Mechatronics for Rehabilitation Engineering: Course Notes

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Chapter 2

Electronics Review

This chapter is a quick review of basic circuit theory and electronics. Prior knowledge of these topics is assumed, along with basic understanding of linear time invariant systems - Fourier/Laplace transforms, linear constant coefficient differential equations, transfer functions, and frequency response. The reader is encouraged to review the material in this chapter before proceeding to the next chapter.

2.1 Review of Basic Circuit Theory

The interaction of electromagnetic fields with matters is the basis of all electrical and electronic devices. These interactions are often described, analyzed and synthesized through the abstractions of electrical circuit theory. The following are the most basic circuit two-terminal elements we will need for now. We will introduce new ones as and when they are required. Each circuit element has a unique voltage-current relationship, and it is *important* that you know these by heart.

Independent Voltage source. An idea voltage source provides a fixed voltage V between its two terminals, and can provide any amount of current. Notice that the voltage V can be fixed or time varying. For example, for a DC vortlage source with V=5V, the voltage across the two terminals will be 5V for all time. But for a time varying AC source, $V=5\sin(100\pi t)$, the voltage across its terminal will vary with time. We will often drop the adjective "independent"

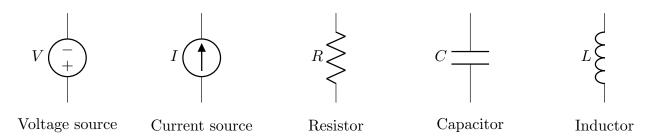


Figure 2.1: Basic circuit elements: voltage source, current source, resistor, capacitor, and inductor.

when we are sure that the context is clear. We will look at dependent sources later, and we will always use the adjective "dependent" to refer to them.

Independent Current source. An ideal cuyrrent source provide a fixed amount of current to flow through its terminals (out through one and in through the other), irrespective of the voltage across its terminals. Current sources can also be time-varying.

Resistor. A passive element where the current i_R flowing through the element is proportional to the voltage v_R across its terminals.

$$V_R \propto i_R$$

In the case of linear resistors, the proportionality factor is constant, resulting in Ohm's law,

$$v_R = R i_R \tag{2.1}$$

The units of R are $V.A^{-1}$ or $Omhs(\Omega)$. R is in general positive. The power absorbed by a resistor is given by the product of the voltage across it and the current flowing through the resistor,

$$P = v_R i_R = i_R^2 R = \frac{v_R^2}{R}$$
 (2.2)

This power is dissipated as heat by the resistor. Note that the power absorbed by a resistor is always positive, since R is positive.

We will later see non-linear resistors, where the resistance varies as a function of the applied votlage, temperature and other factors.

Capacitor. A capacitor is another passive element with the following voltage current relationship.

$$i_C = C \frac{dv_C}{dt} \tag{2.3}$$

The current i_C through the capacitor is proportional to the rate of change of voltage across its terminals v_C . The proportionality factor is called the capacitance C, and has units of F (Farads) or $C \cdot V^{-1}$. The voltage across the capacitor is given by at any given time is proportional to the integral of the current flowing through itor the charge stored in the capacitor. The voltage across the capacitor is given by

$$q = C v_C \tag{2.4}$$

The instantaneous power absorbed by the capacitor is given by,

$$P = v_C i_C = C v_C \frac{dv_C}{dt}$$
 (2.5)

The power absorbed by the capacitor can be positive or negative, depending on the direction of current flow. If the current is flowing into the capacitor, then the voltage across it is increasing, and the power absorbed is positive. If the current is flowing out of the capacitor, then the voltage across it is decreasing, and the power absorbed is negative. A capacitor

stores energy in the electric field between its plates. The energy stored in a capacitance at any given time depends on the charge stored in it, and is given by,

$$E = \frac{1}{2}C v_C^2 \tag{2.6}$$

Inductor. An inductor is another passive element with the following voltage current relationship.

