

## Lecture 16

---

### Example 1. Explicit Solution to 1D flow

Consider a 1D reservoir with the following reservoir and fluid properties:

$$\phi = 0.2$$

$$k = 50 \text{ mD}$$

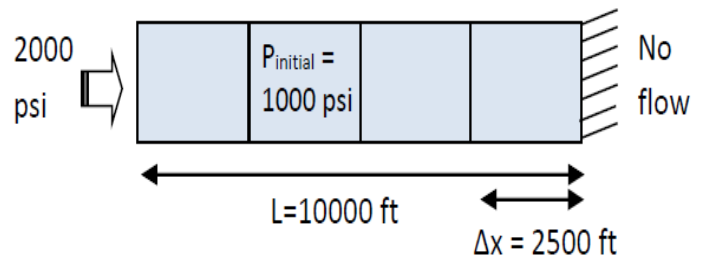
$$\mu = 1 \text{ cp}$$

$$B_w = 1 \text{ rb/stb (assume constant)}$$

$$c_t = 10^{-6} \text{ psi}^{-1}$$

$$L = 10,000 \text{ ft (Reservoir Length)}$$

$$A = 200,000 \text{ ft}^2 \text{ (Cross Sectional Area)}$$



The initial condition is  $P = 1,000 \text{ psi}$ . The boundary conditions are  $P = 2,000 \text{ psi}$  at  $x = 0$  and no flow ( $q = 0$ ) at  $x = L$ . Determine the pressure field in the reservoir using 4 uniform blocks. Use a time step of  $\Delta t = 1.0$  days.

Solution:

The pressure is governed by the 1D diffusivity equation with the following boundary conditions:

$$\frac{1}{\alpha} \frac{\partial p}{\partial t} = \frac{\partial^2 p}{\partial x^2}; \quad \alpha = \left( \frac{k}{\mu \phi c_t} \right)$$

$$IC : p(x, 0) = 1000 \text{ psi}$$

$$BC1 : p(0, t) = 2000 \text{ psi}$$

$$BC2 : \frac{\partial p}{\partial x}(L, t) = 0$$

The explicit finite difference solution is given by the formula:

$$P_i^{n+1} = P_i^n + \frac{\alpha \Delta t}{(\Delta x)^2} [P_{i-1}^n - 2P_i^n + P_{i+1}^n]$$

Where the dimensionless diffusivity is written as:

$$\eta = \alpha \frac{\Delta t}{(\Delta x)^2} = \left( \frac{k}{\mu \phi c_t} \right) \frac{\Delta t}{(\Delta x)^2} = \frac{50 \text{ mD}}{(1 \text{ cp})(0.2)(1.0 \times 10^{-6} \text{ psi}^{-1})} \frac{1.0 \text{ days}}{(2500 \text{ ft})^2} = 40 \frac{\text{mD-days-psi}}{\text{cp-ft}^2} = 0.2532$$

Note that we are guaranteed stability with the explicit method since  $\eta < 0.5$ .

At t = 0 days, the pressure is uniform at 1000 psi.  $P^0 = [1000 \ 1000 \ 1000 \ 1000] \text{ psi}$

At t = 1.0 days (n = 1), we solve explicitly for the pressures

$$\begin{aligned}P_1^1 &= P_1^0 + \eta \left[ P_0^0 - 2P_1^0 + P_2^0 \right] = 1000 + 0.2532 \left[ P_0^0 - 2(1000) + 1000 \right] = \\P_2^1 &= P_2^0 + \eta \left[ P_1^0 - 2P_2^0 + P_3^0 \right] = 1000 + 0.2532 \left[ 1000 - 2(1000) + 1000 \right] = \\P_3^1 &= P_3^0 + \eta \left[ P_2^0 - 2P_3^0 + P_4^0 \right] = 1000 + 0.2532 \left[ 1000 - 2(1000) + 1000 \right] = \\P_4^1 &= P_4^0 + \eta \left[ P_3^0 - 2P_4^0 + P_5^0 \right] = 1000 + 0.2532 \left[ 1000 - 2(1000) + P_5^0 \right] =\end{aligned}$$

There is no “block #0” or “block #5”. We must use our boundary conditions:

