

`ClearAll["Global`*"]`

The energies of the states  $|1\rangle$ ,  $|2\rangle$ , as well as the energy of the interaction are

$$E_1, E_2, E_{\text{int}} = \hbar \Omega \cos[\omega t - \phi];$$



The matrix with the ground state energies is

`MatrixForm[{{E1, 0}, {0, E2}}]`

$$\begin{pmatrix} E_1 & 0 \\ 0 & E_2 \end{pmatrix}$$

These are associated with the states  $c_1$  and  $c_2$  whose phases are  $\xi_1$  and  $\xi_2$ . We choose these phases in interaction picture so that

$$\xi_1 = E_1 t / \hbar$$

$$\xi_2 = E_1 t / \hbar + \omega t$$

$$\frac{t E_1}{\hbar}$$

$$\omega t + \frac{t E_1}{\hbar}$$

Where  $\omega$  is the frequency of the laser and  $\omega_0$  is the difference in frequency between the two energy levels that help us define the detuning  $\Delta$

$$\Delta = \hbar (\omega_0 - \omega)$$

$$\hbar (-\omega + \omega_0)$$

which give a hamiltonian such as

`H = MatrixForm[{{0,  $\Omega (1 + e^{2 i \omega t}) / 2$ }, { $\Omega (1 + e^{2 i \omega t}) / 2$ ,  $\Delta$ }}]`

$$\begin{pmatrix} 0 & \frac{1}{2} (1 + e^{2 i \omega t}) \Omega \\ \frac{1}{2} (1 + e^{2 i \omega t}) \Omega & \hbar (-\omega + \omega_0) \end{pmatrix}$$

In Rotating Wave Approximation RWA  $1 + e^{2 i \omega t} \simeq 1$  when averaged over time

`H = MatrixForm[{{0,  $\Omega / 2$ }, { $\Omega / 2$ ,  $\Delta$ }}]`

$$\begin{pmatrix} 0 & \frac{\Omega}{2} \\ \frac{\Omega}{2} & \hbar (-\omega + \omega_0) \end{pmatrix}$$

So the Schrödinger equation  $c_i' = H * c_i$  gives the following system of equations where we suppose the laser is in resonance with the difference in frequencies of the two levels ( $\Delta=0$ )

$$c_1'[t] == -i \Omega[t] * c_2[t] / 2;$$

$$c_2'[t] == -i \Omega[t] * c_1[t] / 2 - i \Delta c_2[t];$$

Value of the bandwidth( $\tau$ ) and its inverse at half the amplitude are

$$\tau = 5;$$

$$1 / \tau$$

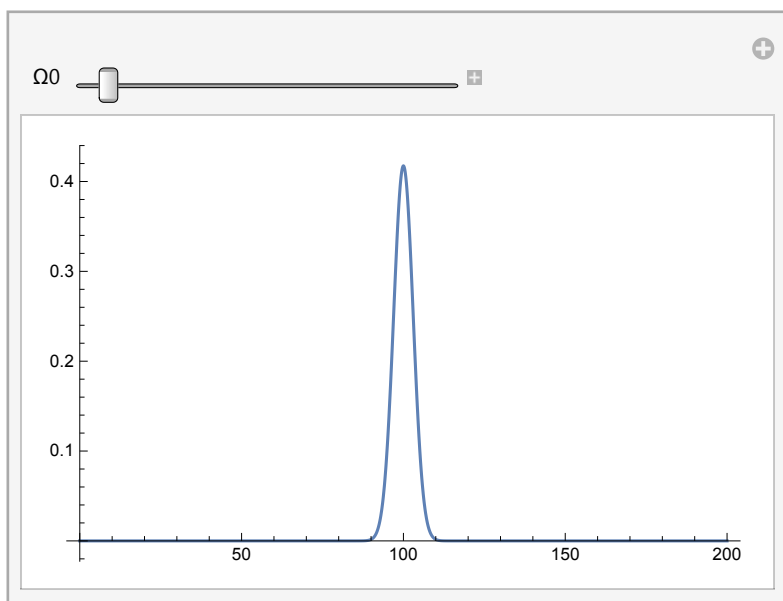
$$\frac{1}{5}$$

Now we plot these values to find the shape of the gaussian in the time span that goes between(all constatnts and variables are in nanoseconds):

```
ti = 0;
tf = 200;
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The frequency when detuned by one peak Rabi is (multiplied by  $\ln(2)$  for the BW we choose)