$$v_L = L \frac{di_L}{dt} \tag{2.7}$$

The voltage v_L across the inductor is proportional to the rate of change of current i_L flowing through it. The proportionality factor is called the inductance L, and has units of H (Henries) or $V \cdot s \cdot A^{-1}$. The current through the inductor is given by at any given time is proportional to the integral of the voltage across it. The current through the inductor is given by

$$i_L = \frac{1}{L} \int v_L \, dt \tag{2.8}$$

The instantaneous power absorbed by the inductor is given by,

$$P = v_L i_L = L i_L \frac{di_L}{dt} \tag{2.9}$$

The power absorbed by the inductor can be positive or negative, depending on the direction of current flow. If the current is flowing into the inductor, then the voltage across it is increasing, and the power absorbed is positive. If the current is flowing out of the inductor, then the voltage across it is decreasing, and the power absorbed is negative. An inductor stores energy in the magnetic field around it. The energy stored in an inductor at any given time depends on the current flowing through it, and is given by,

$$E = \frac{1}{2}L i_L^2 \tag{2.10}$$

2.1.1 Kirchoff's Laws

The five elements alone are not that interesting. But interesting things can be done by connecting these elements together in different ways to form an electrical circuit. The elements are connected together by wires, which are assumed to be perfect conductors, i.e. zero resistance. Consider the following circuit (Figure 2.1),

How do we find out the voltages and currents in the circuit? Kirchoff's laws can be used for analysing such circuits, which are based on the conservation of charge and energy. The circuit in Figure 2.2 is a simple electrical circuit with a voltage source, a current source, and a bunch of resistors. The voltage source provides a fixed voltage V_s between its two terminals, and the current source provides a fixed amount of current I_s to flow through its terminals. The voltages and currents in the rest of the elements will be determined by Kirchoff's laws with the constraints imposed by the voltage and current sources. The two laws are:

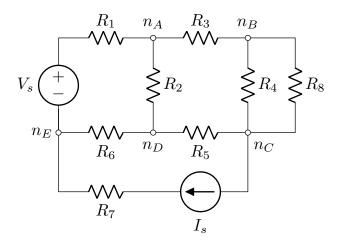


Figure 2.2: A simple electrical circuit with a voltage sourse, a current source, and a bunch of resistors.

1. **Kirchoff's current law (KCL):** The sum of the currents entering a node is equal to the sum of the currents leaving the node. A *node* is a point at which two or more circuit elements are connected together. In Figure 2.2, n_A , n_B , n_C and n_D are examples of nodes where three elements are connected together. There are three other nodes in the circuit, can you identify them?

The sum of the currents at a node is equal to zero. This is based on the conservation of charge, and can be expressed mathematically as:

$$\sum_{i=1}^{n} i_i = 0 (2.11)$$

where i_i is the current flowing into or out of the node, and n is the number of elements connected to the node. The current flowing into a node is positive, and the current flowing out of a node is negative.

2. **Kirchoff's voltage law (KVL):** The sum of the voltages around a closed loop in a circuit is equal to zero. A *closed loop* is a path in the circuit that starts and ends at the same node, and does not cross itself. In Figure 2.2, the path starting from n_A , to n_D , to n_B , and back to n_A is a closed path. This path invludes the resistors R_3 , R_4 , R_5 , and R_2 .

This is based on the conservation of energy, and can be expressed mathematically as:

$$\sum_{i=1}^{n} v_i = 0 \tag{2.12}$$

where v_i is the voltage across each element in the loop, and n is the number of elements in the loop.

The voltage across an element is positive if the current is flowing into the positive terminal of the element, and negative if the current is flowing out of the positive terminal of the element.

Note that the two laws apply for any type of circuit element used in the circuits, independnet or dependent voltage/current sources, resistors, capacitors, inductors, other two, three or four terminal elements.

2.1.2 Series and Parallel Connections

Two elements that share the same voltage across them between a given pair of notes are said to be **parallel** to each other. In Figure 2.2, R_4 and R_8 are parallel to each other. In a single loop, two elements that share the same current are said to be in **series** with each other. In Figure 2.2, V_s and R_1 are in series, I_s and R_7 are in series.

Resistors in series and parallel

When n resistors $R_1, R_2, \dots R_n \ge 0$ are in series, these can be combined to an equivalent resistor with resistance R_{eq} given by the following,

$$R_{eq} = \sum_{i=1}^{n} R_i \implies R_{eq} \ge \max_{1 \le i \le n} R_i. \tag{2.13}$$

Note that in a series connection the equivalent resistance is at least as large as the largest value of R_1 to R_n .

