

Mechatronics for Rehabilitation Engineering: Course Notes

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Chapter 1

Review of Basic Circuit Theory

This chapter is a quick review of basic circuit theory. Prior knowledge of these topics is assumed, along with basic understanding of linear time invariant systems - Fourier/Laplace transforms, linear constant coefficient differential equations, transfer functions, and frequency response. The reader is encouraged to review the material in this chapter before proceeding to the next chapter.

The interaction of electromagnetic fields with matters is the basis of all electrical and electronic devices. These interactions are often described, analyzed and synthesized through the abstractions of electrical circuit theory. The following are the most basic circuit two-terminal elements we will need for now. We will introduce new ones as and when they are required. Each circuit element has a unique voltage-current relationship, and it is *important* that you know these by heart.

1.1 Basic

Independent Voltage source. An Ideal voltage source provides a fixed voltage V between its two terminals, and can provide any amount of current. Notice that the voltage V can be fixed or time varying. For example, for a DC voltage source with $V = 5V$, the voltage across the two terminals will be $5V$ for all time. But for a time varying AC source, $V = 5 \sin(100\pi t)$, the voltage across its terminal will vary with time. We will often drop the adjective "independent" when we are sure that the context is

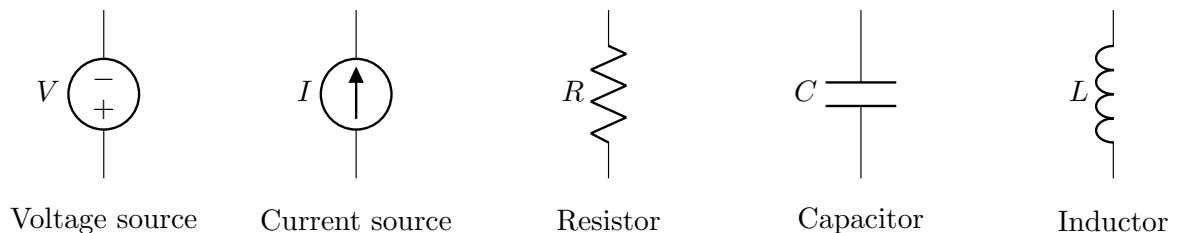


Fig. 1.1: Basic circuit elements: voltage source, current source, resistor, capacitor, and inductor.

clear. We will look at dependent sources later, and we will always use the adjective "dependent" to refer to them.

Independent Current source. An ideal current source provides a fixed amount of current to flow through its terminals (out through one and in through the other), irrespective of the voltage across its terminals. Current sources can also be time-varying.

Resistor. A passive element where the current i_R flowing through the element is proportional to the voltage v_R across its terminals.

$$v_R \propto i_R$$

In the case of linear resistors, the proportionality factor is constant, resulting in Ohm's law,

$$v_R = R i_R \quad (1.1)$$

The units of R are $V \cdot A^{-1}$ or *Ohms* (Ω). R in general is positive. The power absorbed by a resistor is given by the product of the voltage across it and the current flowing through it,

$$P = v_R i_R = i_R^2 R = \frac{v_R^2}{R} \quad (1.2)$$

This power is dissipated as heat by the resistor. Note that the power absorbed by a resistor is always positive, since R is positive.

We will later see non-linear resistors, where the resistance varies as a function of the applied voltage, temperature and other factors.

Capacitor. A capacitor is another passive element with the following voltage current relationship.

$$i_C = C \frac{dv_C}{dt} \quad (1.3)$$

The current i_C through the capacitor is proportional to the rate of change of voltage across its terminals v_C . The proportionality factor is called the capacitance C , and has units of F (Farads) or $C \cdot V^{-1}$. The voltage across the capacitor at any given time is proportional to the integral of the current flowing through it or the charge stored in the capacitor. The voltage across the capacitor is given by

$$v_C = \frac{q}{C} = \frac{1}{C} \int i_C dt \quad (1.4)$$

The instantaneous power absorbed by the capacitor is given by,

$$P = v_C i_C = C v_C \frac{dv_C}{dt} \quad (1.5)$$

The power absorbed by the capacitor can be positive or negative, depending on the direction of current flow. If the current is flowing into the capacitor, then the voltage across it is increasing, and the power absorbed is positive. If the current is flowing out of the capacitor, then the voltage across it is decreasing, and the power absorbed is

negative. A capacitor stores energy in the form of electric field between its plates. The energy stored in a capacitor at any given time depends on the charge stored in it, and is given by,

$$E = \frac{1}{2} C v_C^2 \quad (1.6)$$

Inductor. An inductor is another passive element with the following voltage current relationship.

$$v_L = L \frac{di_L}{dt} \quad (1.7)$$

The voltage v_L across the inductor is proportional to the rate of change of current i_L flowing through it. The proportionality factor is called the inductance L , and has units of H (Henries) or $V \cdot s \cdot A^{-1}$. The current through the inductor at any given time is proportional to the integral of the voltage across it. The current through the inductor is given by

$$i_L = \frac{1}{L} \int v_L dt \quad (1.8)$$

The instantaneous power absorbed by the inductor is given by,

$$P = v_L i_L = L i_L \frac{di_L}{dt} \quad (1.9)$$

The power absorbed by the inductor can be positive or negative, depending on the direction of current flow. If the current is flowing into the inductor, then the voltage across it is increasing, and the power absorbed is positive. If the current is flowing out of the inductor, then the voltage across it is decreasing, and the power absorbed is negative. An inductor stores energy in the form of magnetic field around it. The energy stored in an inductor at any given time depends on the current flowing through it, and is given by,

$$E = \frac{1}{2} L i_L^2 \quad (1.10)$$

1.2 Kirchoff's Laws

The five elements alone are not that interesting. But interesting things can be done by connecting these elements together in different ways to form an electrical circuit. The elements are connected together by wires, which are assumed to be perfect conductors, i.e. zero resistance. Consider the following circuit (Figure 1.2),

How do we find out the voltages and currents in the circuit? Kirchoff's laws can be used for analysing such circuits, which are based on the conservation of charge and energy. The circuit in Figure 1.2 is a simple electrical circuit with a voltage source, a current source, and a bunch of resistors. The voltage source provides a fixed voltage V_s between its two terminals, and the current source provides a fixed amount of current I_s to flow through its terminals. The voltages and currents in the rest of the elements

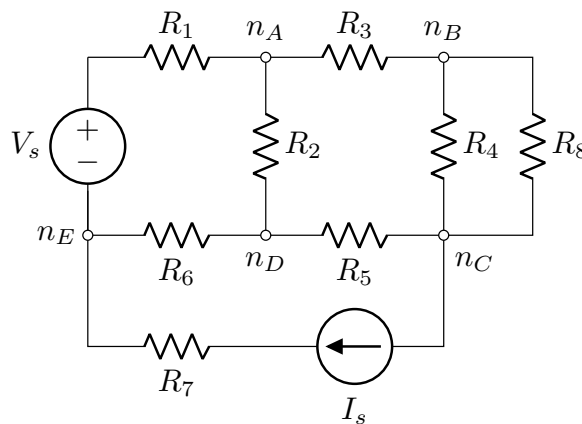


Fig. 1.2: A simple electrical circuit with a voltage source, a current source, and a bunch of resistors.

will be determined by Kirchoff's laws with the constraints imposed by the voltage and current sources. The two laws are:

1. **Kirchoff's current law (KCL):** The sum of the currents entering a node is equal to the sum of the currents leaving the node. A *node* is a point at which two or more circuit elements are connected together. In Figure 1.2, n_A , n_B , n_C , n_D and n_E are examples of nodes where three elements are connected together. There are two other nodes in the circuit, can you identify them?

The sum of the currents at a node is equal to zero. This is based on the conservation of charge, and can be expressed mathematically as:

$$\sum_{i=1}^n i_i = 0 \quad (1.11)$$

where i_i is the current flowing into or out of the node, and n is the number of elements connected to the node. The current flowing into a node is positive, and the current flowing out of a node is negative.

2. **Kirchoff's voltage law (KVL):** The sum of the voltages around a closed loop in a circuit is equal to zero. A *closed loop* is a path in the circuit that starts and ends at the same node, and does not cross itself. In Figure 1.2, the path starting from n_A , to n_D , to n_B , and back to n_A is a closed path. This path includes the resistors R_2 , R_5 , R_4 , and R_3 .

This is based on the conservation of energy, and can be expressed mathematically as:

$$\sum_{i=1}^n v_i = 0 \quad (1.12)$$

where v_i is the voltage across each element in the loop, and n is the number of elements in the loop.

The voltage across an element is positive if the current is flowing into the positive terminal of the element, and negative if the current is flowing out of the positive terminal of the element.

Note that the two laws apply for any type of circuit element used in the circuits, independent or dependent voltage/current sources, resistors, capacitors, inductors, either two, three or four terminal elements.

1.3 Series and Parallel Connections

Two elements that share the same voltage across them between a given pair of nodes are said to be **parallel** to each other. In Figure 1.2, R_4 and R_8 are parallel to each other. In a single loop, two elements that share the same current are said to be in **series** with each other. In Figure 1.2, V_s and R_1 are in series, I_s and R_7 are in series.

1.3.1 Resistors in series and parallel

When n resistors $R_1, R_2, \dots, R_n \geq 0$ are in series, these can be combined to an equivalent resistor with resistance R_{eq} given by the following,

$$R_{eq} = \sum_{i=1}^n R_i \implies R_{eq} \geq \max_{1 \leq i \leq n} R_i. \quad (1.13)$$

Note that in a series connection the equivalent resistance is at least as large as the largest value of R_1 to R_n .

When n resistors $R_1, R_2, \dots, R_n \geq 0$ are in parallel, these can be combined to an equivalent resistor with resistance R_{eq} given by the following,

$$\frac{1}{R_{eq}} = \sum_{i=1}^n \frac{1}{R_i} \implies R_{eq} = \frac{R_1 R_2 \cdots R_n}{R_1 + R_2 + \cdots + R_n} \implies R_{eq} \leq \min_{1 \leq i \leq n} R_i \quad (1.14)$$

Note that in a parallel connection, the equivalent resistance cannot be larger than the smallest value of R_1 to R_n .

1.3.2 Capacitors in series and parallel

Series connection of n capacitors $C_1, C_2, \dots, C_n \geq 0$

$$C_{eq} = \frac{C_1 C_2 \cdots C_n}{C_1 + C_2 + \cdots + C_n} \quad (1.15)$$

Parallel connection of n capacitors $C_1, C_2, \dots, C_n \geq 0$

$$C_{eq} = C_1 + C_2 + \cdots + C_n \quad (1.16)$$

1.3.3 Inductors in series and parallel

Series connection of n inductors $L_1, L_2, \dots, L_n \geq 0$

$$L_{eq} = L_1 + L_2 + \dots + L_n \quad (1.17)$$

Parallel connection of n inductors $L_1, L_2, \dots, L_n \geq 0$

$$L_{eq} = \frac{L_1 L_2 \dots L_n}{L_1 + L_2 + \dots + L_n} \quad (1.18)$$

Its left as an exercise for you to verify these expressions.

What does the equivalent resistance actually mean? The equivalent resistor with resistance R_{eq} has the same voltage-current relationship as the individual elements in series or parallel connection. We can replace the series or parallel connection of the individual resistors R_1 to R_n by a single resistor with value R_{eq} without changing the voltage current relationships in the circuit. The same argument applies for equivalent capacitors and inductors.

1.3.4 Voltage sources in series and parallel

Series connection of n voltage sources V_1, V_2, \dots, V_n will result in an equivalent voltage source V_{eq} given by

$$V_{eq} = V_1 + V_2 + \dots + V_n \quad (1.19)$$

Voltage sources should not be connected in parallel, as this will result in a short circuit. Ideally, an infinite current will flow through the connection, because the voltage sources force a potential difference between the two ends of the wire to be zero and it has zero resistance. Parallel connections are allowed only when the two sources have the same voltage and polarity.

1.3.5 Current sources in series and parallel

Parallel connection of n current sources I_1, I_2, \dots, I_n will result in an equivalent current source I_{eq} given by

$$I_{eq} = I_1 + I_2 + \dots + I_n \quad (1.20)$$

Current sources should not be connected in series; series connections are allowed only when the two sources have the same current and polarity.

