

REPRESENTATIONS OF $\mathfrak{sl}_2\mathbb{C}$

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The representation theory of the Lie algebra $\mathfrak{sl}_2\mathbb{C}$ is extremely important in the representation theory of semisimple Lie algebras, and consequently, it is also important in the representation theory of Lie algebras. It has an explicit and relatively simple description of all the irreducible representations, which we will lay out here.

The *special linear group* $SL_n\mathbb{C}$ is the group of complex $n \times n$ matrices with determinant 1. By differentiating the determinant map at the identity, we get that its Lie algebra, denoted $\mathfrak{sl}_n(\mathbb{C})$ is the space of complex $n \times n$ matrices with trace 0, equipped with the commutator bracket. It admits a basis given by the matrices

$$H = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad X = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \quad Y = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

It's easy to verify that these satisfy the commutation relations

$$[H, X] = 2X \quad [H, Y] = -2Y \quad [X, Y] = H$$

As it turns out, this basis will make the analysis of irreducible representations quite simple. Let V be an irreducible representation of $\mathfrak{sl}_2\mathbb{C}$. From the semisimplicity of $\mathfrak{sl}_2\mathbb{C}$, we have that the representation preserves the Jordan decomposition, and since H is diagonalizable, its action on V will also be diagonalizable, giving us a direct sum decomposition $V = \bigoplus_{\lambda} V_{\lambda}$ of eigenspaces V_{λ} of H . Let $v \in V_{\lambda}$ be an eigenvector of H with eigenvalue λ . To analyze the action of X on v , we note a few helpful identities

$$\begin{aligned} HX &= XH + HX - XH = XH + [H, X] \\ HY &= YH + HY - YH = YH - [H, Y] \end{aligned}$$

This then gives us the identities

$$\begin{aligned} HXv &= XHv + [H, X]v = \lambda Xv + 2Xv = (\lambda + 2)Xv \\ HYv &= YHv + [H, Y]v = \lambda Yv - 2Yv = (\lambda - 2)Yv \end{aligned}$$

So the action of X raises the eigenvalue by 2, while the action of Y lowers the eigenvalue by 2. Furthermore, all the eigenvalues must differ by some multiple of 2. If we fix an eigenvalue λ , then the subspace

$$\bigoplus_{k \in \mathbb{Z}} V_{\lambda+2k}$$

is fixed by the actions of H , X , and Y , so it is invariant under the action of $\mathfrak{sl}_2\mathbb{C}$. Then by irreducibility of V , this must be the whole space. Since V is finite dimensional, there exists some maximal eigenvalue n of V . Let $v \in V_n$. Since $V_{n+2} = 0$, $Xv = 0$, so we now want to observe the action of Y on v . We claim that the set

$$S = \{Y^m v : m \in \mathbb{Z}^{\geq 0}\}$$

spans all of V . Due to the irreducibility of V , it suffices to show that the span of these vectors is fixed by the action of H , X , and Y . It is clearly fixed by Y , and since all elements of S are eigenvectors of H , it is clear that the subspace is preserved by H as well. The final thing to check is the action of X . We know $Xv = 0$, and for XY , we compute

$$\begin{aligned} XYv &= YXv - [Y, X]v \\ &= 0 - (-Hv) \\ &= nv \end{aligned}$$

Proposition 1.1.

$$XY^m v = m(n - m + 1)Y^{m-1}v$$

Proof. We do this by induction. The base case $m = 1$ is verified above. Then assuming the proposition for m , we wish to verify the identity for Y^{m+1} . We compute

$$\begin{aligned}
 XY^{m+1}v &= [X, Y](Y^m v) + YX(Y^m v) \\
 &= HY^m v + Y(XY^m v) \\
 &= (n - 2m)Y^m v + Y(m(n - m + 1)Y^{m-1}v) \\
 &= ((n - 2m) + (m(n - m + 1)))Y^m v
 \end{aligned}$$

Expanding out the coefficient and the expression $(m + 1)(n - (m + 1) + 1)$ then verifies this identity. \blacksquare

Therefore X fixes this subspace, showing that the span of S must be all of V . We also note that this gives a basis of eigenvectors for V , with each eigenvector $Y^m v$ having a different eigenvalue, which implies that it spans the eigenspace with eigenvalue $n - 2m$. Finally, we note that since V is finite dimensional, there must exist some smallest integer m such that $Y^m v = 0$. Then from our proposition above, we get

$$0 = Y^m v = m(n - m + 1)Y^{m-1}v \implies m(n - m + 1) = 0 \implies n - m + 1 = 0$$

so n is a nonnegative integer, and the fact that $n - m + 1$ implies that the eigenvalues of H must be symmetric about 0, and the dimension of V must be $n + 1$. Since the representation is determined by this maximal eigenvalue n , we let V^n denote the irreducible representation in which the maximal eigenvalue of H is n .

Consider the standard representation of $\mathfrak{sl}_2\mathbb{C}$, where the matrices act on \mathbb{C}^2 in the standard way. Then the standard basis vectors e_1 and e_2 are eigenvectors of H with eigenvalues 1 and -1 respectively. We then note that $Xe_1 = Ye_2 = 0$, and $Xe_2 = e_1$ and $Ye_1 = e_2$, so their action of \mathbb{C}^2 satisfies the exact same relations as the irreducible representation V^1 . we can then take the symmetric power $\text{Sym}^2 \mathbb{C}^2$, which has basis $\{e_1^2, e_1 e_2, e_2^2\}$. Using the standard rules for symmetric powers of a Lie algebra representation, we compute

$$\begin{aligned}
 H(e_1^2) &= e_1 H e_1 + (H e_1) e_1 = 2e_1^2 \\
 H(e_1 e_2) &= e_1 H e_2 + (H e_1) e_2 = -e_1 e_2 + e_1 e_2 = 0 \\
 H(e_2^2) &= e_2 H e_2 + (H e_2) e_2 = -e_2^2 - e_2^2 = -2e_2^2
 \end{aligned}$$

Therefore, $\text{Sym}^2 \mathbb{C}^2$ is the irreducible representation V^2 . As it turns out, the symmetric powers of the standard representation on \mathbb{C}^2 form the complete set of irreducible representations of $\mathfrak{sl}_2\mathbb{C}$. The symmetric power $\text{Sym}^n \mathbb{C}^2$ has basis $\{e_1^{n-k} e_2^k : 0 \leq k \leq n\}$. A computation of the action of H on these vectors to be

$$\begin{aligned}
 H(e_1^{n-k} e_2^k) &= e_1^{n-k} H(e_2^k) + H(e_1^{n-k}) e_2^k \\
 &= -k e_1^{n-k} e_2^k + (n - k) e_1^{n-k} e_2^k \\
 &= (n - 2k) e_1^{n-k} e_2^k
 \end{aligned}$$

Giving us that the representation is the irreducible representation V^n . This tells us that all the irreducible representations of $\mathfrak{sl}_2\mathbb{C}$ are given by symmetric powers of the standard representation of $\mathfrak{sl}_2\mathbb{C}$ acting on \mathbb{C}^2 .