When n resistors $R_1, R_2, \dots R_n \ge 0$ are in parallel, these can be combined to an equivalent resistor with resistance R_{eq} given by the following,

$$\frac{1}{R_{eq}} = \sum_{i=1}^{n} \frac{1}{R_i} \implies R_{eq} = \frac{R_1 R_2 \cdots R_n}{R_1 + R_2 + \cdots + R_n} \implies R_{eq} \le \min_{1 \le i \le n} R_i$$
 (2.14)

Note that in a parallel connection, the equivalent resistance cannot be larger than the smallest value of R_1 to R_n .

Capacitors in series and parallel

Series connection of *n* capacitors $C_1, C_2, \dots C_n \ge 0$

$$C_{eq} = \frac{C_1 C_2 \cdots C_n}{C_1 + C_2 + \cdots + C_n} \tag{2.15}$$

Parallel connection of *n* capacitors $C_1, C_2, \cdots C_n \ge 0$

$$C_{eq} = C_1 + C_2 + \dots + C_n \tag{2.16}$$

Inductors in series and parallel

Series connection of *n* capacitors $L_1, L_2, \dots L_n \ge 0$

$$L_{eq} = L_1 + L_2 + \dots + L_n \tag{2.17}$$

Parallel connection of *n* capacitors $L_1, L_2, \dots L_n \geq 0$

$$L_{eq} = \frac{L_1 L_2 \cdots L_n}{L_1 + L_2 + \cdots + L_n} \tag{2.18}$$

Its left as an exercise for you to verify these expressions.

What does the equivalent resistance actually mean? The equilvant resistor with resistance R_{eq} has the same voltage-current relationship as the individual elements in series or parallel connection. We can replace the series or parallel connection of the individual resistors R_1 to R_n by a single resistor with value R_{eq} without changing the voltage current relationships in the circuit. The same argument applies for equivalent capacitors and inductors.

Voltage sources in series and parallel

Series connection of n voltage sources $V_1, V_2, \dots V_n$ will result in an equivalent voltage source V_{eq} given by

$$V_{eq} = V_1 + V_2 + \dots + V_n \tag{2.19}$$

Voltage soruces should not be connected in parallel, as this will result in a short circuit. Ideally, an infinite currelt will flow through the connection, because the volaage sources force a potential difference between the two end of the wires and it has zero resistance. Parallel connections are allowed only when the two sources have the same voltage and polarity.

Current sources in series and parallel

Parallel connection of n current sources $I_1, I_2, \dots I_n$ will result in an equivalent current source I_{eq} given by

$$I_{eq} = I_1 + I_2 + \dots + I_n \tag{2.20}$$

Current sources should not be connected in series; series connections are allowed only when the two sources have the same current and polarity.

2.1.3 Superposition Principle

Linear approximations of circuits are often employed as first order approximations when analysing circuits. A linear circuit is one that consists of linear passive elements, independent sources and linear dependent sources. Linear circuits follow the superposition principle, which states that the response of a linear circuit to a linear combination of inputs is equal to the corresponding linear combination of the responses to each input applied separately.

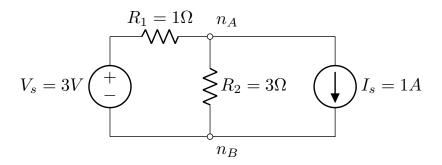


Figure 2.3: A simple circuit with two sourcers

Solving the following circuit (Figure 2.3) should make this concept clear. For the circuit in Figure 2.3, perform the following calculation and compare your results.

Step 1. Solve for the voltage across and the current through the resistors R_1 and R_2 ; we will refer to these are v_{R1} , i_{R1} and v_{R2} , i_{R2} , respectively. You should use both Kirchoff's current and voltages laws to compute these variables.

Let's now consider of the sources in the circuit seperately. This would mean making the source values "zero". This corresponds to two different operations on the circuit. Zeroing a voltage source corresponds to replacing it with a wire (a short circuit), while zeroing a current source corresponds to simple removing the current source (an open circuit).