- 1) At  $x = 0$ , we have a constant pressure boundary  $P = P_{in} = 2000$  psi

$$P_{in} = \frac{P_0 + P_1}{2} \Rightarrow P_0 = 2P_{in} - P_1$$

- 2) At  $x = L$ , we have a “no-flow” condition

$$\frac{P_4 - P_5}{\Delta x} = 0 \Rightarrow P_5 = P_4$$

Substitute these values into the pressure equations:

$$P_1^1 = P_1^0 + \eta \left[ (2P_m - P_1^0) - 2P_1^0 + P_2^0 \right] = 1000 + 0.2532 \left[ (2 \cdot 2000 - 1000) - 2(1000) + 1000 \right] = 1506.4 \text{ psi}$$

$$P_2^1 = P_2^0 + \eta \left[ P_1^0 - 2P_2^0 + P_3^0 \right] = 1000 + 0.2532 \left[ 1000 - 2(1000) + 1000 \right] = 1000 \text{ psi}$$

$$P_3^1 = P_3^0 + \eta \left[ P_2^0 - 2P_3^0 + P_4^0 \right] = 1000 + 0.2532 \left[ 1000 - 2(1000) + 1000 \right] = 1000 \text{ psi}$$

$$P_4^1 = P_4^0 + \eta \left[ P_3^0 - 2P_4^0 + P_4^0 \right] = 1000 + 0.2532 \left[ 1000 - 2(1000) + 1000 \right] = 1000 \text{ psi}$$

At t = 2.0 days (n = 2)

$$P_1^2 = P_1^1 + \eta \left[ (2P_m - P_1^1) - 2P_1^1 + P_2^1 \right] = 1506.4 + 0.2532 \left[ (2 \cdot 2000 - 1506.4) - 2(1506.4) + 1000 \right] = 1628.1 \text{ psi}$$

$$P_2^2 = P_2^1 + \eta \left[ P_1^1 - 2P_2^1 + P_3^1 \right] = 1000 + 0.2532 \left[ 1506.4 - 2(1000) + 1000 \right] = 1128.2 \text{ psi}$$

$$P_3^2 = P_3^1 + \eta \left[ P_2^1 - 2P_3^1 + P_4^1 \right] = 1000 + 0.2532 \left[ 1000 - 2(1000) + 1000 \right] = 1000 \text{ psi}$$

$$P_4^2 = P_4^1 + \eta \left[ P_3^1 - 2P_4^1 + P_4^1 \right] = 1000 + 0.2532 \left[ 1000 - 2(1000) + 1000 \right] = 1000 \text{ psi}$$

At t = 3.0 days (n = 3)

$$P_1^3 = P_1^2 + \eta \left[ (2P_m - P_1^2) - 2P_1^2 + P_2^2 \right] = 1628.1 + 0.2532 \left[ (2 \cdot 2000 - 1628.1) - 2(1628.1) + 1128.2 \right] = 1689.9 \text{ psi}$$

$$P_2^3 = P_2^2 + \eta \left[ P_1^2 - 2P_2^2 + P_3^2 \right] = 1128.2 + 0.2532 \left[ 1628.1 - 2(1128.2) + 1000 \right] = 1223.3 \text{ psi}$$

$$P_3^3 = P_3^2 + \eta \left[ P_2^2 - 2P_3^2 + P_4^2 \right] = 1000 + 0.2532 \left[ 1128.2 - 2(1000) + 1000 \right] = 1032.5 \text{ psi}$$

$$P_4^3 = P_4^2 + \eta \left[ P_3^2 - 2P_4^2 + P_4^2 \right] = 1000 + 0.2532 \left[ 1000 - 2(1000) + 1000 \right] = 1000 \text{ psi}$$

At  $t = \infty$ , the reservoir comes to equilibrium at  $P = 2000$  psi in all blocks

Table 1.1 Summary of block pressures using the explicit method for example 1.

	<b>Block #1 (psi) 1250 ft</b>	<b>Block #2 (psi) 3750 ft</b>	<b>Block #3 (psi) 6250 ft</b>	<b>Block #4 (psi) 8750 ft</b>
0 days	1000	1000	1000	1000
1 days	1506.4	1000	1000	1000
2 days	1628.1	1128.2	1000	1000
3 days	1689.9	1222.3	1032.5	1000
$\infty$	2000	2000	2000	2000

The solution makes physical sense. The pressure in the reservoir increases over time as a result of the dirichlet (constant pressure) boundary condition of  $P = 2000$  psi, until the reservoir comes to steady state at 2000 psi. Pressure increases fastest in block #1 because it is near the boundary and slowest in block #4 because it is far from the boundary.