1.4 Superposition Principle

Linear approximations of circuits are often employed as first order approximations when analysing circuits. A linear circuit is one that consists of linear passive elements, independent sources and linear dependent sources. Linear circuits follow the superposition

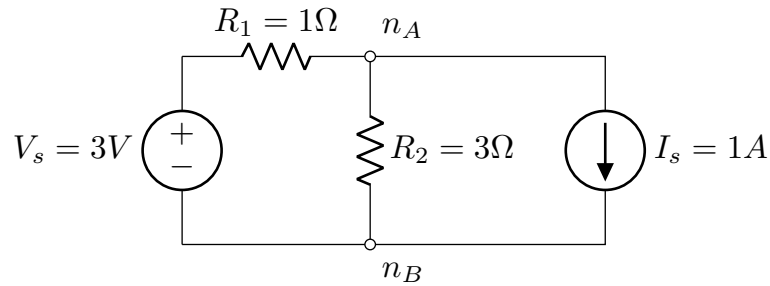


Fig. 1.3: A simple circuit with two sources

principle, which states that the response of a linear circuit to a linear combination of inputs is equal to the corresponding linear combination of the responses to each input applied separately. Solving the following circuit (Figure 1.3) should make this concept clear. For the circuit in Figure 1.3, perform the following calculation and compare your results.

Step 1. Solve for the voltage across and the current through the resistors R_1 and R_2 ; we will refer to these as v_{R1} , i_{R1} and v_{R2} , i_{R2} , respectively. You should use both Kirchhoff's current and voltages laws to compute these variables.

Let's now consider of the sources in the circuit separately. This would mean making the source values "zero". This corresponds to two different operations on the circuit. Zeroing a voltage source corresponds to replacing it with a wire (a short circuit), while zeroing a current source corresponds to simply removing the current source (an open circuit).

Step 2. Zero the voltage source V_s and compute the voltages and currents associated with the two resistors. We will refer to these as $v_{R1, V_s=0}$, $v_{R2, V_s=0}$, $i_{R1, V_s=0}$, and $i_{R2, V_s=0}$.

Step 3. Zero the current source I_s and compute the voltages and currents associated with the two resistors. We will refer to these as $v_{R1, I_s=0}$, $v_{R2, I_s=0}$, $i_{R1, I_s=0}$, and $i_{R2, I_s=0}$.

Step 4. For a linear circuit, shown in Figure 1.3, the following will always be true.

$$\begin{aligned}
 v_{R1} &= v_{R1, V_s=0} + v_{R1, I_s=0} \\
 v_{R2} &= v_{R2, V_s=0} + v_{R2, I_s=0} \\
 i_{R1} &= i_{R1, V_s=0} + i_{R1, I_s=0} \\
 i_{R2} &= i_{R2, V_s=0} + i_{R2, I_s=0}
 \end{aligned} \tag{1.21}$$

Let's assume that I am only interested in i_{R2} . Can I use $i_{R2, V_s=0}$ and $i_{R2, I_s=0}$ to compute the current i_{R2} if $V_s = 1V$ and $I_s = -2A$?

1.5 Practical Voltage and Current Sources

The independent voltage and current sources we have discussed so far are "ideal" sources. Practical or real sources do not behave like them - a battery cannot pro-

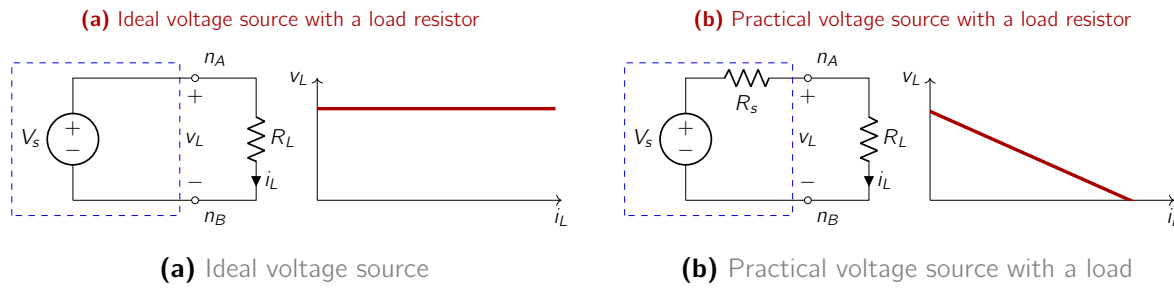


Fig. 1.4: Comparison of the voltage-current relationship of an ideal and a practical voltage source.

vide any amount of current for a load without any changes to the voltage across its terminals.

A good model of practical voltage source is an ideal voltage source V_s in series with a *internal, source or output* resistor R_s . And for a practical current source, it is an ideal current source I_s in parallel with a resistor R_s . The voltage across the terminal of a voltage source as a function of the current drawn from it is depicted for an ideal and practical voltage source in Figure 1.4. For the ideal source (Figure 1.4a), the voltage v_L is independent of the current i_L drawn from it. For the practical voltage source (Figure 1.4b), the voltage across the terminals is a function of the current drawn from it,

$$v_L = V_s - R_s i_L \quad (1.22)$$

When $R_L = \infty$ ($i_L = 0$), the voltage across the terminals is equal to the voltage of the source V_s . This is maximum voltage the practical source can provide. This is also know as the *open circuit voltage* v_{oc} of the source. When $R_L = 0$ ($v_L = 0$), the current drawn from the source $i_L = \frac{V_s}{R_s}$. This is the maximum current the voltage source can provide. This is also known as the *short circuit current* i_{sc} of the source.

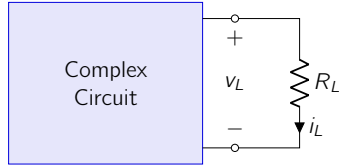
Problem 1.1. Plot the voltage-current relationship of a practical current source with $I_s = 2A$ and $R_s = 10\Omega$. What are v_{oc} and i_{sc} ?

1.6 Thevenin's and Norton's Theorems

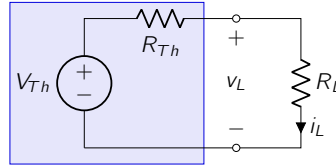
Thevenin's and Norton's theorems are two important theorems in circuit theory that allow us to simplify complex circuits into simpler equivalent circuits. These theorems are based on the superposition principle, and can be used to analyze linear circuits with independent and dependent sources.

Thevenin's Theorem. Thevenin's theorem states that any linear circuit with independent and dependent sources can be replaced by an equivalent circuit with a single voltage source V_{th} in series with a resistor R_{th} , connected to the load resistor R_L .

(a) Complex circuit with a load resistor



(b) Thevenin equivalent circuit



(c) Norton equivalent circuit

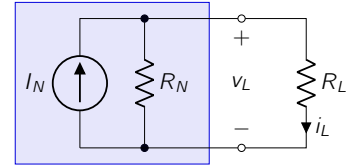


Fig. 1.5: Thevenin's and Norton's circuits for a complex linear circuit. The Thevenin and Norton equivalent circuits have the same voltage-current relationship for any load.

Norton's Theorem. Norton's theorem states that any linear circuit with independent and dependent sources can be replaced by an equivalent circuit with a single current source I_N in parallel with a resistor R_N , connected to the load resistor R_L .

Computing the Thevenin and Norton equivalent circuits. The Thevenin and Norton equivalent circuits can be computed using the following steps:

1. Remove the load resistor R_L from the circuit.
2. Compute the open circuit voltage v_{oc} across the terminals of the load resistor. This is the Thevenin voltage V_{th} . This can be done using the superposition principle, where we zero all the independent sources except one and compute the open circuit voltage. The overall open circuit voltage is the sum of the open circuit voltages when all other sources are zeroed.
3. Compute the short circuit current i_{sc} through the terminals of the load resistor. This is the Norton current I_N . This too can be calculated using the superposition principle.
4. Compute the Thevenin resistance R_{th} by zeroing all independent sources in the circuit and computing the equivalent resistance seen from the terminals of the load resistor (without the load resistor). This is also equal to the Norton resistance R_N .
5. The Thevenin and Norton equivalent circuits are then given by:

$$\begin{aligned} V_{th} &= v_{oc} \\ I_N &= i_{sc} \\ R_{th} &= R_N \end{aligned} \tag{1.23}$$

6. The load resistor R_L can be connected to either the Thevenin or Norton equivalent circuit, and the voltage and current across it can be computed using the voltage-current relationship of the equivalent circuit.

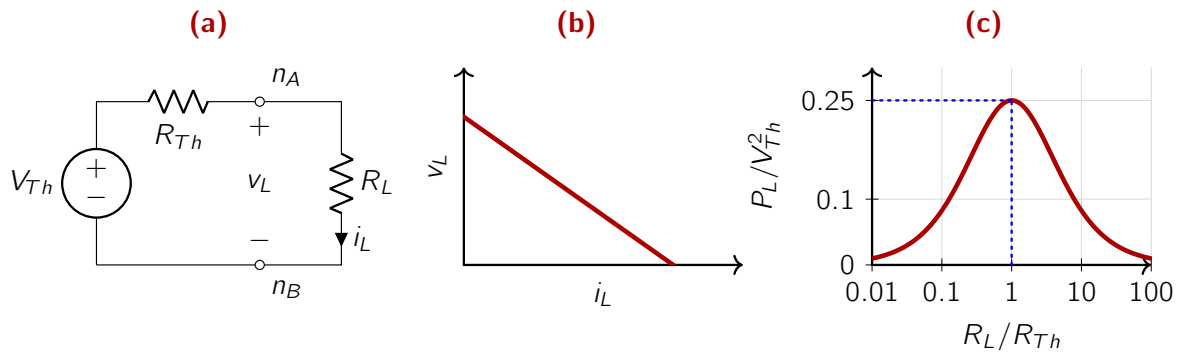


Fig. 1.6: Maximum power transfer theorem.

Problem 1.2. Compute the Thevenin and Norton equivalent circuits for the circuit shown in Figure 1.2 assuming the following as the load resistor: (a) R_8 ; (b) R_2 ; and (c) R_7 .

1.7 Maximum Power Transfer Theorem

The maximum power transfer theorem states that the maximum power is transferred to the load resistor R_L when the load resistance is equal to the Thevenin resistance R_{th} of the circuit. Consider the following circuit (Figure 1.6). The power absorbed by the load resistor R_L is given by,

$$P_{R_L} = \frac{R_L}{(R_{th} + R_L)^2} V_{th}^2 \quad (1.24)$$

Its easy to check that the optimal value of R_L that maximizes the power absorbed by the load resistor is given by,

$$\begin{aligned} R_L^* &= \arg \max_{R_L} \frac{R_L}{(R_{th} + R_L)^2} V_{th}^2 = R_{th} \\ P_L^* &= \max_{R_L} \frac{R_L}{(R_{th} + R_L)^2} V_{th}^2 = \frac{1}{4} \left(\frac{V_{th}^2}{R_{th}} \right) \end{aligned} \quad (1.25)$$

Problem 1.3. Prove the statements in Eq. 1.25. (*Hint:* Use the first order condition for maximization of a continuous function.)

1.8 RC, RL, and RLC Circuits

Unlike resistors, that neither remember the past values of its votlage or current, capacitors and indutors retain memory of the past current and past voltage, respectively. This

can be easily seen from their respective voltage-current relationships given by Eq. 1.3 and Eq. 1.7, respectively.

Capacitor:

$$i_C = C \frac{dv_C}{dt} \implies v_C(t) = \frac{1}{C} \int_0^t i_C(\tau) d\tau + v_C(0) \quad (1.26)$$

The instantaneous voltage across a capacitor contains some information about the past history the current that has flown through the capacitor. This voltage $v_C(t)$ is determined by the charge on the capacitor at time t . Note, that the voltage across the capacitor cannot change instantaneously, as this would require an infinite current to flow through the capacitor. Theoretically, however, an impulse (Dirac delta function) current applied to the capacitor can produce instantaneous change in the capacitor's voltage.

Inductor:

$$v_L = L \frac{di_L}{dt} \implies i_L(t) = \frac{1}{L} \int_0^t v_L(\tau) d\tau + i_L(0) \quad (1.27)$$

Similarly, the instantaneous current through the inductor contains some information about the history of voltage applied across the inductor. Current through an inductor cannot change instantaneously, as this would require an infinite voltage to be applied across the inductor. An impulse voltage applied to the inductor can produce instantaneous change in the inductor's current.