Step 2. Zero the voltage source V_s and compute the voltages and currents associated with the two resistors. We will refer to these as $v_{R1,V_s=0}$, $v_{R2,V_s=0}$, $i_{R1,V_s=0}$, and $i_{R2,V_s=0}$.

Step 3. Zero the current source I_s and compute the voltages and currents associated with the two resistors. We will refer to these as $v_{R1,I_s=0}$, $v_{R2,I_s=0}$, $i_{R1,I_s=0}$, and $i_{R2,I_s=0}$.

Step 4. For a linear circuit, shown in Figure 2.3, the following will always be true.

$$v_{R1} = v_{R1,V_s=0} + v_{R1,I_s=0}$$

$$v_{R2} = v_{R2,V_s=0} + v_{R2,I_s=0}$$

$$i_{R1} = i_{R1,i_s=0} + i_{R1,I_s=0}$$

$$i_{R2} = i_{R2,i_s=0} + i_{R2,I_s=0}$$
(2.21)

Let's assume that I am only interested in i_{R2} . Can I use $i_{R2,V_s=0}$ and $i_{R2,I_s=0}$ to compute the current i_{R2} if $V_s=1V$ and $I_s=-2A$?

2.1.4 Practical Voltage and Current Sources

The independent voltage and current sources we have discussed so far are "ideal" sources. Practical or real sources do not behave like them - a battery cannot provide any amount of current for a load without any changes to the voltage across its terminals.

A good model of practical voltage source is an ideal voltage source V_s in series with a internal, source or output resistor R_s . And for a practical current source, it is an ideal current source I_s in parallel with a resistor R_s . The voltage across the terminal of a voltage source

(a) Ideal voltage source with a load resistor

(b) Practical voltage source with a load resistor

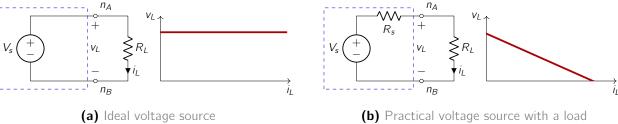


Figure 2.4: Comparison of the voltage-current relationship of an ideal and a practical voltage source.

as function of the current draw from it is depicted for an ideal and practical voltage source in Figure 2.4. For the ideal source (Figure 2.4a), the voltage v_L is independent of the current i_L drawn from it. For the practical voltage source (Figure 2.4b), the voltage across the terminals is a function of the current drawn from it,

$$v_L = V_s - R_s i_L \tag{2.22}$$

When $R_L = \infty$ ($i_L = 0$), the voltage across the terminals is equal to the voltage of the source V_s . This is maximum voltage the practical source can provide. This is also know as the *open circuit voltage* v_{oc} of the source. When $R_L = 0$ ($v_L = 0$), the current drawn from the source $i_L = \frac{V_s}{R_s}$. This is the maximum current the voltage source can provide. This is also known as the *short circuit current* i_{sc} of the source.

Problem 2.1. Plot the voltage-current relationship of a practical current source with $I_s = 2A$ and $R_s = 10\Omega$. What are v_{oc} and i_{sc} ?

2.1.5 Thevenin's and Norton's Theorems

Thevenin's and Norton's theorems are two important theorems in circuit theory that allow us to simplify complex circuits into simpler equivalent circuits. These theorems are based on the superposition principle, and can be used to analyze linear circuits with independent and dependent sources.

Thevenin's Theorem. Thevenin's theorem states that any linear circuit with independent and dependent sources can be replaced by an equivalent circuit with a single voltage source V_{th} in series with a resistor R_{th} , connected to the load resistor R_{L} .

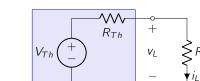
Norton's Theorem. Norton's theorem states that any linear circuit with independent and dependent sources can be replaced by an equivalent circuit with a single current source I_N in parallel with a resistor R_N , connected to the load resistor R_L .

