

Relativity Lecture 1 Notes.

Suppose we have 2 coordinate systems:

$\rightarrow A(x, t)$ B is moving with velocity v relative to A.

$\rightarrow B(x', t')$ In B's perspective, A is moving with velocity $-v$ w.r.t B.

Assuming the (x, t) & (x', t') coordinate system is linearly related, we have the following.

$$\begin{cases} x' = \gamma(v) \cdot (x - vt) \quad \text{①} \\ x = \gamma(v) \cdot (x' + vt') \quad \text{②} \end{cases} \rightarrow \frac{x}{\gamma} - vt' = x' \quad \frac{x}{\gamma} - vt' = \gamma(x - vt)$$

From ①: $\frac{x'}{\gamma} + vt = x$

$$\frac{x}{\gamma} - \gamma x + \gamma vt = vt'$$

$$\text{③} = \text{②} \quad \frac{x'}{\gamma} + vt = \gamma(x' + vt')$$

$$\frac{x'}{\gamma^2} + \frac{v}{\gamma} t = x' + vt'$$

$$\frac{v}{\gamma} t = \left(1 - \frac{1}{\gamma^2}\right)x' + vt'$$

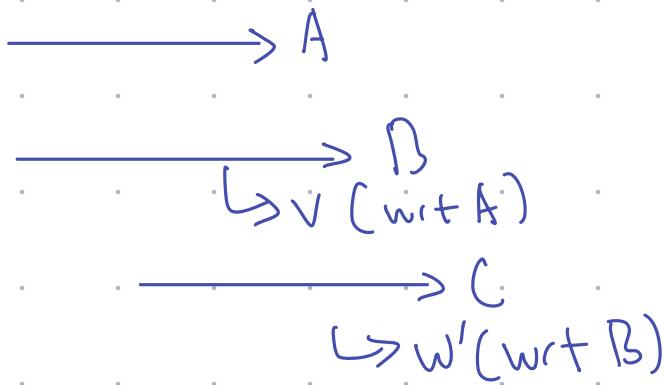
$$t = \frac{\gamma}{v} \left(1 - \frac{1}{\gamma^2}\right)x' + \frac{\gamma}{\gamma} vt'$$

$$= \frac{\gamma}{v} \left(1 - \frac{1}{\gamma^2}\right)x' + \gamma t'$$

$$t = \gamma \left[\frac{1}{v} \left(1 - \frac{1}{\gamma^2}\right)x' + t' \right]$$

$$\frac{\gamma^2 - 1}{\gamma^2}$$

$$\frac{\frac{1}{A^2} - 1}{\frac{1}{A^2}} = \frac{1 - A^2}{A^2}$$



① $x' = \omega' t'$ where does that come from?

② $\omega = \frac{x}{t} = \frac{x' + vt'}{t'}$

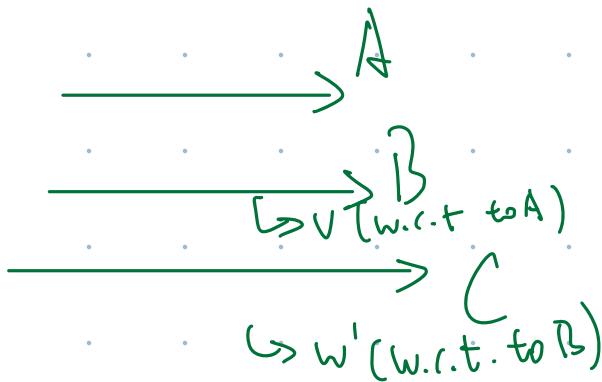
$$= \frac{\omega' t' + vt'}{\frac{1}{v} \left(1 - \frac{1}{\gamma^2}\right) \omega' t' + t'}$$

$$= \frac{\omega' + v}{\frac{1}{v} \left(1 - \frac{1}{\gamma^2}\right) \omega' + 1}$$

Now I am going to rewrite γ as $\gamma(v)$ to specify that it depends on v .

$$\omega = \frac{\omega' + v}{\frac{1}{v} \left(1 - \frac{1}{\gamma^2(v)}\right) \omega' + 1}$$

Now we want to see everything from observer C's perspective.