RC Circuit. Consider a simple RC circuit shown in Figure 1.7a. We can write Kirchhoff's voltage law for the circuit as follows,

$$RC \frac{dv_C}{dt} + v_C = V_s \implies v_C(t) = e^{-t/RC} \left[v_C(0) + \int_0^t \frac{1}{RC} e^{\tau/RC} V_s(\tau) d\tau \right] \quad (1.28)$$

The above equation gives the general solution for the voltage across the capacitor. We can derive all other variables of interest from v_C . The response to an impulse input, step input, sinusoidal input, or any arbitrary V_s can be computed using the above equation. The response for a step input obtained using a fixed sources and a switch is given in Figure 1.7a. RC is the time constant of the circuit, and is a measure of how fast the capacitor charges or discharges and it has the units of time.

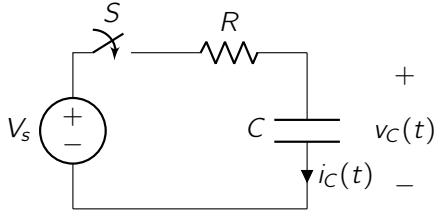
Using Laplace transform to analyze the circuit, we can write the following expression for the voltage across the capacitor,

$$V_C(s) = \frac{1}{sRC + 1} V_s(s) + \frac{v_C(0)}{sRC + 1} \quad (1.29)$$

$V_C(s)$ and $V_s(s)$ are the Laplace transforms of $v_C(t)$ and $V_s(t)$. Replacing $s = j\omega$ will give us the frequency response of the system with $V_C(s)$ as the output and $V_s(s)$ as the input.

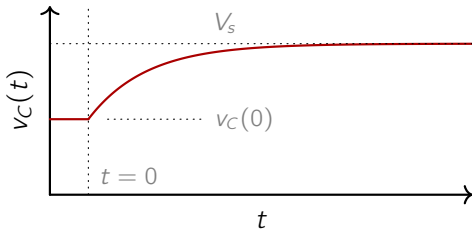
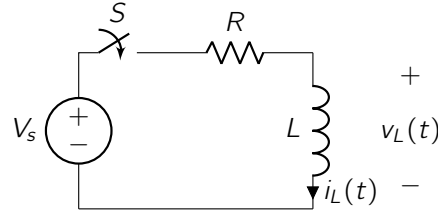
RL Circuit. Consider a simple RL circuit shown in Figure 1.7b. Kirchhoff's voltage law for the circuit is as follows,

$$\frac{L}{R} \frac{di_L}{dt} + i_L = \frac{1}{R} V_s \implies i_L(t) = e^{-tR/L} \left[i_L(0) + \int_0^t \frac{1}{L} e^{\tau R/L} V_s(\tau) d\tau \right] \quad (1.30)$$

(a) A Simple RC Circuit

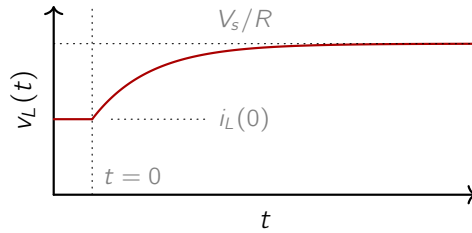
Assuming the switch S is closed at $t = 0$, the voltage across the capacitor $v_C(t)$ is given by,

$$v_C(t) = V_s \cdot \left(1 - e^{-\frac{t}{RC}}\right) + v_C(0)e^{-\frac{t}{RC}}, \quad t \geq 0$$

**(a)** Ideal voltage source**(b) A Simple RL Circuit**

Assuming the switch S is closed at $t = 0$, the current through the inductor $i_L(t)$ is given by,

$$i_L(t) = \frac{V_s}{R} \cdot \left(1 - e^{-\frac{t}{L/R}}\right) + i_L(0)e^{-\frac{t}{L/R}}, \quad t \geq 0$$

**(b)** Practical voltage source with a load**Fig. 1.7:** Simple RC and RL circuits and their transient responses.

The above equation gives the general solution for the current through the inductor. The response to an impulse input, step input, sinusoidal input, or any arbitrary V_s can be computed using the above equation. The response for a step input obtained using a fixed source and a switch is given in Figure 1.7a. $\frac{L}{R}$ is the time constant of the circuit, and is a measure of how fast the current through the inductor can change; it has the units of time.

Similarly, applying the Laplace transform, we have,

$$I_L(s) = \frac{1}{sL + R}V_s(s) + \frac{i_L(0)}{sL + R} \quad (1.31)$$

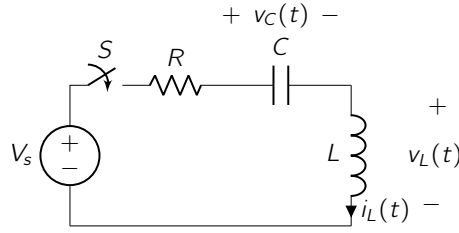
RLC Circuit. Consider a simple series RLC circuit shown in Figure 1.8. Kirchhoff's voltage law for the circuit is as follows,

$$LC \frac{d^2 v_C}{dt^2} + RC \frac{dv_C}{dt} + v_C = V_s \quad (1.32)$$

The response of the circuit when the switch is closed at $t = 0$ is qualitatively different depending on the values of R , L and C . This is better understood in the Laplace domain. The Laplace transform of the response $v_C(t)$ is given by,

$$V_C(s) = \frac{V_s(s)}{LCs^2 + RCs + 1} + v_C(0) \frac{C(Ls + R)}{LCs^2 + RCs + 1} + i_L(0) \frac{RC}{LCs^2 + RCs + 1} \quad (1.33)$$

(a) A Simple RLC Circuit



(b) Four different step responses

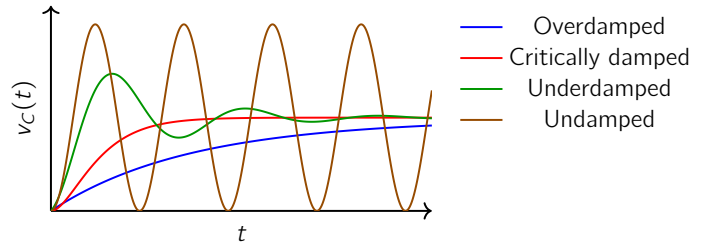


Fig. 1.8: A simple RLC circuit and the different responses of the circuit when the switch S is closed at time $t = 0$, with zero initial capacitor voltage and inductor current when $t = 0$.

Note that $i_L(0) = \frac{dv_C}{dt}(0)$. The denominator of the above equation is a second order polynomial in s , and thus the response of the circuit will depend on the roots of the polynomial. The four different responses are shown in Figure 1.8.

1.9 Steady State Sinusoidal Analysis

The previous subsection looked at the transient response of the some simple circuits involving R , L , and C . However, we are often interested in the steady state response of these circuits to sinusoidal excitations. This is because sinusoidal signals are eigenfunctions of linear systems, such as the linear circuits we have discussed so far. With sinusoidal excitations, all currents and voltages in the circuit will also be sinusoidal with the same frequency as the input excitations, although with different amplitudes and phases. Complex numbers are used in this analysis for compact representation of amplitude and phase of signals.

For steady state frequency analysis for circuits involving R , L , and C , we extend the idea of resistance to impedance. Fourier transforms provide a natural way to analyze the steady state response of linear circuits to sinusoidal excitations. We first recast the voltage-current relationships of the circuit elements in the Fourier domain.

$$\begin{aligned}
 \textbf{Resistor:} \quad v_R &= R i_R \implies V_R(j\omega) = R \cdot I_R(j\omega) \\
 \textbf{Capacitor:} \quad i_C &= C \frac{dv_C}{dt} \implies V_C(j\omega) = \frac{1}{j\omega C} \cdot I(j\omega) \\
 \textbf{Inductor:} \quad v_L &= L \frac{di_L}{dt} \implies V_L(j\omega) = j\omega L \cdot I_L(j\omega)
 \end{aligned} \tag{1.34}$$

The Fourier transformed voltage-current relationships of the circuit elements look like Ohms law, with the concept of resistance extended to impedance. The impedance of a capacitor and inductor defined as the ratio of the Fourier transformed voltage to the Fourier transformed current. They are functions of the frequency of the voltage/current.

The impedance of a resistor, capacitor, and inductor are given by,

$$\begin{aligned}
 \text{Resistor: } Z_R &= R \\
 \text{Capacitor: } Z_C &= \frac{1}{j\omega C} \\
 \text{Inductor: } Z_L &= j\omega L
 \end{aligned} \tag{1.35}$$

Equivalent Impedance. The equivalent impedance of a circuit with impedances Z_1, Z_2, \dots, Z_n in series is given by $Z_{eq} = Z_1 + Z_2 + \dots + Z_n$. When the impedances Z_1, Z_2, \dots, Z_n are in parallel, the equivalent impedances is given by $\frac{1}{Z_{eq}} = \frac{1}{Z_1} + \frac{1}{Z_2} + \dots + \frac{1}{Z_n}$.

Amplitude and phase modification by an impedance. If a sinusoidal current excitation $i(t) = I_o \sin(\omega t)$ is passed through an impedance Z . Then the voltage across the impedance is given by,

$$v(t) = |Z|I_o \sin(\omega t + \arg Z) \tag{1.36}$$

The amplitude of the sinusoidal voltage is $|Z|I_o$, and the phase of the sinusoidal voltage is $\arg Z$, with respect to the current $i(t)$.

1.10 Exercise

- Plot the current through a resistor $R = 10\Omega$, capacitor $C = 5\mu F$, and inductor $L = 2mH$ for the following voltage inputs.

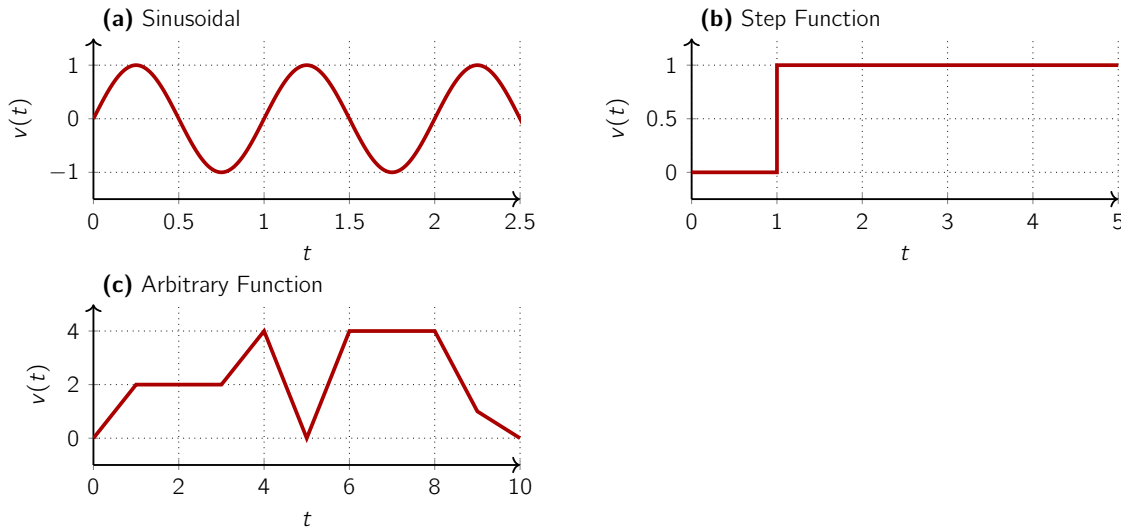


Fig. 1.9: [Exercise 1] Voltage inputs applied to a resistor, capacitor, or an inductor. The units of voltage are in volts, and the time is in seconds.

2. Find the Thevenin and Norton equivalent circuits of the following circuits (Figure 1.10). Note that one of the problems has a dependent source. Dependent sources are voltage or current sources (diamond shaped elements), whose voltage or current, respectively, depend on the voltage or current across another element in the circuit.

