

# Boiling Study for H, H2, H3 and He3 Targets

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## Abstract

When the beam passes through a cryo-target, the local temperature fluctuations could cause a variation in the target density. This density variation is called *Boiling*, and it increases with increasing current. The density fluctuation due to passing beam, or the boiling effect, for jlab Tritium experiments has been studied for Hydrogen (H), Deuterium (H2), Tritium (H3) and Helium (He3) targets.

## 1 Method Overview

One way to study the boiling effect is by using the **Yield Analysis**, where the charge normalized yield is calculated versus current. A decrease in the charged normalized yield at higher currents, indicates that the target's density is decreasing. In another words, the target is *boiling*.

### Steps for Yield Analysis:

1. At each current, the total charge is calculated using:

$$Q = a \times \Delta Counts + b \times \Delta Time \quad (1)$$

wheren  $Q$  is the total charge,  $a$  and  $b$  are the slop and offset identified by calibrating the Beam Charge Monitors "BCM's" frequency *vs.* unser current . One can get the *Counts* and *Time* from the BCM and Clock scalers respectively. For this study, the dnevr BCM is used, and the calibration constants are:  $a = 0.0003264$  and  $b = 0.1055$  (see Fig.1)

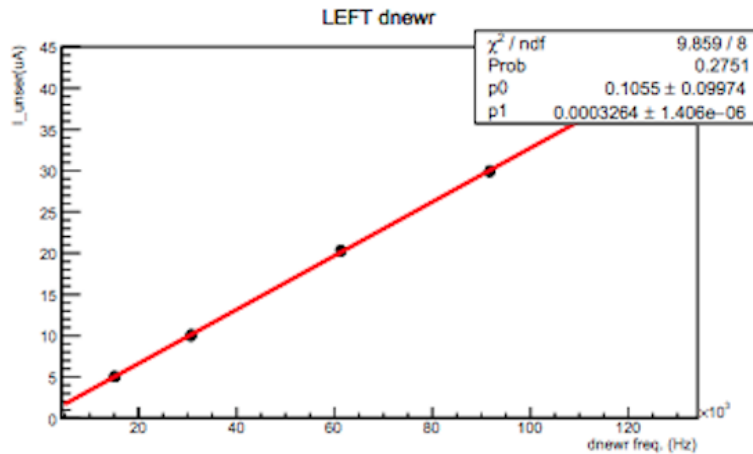


Figure 1: dnevr Calibration by: Nathaly Santiesteban

2. The charge yield of a current  $I$ ,  $Yield(I)$ , is then calculated using the following:

$$Yield(I) = \frac{N_{good} \times PS}{Q \times efficiencies \times LT} \quad (2)$$

where  $N_{good}$  is the number of good electrons at each current. Determining  $N_{good}$ , *i.e.* selecting good electrons, is discussed in the following sections. PS and LT are the pre-scale and live-time of each run respectively. The *efficiencies* include both detectors and trigger efficiencies.

3. The charge normalized yield is calculated by normalizing the charge yield for a current  $I$ , eq.2, over the charge yield when there is no current.

$$Y_{norm}(I) = \frac{Yield(I)}{Yield(I=0)} \quad (3)$$

Since we don't really know what the charge yield should be at zero current, we use the charge yield at the lowest current available ( $2.5 \mu\text{A}$  for this study) as a normalization factor. Then fit the results so that the normalized yield will be equal to 1 at  $I = 0$ .

## 2 Electron Selection:

In order to extract a good electron sample, several cuts were applied on the recorded data:

1. **Selecting Events from Pion Rejector Layers PRL1&PRL2 and Gas Čerenkov GC (PID cut):**

To identify the good electrons from the unwanted particles, e.g. pions, only events that leave a pre-specified range of energy on pion rejector layers and Gas Čerenkov will be selected. We refer to this way of Particle IDentification as a **PID cut**. In this analysis, the following PID cuts were applied (see Fig.2 &3):

```
L.prl1.e>600
&&L.prl2.e>300
(L.prl1.e+L.prl2.e)>1700
(L.prl1.e+L.prl2.e)<2300
L.cer.asum_c>1000
```

