

Advanced Plasma - Plasma Physics for Fusion: Assignment - James Davies

09/11/2017

Question 1

Consider a current, I , being carried by a copper rod in the z -direction. This will produce a magnetic field

$$B = \frac{\mu_0 I}{2\pi R} \hat{\phi}$$

in the toroidal direction. This system is shown in figure 1:

Info for figure 1: a.) A current in the \hat{z} direction at $R=0$ leads to a toroidal field. b.) Illustration of the toroidal field strength dependence on R

a.) Suppose 1MA of current flows through the central rod, find the magnetic field strength at $R=1m$ and use this to calculate the Larmor radius, r_L , and the gyro-frequency, Ω , for an electron with temperature 10keV and an alpha with temperature 3.5MeV. You may assume $V_L = V_{th} = \sqrt{2k_B T_e / m}$ and $m_e = 1/m_p$, where m_p is the proton mass.

$$B = \frac{1.257 \times 10^{-6} \times 1 \times 10^6}{2\pi \times 1} = \underline{\underline{0.2 T}}$$

$$\Omega = \frac{ZeB}{m}$$

$$\text{For electron: } \Omega = \frac{1.6 \times 10^{-19} \times 0.2}{9.11 \times 10^{-31}} = 3.51 \times 10^{10} \text{ Hz} = \underline{\underline{35.1 GHz}}$$

$$\text{For } \alpha\text{-particle } \Omega = \frac{2 \times 1.6 \times 10^{-19} \times 0.2}{4 \times 1.67 \times 10^{-27}} = 9.58 \times 10^6 \text{ Hz} = \underline{\underline{9.58 MHz}}$$

$$V_L = V_{th} = \sqrt{2k_B T_e / m}$$

$$T(\text{keV}) = 8.62 \times 10^{-8} T(\text{K})$$

for electron:

$$V_L = \sqrt{2(1.38 \times 10^{-23}) \left(\frac{10}{8.62 \times 10^{-8}} \right) / 9.11 \times 10^{-31}} = 59.3 \times 10^6 \text{ ms}^{-1}$$

for α -particle $V_L = \sqrt{2(1.38 \times 10^{-23}) \left(\frac{3500}{8.62 \times 10^{-8}} \right) / (4 \times 1.67 \times 10^{-27})} = 13.0 \times 10^6 \text{ ms}^{-1}$

$$r_L = \frac{V_L}{\Omega}$$

for electron:

$$r_L = \frac{59.3 \times 10^6}{3.51 \times 10^{10}} = \underline{\underline{1.69 \times 10^{-3} \text{ m}}}$$

for ~~electron~~ α -particle:

$$r_L = \frac{13.0 \times 10^6}{9.58 \times 10^6} = \underline{\underline{1.36 \text{ m}}}$$

b.) In this system the ∇B and curvature drift velocities may be combined and written as:

$$\underline{V_D} = \frac{m}{qZeB} \left(\underline{V_H}^2 + \frac{V^2}{2} \right) \underline{e_z}$$

This combined drift, V_D , is often referred to as simply the magnetic drift.

i.) Assuming $v_{\parallel}^2 + v_{\perp}^2/2 = v_{th}^2$, calculate the drift velocity for electrons and Deuterium ions ($m_D = 2m_p$) at $R=1m$ with 10keV thermal energy.

$$V_D = \frac{m}{RZeB} \left(v_{th}^2 \right)$$

for electrons:

$$V_D = \frac{9.11 \times 10^{-31}}{1 \times 1.6 \times 10^{-19} \times 0.2} \left((59.3 \times 10^6)^2 \right) = \underline{\underline{1.00 \times 10^5 \text{ ms}^{-1}}}$$

for ions:

$$v_{th}^2 = \left(2 (1.38 \times 10^{-23}) \left(\frac{10}{8.62 \times 10^{-8}} \right) \right) / (2 \times 1.67 \times 10^{-27})$$

$$= 9.59 \times 10^{11}$$

$$V_D = \frac{2 \times 1.67 \times 10^{-27}}{1 \times 2 \times 1.6 \times 10^{-19} \times 0.2} \left(9.59 \times 10^{11} \right) = \underline{\underline{5.00 \times 10^4 \text{ ms}^{-1}}}$$

ii.) In a terrestrial system, such as the one under consideration here, there will be a gravitational force, mg , acting on the plasma species in the $-e_z$ direction. This force will lead to a time independent particles drift. Starting from the equation of motion:

$$m \frac{dv}{dt} = -mg \underline{e_z} + Ze(v \times B)$$

derive an expression for the perpendicular velocity of the drift due to the gravitational force, $V_{g\perp}$. You may find the following vector identity useful:

$$(A \times B) \times C = (A \cdot C)B - (B \cdot C)A$$

$$m \frac{d\mathbf{v}}{dt} = -mg\mathbf{e}_z + Ze(\mathbf{v} \times \mathbf{B})$$