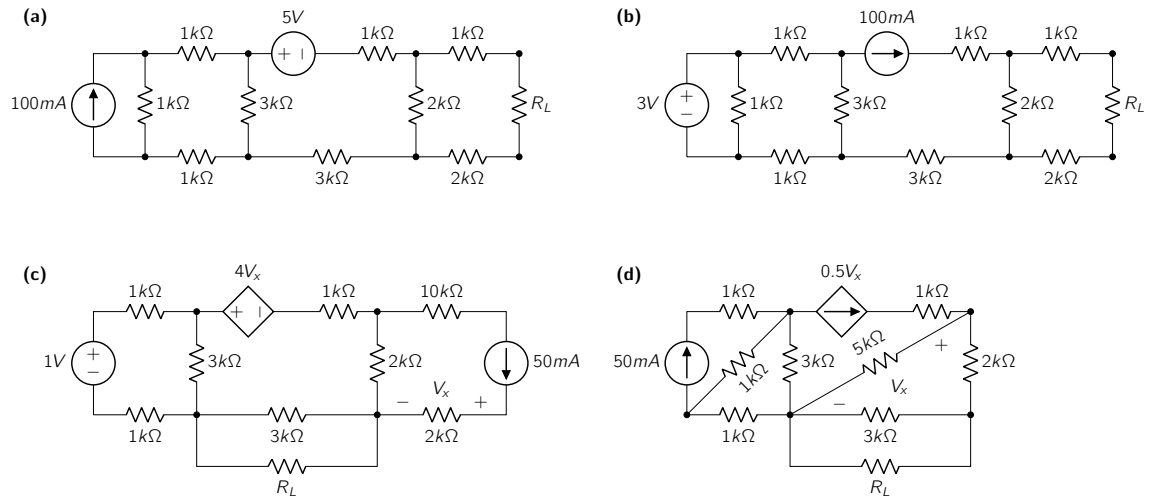


Fig. 1.10: [Exercise 2] Find the equivalent impedance of the circuit.

3. A certain red LED has a maximum current rating of 35mA , and if this value is exceeded, overheating and catastrophic failure will result. The resistance of the LED is a nonlinear function of its current, but the manufacturer warrants a minimum resistance of 47Ω and a maximum resistance of 117Ω . Only 9V batteries are available to power the LED. Design a suitable circuit to deliver the maximum power possible to the LED without damaging it. Use only combinations of the standard resistor values. (*This problem is from Engineering Circuit Analysis by Hayt Jr. et al.*)
4. The load resistor in Figure 1.11 can safely dissipate up to 1W before overheating and bursting into flame. The lamp can be treated as a 10.6Ω resistor if less than 1A flows through it and a 15Ω resistor if more than 1A flows through it. What is the maximum permissible value of I_s ? Verify your answer with an appropriate computer simulation. (*This problem is from Engineering Circuit Analysis by Hayt Jr. et al.*)
5. In the following circuit (Figure 1.12), find the voltage across the capacitors C_1 and C_2 . The single pole double throw switch S is connected to the top terminal before time $t = 0$. At time $t = 0$, switch is flipped to the bottom terminal. Assume that before time $t = 0$, the voltage across the capacitor C_2 , $v_{C2} = 1\text{V}$. Draw the plot of the voltage across the capacitors C_1 and C_2 as function of time.

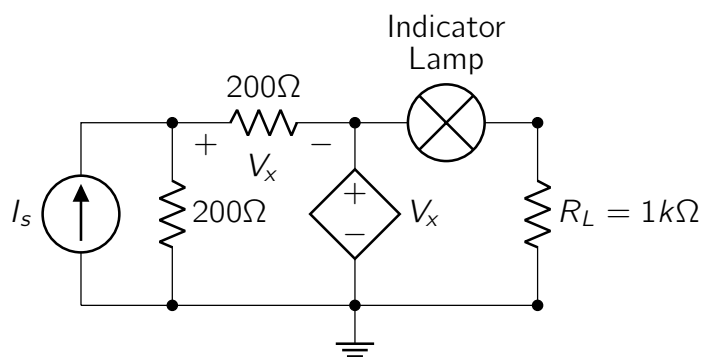


Fig. 1.11: [Exercise 4] Find the equivalent impedance of the circuit.

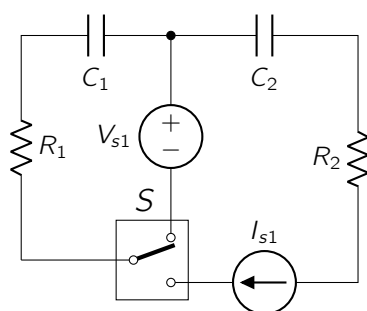


Fig. 1.12: [Exercise 5] Find the equivalent impedance of the circuit.

6. Figure 1.13 shows a voltage divider circuit with a sinusoidal voltage source $V_s(t) = 10 \sin(1000\pi t)$. The impedance $Z_1 = 3 + j4\Omega$. What should the impedance Z_2 be if the voltage across Z_2 has one-fourth the amplitude of V_s and with a phase difference of 45° ?

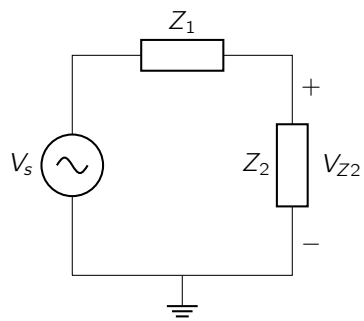


Fig. 1.13: [Exercise 5] Find the equivalent impedance of the circuit.

Chapter 2

Electronics Review

This chapter is a quick review of basic electronics. The previous chapter involved simple, linear passive two terminal elements, which are essential for building most electrical circuits. This chapter will focus on some essential non-linear, two/three terminal electronics components which form the basis for more useful electronics circuits such as amplifiers, oscillators, filters, power supplies, switches, etc. We will review three important electronic components and their circuits in this chapter: *diode*, *bipolar junction transistor (BJT)*, and *metal oxide field effect transistor*.

2.1 Basic semiconductor concepts

Conductivity of a material is proportional to the concentration of free electrons.

$$\begin{aligned}\textbf{Conductor: } n_{cond} &\approx 10^{28} \text{electrons}/\text{m}^3 \\ \textbf{Insulator: } n_{ins} &\approx 10^7 \text{electrons}/\text{m}^3 \\ \textbf{Semiconductor: } n_{ins} &< n_{sem} < n_{cond}\end{aligned}\tag{2.1}$$

2.1.1 Intrinsic semiconductors

Silicon or Germanium have crystalline structure with four covalent bonds between neighbouring atoms. At 0°K semiconductors all covalent bonds are in place tightly holding on

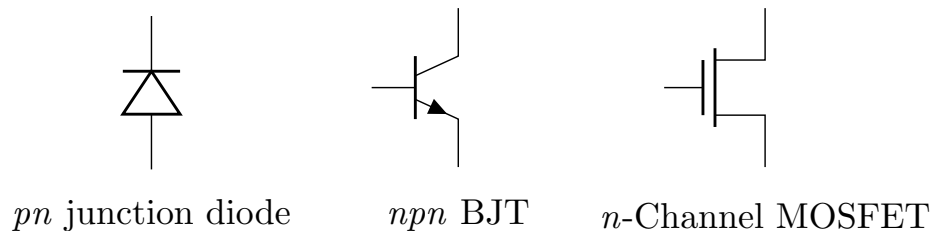


Fig. 2.1: Three most essential electronic components reviewed in this chapter.

to electrons in the bonds. Thus these behave as perfect insulators at low temperatures. With increase in temperature, covalent bonds are ruptured, releasing *free electrons* to roam around in the crystal and become available for conduction when there is an externally applied electric field. The missing electron in the ruptured covalent bond is called a hole, which too act as carrier of electricity. When an electric field is applied, electrons from neighbouring covalent bonds jump into a hole creating a hole in their previous bond. This results in the hole effectively moving the direction opposite to the jumping electrons from the covalent bonds. Thus, unlike conductors, semiconductors can have two types of charge carriers: *electrons* and *holes*. The concentration of free electrons and holes in a semiconductor determines its conductivity, which is controlled by temperature. In intrinsic semiconductors, the concentration of free electrons and holes is equal, i.e. $n_{e^-} = n_{h^+} = n_i$. This concentration is given by,

$$n_i = B T^{3/2} e^{-\frac{E_g}{k_B T}} \quad (2.2)$$

where B is a material constant ($7.3 \times 10^{15} \text{cm}^{-3} \text{K}^{-3/2}$ for Si), E_g is the band gap energy, k_B is the Boltzmann constant, and T is the absolute temperature in Kelvin. The band gap energy is the energy required to break a covalent bond and create a free electron-hole pair. It should be noted that the conductivity of semiconductors increases with temperature, making them suitable for thermometry.

2.1.2 Doped Semiconductors

With intrinsic semiconductors, the conductivity is still quite small at room temperature. Another precise and controlled way to change a semiconductor's conductivity is by adding impurities, called *doping*. The two types of doping are: ***n-type*** and ***p-type*** doping.

In a ***n-type*** semiconductor, a small amount of pentavalent atoms (e.g. Phosphorus, Arsenic) are added to the semiconductor. These atoms have five valence electrons, and when added to the silicon crystal, four of these electrons form covalent bonds with the neighbouring silicon atoms, while the fifth electron is free to roam around in the crystal. This results in an increase in the concentration of free electrons, making the semiconductor more conductive. The concentration of free electrons in an *n-type* semiconductor is given by,

$$n_{e^-} = n_i + N_D \quad n_{h^+} = n_i \quad (2.3)$$

where, N_D is concentration of the doping element.

In a ***p-type*** semiconductor, a small amount of trivalent atoms (e.g. Boron) is added to the silicon crystal structure. This leaves the doped elements with four covalent bonds with one of the bonds lacking an electron or a hole. This hole acts as a charge carrier. The concentration of the holes in a *p-type* semiconductor is given by,

$$n_{h^+} = n_i + N_A \quad n_{e^-} = n_i \quad (2.4)$$

where, N_A is concentration of the doping element in the p -type semiconductor.

The doped semiconductors will have higher conductivity than the intrinsic semiconductors depending on the concentration of the impurities added to the silicon crystal.

2.1.3 Flow of current in semiconductors

There are two types of current flow mechanisms in semiconductors unlike pure conductors – (a) **drift current** and (b) **diffusion current**. Both are important for understanding the operation of the basic electronic components.

Drift current. Drift current is established by an external electric field, for example when a semiconductor is connected to a battery. Holes will accelerate in the direction of the field, while the electrons accelerate in the opposite direction. Bumping into the atoms of the crystal structure, the holes and electrons acquire an average drift velocity given by,

$$\nu_{p-drift} = \mu_p E \quad \& \quad \nu_{n-drift} = -\mu_n E \quad (2.5)$$

This constitutes a drift current through the semiconductor given by the following,

$$I_{drift} \propto q (n_h \mu_p + n_e \mu_n) E \quad (2.6)$$

where, q is the magnitude of electron charge.

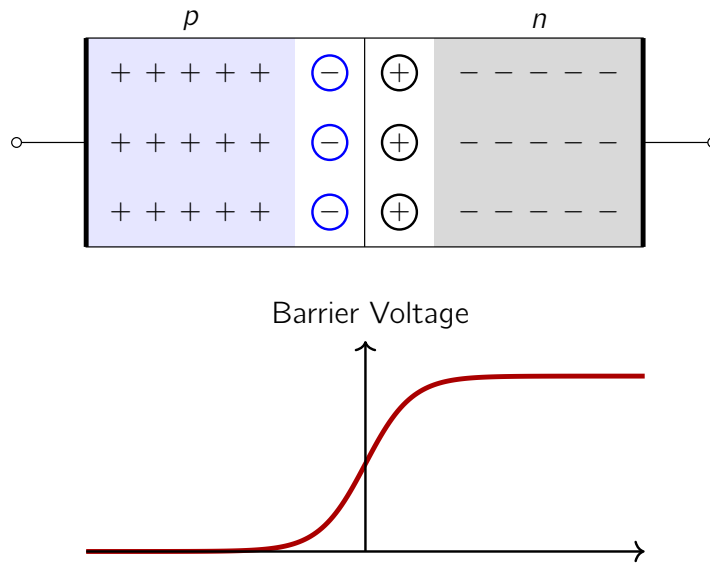
Diffusion current. Diffusion currents result when a concentration gradient exists across a semiconductor.

$$I_{diff} \propto -q D_q \frac{dp(x)}{dx} \quad (2.7)$$

where, $p(x)$ is the concentration of the charge carrier along the x direction. Diffusion currents play an important role in the functioning of the BJT.