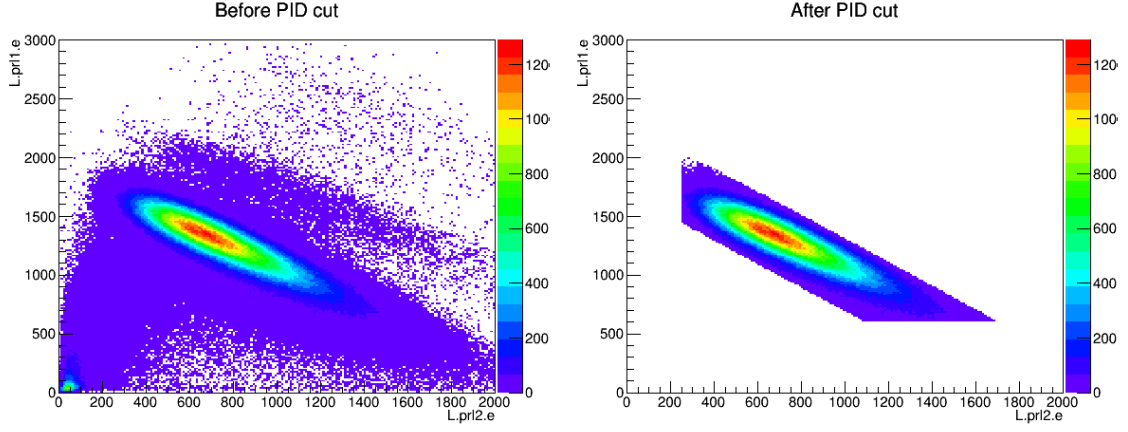


Figure 2: PID cut applied on the calorimeters PRL1 & PRL2 for LHRS

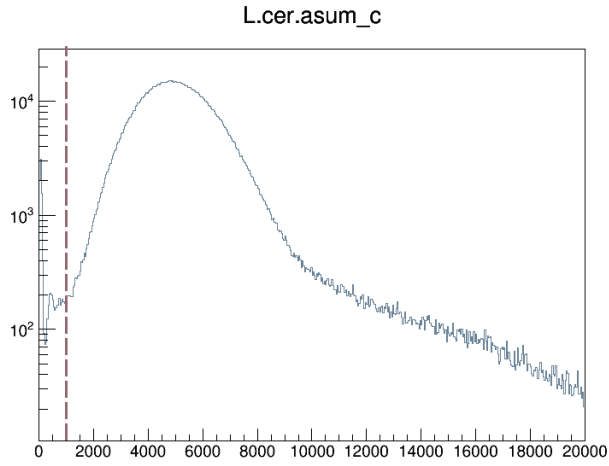


Figure 3: PID cut applied on Gas Čerenkov ADC sum

## 2. Selecting events based on the number of tracks (One-track Cut):

Good events are characterized of leaving only one track on the Vertical Drift Chambers VDC. Therefore, only events with one-track were selected, and we call this selection **One-track Cut**

$$\text{L.tr.n}=1$$

## 3. Selecting events based on the trigger (Trigger Cut):

Three triggers were used during the experiment for LHRS. T1 [S0&&S2], T2 [S0&&S2 && GC] and T3 [(S0 || S2) && GC]. T2 is the trigger used through this study. The selection of events that are recorded by a specific trigger is called **Trigger Cut**

$$\text{DL.evtypebits}\gg 2\&1$$

## 4. Selecting events based on the spectrometer acceptance (Acceptance Cut):

Electrons that scattered within a specific range of  $\theta$ ,  $\phi$  and  $\delta$  were selected and let's call this selection the **Acceptance Cuts**. This cut is to make sure that the events are chosen within the acceptance range of the spectrometer.

$$\begin{aligned} \text{abs(L.tr.tg\_th)} &< 0.03 \\ \text{abs(L.tr.tg\_ph)} &< 0.04 \\ \text{abs(L.tr.tg\_dp)} &< 0.05 \end{aligned}$$

#### 5. Removing events from target end-caps (End-cap Cut):

Each cell has two end-caps made of Aluminum. When selecting the good electrons, the events that come from the cell's end-cap need to be excluded. Most of those events can be removed by cutting out the peaks of both end-caps in the  $z_{react}$  (or  $y_{tar}/\sin\theta$ ) distribution. The following cut on the target length was applied (See Fig.4):

$$\text{abs(L.tr.tg\_y}/\sin\theta) < 0.075 \text{ m}$$

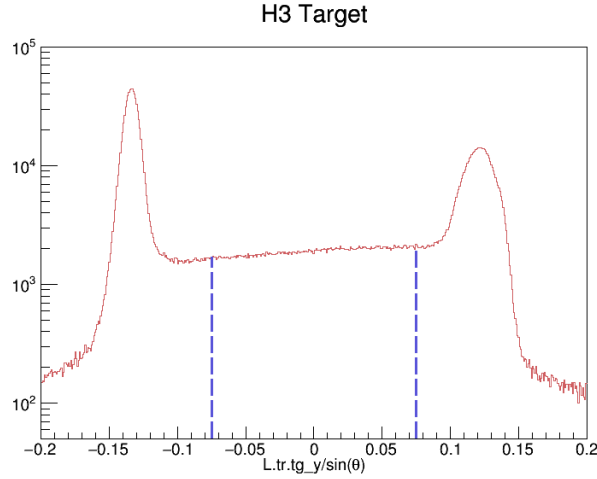


Figure 4: Cutting on the target length to get rid of end-cap events

However, cutting the end-cap peaks will not remove all of the end-cap events. There are still some Aluminum events mixed with the events from the cryo-target. These end-cap events are called "end-cap background". In order to estimate this background, the following steps were taken:

- Estimating the end-cap background *vs.* current:

In order to check how end-cap background changes with increasing current, a comparison between the background at low and high current was estimated for an empty cell (Fig.5). Each run was normalized over the total charge yield, and the ratio of the events at high current to low current was estimated. The ratio was found to be  $\sim 1.006$ , which indicates that the end-cap background will *not* increase with increasing current and it is a constant number. Therefore, for the charge normalized yield this value will cancel out. However, when calculating the total charge yield, the background events need to be extracted from cryo-target events.

