

# Conformal twistor KIDs

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## Abstract

Write-up of Edgar's twistor KID calculations from the new notebook.

## 1 Definitions

We are interested here in the twistor equation:

$$\nabla_{A'}(A\kappa_B) = 0,$$

encoded in the vanishing of the *zero quantity*  $H_{A'AB} := 2\nabla_{A'}(A\kappa_B)$ . It will prove useful to define the auxiliary quantity

$$\xi_{A'} := \nabla^A{}_{A'}\kappa_A.$$

We also define the zero quantity  $S_{A'B'A} := \nabla_{QA'}H_{B'A}{}^Q$ . This may be expressed alternatively in terms of the auxiliary field<sup>•1</sup>  $\xi_{A'}$ , by an easy computation, as follows:

$$S_{A'B'A} = -\nabla_{AB'}\xi_{A'} + 2\Lambda\epsilon_{A'B'}\kappa_A - 2\Phi_{AQA'B'}\kappa^Q. \quad (1)$$

•1: I think we should change the sign in either the definition of  $S$  or  $\xi$ , to be more consistent.

The other (symmetrised) contraction yields the *Buchdahl constraint*:

$$0 = \nabla_{(A}{}^{A'}H_{|A'|BC)} = \Psi_{ABCD}\kappa^D,$$

though this won't feature in the calculations here.

## 2 Wave equations

### 2.1 For the twistor fields

A short computation shows that

$$\square\kappa_B = -2\Lambda\kappa_B + \frac{2}{3}\nabla^{AA'}H_{A'AB}.$$

Hence, if  $\kappa_B$  solves the twistor equation, then it necessarily satisfies the wave equation

$$\square\kappa_B + 2\Lambda\kappa_B = 0. \quad (2)$$

We will choose to propagate a twistor candidate according to this equation. It is maybe worth noticing that  $S_{A'}{}^{A'}{}_B = \nabla^{AA'}H_{A'AB}$ . Hence, if  $\kappa_A$  satisfies (2), then necessarily  $S_{A'B'A} =$

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$S_{(A'B')A}$ , since it is assured that  $\nabla^{A'} H_{A'AB} = 0$ .

We will also need a wave equation for the auxiliary field  $\xi_{A'}$ . To get this, take the contracted derivative of (2) and commute derivatives, finally resulting in

$$\square \xi_{A'} + 2\Lambda \xi_{A'} - 8(\nabla_{AA'} \Lambda) \kappa^A = 0. \quad (3)$$

Again, if  $\kappa_A$  is a twistor, then the resulting  $\xi_{A'}$  solves (3). Given a twistor candidate and its corresponding  $\xi_{A'}$ , we choose to propagate the latter according to this wave equation.

## 2.2 For the zero quantities

In order to derive wave equations for the zero quantities, we assume that twistor candidate,  $\kappa_A$  and its auxiliary spinor  $\xi_{A'}$  satisfy the wave equations (2)–(3); at the end of the day, these will be satisfied by construction, since the candidate quantities will be propagated off the initial hypersurface using these wave equations.

To get a wave equation for  $H_{A'AB}$ , simply take the definition of  $\mathbf{S}$  in terms of  $\mathbf{H}$  and take a contracted derivative —i.e. consider  $\nabla_A{}^{D'} S_{D'A'B}$ . Ultimately, we get

$$\square H_{A'AB} = 8\Lambda H_{A'AB} - 2\Psi_{ABCD} H_{A'}{}^{CD} - 2\Phi_{ADA'D'} H^{D'D}{}_B - 2\nabla_{AD'} S^{D'}{}_{A'B}.$$

It is important to note that this is expressible (in a regular way) in terms of the rescaled Weyl spinor  $\phi_{ABCD} = \Theta^{-1} \Psi_{ABCD}$ :

$$\square H_{A'AB} = 8\Lambda H_{A'AB} - 2\Theta \phi_{ABCD} H_{A'}{}^{CD} - 2\Phi_{ADA'D'} H^{D'D}{}_B - 2\nabla_{AD'} S^{D'}{}_{A'B}.$$