2.2 Diode

When a p -type and n -type semiconductors are brought in close contact to each other, there is a big concentration difference between the charge carriers across the interface. A pn -junction diode made by bringing a n -type and p -type semiconductors together as shown in Figure 2.2. Due to the concentration difference, electrons from the n -type semiconductor diffuse into the p -type semiconductor, while holes from the p -type semiconductor will diffuse into the n -type semiconductor. Note that this is the diffusion current I_D , which flows even when the diode is open circuited. This diffusion current sets up the depletion region at the interface, where there are no free charge carriers. This region consists of fixed positive and negative charges, which creates an electric field across the junction. This electric field opposes further diffusion of charge carriers across the junction, and the potential or *barrier voltage* V_0 due to the electric field must

Fig. 2.2: *pn* junction diode

be overcome for the diffusion current to flow across the junction. Note that the diffusion current I_D is due to the majority charge carriers – holes from the *p* side and *electrons* from the *n*-side.

There is a drift current I_S that flows across the junction due to the minority charge carriers – the *electrons* from the *p*-side and the *holes* from the *n*-side. These thermally generated minority charge carriers, will be swept across and transported across the junction due to the electric field, which tends to reduce the barrier voltage V_0 . Thus, when the diode is open circuited, we have an equilibrium $I_D = I_S$.

Barrier voltage of a diode depends on several factors, the amount of doping of the *n* and *p* semiconductors, the temperature, etc. The exact relationship is given by,

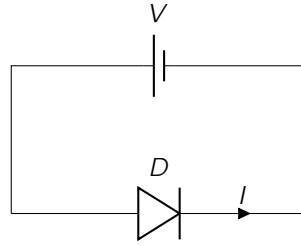
$$V_0 = V_T \ln \left(\frac{N_A N_D}{n_i^2} \right) \quad (2.8)$$

where, V_T is the thermal voltage given by $V_T = \frac{k_B T}{q}$, where k_B is the Boltzmann constant, T is the absolute temperature in Kelvin. For silicon diodes, this is approximately $V_T \approx 26\text{mV}$ at room temperature. The barrier voltage V_0 is typically around 0.7V for silicon diodes, and around 0.3V for germanium diodes.

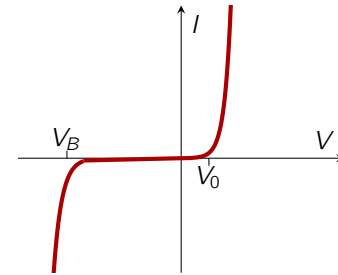
2.2.1 Applying a voltage across a diode

When we connect a voltage source across the diode (Figure 2.3), we can either apply a forward bias or a reverse bias. When the positive terminal of the voltage source is connected to the *p*-side and the negative terminal is connected to the *n*-side, we have a **forward bias**. This reduces the barrier voltage V_0 , and allows the diffusion current I_D to flow across the junction. The diode is said to be *on* in this case, and the current I

Diode with an external voltage source



Diode I-V Characteristic

**Fig. 2.3:** V-I characteristics of a diode

flowing through the diode is called the *forward current*. The voltage-current relationship between the forward current I and the voltage V across the diode is given by,

$$I = I_S \left(e^{\frac{V}{V_T}} - 1 \right) \quad (2.9)$$

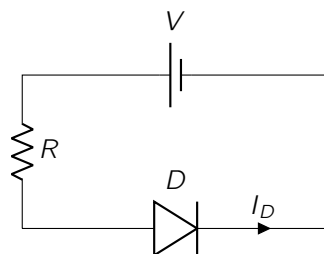
where, I_S is the saturation current, which is the reverse current that flows when the diode is reverse biased, and V_T is the thermal voltage, and q is the magnitude of electron charge.

When the negative terminal of the voltage source is connected to the p -side and the positive terminal is connected to the n -side, we have a **reverse bias**. This increases the barrier voltage V_0 , and prevents the diffusion current I_D from flowing across the junction. The diode is said to be *off* in this case, and only a small reverse saturation current I_S flows through the diode. When the reverse biased voltage is increased beyond a certain threshold, called the *breakdown voltage* V_B , the diode will start conducting in the reverse direction. This is called *avalanche breakdown*, and the diode can be damaged if the current is not limited by a resistor or some other means. The avalanche breakdown is the reason for the sharp rise in the current at the V_B .

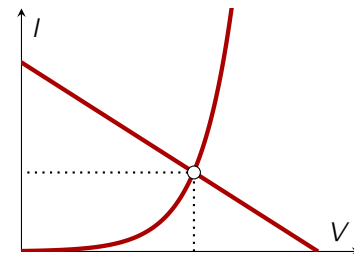
2.2.2 A simple diode circuit

In the circuit shown in Figure 2.3, the moment the forward bias-voltage gets close to 0.6, the current starts to rise dramatically. Slight changes in the voltage will lead to

Diode in series with a resistor and a voltage source



Diode I-V Characteristic

**Fig. 2.4:** Diode voltage and current in a simple diode circuit with a series resistor.

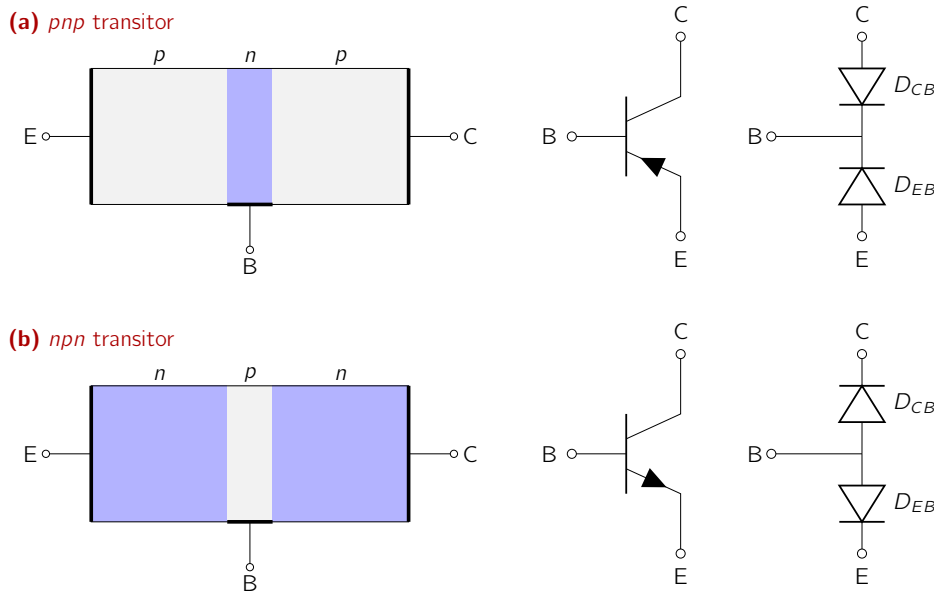


Fig. 2.5: Bipolar junction transistor structure and symbol

large changes in the current if it is not limited by a series resistor. Without a resistor, we can easily burn the diode due to large power dissipation across the diode. Consider a more practical circuit with a voltage source, resistor, and a diode, as shown in Figure 2.4. The voltage across the diode V_D and the current through the diode I_D by writing the Kirchhoff voltage law (KVL) around the loop,

$$V - I_D R - V_D = 0 \quad (2.10)$$

The current I_D and V_D from Eq. 2.9 can be substituted into Eq. 2.10 to solve for V_D and I_D ,

$$V - R I_S \left(e^{\frac{V_D}{V_T}} - 1 \right) - V_D = 0 \quad (2.11)$$

Unfortunately, this is a non-linear solution will need to be solved to find the voltage across the diode V_D and the current through the diode I_D .

Let's assume that $V = 5V$ and $R = 1k\Omega$ in Figure 2.4. A first order approximation can be made by assuming that the diode is *switched off*, i.e. no current is flow through it and see if it leads to a contradiction. If we assume the diode is off and no current is flowing $I_D = 0mA$ through it, then $V_D = 5V$ The diode cannot be off if $V_D \geq 0.7V$. So, the diode must be switched on. We will assume that the diode is on and then find out the current flowing through it and check that it is not close to 0. In the current circuit, have

$$I_D \approx \frac{V - 0.7}{R} = 4.3mA \quad (2.12)$$

I_D is large enough for to flow through the diode when its switched on. Thus, $V_D = 0.7V$ and $I_D = 4.3mA$ is a reasonable first approximation of the diode's voltage and current in the current circuit.

2.3 Bipolar junction transistor (BJT)

The bipolar junction transistor (BJT) is the first of some of the three terminal devices we will come across in this course. There are two types of BJTs – *npn* and *pnp*. BJTs are formed by sandwiching a thin layer of *p*-type semiconductor between two *n*-type semiconductors or vice versa. This is depicted in Figure 2.5. Metal contacts are made with these three semiconductor regions to form the three terminals of the BJT. In an *npn* BJT, one of the *n* regions is the *emitter* which is heavily doped, and the other *n* region is the *collector*. The *p* region is the *base*, which is lightly doped and very thin compared to the widths of the emitter and collector regions. The collector is moderately doped. Three terminal devices like the BJT and MOSFET can be used as an amplifier or a switch, depending on how it is biased. We will only discuss *npn* transistors for the rest of the section; the *pnp* transistor is similar, but with the polarities of the voltages and currents reversed. A BJT is essentially two diodes connected back to back as shown in Figure 2.5

A BJT can be operated in three different modes, depending on the biasing of the terminals. The three modes are: (a) **Cut-off mode**, (b) **Active mode**, and (c) **Saturation mode**. The difference between the three modes are summarized in Table 2.1.

2.3.1 Cut-off mode

In this mode, the emitter-base (EB) and the collector-base (CB) junctions are reverse biased, and no current flows through the transistor.

$$I_E = 0 \quad I_B = 0 \quad I_C = 0 \quad (2.13)$$

2.3.2 Active mode

In this mode, the EB junction is forward biased, while the CB junction is reverse biased. Consider the circuit in Fig. 2.6. When the *npn* BJT's EB junction is forward biased ($V_{BE} \approx 0.7V$), electrons (the majority carrier) from the highly doped emitter region are injected just across the EB junction. The electrons diffuse into the base region, which is lightly doped and very thin. Because $V_C > V_B$, the injected electrons are swept across the base region by the electric field created by the reverse biased CB junction. A small percentage of these injected electrons recombine with the holes in the base region, which contributes to the small base current. The collector current will be proportional to the base current; the current gain is given by the parameter β .

$$I_C = \beta I_B \quad \text{where, } \beta \approx 100 - 1000 \quad (2.14)$$

An equivalent circuit of the BJT in the active model is shown in the following figure (Fig. 2.7). In this circuit, assume that the supply voltages V_{BB} , V_{CC} , and the resistors

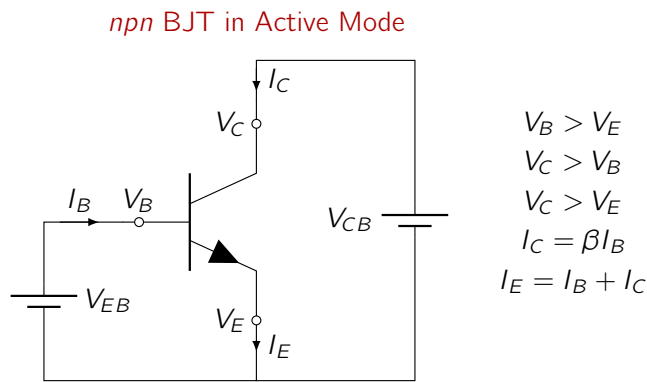


Fig. 2.6: *npn* BJT in active mode.