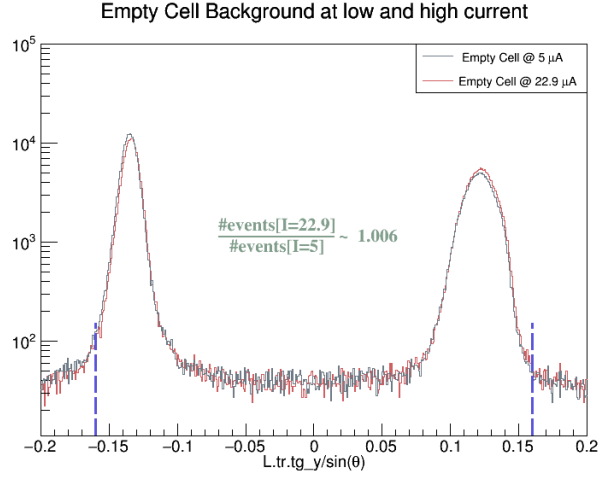


Figure 5: Ratio of end-cap background at high current to low current

- Estimating the end-cap background for each cryo-target:

Even though the end-cap background should be the same for low and high current, this background value might be different for each target. To check that, an empty cell run at low current ( $\sim 5 \mu\text{A}$ ) was compared to each cryo-target at the same current. Since the windows thickness is different slightly from cell to cell, normalizing over the charge yield did not seem accurate. Calculating the charge yield of two cells by applying same cut on the target length, does not mean necessarily that we picked the same thickness of both windows. We might end up normalizing over the charge yield that came out of two different target thickness. To avoid that, and knowing that the peak at the cell window is proportional to the window's thickness, the upstream windows were scaled to match each other. The two runs, *i.e.* the empty cell run and the cryo-target run, were normalized so that the maximum number of events at the upstream windows are the same for both runs. To determine the maximum number of events, a gaussian fit about  $1 \sigma$  was made at the peak (see Fig.6). The y-value at the mean of the fit is considered the maximum number of events for that window. Then, the two windows were normalized so that the peak will match.

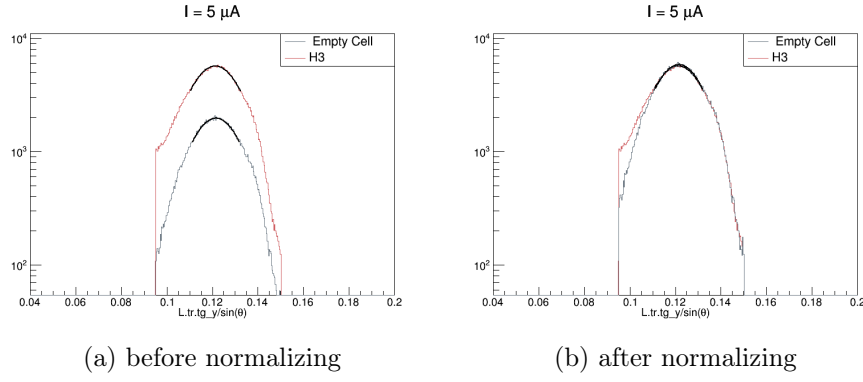


Figure 6: Normalizing the upstream window for two runs

After normalizing the upstream windows, the end-cap contamination was calculated as:

$$\text{End-cap Contamination} = \frac{\text{\#of good events come from empty cell from } -0.075 \text{ to } 0.075}{\text{\#of good events come from cryo-cell from } -0.075 \text{ to } 0.075}$$

This way, one can know how many events coming from end-caps *and* that will trick us to be as good electrons. That is the number that will affect the cryo-target charge yield directly. The other events coming from the end-caps, will be cut out with the good electrons cut and shouldn't be extracted from the number of good events used to measure the charge yield.

The above method was made for H3, He3, D2 and H and the results are shown in Fig.7

