

## Relativity lecture 4.

• Relativistic doppler effect

e → emitter  
r → receiver

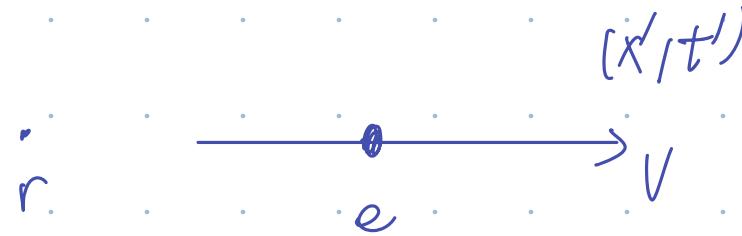
$$f_r = \frac{c - v_r}{c + v_e} f_e$$

↑  
speed of sound

there is a  
preferential.

$$t = \gamma(t' + vx)$$

$$x = \gamma(x' + vt)$$



2 pings  $p_1'(l, t)$

$p_2'(l, t + T_e)$

$$p_1 = (\gamma(l + vt'), \gamma(t' + vl))$$

$$p_2 = (\gamma(l + v(t' + T_e)), \gamma(t' + T_e + vl))$$

time it takes for light to travel,

$p_1$  received at

$$\gamma(t + vl) + \gamma(l + vt)/c$$

$c \approx$

$$p_2 - \gamma(t + T_e + vl) + \gamma(l + v(t + T_e))/c$$

$$T_r = p_2 - p_1$$

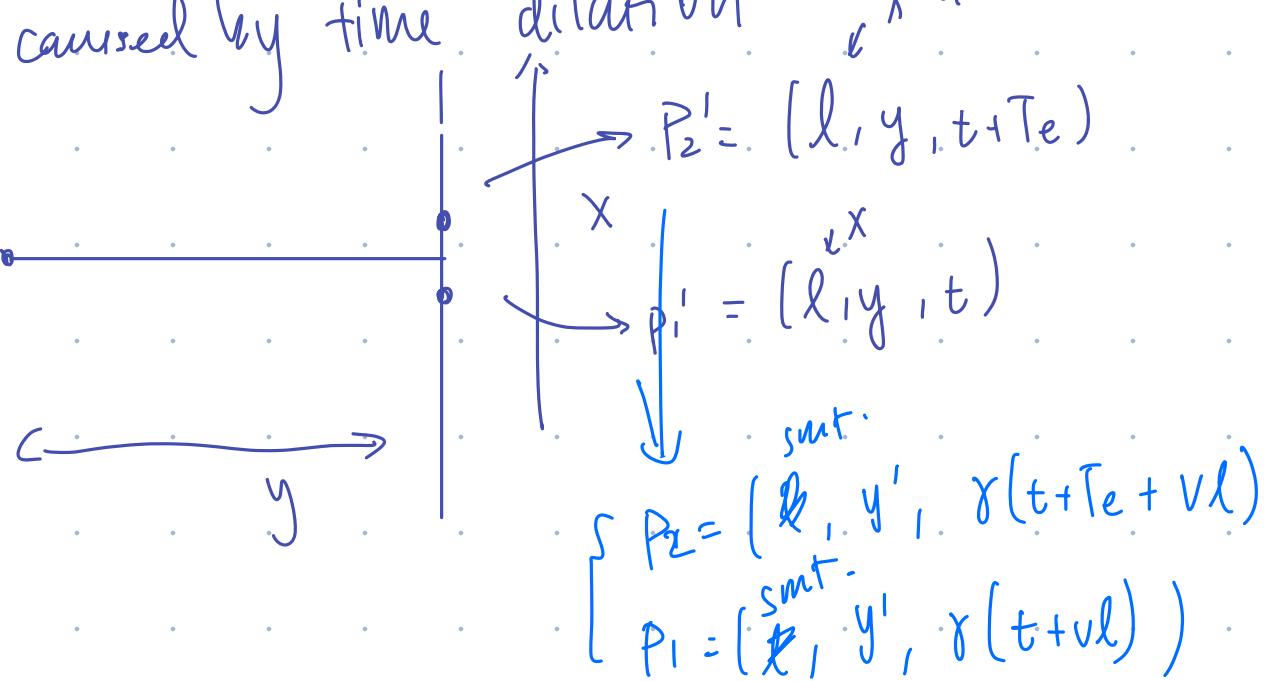
$$= \gamma T_e + \gamma v T_e$$

$$= \gamma(1+v) T_e = \frac{1+v}{\sqrt{1-v^2}} T_e$$

$$= \frac{\sqrt{1+v} \sqrt{1+v}}{\sqrt{1+v} \sqrt{1-v}} T_e$$

$$f_r = \frac{\sqrt{1-v}}{\sqrt{1+v}} f_e \rightarrow \text{longitudinal doppler effect.}$$