R_B and R_C are chosen appropriately to place the *npn* transistor in the active mode. The AC source v_s of sufficiently small in amplitude (in *mV*) will result in small changes in I_B , which lead to corresponding change in I_C . The voltage at the collector V_C have a larger amplitude variations riding on a DC voltage, given by the following,

$$V_C = V_{CC} - \beta \frac{R_C}{R_B} (V_{BB} - 0.7) - \beta \frac{R_C}{R_B} v_s \quad (2.15)$$

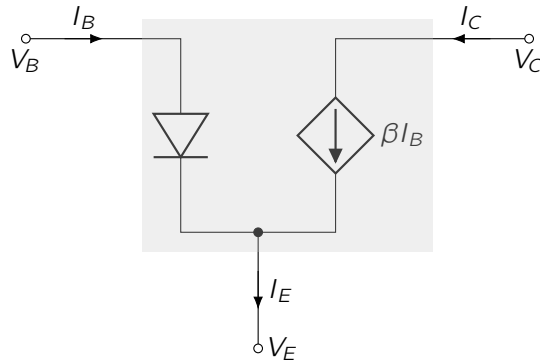
This is easy to derive and is left as an exercise for the reader. The input voltage v_s get amplified by a factor of $\beta \frac{R_C}{R_B}$, which is the voltage gain of this the common emitter amplifier circuit. We will no disucss such circuits any further in the course. This was just to illustrate the utility of the BJT in amplifying small AC signals when operated in the active mode.

Table 2.1: Modes of operation of a bipolar junction transistor (BJT)

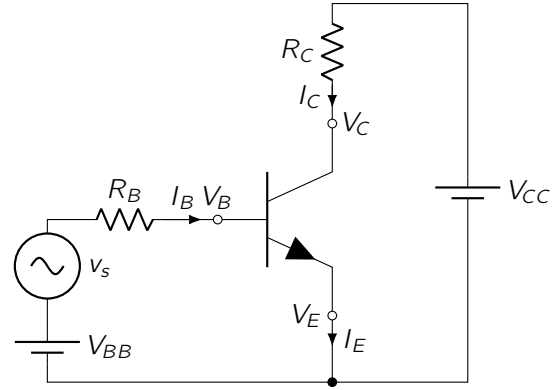
Mode	Emitter-Base Junction	Collector-Base Junction	Operation
Cut-off	Reverse biased	Reverse biased	Transistor is off; no current flows.
Active	Forward biased	Reverse biased	I_C is proportional to I_B
Saturation	Foward biased	Forward biased	I_C and V_{CE} saturate.

2.3.3 Saturation mode

Consider the circuit in Fig. 2.8. If for fixed V_{BB} , we reduce R_B , this will increase I_B , which would result in progressively increasing I_C . However, I_C cannot continue to increase forever. As I_C increases, $V_C = V_{CC} - I_C R_C$ decreases. Once its value falls sufficiently below the base voltage V_B , the collector-base (CB) junction becomes forward biased. This is called the **saturation mode** of operation of the BJT. In this mode, the BJT behaves like a closed switch, and the collector current I_C saturates at a value

(a) *npn* BJT Active Mode Equivalent Circuit

(b) A simple Common-Emitter Amplifier Circuit

**Fig. 2.7:** *npn* BJT active mode equivalent circuit, and a simple common emitter voltage amplifier circuit.

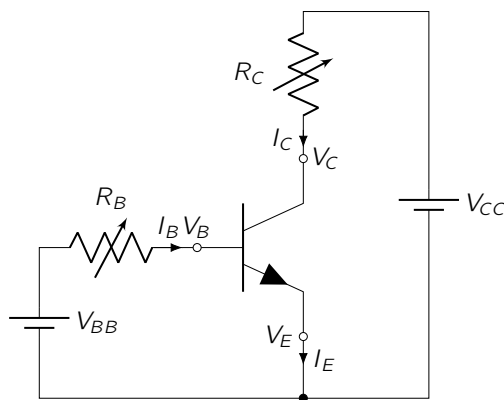
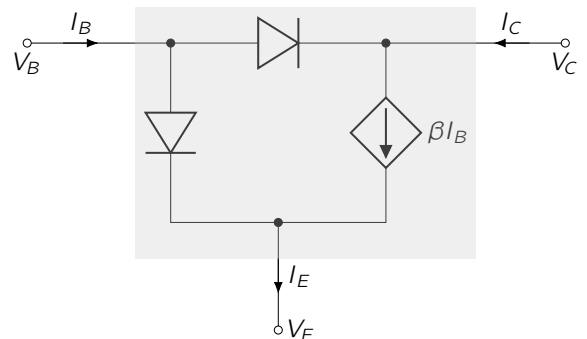
that is independent of the base current I_B . The collector-emitter voltage saturates at $V_{CE(sat)} \approx 0.2V$ for silicon BJTs, making the collector current independent of the base current.

$$I_C \approx \frac{V_{CC} - V_{CE(sat)}}{R_C} \quad \text{and} \quad V_{CE(sat)} \approx 0.2V \quad (2.16)$$

The saturation mode is used in switching applications, where the BJT is used as a switch to turn on or off a load connected to the collector; think of the collector resistor as the load. We will discuss some applications later in the chapter.

2.4 Metal oxide field effect transistor (MOSFET)

While the BJT might be a popular choice for many applications where circuits are made on a breadboard or a PCB. However, when it comes to integrated circuits, the MOSFET is the most popular choice. MOSFETS can be made quite small, are easier

Equivalent circuit of *npn* BJT in Saturation**Fig. 2.8:** *npn* BJT active mode equivalent circuit, and a simple common emitter voltage amplifier circuit.

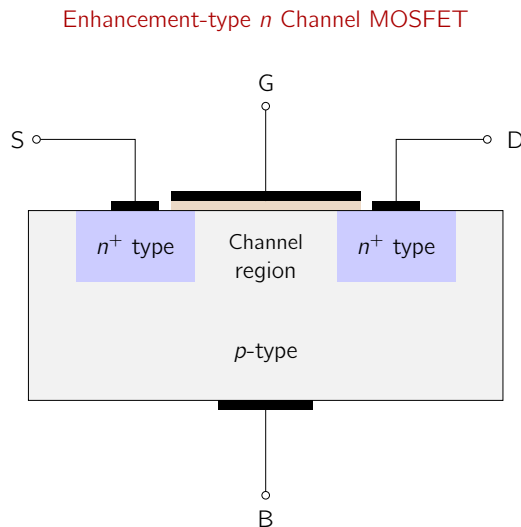
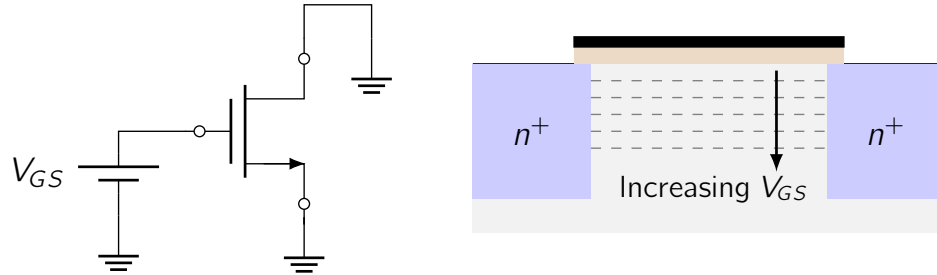


Fig. 2.9: n Channel Enhancement MOSFET.

to fabricate, and consume less power than BJTs. They can be used for both digital and analog integrated circuits. There are four types of MOSFETs, n -channel or p -channel enhancement or *depletion* type MOSFETs. In this section, we will only focus on the enhancement-type n -channel MOSFET to understand the fundamental principle of operation of the MOSFETs.

The physical structure of an enhancement-type n -channel MOSFET is shown in Figure 2.9. The MOSFET in this figure has four terminals - source (S), drain (D), gate (G), and body (B). The source and drain are heavily doped n -type regions, while the channel is lightly doped p -type region. The region between drain and the source is called the *channel region*. The gate is insulated from the channel by a thin layer of oxide, which acts as a dielectric. The gate terminal is used to control the formation of a conductive channel between the source and drain terminals. The body terminal is connected to the source terminal, and is usually shorted to the source in most applications. The MOSFET is a four terminal device, but it can be treated as a three terminal device by shorting the body and source terminals.

The drain is always kept at a higher potential than the source. The connection between the drain to the source is essentially two diodes connected back-to-back, implying that applying a voltage between the drain and source does not result in any current because one of the diodes is reverse biased. However, when a positive voltage is applied to the gate with respect to the body/source, things get interesting as shown in Figure 2.10. Applying a positive gate voltage V_{GS} with respect to the source terminal creates an electric field across the oxide layer, which penetrates into the channel region. This electric field attracts electrons into the channel region right below the gate oxide layer. As the voltage is increased beyond a threshold V_t , a conductive channel is formed between the source and drain terminals, which can now conduct current between the

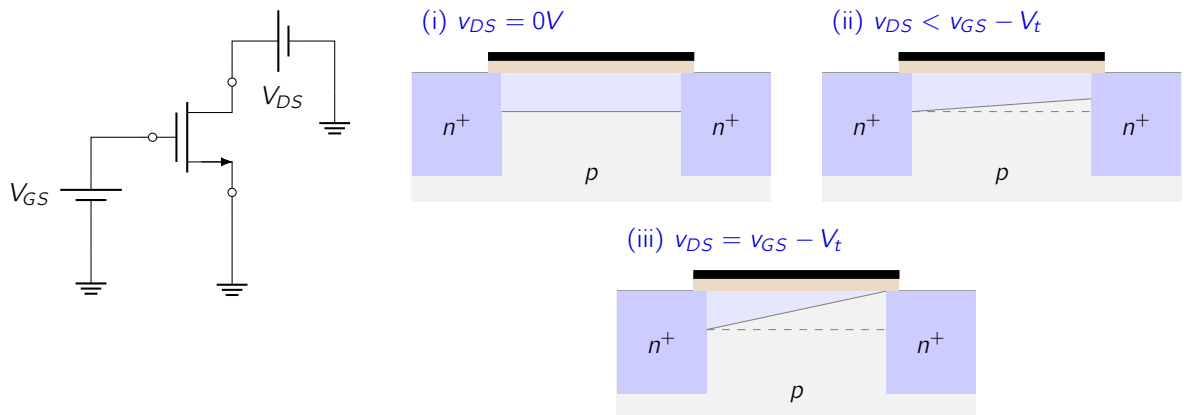
Effect varying V_{GS} on the MOSFET channel**Fig. 2.10:** n -Channel MOSFET V_{GS} - V_{DS} Characteristics.

source and the drain. The threshold voltage is typically between $0.3V - 1.0V$. Any v_{GS} applied beyond V_t is called the overdrive voltage or *effective voltage* v_{OV} .

In Figure 2.10 no current flows because both the source and drain are grounded. When a drain source voltage is applied, as shown in Figure ??, the MOSFET can be turned on by applying a positive gate voltage V_{GS} with respect to the source. The current flowing through the MOSFET is called the *drain current* I_D . The relationship between the drain current i_D , the gate-source voltage V_{GS} , and the drain-source voltage V_{DS} is given by,

$$i_D = k_n \left((v_{GS} - V_t) v_{DS} - \frac{1}{2} v_{DS}^2 \right) \quad (2.17)$$

where, k_n is the process transconductance parameter, and V_t is the threshold voltage. This expression holds only when $V_{GS} > V_t$ and $0 \leq V_{DS} < V_{GS} - V_t$. Something very interesting happens to the channel when a drain-source voltage is applied and its magnitude is increased, as shown in Figure 2.11. The channel becomes trapezoidal in shape as v_{DS} is increased due to different v_{GS} and v_{GD} . This causes the resistance of the channel to increase, tapering off the current flowing through it. When the v_{DS} is increased further, the channel gets pinched off and the current saturates.

Effect varying V_{DS} on the MOSFET channel for a fixed V_{GS} **Fig. 2.11:** n -Channel MOSFET V_{DS} Characteristics.