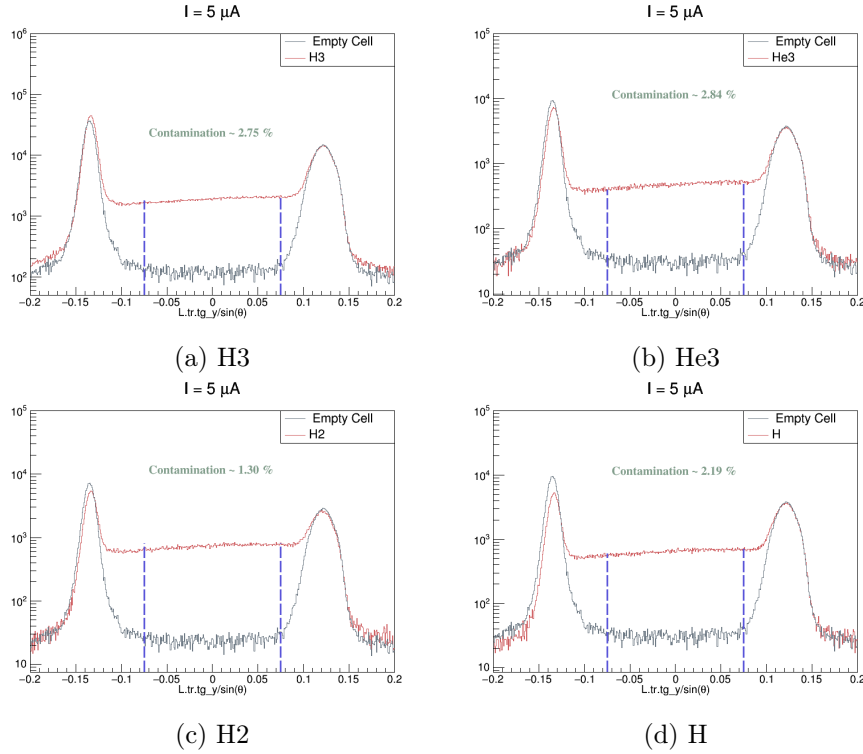


Figure 7: End-cap Background for each target

When calculating the total charge yield, the end-cap background should be extracted from the total number of good events of each target. However, since

this background is assumed to be almost the same with increasing current, the *Normalized* charge yield -which is the main goal of this study- will be the same, i.e. the background will cancel out. The number of charge yield before and after extracting end-cap background will be listed for each target.

## 6. Selecting events from a stable beam (Beam Trip Cut):

If the beam is stable during the run (see Fig.8a), the total charge can be calculated using Eq.1. Both  $\Delta Counts$  and  $\Delta Time$  can be extracted from the scalers at the beginning and the end of the run. If the beam tripped for a small period of time, i.e. at the beginning or end of the run, one can just cut those periods of the run when the beam was tripping, and handle only the period of the run when the beam was stable. However, sometimes the beam trips significantly from time to time during the run (see Fig.8b), which makes selecting a period of the run with a stable beam hard and maybe not accurate. It is crucial when selecting the good electrons to make sure that those electrons survived the beam tripping, and we call this selection **Beam Trip Cut**.

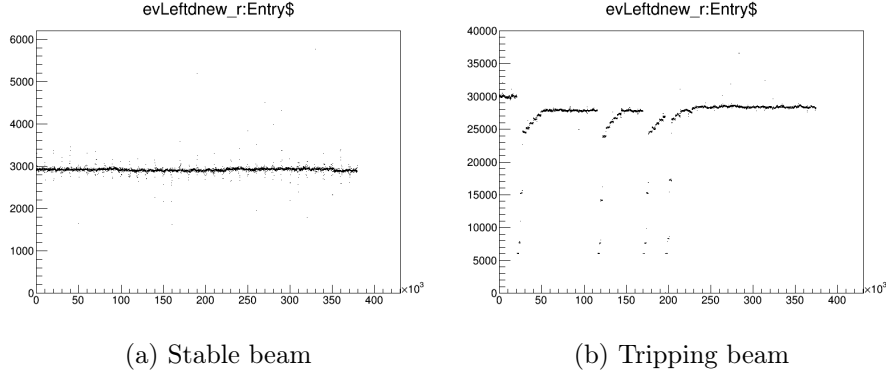


Figure 8: Examples for runs with stable beam and tripping beam

In this analysis, event-by-event check was applied. If the event passed all of the above cuts, i.e. PID, trigger, one track, end-cap and acceptance cuts, then the current that produced that event will be checked. If the current was within  $\pm 2\mu A$  about the current requested for the run, that event will be considered as a *good electron*. Otherwise, the event will be neglected. Fig.8 illustrates how the beam trip cut was applied for this study.

$$\epsilon_{cer} = \frac{N_e^{GC}}{N_e^{sample}} \quad (4)$$

where  $N_e^{sample}$  is the number of electrons obtained after applying PRL, trigger, one-track, end-cap, beam trip acceptance cuts to get good electron sample.  $N_e^{GC}$  is the number of electrons obtained after applying a Čerenkov cut in addition to the same cuts applied to obtain  $N_e^{sample}$ .

## 7. VDC Efficiency ( $\epsilon_{trk=1}$ ):

The efficiency of the VDC can be obtained by checking the efficiency of the one-track

cut as follows:

$$\epsilon_{trk=1} = \frac{N_{trk=1}}{N_{\text{all tracks}}} \quad (5)$$

where  $N_{\text{all tracks}}$  is the number of electrons that passes PID (Cer and PRL's) , trigger, end-cap, acceptance and beam trip cuts, that left any track on the VDC's including zero tracks.  $N_{trk=1}$  is the number of electrons that passed all previous cuts *but* left only one track on the VDC's.