To close the system, we need a wave equation for  $S_{A'B'A}$ . To get this, we will apply the D'Alembertian to (1), commute derivatives, and substitute the wave equation for  $\xi_{A'}$ , equation (3). Finally, we arrive at<sup>•2</sup>

$$\begin{aligned} \square S_{A'B'A} = & 6\Lambda S_{A'B'A} - 4\Phi_{ABC'(A'} S_{B')}{}^{C'B} - 2\Theta \bar{\phi}_{A'B'C'D'} S^{C'D'}{}_A \\ & - \frac{2}{3} \Phi_{BCA'B'} (\nabla_{AC'} H^{C'BC} + 2\nabla^C{}_{C'} H^{C'}{}_A{}^B) \\ & + 4H_{(A'|AB|} \nabla^B{}_{B'} \Lambda - 2(\nabla_{CC'} \Phi_{ABA'B'}) H^{C'BC}. \end{aligned} \quad (4)$$

•2: Is it possible to simplify this a bit more?

Note that the terms on the right-hand-side are homogeneous in  $\mathbf{S}$ ,  $\mathbf{H}$  and  $\nabla \mathbf{H}$ .

## 3 What goes wrong in the higher-valence case?

Define also the “Buchdahl zero quantity”:

$$B_{ABCD} = \phi_{F(ABC} \kappa_{D)}{}^F.$$

Note that

$$\nabla_{(A}{}^{A'} H_{|BCD)} = 6\Theta B_{ABCD}$$

Can derive equations of the form

$$\square H_{A'ABC} = (\mathbf{H}, \nabla \mathbf{S}), \quad (5)$$

$$\square H_{A'ABC} = (\mathbf{H}, \mathbf{B}, \nabla \mathbf{B}), \quad (6)$$

$$\square S_{AA'BB'} = (\mathbf{H}, \mathbf{S}, \Theta \mathbf{B}, \nabla \Theta \cdot \nabla \mathbf{B}) = (\mathbf{H}, \mathbf{S}, \nabla \mathbf{H}, \nabla \Theta \cdot \nabla \mathbf{B}), \quad (7)$$

$$\square B_{ABCD} = (\mathbf{H}, \mathbf{B}) + \frac{2}{3} \nabla_{\xi} \phi_{ABCD} = (\mathbf{H}, \mathbf{B}) + \frac{2}{3} \mathcal{L}_{\xi} \phi_{ABCD}. \quad (8)$$

The final equality follows from the fact that  $\nabla_{(A}{}^{A'} \xi_{B)A'} = 0$ , as a consequence of the assumed wave equation for  $\kappa_{AB}$ .<sup>•3</sup>

•3: Generally speaking, do we need to worry about the fact that  $\kappa_{AB}, \xi_{AA'}$  are being propagated independently (though consistently)?

**Note:** In the equation for  $\mathcal{S}$ , we couldn't replace the  $\nabla \mathbf{B}$  with  $\nabla \nabla \mathbf{H}$  terms even if we wanted to, because then we would have  $\Theta^{-1}$  factors appearing.

**Proposal:** Define  $F_{A'BCD} := \nabla^A_{A'} B_{ABCD}$ . Then we get a wave equation for  $B$  trivially:

$$\square B_{ABCD} \propto \nabla_{(A}{}^{A'} F_{|A'|BCD)} + \text{curv.} \times \mathbf{B}.$$

To get the remaining equation for  $F_{A'BCD}$  take a contracted derivative of (8), commute derivatives on the  $\nabla_{\xi} \phi$  term and use the fact that  $\nabla^A_{A'} \phi_{ABCD} = 0$ .