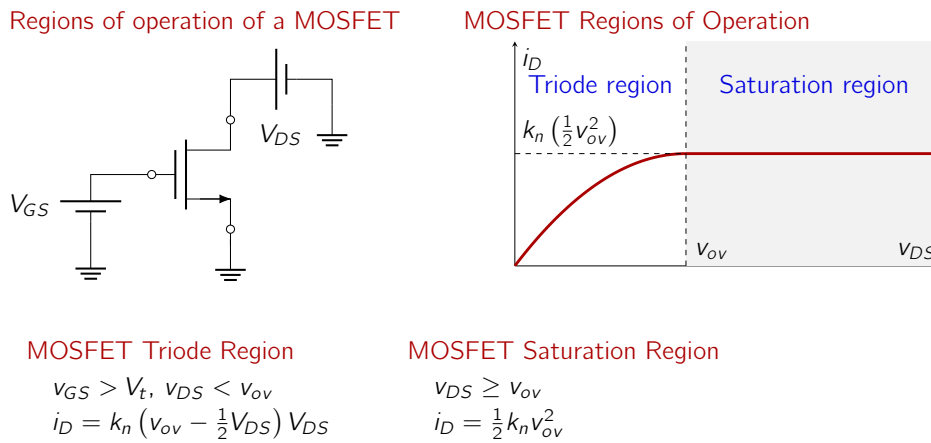


Fig. 2.12: Regions of operation of a MOSFET.

MOSFET regions of operation

The MOSFET can be operated in three different regions, depending on the values of V_{GS} and V_{DS} . The three regions are: (a) **Cut-off region**, (b) **Triode region**, and (c) **Saturation region**. The difference between the three regions are summarized in Table 2.2.

Table 2.2: Modes of operation of a MOSFET

Mode	Gate-Source Voltage	Drain-Source Voltage	Operation
Cut-off	$V_{GS} < V_t$	$V_{DS} = 0$	Transistor is off; no current flows.
Triode	$V_{GS} > V_t$	$V_{DS} < V_{GS} - V_t$	I_D is proportional to V_{DS} .
Saturation	$V_{GS} > V_t$	$V_{DS} \geq V_{GS} - V_t$	I_D is constant.

Figure 2.12 shows i_D versus v_{DS} for a fixed value of $v_{GS} > V_t$. When v_{DS} is increased, the relationship is approximately linear for small values of v_{DS} , while increasing v_{DS} slowly constricts the channel, until it is pinched off, while i_D saturates. When $v_{DS} < V_{GS} - V_t$, the MOSFET is in the *triode region*, where the drain current I_D is proportional to the drain-source voltage V_{DS} . In this region, the MOSFET behaves like a variable resistor. When v_{DS} is increased beyond $V_{GS} - V_t$, the MOSFET enters the *saturation region*, where the drain current I_D becomes constant and independent of the drain-source voltage V_{DS} . In the saturation mode, the MOSFET behaves like a voltage controlled current source, where the drain current I_D is controlled by the gate-source voltage V_{GS} . When $V_{GS} < V_t$, the MOSFET is in the *cut-off region*, where no current flows through the MOSFET. The current versus voltage characteristics of a MOSFET is shown in Figure 2.13. When v_{DS} is large enough to keep the MOSFET in the saturation region, the drain current I_D

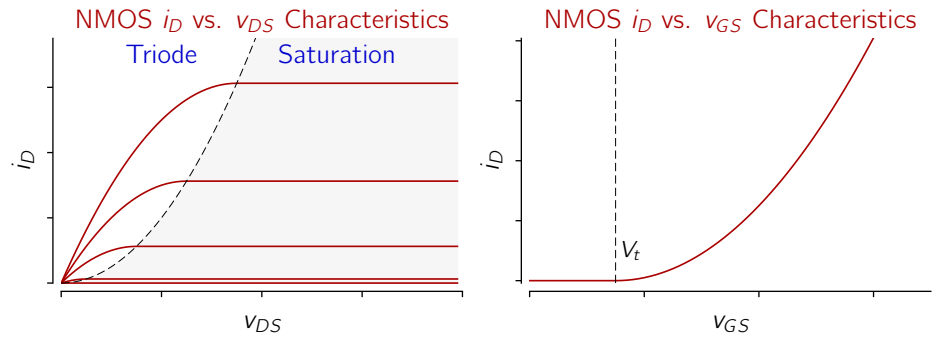


Fig. 2.13: Current vs. voltage characteristics of a MOSFET.

is given by,

$$i_D = \frac{1}{2} k_n (v_{GS} - V_t)^2 \quad (2.18)$$

2.5 Some Diode Circuits

The diode is used in many circuits, such as rectifiers, voltage regulators, clipper, clammers, etc. In this section, we will briefly presents some common and basic diode circuits. The analysis of these circuits are left as an exercise for the reader. The reader is encouraged to simulate these circuits using a circuit simulation software such as LTSpice to gain a better understanding of the operation of these circuits.

2.5.1 Rectifiers

Figure 2.14 shows a half-wave and full-wave rectifier circuit. Assume that the input signal $v_{in}(t)$ is sinusoidal, i.e., $v_{in}(t) = v_o \sin(\omega t)$. The output signal, i.e., the voltage across the load resistance R_L can be computed with different levels of accuracy based on the model we assume for the diode. The output of the half-wave rectifier circuit is

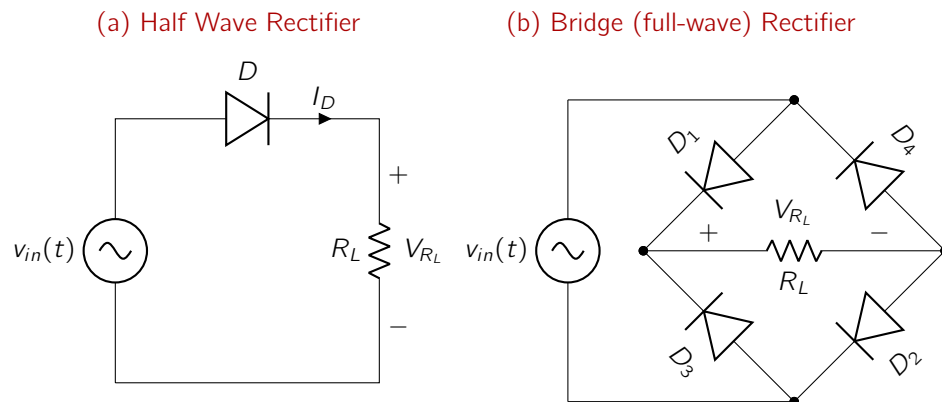


Fig. 2.14: Half-wave and full-wave rectifier circuits.

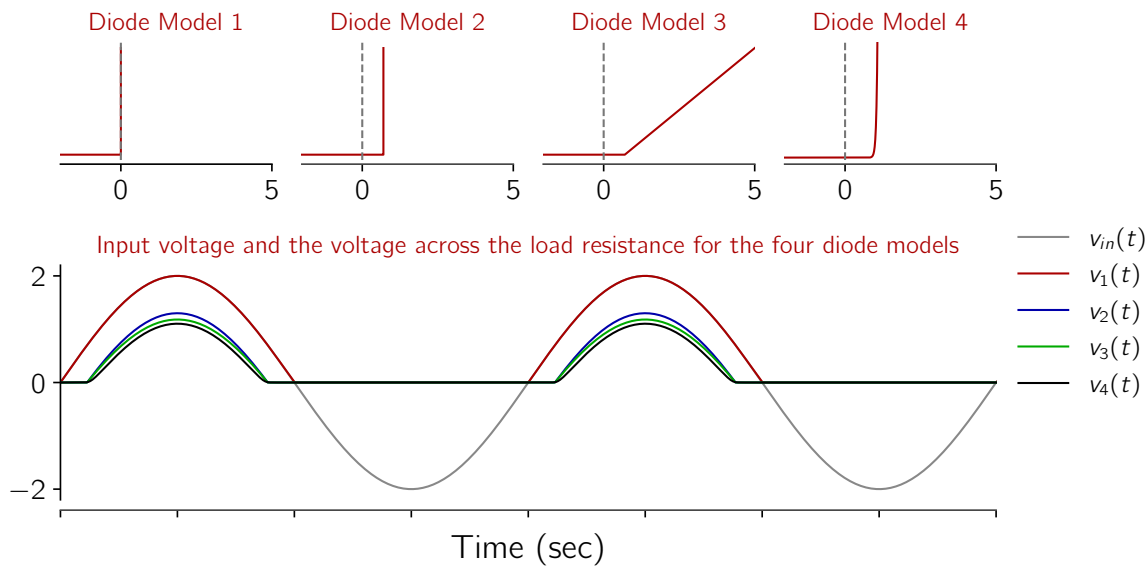


Fig. 2.15: The top row plots show the current-voltage relationship for the assumed diode model. Model 1 is an ideal diode model, while model 4 is a one described in Eq. 2.9. The plot in the bottom row shows the input voltage $v_{in}(t) = 2 \sin(2\pi t)$ in gray, and the voltage across the load resistance $R_L = 100\Omega$ for the different assumed model.

shown in Figure 2.15 for four different models, starting with the simplest idea model (model 1) to a more realistic model (model 4).

Problem 2.1. Can you explain the four output graphs in Figure 2.15 using the four models?

The problem of generating a similar output for the full-wave bridge rectifier is left as an exercise at the end of the chapter.

2.5.2 Voltage regulator

A voltage regulator is a circuit that maintains a constant voltage across its output terminals, even when the input voltage or load current changes. Voltage regulators are widely used in power supplies to provide a stable output voltage. Diodes can be used for building simple voltage regulators. Consider the simple circuit in Figure 2.16. Here, we want a stable 2.8V supply for a circuit and we have a voltage supply V_s that is higher than 2.8V. We can build a stable 2.8V using the arrangement shown in this figure ($R_1 = 1k\Omega$). The variation of V_{out} with variations in $V_{in} \in [5, 15]V$ are shown in the adjacent plots with ($R_L = \infty\Omega$) and without a load resistance (top plot; $R_L = 1k\Omega$). The bottom plot shows the change in the V_{out} as a function of the load resistance $R_L \in [0.1k, 10k]\Omega$. Both of these plots show that that V_{out} is fairly robust to large variations in V_{in} and R_L ; for very low values of V_{in} and R_L , the circuit fails as expected.

A simple voltage regulator

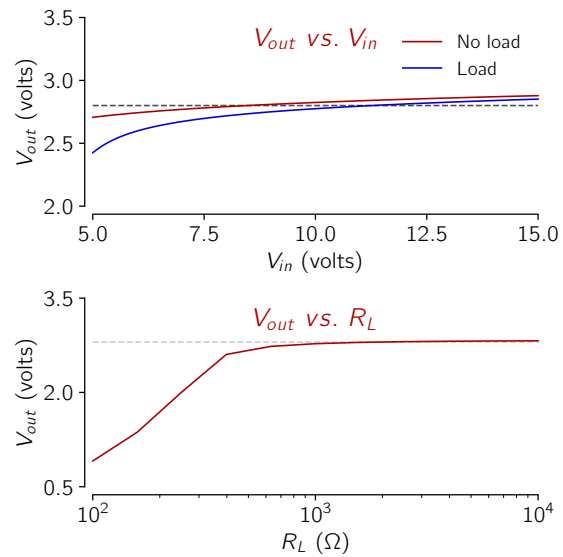
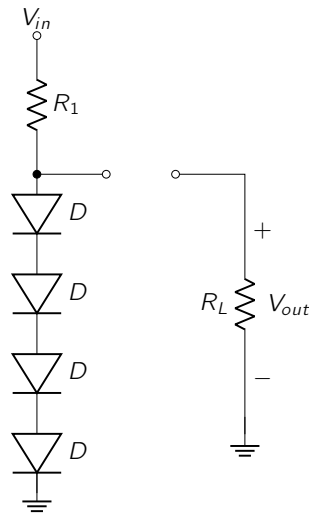
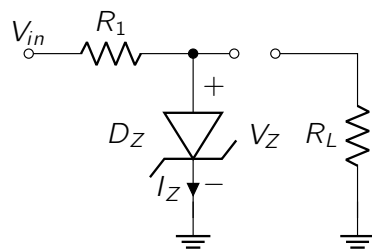


Fig. 2.16: A simple voltage regulator circuit using 4 diodes to obtain a stable voltage of 2.8V. The plots shown in the figure were obtained by simulating this circuit in LTSpice.

