## Fluid dynamics Assignment 01

### N'Dah Jean KOUAGOU

January 14, 2019

### 1. Sketching contours

(i) 
$$f(x,y) = x^2 + y^2$$

We draw the curves  $f(x,y) = constant \in \mathbb{R}_0^+$ .

If constant = 0 then we get the origin O(0,0).

If constant > 0 then we have a circle.

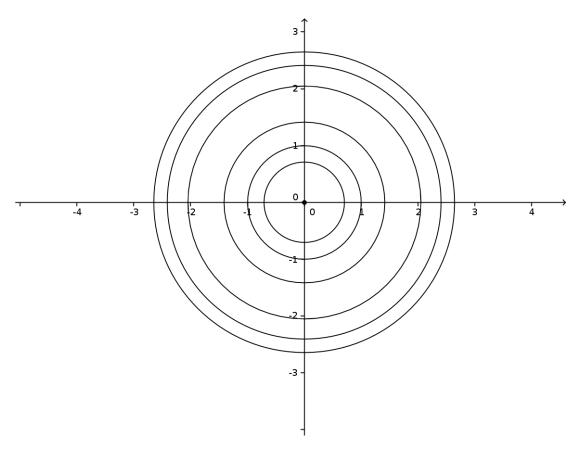


Figure 1: Contours of  $f(x, y) = x^2 + y^2$ 

(ii) 
$$f(x,y) = x^2 - y^2$$

Here again we draw the curves f(x,y) = constant

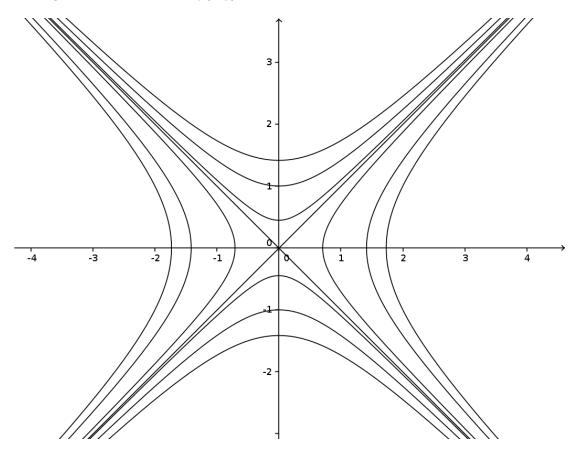


Figure 2: Contours of  $f(x, y) = x^2 - y^2$ 

(iii) 
$$f(x,y) = xy$$

Here again we draw the curves f(x,y) = constant

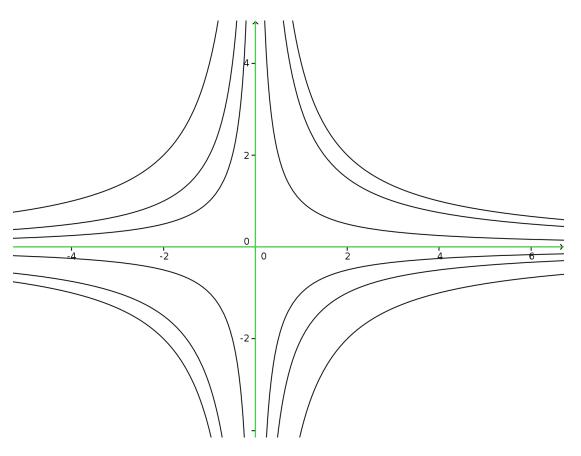


Figure 3: Contours of f(x, y) = xy

# 2. Gradient and directional derivative of $f(x,y) = 3x^2 - 2xy + y^3$

$$f(x,y) = 3x^2 - 2xy + y^3$$

•  $\nabla f(x,y)$ We have:

$$\nabla f(x,y) = \begin{pmatrix} \frac{\partial f}{\partial x}(x,y) \\ \frac{\partial f}{\partial y}(x,y) \end{pmatrix}$$
$$= \begin{pmatrix} 6x - 2y \\ -2x + 3y^2 \end{pmatrix}$$
$$\nabla f(x,y) = \begin{pmatrix} 6x - 2y \\ -2x + 3y^2 \end{pmatrix}$$

• Directional derivative of f in the direction of the vector  $u = \begin{pmatrix} 3 \\ 4 \end{pmatrix}$  at the point x = (1,1). We denote by  $D_u f(1,1)$  that derivative.