In the perpendicular direction, the doppler effect  
 > caused by time dilation  $\gamma^{x\text{-dir}}$



$$T_p = \gamma T_e$$

$$f_r = \sqrt{1-v^2} f_e$$

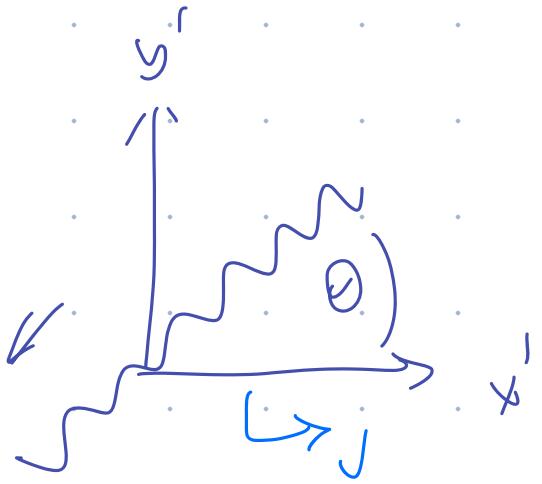
$$\cos(\vec{k} \cdot \vec{r} + \omega t)$$

$$n = \frac{\vec{k} \cdot \vec{r}}{|\vec{k}|} \quad \cos(kr + \omega t)$$

$$= \cos k \left( r + \frac{\omega}{k} t \right)$$

wave velocity

moves to the left in the  $\frac{\vec{k}}{|\vec{k}|}$  direction.



$$\text{For light } \frac{c}{k} = c = 1$$

$$\text{so } |\vec{k}| = \omega$$

emitted wave

$$\cos[\omega_e [(\cos\theta_e x' + \sin\theta_e y') + t']]$$



$$= \cos[\omega_e [r(x-vt) \cos\theta_e + y \sin\theta_e + \gamma(t-vx)]]$$

$$= \cos[\omega_e [\gamma(\cos\theta_e - v)x + y \sin\theta_e + \gamma(1 - v \cos\theta_e)t]]$$

$$\omega_r = \gamma(1 - v \cos\theta_e)t \cdot \omega_e$$

$$f_r = f_e \gamma(1 - v \cos\theta_e)t$$

Knowing the angle is not very useful.

we want to trade for  $\theta_{\text{rel}}$ .

$$\cos \theta_r = \frac{\gamma (\cos \theta_e - v)}{\sqrt{\gamma^2 (\cos \theta_e - v)^2 + \sin^2 \theta_e}}$$

$$= \frac{\frac{1}{\sqrt{1-v^2}} (\cos \theta_e - v)}{\sqrt{\frac{1}{1-v^2} (\cos \theta_e - v)^2 + \sin^2 \theta_e}}$$

$$= \frac{\cos \theta_e - v}{\sqrt{(\cos \theta_e - v)^2 + (1-v^2) \sin^2 \theta_e}}$$

$$= \frac{\cos \theta_e - v}{\sqrt{\cos^2 \theta_e + v^2 - 2v \cos \theta_e + \sin^2 \theta_e - v^2 \sin^2 \theta_e}}$$

$$= \frac{\cos \theta_e - v}{\sqrt{v^2 + 1 - v^2 \sin^2 \theta_e - 2v \cos \theta_e}}$$

$$= \frac{\cos \theta_e - v}{\sqrt{1 + v^2 \cos^2 \theta_e - 2v \cos \theta_e}}$$

$$= \frac{\cos \theta_e - v}{\sqrt{(1 - v \cos \theta_e)^2}} = \frac{\cos \theta_e - v}{1 - v \cos \theta_e} = \cos \theta_r$$