To observer B, C is moving with vel w'

$$\left\{ \begin{array}{l} x'' = \gamma(w') [x' - \omega' t'] \\ x' = \gamma(w') [x'' + \omega^2 t''] \end{array} \right. \begin{array}{l} \textcircled{1} \\ \textcircled{2} \end{array}$$

$$\textcircled{3} = \frac{x'}{\gamma} - \omega' t'' = x''$$

~~$\textcircled{2}$~~ $\textcircled{1} = \textcircled{3}$

also $y' = -vt'$
how did that come about?

$$\frac{x'}{\gamma} - \omega' t'' = \gamma (x' - \omega' t')$$

A moving at
-v w.r.t to
B
a point x' moves

$$\frac{x'}{\gamma} - \gamma x' + \gamma \omega' t' = \omega' t''$$

$$\gamma \left[\left(\frac{1}{\gamma^2} - 1 \right) x' + \omega' t' \right] = \omega' t''$$

$$\frac{1}{\omega'} \left[\left(\frac{1}{\gamma^2} - 1 \right) x' + t' \right] = t''$$

$$-\omega = \frac{x''}{t''} = \frac{\gamma [x' - \omega' t']}{\gamma \left[\frac{1}{\omega'} \left(\frac{1}{\gamma^2} - 1 \right) x' + t' \right]}$$

$$= \frac{-vt' - \omega' t'}{\frac{1}{\omega'} \left(\frac{1}{\gamma^2} - 1 \right) (-vt') + t'}$$

$$= \frac{-(v + \omega')}{\frac{1}{\omega'} \left(1 - \frac{1}{\gamma^2} \right) v + 1}$$

so

$$\omega = \frac{v + \omega'}{\frac{1}{\omega'} \left(1 - \frac{1}{\gamma^2} \right) v + 1}$$

Now, to specify that γ depends on w' .

|| rewrite w as

$$w = \frac{v+w'}{\frac{1}{w'} \left(1 - \frac{1}{\gamma^2(w')} \right) v + 1}$$

Since $w = \frac{v+w'}{\frac{1}{w'} \left(1 - \frac{1}{\gamma^2(w')} \right) v + 1} = \frac{v+w'}{\frac{1}{v} \left(1 - \frac{1}{\gamma^2(v)} \right) w + 1}$

that $\frac{1}{w'} \left(1 - \frac{1}{\gamma^2(w')} \right) = \frac{w'}{v} \left(1 - \frac{1}{\gamma^2(v)} \right)$

$$\frac{\cancel{w'}}{w'^2} \left(1 - \frac{1}{\gamma^2(w')} \right) = \frac{\cancel{w'}}{v^2} \left(1 - \frac{1}{\gamma^2(v)} \right)$$

$$\underbrace{\frac{1}{w'^2} \left(1 - \frac{1}{\gamma^2(w')} \right)}_{\textcircled{1}} = \underbrace{\frac{1}{v^2} \left(1 - \frac{1}{\gamma^2(v)} \right)}_{\textcircled{2}} = B.$$

That means the two can't depend on time, so we can set it to be a constant.



$$\omega = \frac{v + \omega'}{B \omega' v + 1}$$

we could have $\begin{cases} B = 0 : \text{galilean-} \\ B > 1 : \text{relativistic} \end{cases}$

$$B = \frac{1}{v^2} \left(1 - \frac{1}{\gamma^2} \right)$$

$$v^2 B = 1 - \frac{1}{\gamma^2}$$

$$\gamma^2 v^2 B = \gamma^2 - 1$$

$$\gamma^2 [v^2 B - 1] = -1$$

$$\gamma^2 = \frac{1}{1 - v^2 B}$$

$$\gamma = \frac{1}{\sqrt{1 - v^2 B}}$$

turns out that $B = \frac{1}{c^2}$

$$\text{so } \gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Fitcau's
experiment
measured
yields $B = 1.1 \times 10^{-17} \text{ N}^2 \text{ m}^2$

Here we have the Lorentz transformation:

$$\left\{
 \begin{array}{l}
 x' = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} (x - vt) \quad 1 - \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \\
 x = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} (x' + vt) \quad x - \frac{v^2}{c^2} - x \\
 t = \gamma \left[\frac{1}{v} \left(1 - \frac{1}{\gamma^2} \right) x' + t' \right] \\
 = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \left[\frac{1}{v} \left[x - x + \frac{v^2}{c^2} \right] x' + t' \right] \\
 = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \left[\frac{v}{c^2} x' + t' \right] \\
 t' = \gamma \left[\frac{1}{v} \left(\frac{1}{\gamma^2} - 1 \right) x + t \right] \\
 = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \left[\frac{1}{v} \left[1 - \frac{v^2}{c^2} \right] x + t \right] \\
 = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \left(- \frac{v}{c^2} x + t \right).
 \end{array}
 \right.$$

□ Lorentz Transformation