#### 8. Trigger Efficiency ( $\epsilon_{trig}$ ):

As mentioned above, T2 was used to select electrons in this study, and T1 was used to check T2 efficiency as following:

$$\epsilon_{trig} = \frac{PS_{T_2} \times N_{T_2}}{PS_{T_1} \times N_{T_1}} \quad (6)$$

where  $PS_{T_i}$  and  $N_{T_i}$  are the pre-scale and number of events recorded for trigger type  $i$  respectively. The events selected here must pass the PID, one-track, end-cap, acceptance and beam trip cuts.

### 3 Solid Target Boiling "Check":

Since the solid target does not change its density with increasing current, the normalized yield should be  $\tilde{1}$  for any solid target. The upstream and downstream windows of the Tritium cell were selected for this test. The results are shown in Fig.9

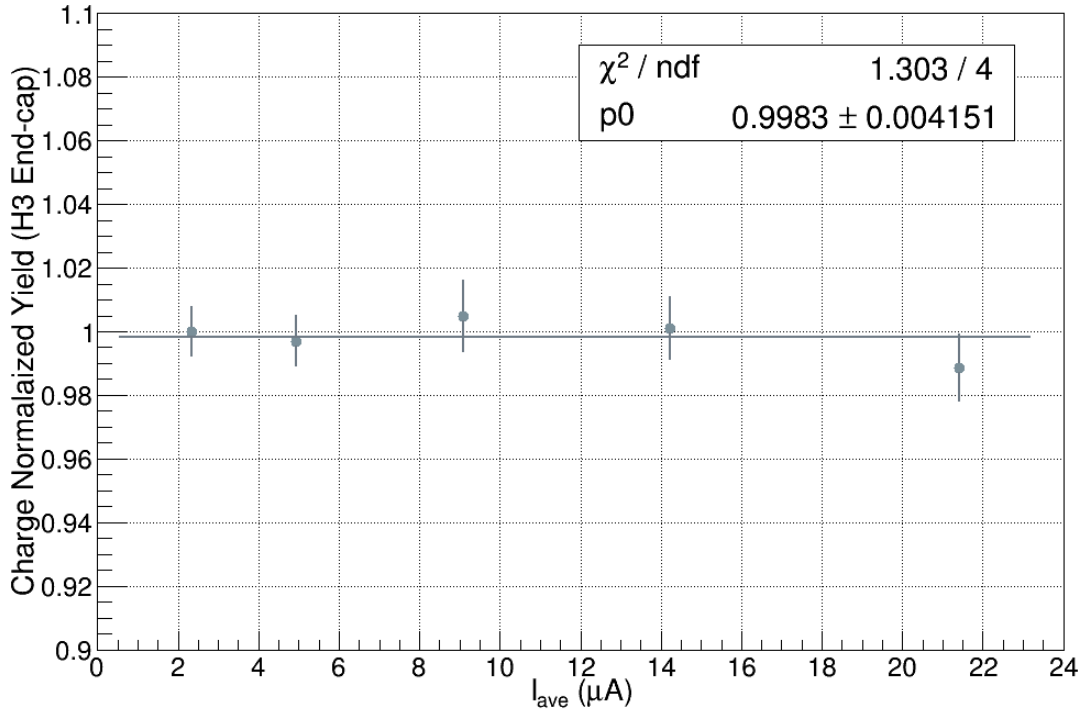


Figure 9: Charge Normalized Yield for Solid Target (Upstream End-cap of H3 cell)