jolie fir  
cos  $\theta_e$

$$(1 - v \cos \theta_e)(1 + v \cos \theta_r)$$

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$$= 1 - v^2$$

$$1 - v \cos \theta_e = \frac{1 - v^2}{1 + v \cos \theta_r}$$

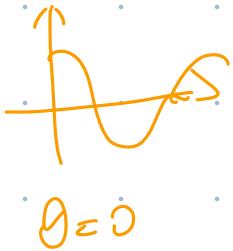
Plug into

$$f_r = f_e \gamma (1 - v \cos \theta_r)$$

$$= f_e \frac{1}{\sqrt{1-v^2}} \frac{1-v^2}{1+v \omega_s r}$$

$$= f_e \sqrt{1-v^2} \frac{1}{1+v \cos \theta_r}$$

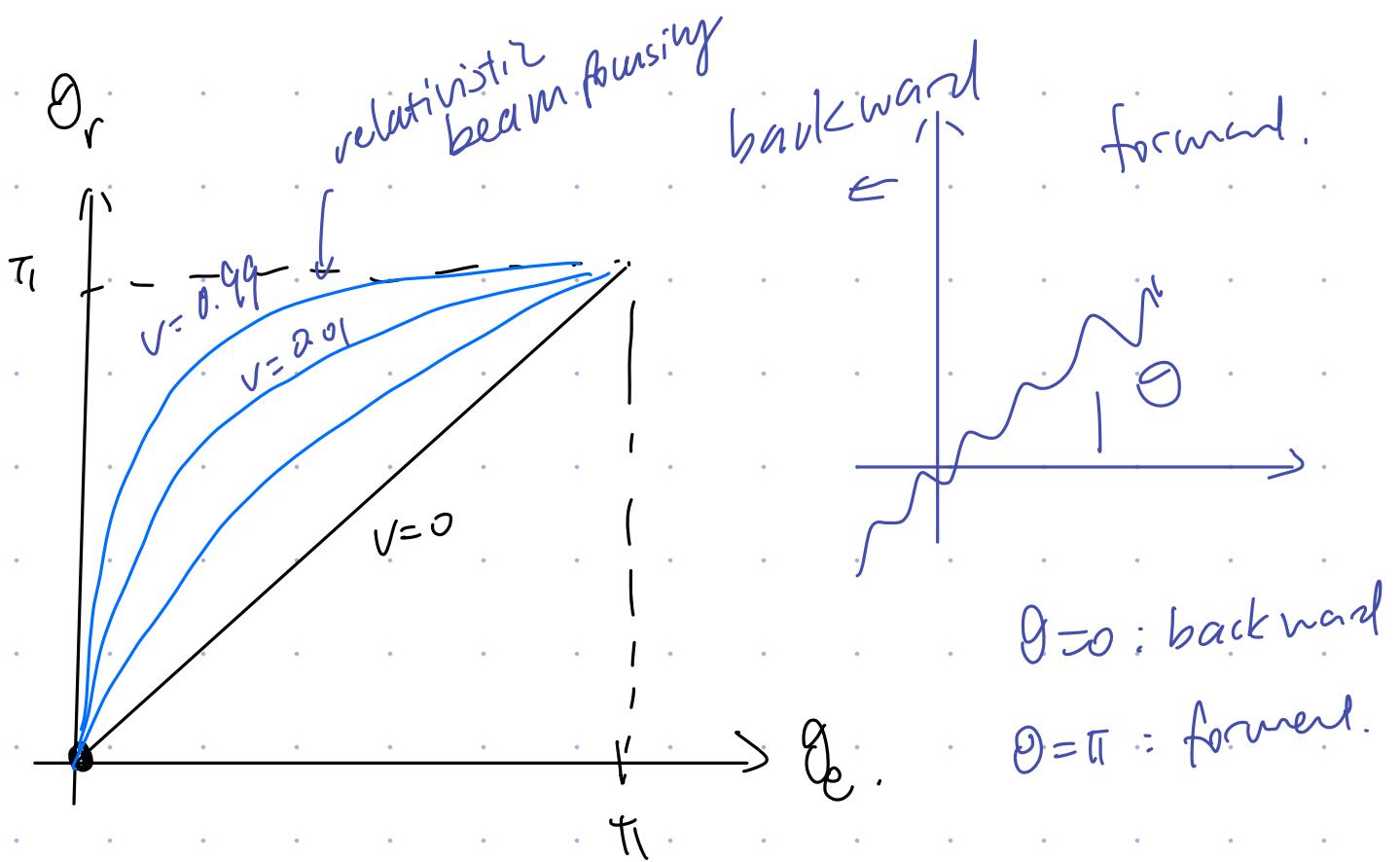
if  $\cos \theta_r = 1 \Rightarrow f_r = f_e \sqrt{1-v^2} \frac{1}{\sqrt{1+v} \sqrt{1+v}}$



$$= f_e \frac{\sqrt{1-v}}{\sqrt{1+v}} \frac{\sqrt{1+v}}{\sqrt{1+v}}$$

$$= \frac{\sqrt{1-v}}{\sqrt{1+v}} f_e$$

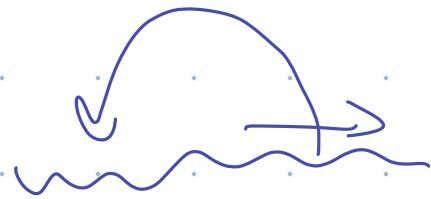
corresponds to the longitudinal effect.



$\theta = 0$ : backward

$\theta = \pi$ : forward.

$$\frac{c \omega \theta_e - v}{1 - v c \omega \theta_e}$$



## Relativistic Momentum

$\vec{P} = m \vec{v} f(\vec{v})$  depend on  $\vec{v}^2$  because of rotational symmetry.

$$\vec{r}' = \vec{r}_0' + (\omega_x', \omega_y', \omega_z') t'$$

$$\frac{d\vec{r}'}{dt'} = \vec{\omega}' = (\omega_x', \omega_y', \omega_z')$$

$$\omega_x = \frac{dx}{dt} = \frac{d\gamma(x' + vt')}{dt'} \frac{dt'}{dt} = \frac{1}{\frac{dt}{dt'}} \cdot \frac{d\gamma(x' + vt')}{dt}$$

$$= \frac{\gamma(\omega_x' + v)}{\gamma(1 + v\omega_x')}$$

$$= \frac{v + \omega_x'}{1 + v\omega_x'}$$

$$\omega_y = \frac{dy}{dt} = \frac{\frac{dy'}{dt'}}{\frac{dt}{dt'}} = \frac{\omega_y'}{\gamma(1 + v\omega_x')}$$

$$\omega_z = \frac{\omega_z'}{\gamma(1 + v\omega_z')}$$