We have

$$D_{u}f(1,1) = \lim_{h \to 0} \frac{f((1,1) + hu) - f(x,y)}{h|u|}$$

$$= \nabla f(1,1) \cdot \frac{u}{|u|} \text{ because } f \text{ is differentiable at } (1,1)$$

$$= \frac{1}{5} \left(4,1\right) \cdot (3,4)$$

$$= \frac{1}{5} (4 \times 3 + 1 \times 4)$$

$$= \frac{1}{5} \times 16$$

$$= \frac{16}{5}$$

$$D_{u}f(1,1) = \frac{16}{5}$$

### 3. Vector velocity field sketch

 $\bullet \ u = (x, y)$ 

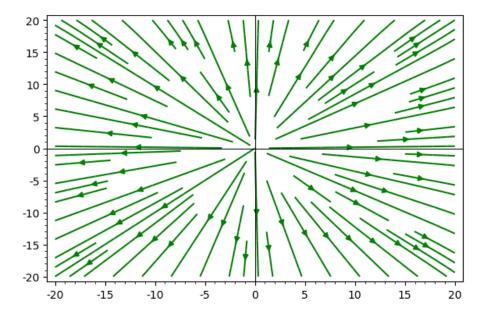


Figure 4: velocity field for u = (x, y)

 $\bullet \ u = (x, -y)$ 

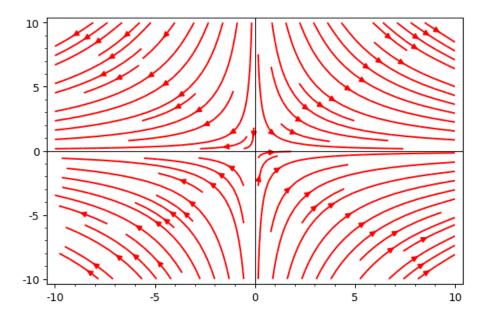


Figure 5: velocity field for u=(x,-y)

 $\bullet \ u = (-y, x)$ 

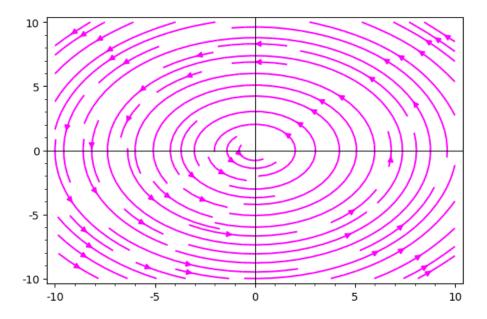


Figure 6: velocity field for u = (-y, x)

#### 4.

## 4.a) Sketch of the streamline passing through the point (1,1) at t=0.

The streamline is a curve instantaneously tangent to the velocity vector of the flow.

Let S(t) = (x(t), y(t)) be a streamline.

We have  $det\left(S'(t),U(t)\right)=0$  where  $U(t)=\left(\frac{1}{t+1},1\right)$  is the velocity.

$$det\left(S'(t), U(t)\right) = 0$$

$$\Leftrightarrow det\left(\begin{matrix} x' & \frac{1}{t+1} \\ y' & 1 \end{matrix}\right) = 0$$

$$\Leftrightarrow x' - \frac{1}{t+1}y' = 0$$

$$\Leftrightarrow \frac{dx}{dt} - \frac{1}{t+1}\frac{dy}{dt} = 0$$

$$\Leftrightarrow (t+1)dx - dy = 0 \text{ multiplying by } (1+t)dt$$

$$\Leftrightarrow \frac{dy}{dx} = 1 + t$$

$$\Leftrightarrow y = (1+t)x + c, \ c \in \mathbb{R}$$

Since at t = 0, x = y = 1, then we have  $1 = 1 + c \Rightarrow c = 0$ .

Now we choose t = 0 to find the streamline passing through the point (1,1) at t = 0 which is the straight line of equation y = x.