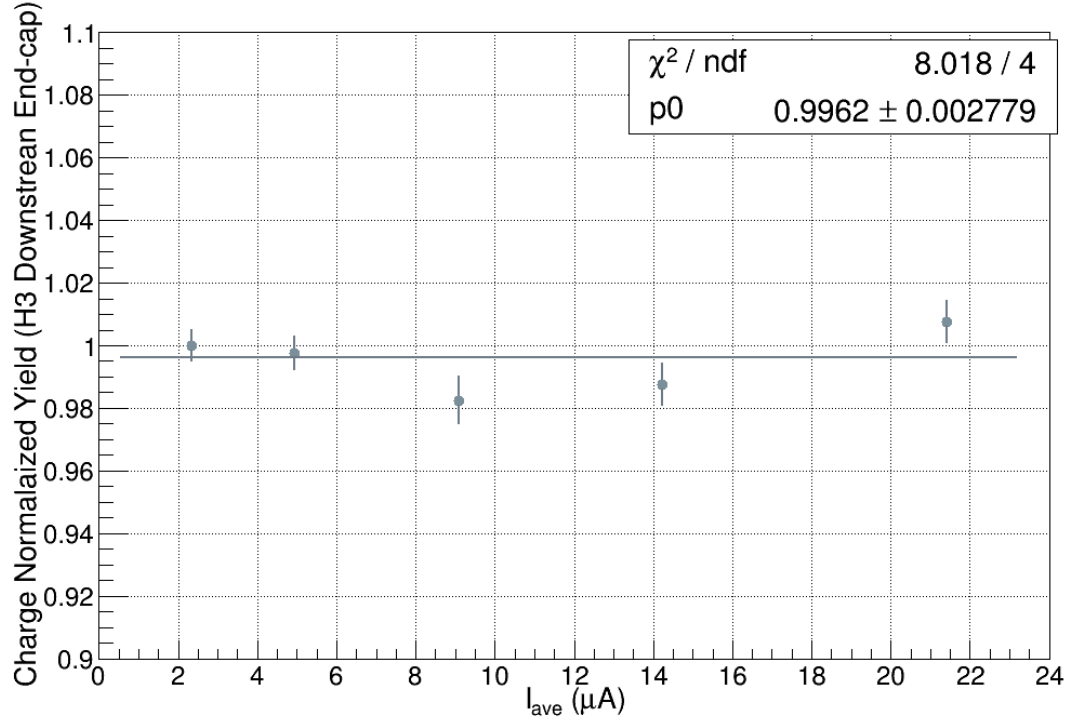


Figure 10: Charge Normalized Yield for Solid Target (Downstream End-cap of H3 cell)

## 4 Tritium Target Boiling:

Table 1: Efficiencies for H3 boiling runs

Run	Cerenkov eff. (%)	one track eff. (%)	trigger [T2] eff (%)	DT (%)
902	99.9828	98.6125	98.9123	4.36
905	99.9879	98.4400	99.4147	4.43
9091	99.9845	98.2996	100.5	2.91
910	99.9925	98.0794	100.	4.54
911	99.9916	97.8293	99.8825	3.43

Table 2: Normalized Yield H3 Target

Run	Current (μA)	#good events	Total Charge (μC)	Yield	Charge Normalized Yield	error
902	2.343	67998	1964.302	111.350	0.983	0.002
905	4.809	64243	4382.545	109.720	0.968	0.002
909	9.094	31366	6141.763	106.410	0.939	0.003
910	14.232	39032	8009.334	104.013	0.918	0.002
911	21.407	34709	14567.098	101.008	0.891	0.003

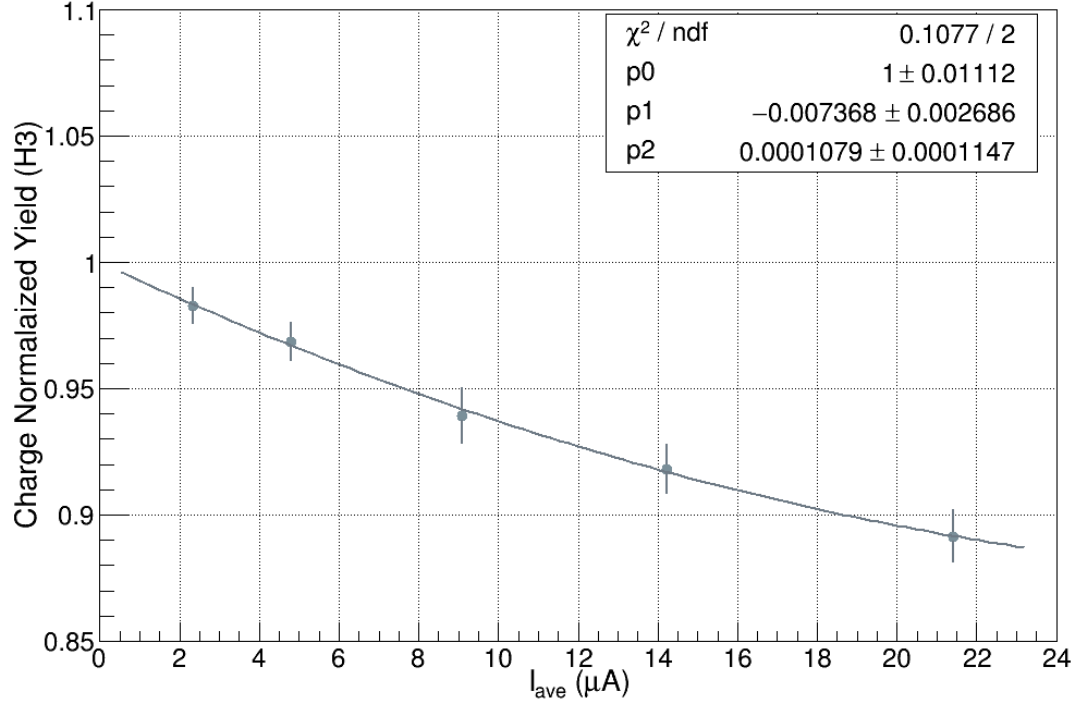


Figure 11: Charge Normalized Yield for Tritium

## 5 Helium-3 Target Boiling:

Table 3: Efficiencies for He3 boiling runs

Run	Cerenkov eff. (%)	one track eff. (%)	trigger [T2] eff (%)	DT (%)
916	99.9812	98.7991	100.	3.70
915	99.9827	98.7315	101.5	3.14
914	99.9927	98.5244	102.2	3.82
913	99.9907	98.3497	101.34	3.93
912	99.9851	98.0509	102.7	2.83