Computing the Thevenin and Norton equivalent circuits. The Thevenin and Norton equivalent circuits can be computed using the following steps:

1. Remove the load resistor R_L from the circuit.

(a) Complex circuit with a load resistor

Complex Circuit



(b) Thevenin equivalent circuit

(c) Norton equivalent circuit

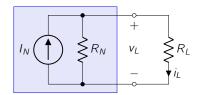


Figure 2.5: Thevenin's and Norton's circuits for a complex linear circuit. The Thevenin and Norton equivalent circuits have the same voltage-current relationship for any load.

- 2. Compute the open circuit voltage v_{oc} across the terminals of the load resistor. This is the Thevenin voltage V_{th} . This can be using the superposition principle, where we zero all the independent sources except one and compute the open circuit voltage. The overall open circuit voltage is the sum of the open circuit voltages when all other sources are zeroed.
- 3. Compute the short circuit current i_{sc} through the terminals of the load resistor. This is the Norton current I_N . This too can be calculated using the superposition principle.
- 4. Compute the Thevenin resistance R_{th} by zeroing all independent sources in the circuit and computing the equivalent resistance seen from the terminals of the load resistor (without the load resistor). This is also equal to the Norton resistance R_N .
- 5. The Thevenin and Norton equivalent circuits are then given by:

$$V_{th} = v_{oc}$$

$$I_N = i_{sc}$$

$$R_{th} = R_N$$
(2.23)

6. The load resistor R_L can be connected to either the Thevenin or Norton equivalent circuit, and the voltage and current across it can be computed using the voltage-current relationship of the equivalent circuit.

Problem 2.2. Compute the Thevenin and Norton equivalent circuits for the circuit shown in Figur 2.2 assuming the following as the load resistor: (a) R_8 ; (b) R_2 ; and (c) R_7 .

2.1.6 Maximum Power Transfer Theorem

The maximum power transfer theorem states that the maximum power is transferred to the load resistor R_L when the load resistance is equal to the Thevenin resistance R_{th} of the circuit.

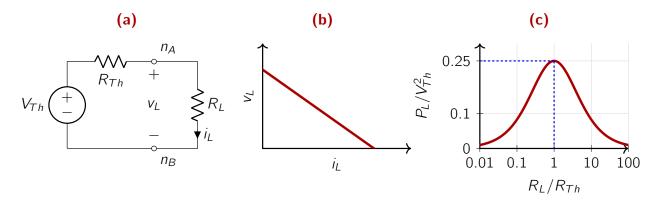


Figure 2.6: Maximum power transfer theorem.

Consider the following circuit (Figure 2.6). The power absorbed by the load resistor R_L is given by,

$$P_{R_L} = \frac{R_L}{(R_{th} + R_L)^2} V_{th}^2 \tag{2.24}$$

Its easy to check that the optimal value of R_L that maximizes the power absorbed by the load resistor is given by,

$$R_{L}^{*} = \arg\max_{R_{L}} \frac{R_{L}}{(R_{th} + R_{L})^{2}} V_{th}^{2} = R_{Th}$$

$$P_{L}^{*} = \max_{R_{L}} \frac{R_{L}}{(R_{th} + R_{L})^{2}} V_{th}^{2} = \frac{1}{4} \left(\frac{V_{Th}^{2}}{R_{Th}} \right)$$
(2.25)

Problem 2.3. Prove the statements in Eq. 2.25. (*Hint:* Use the first order condition for maximization of a continuous function.)

2.1.7 RC, RL, and RLC Circuits

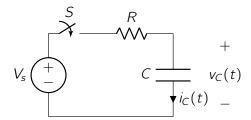
Unlike resistors, that neither remember the past values of its votlage or current, capacitors and indutors retain memory of the past current and past voltage, respectively. This can be easily seen from their respective voltage-current relationships given by Eq. 2.3 and Eq. 2.7, respectively.

Capacitor:

$$i_C = C \frac{dv_C}{dt} \implies v_C(t) = \frac{1}{C} \int_0^t i_C(\tau) d\tau + v_C(0)$$
 (2.26)

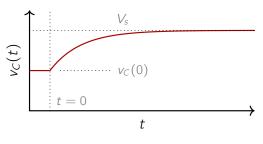
The instantaneous voltage across a capacitor contains some information about the past history the current that has flown through the capacitor. This voltage $v_C(t)$ is determined by the charge on the capacitor at time t. Note, that the voltage across the capacitor cannot change instantaneously, as this would require an infinite current to flow through the capacitor. Theoretically, however, an impulse (Dirac delta function) current applied to the capacitor can produce instantaneous change in the capacitor's voltage.