The sketch of that streamline is below:

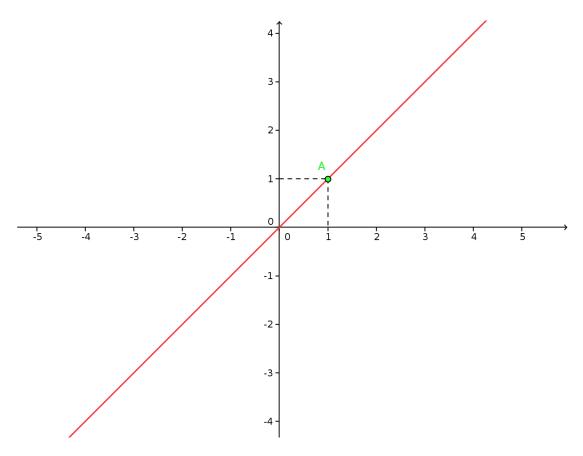


Figure 7: Streamline at t=1

# 0.1 4.b) The path of a fluid particle that is released from (1,1) at t=0.

Let  $X(t) = \left(x(t), y(t)\right)$  be the instantaneous position of the particle. We have:

$$\begin{cases} \frac{dx}{dt}(t) = \frac{1}{1+t} \\ \frac{dy}{dt}(t) = 1 \end{cases} \Rightarrow \begin{cases} x(t) = \ln(1+t) + c_1 \\ y(t) = t + c_2 \end{cases}$$
  
At  $t = 0, x = y = 1$ .

So  $c_1 = 1, c_2 = 1$  and we have

$$\begin{cases} x(t) = 1 + ln(1+t) \\ y(t) = 1 + t \end{cases}$$
 Then  $x(t) = 1 + ln(y(t))$ , that is  $y(t) = \exp(x(t) - 1)$ 

It follows that the path equation is  $y = \exp(x - 1)$  and we draw it below:

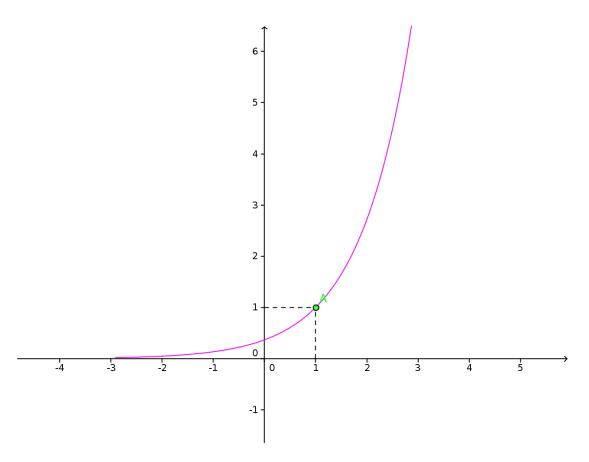


Figure 8: Path of a fluid particle

### The position at time t = 0 of a streak of dye:

For the pathline trajectory, we got

$$\begin{cases} x = \ln(1+t) + 1\\ y = t + 1 \end{cases}$$

Now consider: 
$$\begin{cases} x_{\tau} = \ln(1 + t_{\tau}) + 1 \\ y\tau = t_{\tau} + 1 \end{cases}$$

We have

$$\begin{cases} x_{\tau} - \ln(1 + t_{\tau}) = 1 \\ y_{\tau} - t_{\tau} = 1 \end{cases}$$

$$\Longrightarrow \begin{cases} x = x_{\tau} - \ln(1 + t_{\tau}) - 1 + \ln(1 + t) + 1 \\ y = y_{\tau} - t_{\tau} - 1 + t + 1 \end{cases}$$

$$\Longrightarrow \begin{cases} x = x_{\tau} - \ln(1 + t_{\tau}) + \ln(1 + t) \\ y = y_{\tau} - t_{\tau} + t \end{cases}$$

$$\Longrightarrow \begin{cases} x = x_{\tau} - \ln(1 + t_{\tau}) + \ln(1 + t) \\ t_{\tau} = y_{\tau} + t - y \end{cases}$$

$$\Longrightarrow \begin{cases} x = x_{\tau} - \ln(1 + t_{\tau}) + \ln(1 + t) \\ t_{\tau} = y_{\tau} + t - y \end{cases}$$

Then

$$x = x_{\tau} - \ln(1 + t_{\tau}) + \ln(1 + t) = x_{\tau} - \ln(1 + y_{\tau} + t - y) + \ln(1 + t)$$
 using (\*\*)  
Also  $-1 < t \le 0$  and at  $t = 0$ ,  $(x_{\tau}, y_{\tau}) = (1, 1)$ . As a result,

$$x = 1 - \ln(2 - y) \Longrightarrow 1 - x = \ln(2 - y) \Longrightarrow e^{1 - x} = 2 - y \Longrightarrow y = 2 - e^{1 - x}$$