Split \mathbf{v} into $\mathbf{v}_{\text{gyro}} + \mathbf{v}_g$:

$$\mathbf{v} = \mathbf{v}_{\text{gyro}} + \mathbf{v}_g$$

Looking at \mathbf{v}_{gyro} :

$$m \frac{d\mathbf{v}_{\text{gyro}}}{dt} = Ze \mathbf{v}_{\text{gyro}} \times \mathbf{B}$$

Therefore:

$$m \frac{d(\mathbf{v}_{\text{gyro}} + \mathbf{v}_g)}{dt} = -mg\mathbf{e}_z + Ze \mathbf{v}_{\text{gyro}} \times \mathbf{B} + Ze \mathbf{v}_g \times \mathbf{B}$$

$$m \frac{d\mathbf{v}_g}{dt} = -mg\mathbf{e}_z + Ze \mathbf{v}_g \times \mathbf{B}$$

Since acceleration due to gravity is constant: $\frac{d\mathbf{v}_g}{dt} = 0$

$$\therefore mg\mathbf{e}_z = Ze \mathbf{v}_g \times \mathbf{B}$$

Crossing both sides with \mathbf{B}

$$mg\mathbf{e}_z \times \mathbf{B} = Ze (\mathbf{v}_g \times \mathbf{B}) \times \mathbf{B}$$

Using $(\mathbf{A} \times \mathbf{B}) \times \mathbf{C} = (\mathbf{A} \cdot \mathbf{C})\mathbf{B} - (\mathbf{B} \cdot \mathbf{C})\mathbf{A}$ on $(\mathbf{v}_g \times \mathbf{B}) \times \mathbf{B}$:

$$(\mathbf{v}_g \times \mathbf{B}) \times \mathbf{B} = (\mathbf{v}_g \cdot \mathbf{B})\mathbf{B} - (\mathbf{B} \cdot \mathbf{B})\mathbf{v}_g = B^2 \left(\frac{(\mathbf{v}_g \cdot \mathbf{B})\mathbf{B}}{B^2} - \mathbf{v}_g \right)$$

If \mathbf{v}_g is perpendicular to \mathbf{B} then $(\mathbf{v}_g \cdot \mathbf{B}) = 0$

$$\therefore \mathbf{v}_g = \frac{-mg\mathbf{e}_z \times \mathbf{B}}{Ze B^2}$$

iii.) Calculate the magnitude of the gravitational drift velocity for electrons and deuterium ions on the surface of Earth ($g = 9.81 \text{ m s}^{-2}$). Compare these values to the magnetic drift velocity in part i.) and discuss the importance of gravitational drift.

$$V_g = \frac{-mgB}{ZeB^2} \quad \text{in } Z \text{ direction}$$

$$= - \frac{9.11 \times 10^{-31} \times 9.81 \times 0.2}{1.6 \times 10^{-19} (0.2)^2} = -2.79 \times 10^{-10} \text{ m s}^{-1} \quad \text{for electrons}$$

$$= - \frac{2 \times 1.67 \times 10^{-27} \times 9.81 \times 0.2}{2 \times 1.6 \times 10^{-19} \times (0.2)^2} = -5.12 \times 10^{-7} \text{ m s}^{-1} \quad \text{for ions}$$



Question 2

The system in question 1 uses a purely toroidal magnetic field. In this question we will consider introducing a poloidal component of the magnetic field, which can be achieved by driving a toroidal current, I_T , at $R=1m$. The result of this is that the field lines become helical, orbiting about the location of the current as shown in Fig 2.

As particles travel along the now helical field lines, they also move in R and Z . To describe this it is convenient for us to introduce the poloidal angle, θ , as shown in Fig 2. As the particle moves along the field line it will move in both the toroidal and poloidal angles.

a) Describe why the magnetic drift means a purely toroidal field, such as that in question 1, cannot confine a plasma indefinitely.

A toroidal magnetic field like the one in question 1 would not ~~confine~~ confine a plasma due to ∇B drift. As B is proportional to $1/R$, there is a grad B inwards, this generates a drift perpendicular to B . This results in a vertical charge separation of ions + electrons producing an E field. This then produces an $E \times B$ drift to a large radius, resulting in a loss of confinement!

b) The magnitude of the toroidal component of the magnetic field, B_T , depends on R and as such the total field magnitude, B , will vary along the field line. Let us assume that this variation is described by $B = B_0(1 - a \cos \theta)$, where B_0 and a are some constants and θ is the poloidal angle. Starting from the conservation of energy and ~~magnetic~~ magnetic moment show that a particle at $\theta = 0$ with pitch angle λ (defined through $\sin \lambda = v_{\perp} / v \neq 0$) will be reflected at a poloidal angle θ satisfying

$$\cos \theta = \frac{1}{a} \left[1 - \frac{1-a}{\sin^2 \lambda} \right]$$

Electrostatic potential ϕ may be neglected!

