-Sitter through asympt ID.
- Type D initial data characterisation.

Future directions

- Conds on a spacelikefor simplicity
- Other existence and uniqueness thms for wave eqs →
 - Characteristic IVP
 - IBVP
- → Data on null or timelike ¶



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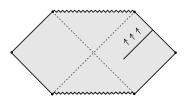
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Many thanks for your attention