(a) A Simple RC Circuit



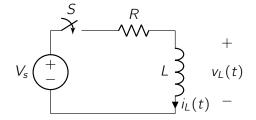
Assuming the switch S is closed at t=0, the voltage across the capacitor $v_C(t)$ is given by,

$$v_C(t) = V_s \cdot \left(1 - e^{-\frac{t}{RC}}\right) + v_C(0)e^{-\frac{t}{RC}}, \ t \ge 0$$



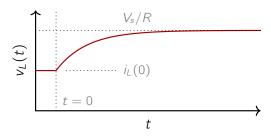
(a) Ideal voltage source

(b) A Simple RL Circuit



Assuming the switch S is closed at t = 0, the current through the inductor $i_L(t)$ is given by,

$$i_L(t) = \frac{V_s}{R} \cdot \left(1 - e^{-\frac{t}{L/R}}\right) + i_L(0)e^{-\frac{t}{L/R}}, \ t \ge 0$$



(b) Practical voltage source with a load

Figure 2.7: Simple RC and RL circuits and their transient responses.

Inductor:

$$v_{L} = L \frac{di_{L}}{dt} \implies i_{L}(t) = \frac{1}{L} \int_{0}^{t} v_{L}(\tau) d\tau + i_{L}(0)$$

$$(2.27)$$

Similarly, the instantaneous current through the inductor contains some information about the history of voltage applied across the inductor. Current through an inductor cannot change instantaneously, as this would require an infinite voltage to be applied across the inductor. An impulse voltage applied to the inductor can produce instantaneous change in the inductor's current.

RC Circuit. Consider a simple RC circuit shown in Figure 2.7a. We can write Kirchoff's voltage law for the circuit as follows,

$$RC\frac{dv_C}{dt} + v_C = V_s \implies v_C(t) = e^{-t/RC} \left[v_C(0) + \int_0^t \frac{1}{RC} e^{\tau/RC} V_s(\tau) d\tau \right]$$
 (2.28)

The above euqation gives the general solution for the voltage across the capactior. We can derive all other variables of interest from v_C . The response to an impulse input, step input, sinusoidal input, or any arbitrary V_s can be computed using the above equation. The response for a step input obtained using a fixed sources and a swtich is given in Figure 2.7a. RC is the time constant of the circuit, and is a measure of how fast the capacitor charges or discharges and it has the units of time.

(a) A Simple RLC Circuit

(b) Four different step responses

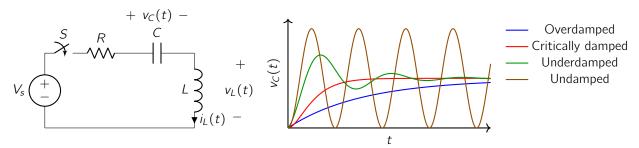


Figure 2.8: A simple RLC circuit and the different respones of the circuit when the switch S is closed at time t=0, with zero initial capacitor voltage and inductor current whent.

Using Laplace transform to analyze the circuit, we can write the following expression for the voltage across the capacitor,

$$V_C(s) = \frac{1}{sRC + 1}V_s(s) + \frac{v_C(0)}{sRC + 1}$$
 (2.29)

 $V_C(s)$ and $V_s(s)$ are the Laplace transforms of $v_C(t)$ and $V_s(t)$. Replacing $s = j\omega$ will give us the frequency response of the system with $V_C(s)$ as the output and $V_s(s)$ as the input.

RL Circuit. Consider a simple RL circuit shown in Figure 2.7b. Kirchoff's voltage law for the circuit is as follows,

$$\frac{L}{R}\frac{di_{L}}{dt} + i_{L} = \frac{1}{R}V_{s} \implies i_{L}(t) = e^{-t/L}\left[i_{L}(0) + \int_{0}^{t} \frac{1}{L}e^{\tau/L}V_{s}(\tau)\,d\tau\right]$$
(2.30)

The above equation gives the general solution for the current through the inductor. The response to an impulse input, step input, sinusoidal input, or any arbitrary V_s can be computed using the above equation. The response for a step input obtained using a fixed sources and a swtich is given in Figure 2.7a. $\frac{L}{R}$ is the time constant of the circuit, and is a measure of how fast the current through the inductor can chamge; it has the units of time.