Looking at the conserved quantities:

$$E = \frac{1}{2}mv^2 + Ze\phi \quad ; \quad \mu = \frac{mV_{\perp}^2}{2B}$$

Since these are conserved and neglecting ϕ

$$\mu = \frac{mV_{\perp,0}^2}{2B_0} = \frac{mV_{\perp}^2}{2B}$$

Then we have: $V_{\perp}^2 = B V_{\perp,0}^2 / B_0$ ①

Also:

$$E/m = V_0^2 = V_{\parallel}^2 + V_{\perp,0}^2 = V_{\parallel}^2 + V_{\perp}^2$$

$$\Rightarrow V_{\parallel}^2 = V_0^2 - V_{\perp}^2$$

Subbing in ①

$$V_{\parallel}^2 = V_0^2 - B V_{\perp,0}^2 / B_0$$

~~Observation~~ Defining pitch angle $\sin \tau = \frac{V_{\perp,0}}{V_0}$

$$\Rightarrow V_{\parallel}^2 = V_0^2 \left[1 - \frac{B}{B_0} \sin^2 \tau \right]$$

So particle to be reflected with θ : $B = \frac{B_0}{\sin^2 \tau} \Rightarrow \sin \tau = \sqrt{B_0/B}$

Variation: $B = B_c (1 - a \cos \theta)$; $B_0 (\theta=0) = B_c (1 - a)$

$$\Rightarrow \sin \tau = \sqrt{(1-a)/(1-a \cos \theta)}$$

Please turn over

$$\sin^2 \lambda = \frac{1-a}{1-a \cos \theta}$$

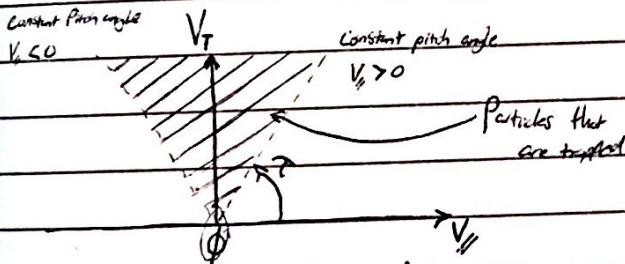
$$1-a \cos \theta = \frac{1-a}{\sin^2 \lambda}$$

$$\Rightarrow \cos \theta = \frac{1}{a} \left[1 - \frac{1-a}{\sin^2 \lambda} \right] \quad \text{As expected}$$

(C.)

A particle is considered trapped if it is reflected at some point along the field line (i.e. somewhere between $\theta=0$ and $\theta=\pi$). Assuming a Maxwellian velocity distribution, use the result of 2b.) with $c \ll 1$ to show that the fraction fraction of particles which are trapped, f_t is given by:

$$f_t = \sqrt{2a}$$



$$N(v) = A \int_0^{2\pi} d\phi \int_0^{\pi/2} \sin \theta d\theta \int_0^{\infty} v^2 \exp(-v^2/v_{th}^2) dv$$

$$= 4\pi A [-\cos \theta]_0^{\pi/2} \int_0^{\infty} v^2 \exp(-v^2/v_{th}^2) dv \quad \text{use } \int_0^{\infty} s^2 \exp(-s^2) ds$$

$$= +4\pi A \cos \pi v_{th} \left(\frac{\sqrt{\pi}}{4} \right) \quad \text{particles trapped}$$

$$N(v) = A \int_0^{2\pi} d\phi \int_0^{\pi} \sin \theta d\theta \int_0^{\infty} v^2 \exp(-v^2/v_{th}^2) dv$$

$$= 2\pi A [-\cos \theta]_0^{\pi} \int_0^{\infty} v^2 \exp(-v^2/v_{th}^2) dv$$

$$\Rightarrow 4\pi A v_{th} \left(\frac{\sqrt{\pi}}{4} \right) \quad \Rightarrow \text{fraction} = \cos \lambda$$

for all particles.

Looking at the critical angle λ :

$$\sin \lambda = \sqrt{\frac{1-a}{1-a \cos \theta}}$$

At the forward limit, critical angle $= \pi$

$$\sin \lambda = \sqrt{\frac{1-a}{1+a}} \approx \sqrt{1-2a} \quad \text{for } a \ll 1$$

$$\therefore \lambda = \sin^{-1}(\sqrt{1-2a})$$

Subbing into fraction \Rightarrow

$$f_{\pm} = \cos(\sin^{-1}(\sqrt{1-2a}))$$

Using trig identity: $\cos(\sin^{-1}(x)) = \sqrt{1-x^2}$

$$f_{\pm} = \sqrt{1-(1-2a)}$$

$$\underline{f_{\pm} = \sqrt{2a}} \quad \text{as expected!}$$