Table 4: Normalized Yield and their errors for He3 runs

Run	Current ( $\mu\text{A}$ )	#good events	Total Charge ( $\mu\text{C}$ )	Yield	Charge Normalized Yield	error
916	2.421	51549	1665.761	96.801	0.991	0.002
915	4.423	39393	2960.469	95.950	0.983	0.002
914	9.472	40088	5761.835	93.396	0.975	0.002
913	14.700	52454	11732.765	93.344	0.956	0.002
912	21.561	38998	17491.905	91.094	0.933	0.002

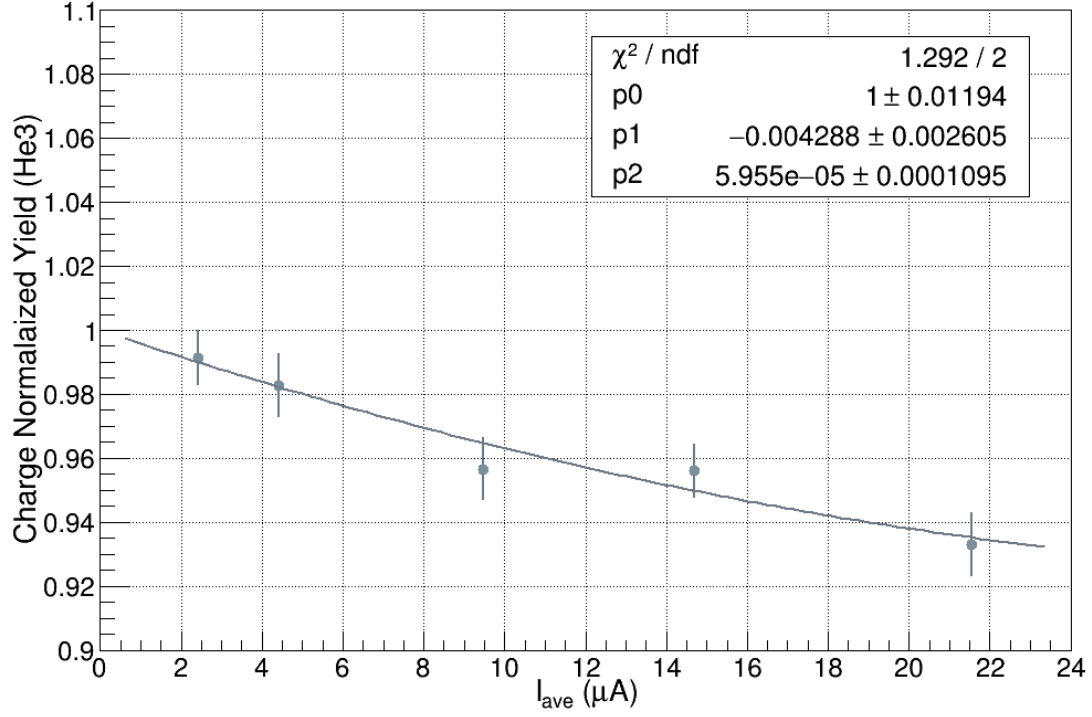


Figure 12: Charge Normalized Yield for He3

## 6 DeuteriumTarget Boiling:

Table 5: Efficiencies for H2 boiling runs

Run	Cerenkov eff. (%)	one track eff. (%)	trigger [T2] eff (%)	DT (%)
919	99.98	98.724	100.3	5.625
918	99.99	98.62	100.5	6.97
920	99.98	98.37	101.2	3.81
921	99.98	98.08	101.6	35.21
923	99.98	97.58	101.1	4.076

Table 6: Normalized Yield and their errors for H2 runs

Run	Current ( $\mu A$ )	#good events	Total Charge ( $\mu C$ )	Yield	Charge Normalized Yield	error
919	2.483	96793	1325.762	234.357	0.979	0.002
918	4.856	115933	2737.609	229.538	0.959	0.002
920	9.784	52131	4902.409	222.121	0.928	0.004
921	14.515	48688	6936.292	217.450	0.909	0.004
923	21.533	52185	10334.375	213.374	0.892	0.004