A plot of the numerical (explicit) solution to the problem is shown in figure 1.1a, 1.1b, and 1.1c using 4, 8, and 12 blocks, respectively. A plot of the actual (analytical) solution is given in figure 1.1d. Recall the analytical solution for this problem is given by:

$$p(x,t) = p_{B1} - \frac{4p_{mit}}{\pi} \sum_{n=0}^{\infty} \frac{(-1)^n}{2n+1} e^{-\frac{\alpha(2n+1)^2 \pi^2 t}{4L^2}} \cos \frac{(2n+1)\pi x}{2L}$$

The numerical solution demonstrates the same trends as the analytical solution, but is not as smooth because only a finite # of grid blocks are used. The numerical solution is particularly poor for only 4 grid blocks, but we could have increased accuracy by using more grids and smaller time steps as long as the stability criterion ( $\eta < 0.5$ ) was still met. Figure 1.1 demonstrates improved accuracy with more grids and smaller timesteps.

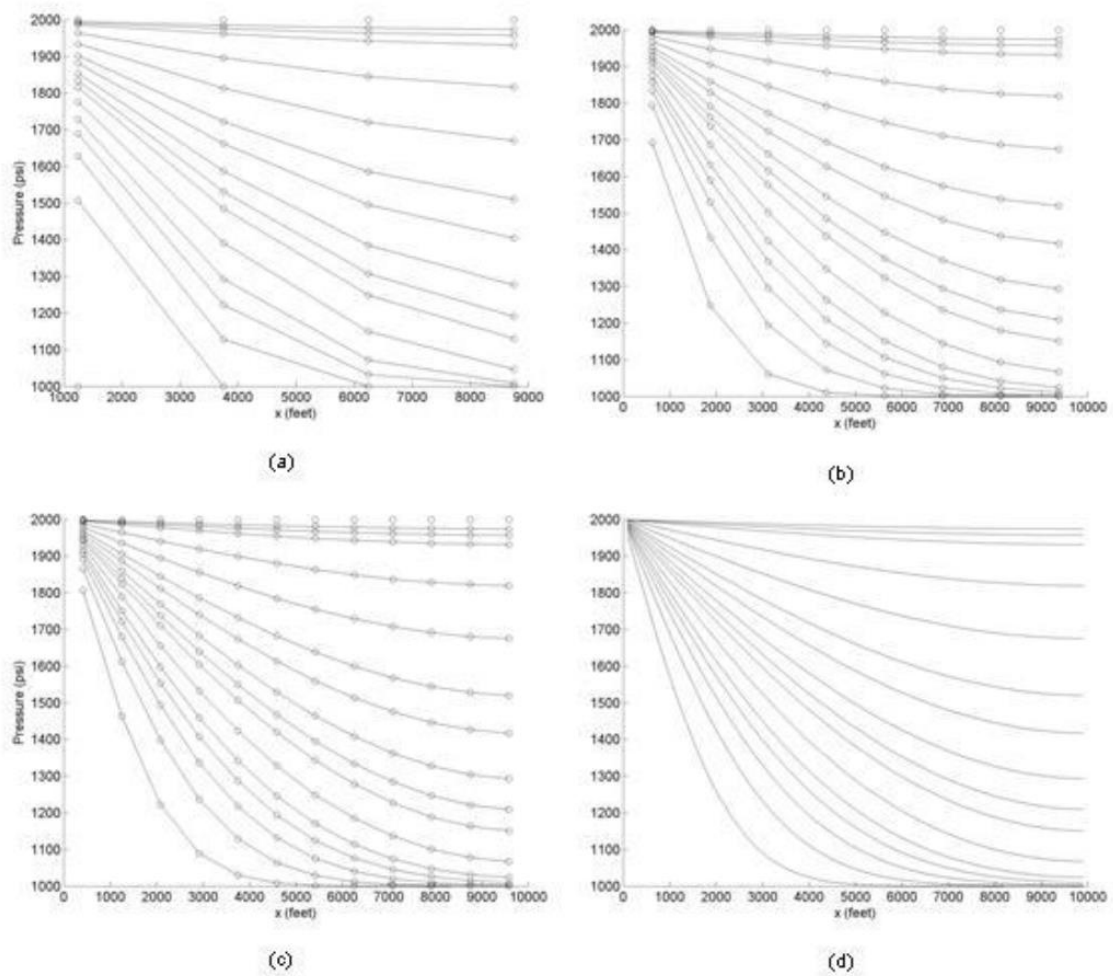


Figure 1.1. Solutions to pressure vs. reservoir distance at various times. (a) 4 grids ( $\Delta x = 2500$  feet),  $\Delta t = 1$  day; (b) 8 grids ( $\Delta x = 1250$  feet),  $\Delta t = 0.1$  day; (c) 12 grids ( $\Delta x = 833$  feet),  $\Delta t = 0.01$  day; (d) analytical solution. For the numerical solutions, the points are the grid pressures and the continuous lines are linearly interpolated values.