We draw that curve below:

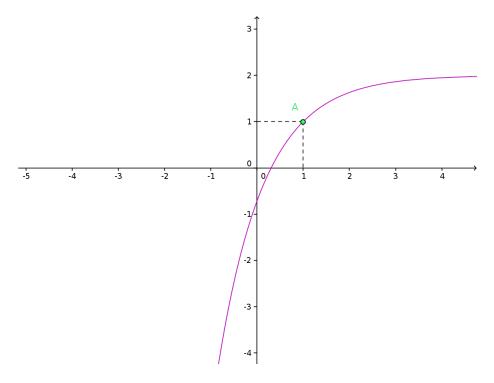


Figure 9: Position at t = 0 of a streak of dye released at the point (1,1)

# 5. Two dimensional flow represented by $\psi(x,y)$ with $u=\frac{\partial \psi}{\partial y}, v=-\frac{\partial \psi}{\partial x}$

Let's show that:

### **5.a)** The streamlines are given by $\psi = const$

Let  $S(t) = \left(x(t), y(t)\right)$  be a streamline. Like we did in 4.a), we have:

 $det \begin{pmatrix} \frac{dx}{dt} & u \\ \frac{y}{dt} & v \end{pmatrix} = 0 \Leftrightarrow det \begin{pmatrix} \frac{dx}{dt} & \frac{\partial \psi}{dy} \\ \frac{dy}{dt} & -\frac{\partial \psi}{\partial x} \end{pmatrix} = 0$  $\Leftrightarrow -\frac{dx}{dt} \frac{\partial \psi}{\partial x} - \frac{dy}{dt} \frac{\partial \psi}{\partial y} = 0$ 

$$\Leftrightarrow \frac{\partial \psi}{\partial x} dx + \frac{\partial \psi}{\partial y} dy = 0 \text{ multiplying by -dt}$$
 
$$\Leftrightarrow d\psi = 0$$

 $\Leftrightarrow \psi = const.$ 

So the streamlines are given by  $\psi = const.$ 

**5.b**) 
$$|{\bf u}| = |\nabla \psi|$$

We have

$$\begin{split} |\mathbf{u}| &= \sqrt{u^2 + v^2} \\ &= \sqrt{\left(\frac{\partial \psi}{\partial y}\right)^2 + \left(-\frac{\partial \psi}{\partial x}\right)^2} \\ &= \sqrt{\left(\frac{\partial \psi}{\partial y}\right)^2 + \left(\frac{\partial \psi}{\partial x}\right)^2} \end{split}$$

Also,

$$\nabla \psi = \begin{pmatrix} \frac{\partial \psi}{\partial x} \\ \frac{\partial \psi}{\partial y} \end{pmatrix}$$
  
So  $|\nabla \psi| = \sqrt{\left(\frac{\partial \psi}{\partial y}\right)^2 + \left(\frac{\partial \psi}{\partial x}\right)^2}$ 

Therefore  $|\nabla \psi| = |\mathbf{u}|$ 

#### 5.c)

Let q be the volume integral crossing a curve from  $x_0$  to  $x_1$ . We have  $q = \int_{x_0}^{x_1} \mathbf{u} \cdot \mathbf{n} ds$  with  $\mathbf{n} ds = \left(dy, -dx\right)$  and  $\mathbf{u} = \left(\frac{\partial \psi}{\partial y}, -\frac{\partial \psi}{\partial x}\right)$ Thus:

$$q = \int_{x_0}^{x_1} \frac{\partial \psi}{\partial x} dx + \frac{\partial \psi}{\partial y} dy$$
$$= \int_{x_0}^{x_1} d\psi$$
$$= \psi(x_1) - \psi(x_0)$$

We conclude that the volume integral crossing a curve from  $x_0$  to  $x_1$ **is:**  $\psi(x_1) - \psi(x_0)$ .

### 5.d) $\psi$ is constant along any fixed boundary

The flow can not cross a boundary and the velocity is O on the boundary. So

$$u = 0i$$
 and  $v = 0j \Rightarrow \begin{cases} \frac{\partial \psi}{\partial y} = 0\\ \frac{\partial \psi}{\partial x} = 0 \end{cases}$ 

 $\Rightarrow \psi$  does not depend on x and does not depend on y.

 $\Rightarrow \psi$  is constant on any fixed boundary.