## 4 A closed system for the Killing spinor case

Zero quantities:

$$H_{A'ABC}, \quad B_{ABCD}, \quad F_{A'BCD} := \nabla^A_{A'} B_{ABCD}$$

From the previous section,  $\nabla_{(A}{}^{A'} H_{|A'|BCD)} = 6\Theta B_{ABCD}$ . Additionally, the wave equation for  $\kappa_{AB}$  is equivalent to  $\nabla^{AA'} H_{A'ABC} = 0$ , hence we have  $\nabla_A{}^{A'} H_{A'BCD} = 6\Theta B_{ABCD}$ . Contracting with  $\nabla^A_{B'}$  we then derive the following wave equation:

$$\square H_{A'ABC} = \dots \quad (9)$$

Similarly, substituting the definition of  $F_{A'ABC}$  in terms of  $B_{ABCD}$ , it is straightforward to verify the following wave equation for  $B_{ABCD}$ :

$$\square B_{ABCD} = 12\Lambda B_{ABCD} - 6\Theta \phi_{(AB}{}^{FG} B_{CD)FG} + 2\nabla_{AA'} F^{A'}{}_{BCD}. \quad (10)$$

The task remaining is to derive a wave equation for  $F_{A'ABC}$ . Let us first consider some useful identities:

$$\begin{aligned} \kappa^{DG} \phi_{(ABC}{}^H \phi_{F)HDG} &= \kappa_A{}^D \phi_{(BC}{}^{GH} \phi_{FD)GH} + \kappa_B{}^D \phi_{(AC}{}^{GH} \phi_{FD)GH} \\ &\quad + \kappa_C{}^D \phi_{(AB}{}^{GH} \phi_{FD)GH} + \kappa_F{}^D \phi_{(BC}{}^{GH} \phi_{AD)GH} \\ &= 2\phi_{(AB}{}^{GH} B_{CF)GH} \end{aligned} \quad (11)$$

Using this identity, we can derive the following alternative (non-homogeneous) wave equation for  $B_{ABCD}$ :

$$\square B_{ABCD} = \frac{2}{3} \nabla_{\xi} \phi_{ABCD} + 8\Lambda B_{ABCD} - 14\Theta \phi_{(AB}{}^{FG} B_{CD)FG} + \frac{2}{3} (\nabla_{FA'} \phi_{G(ABC)} H^{A'}{}_D)^{FG} \quad (12)$$

where, for convenience, we have defined  $\xi_{AA'} := \nabla^B_{A'} \kappa_{AB}$ , as usual. Contrary to previous approaches, however, we won't propagate  $\xi_{AA'}$  independently of  $\kappa_{AB}$ ; it is simply a convenient shorthand.

**Remark 1.** *As an aside, note that (12) expresses the fact that, given a Killing spinor  $\kappa_{AB}$ , the quantity*

$$\mathcal{L}_{\xi} \phi_{ABCD} \equiv \nabla_{\xi} \phi_{ABCD} + \phi_{F(ABC} \nabla_{D)A'} \xi^{FA'} = \nabla_{\xi} \phi_{ABCD}$$

*—the equality being by virtue of the wave equation for  $\kappa_{AB}$ — vanishes, and in particular that in the physical spacetime the Lie derivative of the Weyl curvature along the Killing vector  $\tilde{\xi}$  vanishes, as we know should be the case.*

At this point we note that there are no derivatives of zero quantities appearing on the right-hand-side of (12). This, combined with the fact that  $\nabla^A_{A'} \phi_{ABCD} = 0$ , makes it seem plausible that by applying  $\nabla^A_{A'}$  to (12) we will be able to derive a wave equation for  $F_{A'BCD}$ . We will see that this does indeed work; the only essential difficulty is in showing that  $\nabla^A_{A'} \nabla_{\xi} \phi_{ABCD}$  is expressible in terms of the zero quantities  $\mathbf{H}, \mathbf{B}, \mathbf{F}$  and their first derivatives only. This is done in the following identity:

$$\nabla^A_{A'} \nabla_{\xi} \phi_{ABCD} = \dots \quad (13)$$