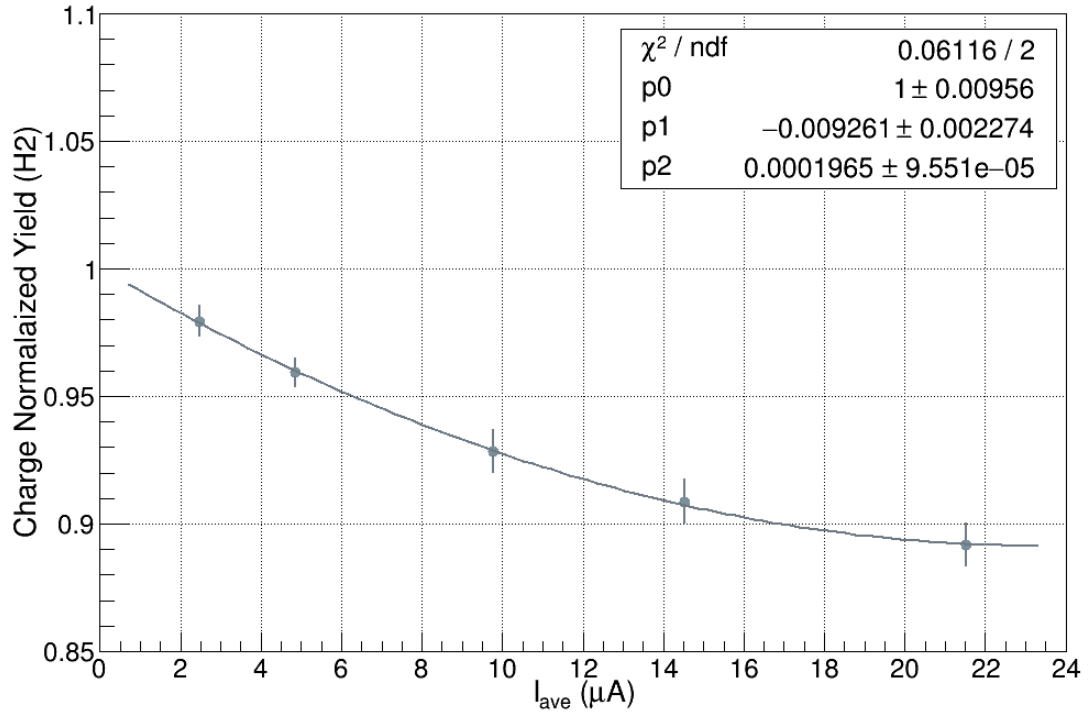


Figure 13: Charge Normalized Yield for H2

## 7 Hydrogen Target Boiling:

Table 7: Efficiencies for H boiling runs

Run	Cerenkov eff. (%)	one track eff. (%)	trigger [T2] eff (%)	DT (%)
901	99.9825	98.8563	100.5	3.08
899	99.9859	98.6726	100.8	3.41
897	99.9846	98.4586	100.9	3.77
896	99.9964	98.2092	100.1	3.68
895	99.987	97.8128	101.4	4.12

Table 8: Normalized Yield and their errors for H runs

Run	Current ( $\mu A$ )	#good events	Total Charge ( $\mu C$ )	Yield	Charge Normalized Yield	error
901	2.011	50379	1264.459	206.816	0.986	0.004
899	4.321	48705	2495.384	203.141	0.968	0.004
897	9.536	57093	6147.207	194.303	0.9326	0.003
896	13.636	54548	8928.621	193.563	0.922	0.003
895	20.937	105313	23805.592	186.091	0.887	0.002

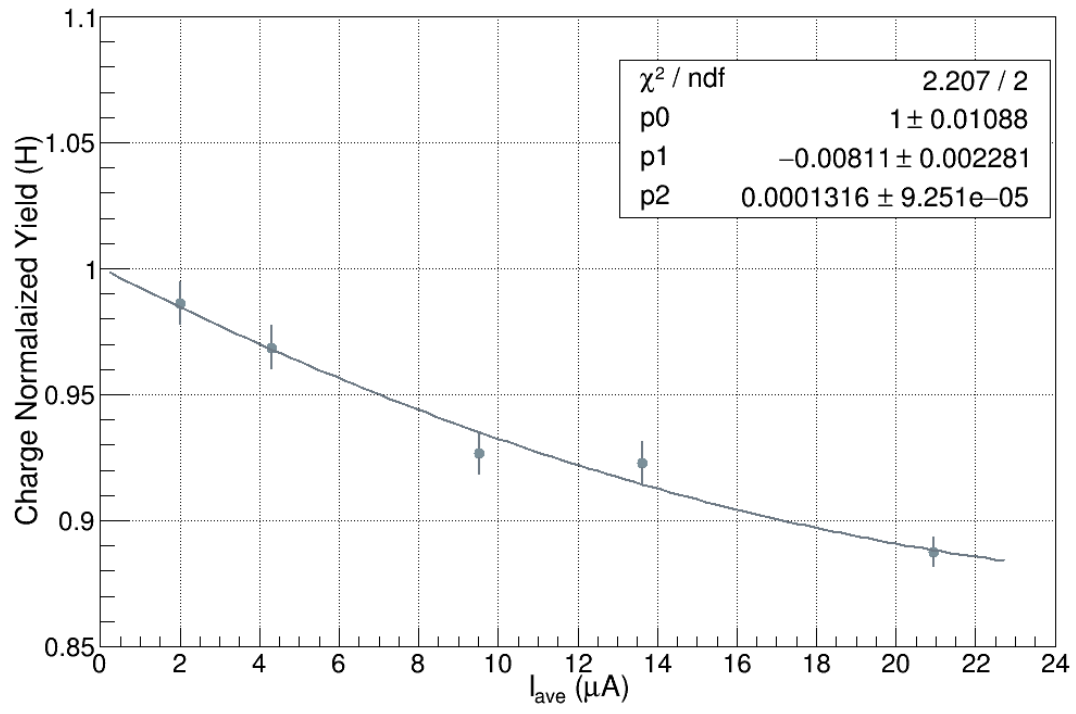


Figure 14: Charge Normalized Yield for H