Zener diode voltage regulator. A Zener diode is a special type of diode that is designed to operate in the reverse breakdown region, with sufficient power dissipation capability. When a Zener diode is reverse biased, it allows current to flow in the reverse direction when the voltage across it exceeds a certain value, called the *Zener voltage* V_Z . This makes it useful for voltage regulation applications. The Zener diode can be used to build a simple voltage regulator circuit as shown in Figure 2.17. The Zener diode is connected in reverse bias in series with a resistor R_1 and the input voltage V_{in} is slowly increased. The current-voltage relationship of the Zener diode (V_Z vs. I_Z) is shown in the figure. For now assume that the load resistance R_L is not connected to the the circuit. In this case, the Zener breakdown voltage is 5.1V, when $V_Z < 5.1V$, hardly any current flows through the Zener diode and $V_Z = V_{in}$. But once V_{in} get closed to 5.1V, the Zener diode starts to conduct and V_Z clamps to 5.1V.

A simple Zener voltage regulator



Zener diode IV characteristics

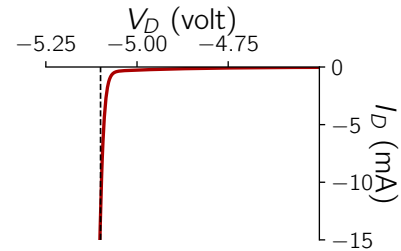


Fig. 2.17: A simple Zener voltage regulator circuit using a Zener diode.

The plot of V_{in} versus V_Z with and without the load resistance is shown in Figure 2.18.

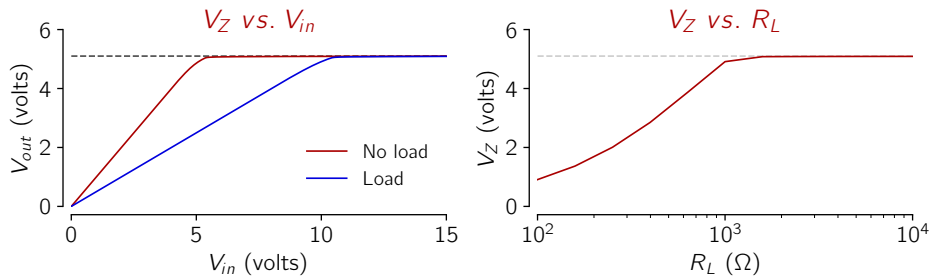


Fig. 2.18: The plot on the left shows the performance of the Zener diode voltage regulator with ($R_L = 1k\Omega$) and without ($R_L = \infty\Omega$) the load. The second plot shows the Zener diode's performance for varying load resistance.

2.6 Some BJT Circuits

These days the most common practical use of the BJT is as a switch. This is what we will focus on; we will also briefly talk about a simple voltage amplifier circuit using the BJT for the sake of demonstration.

2.6.1 BJT as a switch

The BJT can be used as a voltage controlled switch to control the flow of current through different types of loads such as resistors, inductors, LEDs, etc. To act as a switch the BJT is operated in the cut-off or the saturation modes, to turn it off or on, respectively, as shown in Figure 2.19. In the ideal case, it behave alike an open or closed switch between the collector and emitter terminals.

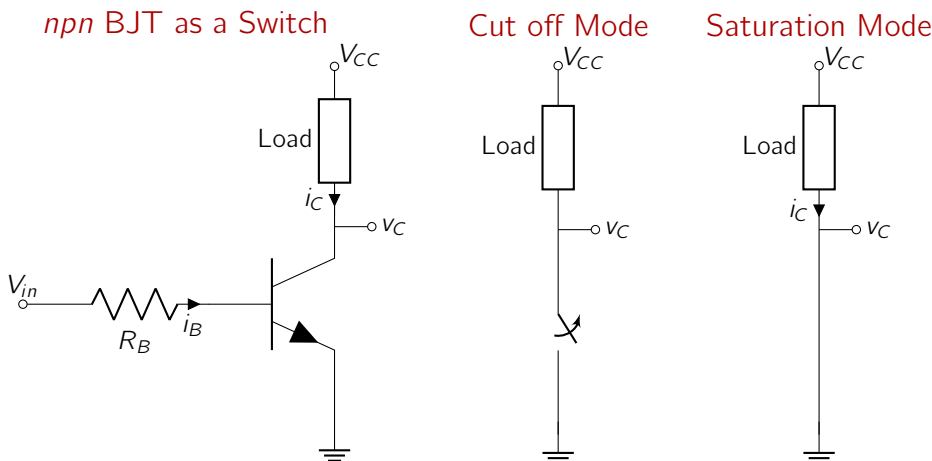


Fig. 2.19: Circuit of a typical *npn* BJT based switch.

Consider a resistive load in Figure 2.19 with load resistance R_L . In the cut-off mode, when the base-emitter junction is reverse biased, i.e. $V_{BE} < 0.7V$, the BJT is off. One way to achieve this is to set $V_{in} = 0V$ for a *npn* BJT. In the cut-off mode, $i_B = 0$, which

means that $i_C = 0$ and $i_E = 0$. The voltage at the collector of the BJT will be equal to the supply voltage V_{CC} , i.e. $v_C = V_{CC} - i_C R_L = V_{CC}$. In saturation mode, the base-emitter and collector-base junctions are both forward biased. In a real *nnp* BJT circuit, the collector-emitter voltage is $v_{CE} \approx 0.2V$. Thus, we have, $i_C = \frac{V_{CC} - v_{CE}}{R_L} = \frac{V_{CC} - 0.2}{R_L}$. For this collector current, the base current must be at least $i_B = \frac{i_C}{\beta}$. Thus, we have $i_B = \frac{V_{in} - 0.7}{R_B}$. This can be achieved by choosing a sufficiently large V_{in} or sufficiently small R_B .

For non-resistive loads, the analysis is similar, but we must consider the load's voltage-current characteristics and appropriately design the circuit. Let's look at a couple of examples.

nnp BJT to switch on/off an LED

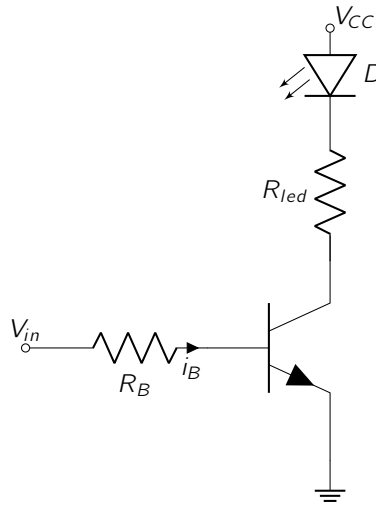


Fig. 2.20: BJT LED control circuit

Example 2.1 (Switch on an LED). We have an LED with a forward voltage drop of $2.0V$ and a forward current of $20mA$. We want to control the LED through a transistor, such that when $3.3V$ is applied to the base of the transistor, the LED switches ON, else it is OFF (Figure 2.20). Assume that the *nnp* transistor has a current gain in the range $\beta = 100 - 200$. We need to choose R_{led} and R_B for this circuit.

We need i_C must be $20mA$ when the LED is turned on, and it has a drop of $2V$. When the transistor is in saturation, we have,

$$V_{CC} - 2V - i_C R_L = 0 \implies R_L = \frac{V_{CC} - 2V}{i_C} = \frac{3.3V - 2V}{20mA} = 65\Omega$$

To get an collector current of $20mA$, we need a base current of $i_B = \frac{i_C}{\beta} = \frac{20mA}{\beta}$. Since we do not have the exact value of β , we will assume the worst case, i.e. lowest possible value for β , which in this case is 100. Thus, we have $i_B = \frac{20mA}{100} = 0.2mA$.

The base current is given by $i_B = \frac{V_{in}-0.7}{R_B}$, where $V_{in} = 3.3V$. Thus, we have, $R_B = \frac{3.3V-0.7V}{i_B} = \frac{2.6V}{0.2mA} = 13k\Omega$.

npn BJT to switch on/off a solenoid

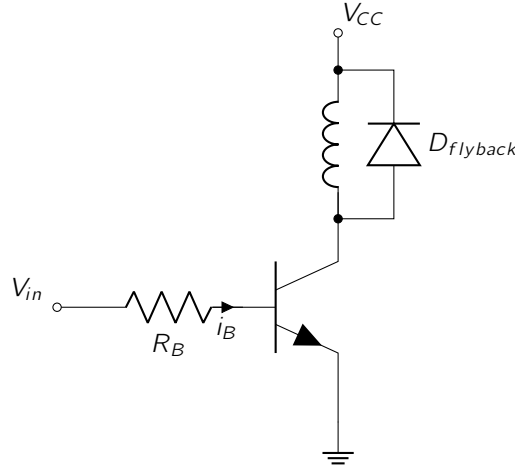


Fig. 2.21: BJT solenoid control circuit

Example 2.2 (Switch on/off a solenoid). Consider a tiny solenoid that switches on at 5V across it coil and a current of 0.8A. We wish to use a *npn* BJT to switch the the solenoid ON and OFF by controlling the BJT's base voltage. The circuit for switching the solenoid ON and OFF is shown in Figure 2.21. We will assume that the BJT has a current gain of $\beta = 500$. We need to choose R_B for this circuit assuming that V_{in} is 0V or 3.3V.

The solenoid has a voltage drop of 5V and a current of 0.8A. When the transistor is in saturation the voltage across the solenoid will be 4.8V, i.e. close to 5V (the solenoid should still work as per the specification). To ensure 0.8A flows through the solenoid, we need to ensure that the base current must be at least $i_B = \frac{i_C}{\beta} = \frac{0.8A}{500} = 1.6mA$. The base current is given by $i_B = \frac{V_{in}-0.7}{R_B}$, where $V_{in} = 3.3V$. Thus, we have, $R_B = \frac{3.3V-0.7V}{1.6mA} = 1.625k\Omega$.

You are probably wondering about the purpose of the $D_{flyback}$ diode. Let's assume that we apply a square pulse as the input V_{in} to the base.

$$V_{in} = \begin{cases} 5V & 0.25s \leq t < 1.25s \\ 0V & \text{otherwise} \end{cases}$$

When the BJT is suddenly turned OFF at time $t = 1.25s$ with 0.8A flowing through it, the current is forced to stop suddenly, which causes a large voltage spike across the solenoid, which will result in a large voltage spike across the BJT, which can

damage it. The flyback diode $D_{flyback}$ is used to suppress this voltage spike by allowing the current to flow through it, and decay down slowly (relatively), thus avoiding a voltage spike across the inductor and the BJT. You will see this in action in the lab exercise, when you are deliberately asked to remove the flyback diode and see what happens to the BJT.

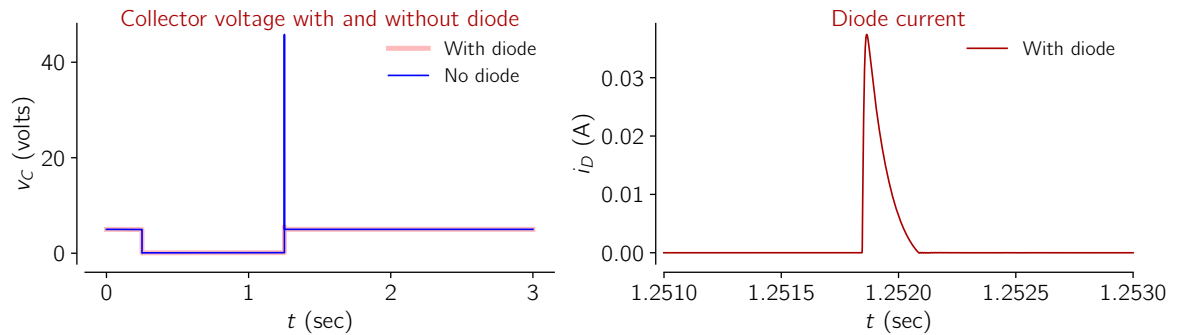


Fig. 2.22: Effect of the flyback diode on the solenoid switching circuit. The left plot shows the collector voltage for a input pulse voltage V_{in} of 5V amplitude over $0.25s \leq t \leq 1.25s$. Without the flyback diode we find a large voltage spike when the collector current is suddenly reduced to zero. With the flyback diode suppressed the voltage spike by allowing the current to flow through it when the BJT is turned off. The right plot shows the transient diode current when the BJT is switched off at $t = 1.25sec$.

2.7 Some MOSFET Circuits