Similarly, applying the Laplace transform, we have,

$$I_{L}(s) = \frac{1}{sL + R} V_{s}(s) + \frac{i_{L}(0)}{sL + R}$$
 (2.31)

RLC Circuit. Consider a simple series RLC circuit shown in Figure 2.8. Kirchoff's voltage law for the circuit is as follows,

$$LC\frac{d^2v_C}{dt^2} + RC\frac{dv_C}{dt} + v_C = V_s$$
 (2.32)

The response of the circuit when the switch is closed at t=0 is qualitatively different depending on the values of R, L and C. This is better understood in the Laplace domain. The Laplace transform of the response $v_c(t)$ is given by,

$$V_C(s) = \frac{V_s(s)}{LCs^2 + RCs + 1} + v_C(0) \frac{C(Ls + R)}{LCs^2 + RCs + 1} + i_L(0) \frac{RC}{LCs^2 + RCs + 1}$$
(2.33)

Note that $i_L(0) = \frac{dv_C}{dt}(0)$. The denominator of the above equation is a second order polynomial in s, and thus the response of the circuit will depend on the roots of the polynomial. The four different responses are shown in Figure 2.8.

2.1.8 Steady State Sinusoidal Analysis

The previous subsection looked at the transient response of the some simple circuits involving R, L, and C. However, we are often interested in the steady state response of these circuits to sinussoidal excitations. This is because sinusoisal signals are eigenfunctions of linear systems, such as the linear circuits we have discussed so far. With sinusoidal excitations, all currents and voltages in the circuit will also be sinusoidal with the same frequency as the input excitations, althought with different amplitudes and phases. Complex numbers are used to in this analysis as for compact representation of amplitude and phase of signals.

For steady state frequency analysis for circuits involving R, L, and C, we extend the idea of resistance to impedance. Fourier transforms provide a natural way to analyze the steady state response of linear circuits to sinusoidal excitations. We first recast the voltage-current relationships of the circuit elements in the Fourier domain.

Resistor:
$$v_R = R i_R \implies V_R (j\omega) = R \cdot I_R (j\omega)$$

Capacitor: $i_C = C \frac{dv_C}{dt} \implies V_C (j\omega) = \frac{1}{j\omega C} \cdot I (j\omega)$
Inductor: $v_L = L \frac{di_L}{dt} \implies V_L (j\omega) = j\omega L \cdot I_L (j\omega)$ (2.34)

The Fourier transformed voltage-current relationships of the circuit elements look like Ohms law, with the concept of resistance extended to impedance. The impedance of a capacitor and inductor defined as the are defined as the ratio of the Fourier transformed voltage to the Fourier transformed current. They are functions of the frequency of the voltage/current. The impedance of a resistor, capacitor, and inductor are given by,

Resistor:
$$Z_R = R$$

Capacitor: $Z_C = \frac{1}{j\omega C}$ (2.35)
Inductor: $Z_L = j\omega L$

Equivalent Impedance. The equivalent impedance of a circuit with impedances $Z_1, Z_2, \cdots Z_n$ in series is given by $Z_{eq} = Z_1 + Z_2 + \cdots + Z_n$. When the impedances $Z_1, Z_2, \cdots Z_n$ are in parallel, the equivalent impedances is given by $\frac{1}{Z_{eq}} = \frac{1}{Z_1} + \frac{1}{Z_2} + \cdots + \frac{1}{Z_n}$.

Amplitude and phase modification by an impedance. If a sinusoidal current excitation $i(t) = I_o \sin(\omega t)$ is passed through an impedance Z. Then the voltage across the impedance is given by,

$$v(t) = |Z|I_o \sin(\omega t + \arg Z)$$
 (2.36)

The amplitude of the sinusoidal voltage is $|Z|I_o$, and the phase of the sinusoidal voltage is arg Z, with respect to the current i(t).