```
Manipulate[
  Plot[ $\Omega_0 * \text{Exp}[-2 \text{Log}[2] (t - (tf - ti) / 2)^2 / \tau^2]$ , {t, ti, tf}, PlotRange -> All],
  {{ $\Omega_0$ , Sqrt[ $\pi \text{Log}[4]$ ] /  $\tau$ }, 0, 10, 2 * Sqrt[ $\pi \text{Log}[4]$ ] /  $\tau$ }]
```



For operating purposes we write down the Rabi frequency relative to its peak

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 $\Omega_{rel}[t_] := \text{Exp}[-2 \text{Log}[2] (t - (tf - ti) / 2)^2 / \tau^2];$ 
```

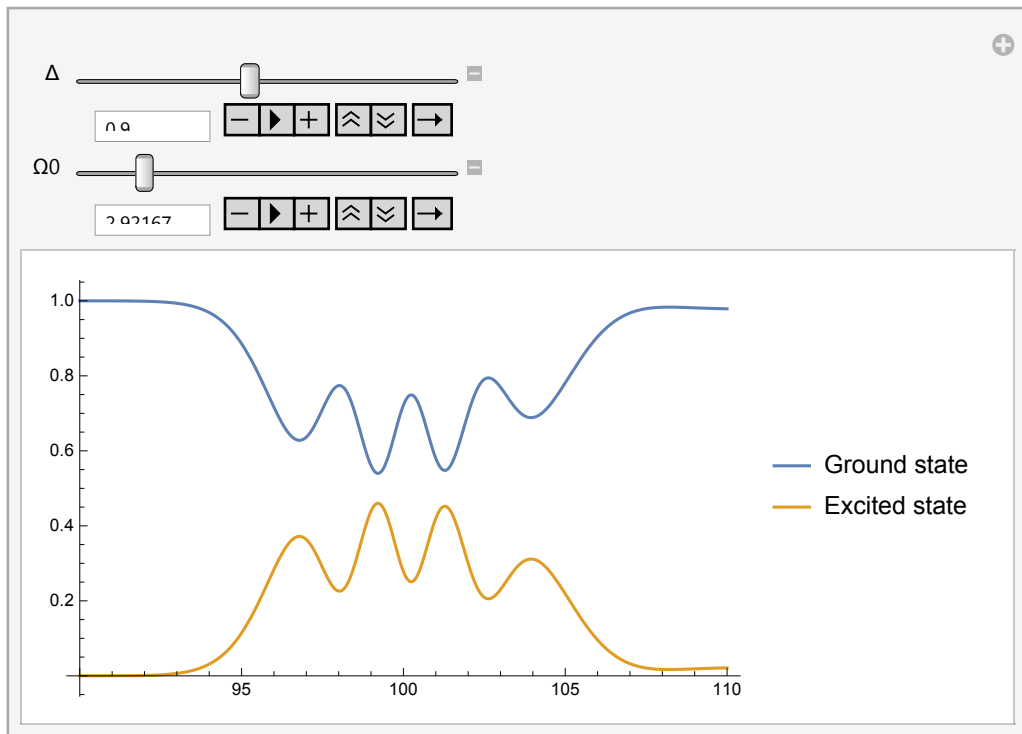
As well as the representation of the population for varying detuning  $\Delta$  and peak Rabi  $\Omega_0$

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Manipulate[Block[{c1Solved, c2Solved, c1, c2, t, s},

  s = NDSolve[{c1'[t] == -i Ω0 * Ωrel[t] * c2[t] / 2,
    c2'[t] == -i Ω0 * Ωrel[t] * c1[t] / 2 - i Δ c2[t],
    c1[ti] == 1, c2[ti] == 0}, {c1[t], c2[t]}, {t, ti, tf}];
  c1Solved = c1[t] /. s[[1, 1]];
  c2Solved = c2[t] /. s[[1, 2]];
  Plot[{c1Solved * Conjugate[c1Solved], c2Solved * Conjugate[c2Solved]}, {t, 90,
    110}, PlotRange → All, PlotLegends → {"Ground state", "Excited state"}]]
, {{Δ, 0}, 0, 2, 0.1}, {{Ω0, Sqrt[π Log[4]] / τ}, 0, 20, Sqrt[π Log[4]] / τ}]

```



The effective Rabi frequency is  $\Omega_{\text{eff}}$

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Manipulate[{Ω0^2 / Sqrt[Ω0^2 + Δ^2]}, {{Δ, 0}, 0, 100, 5}, {{Ω0, 1}, 0, 10, 0.5}]

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