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# Time-independent Schrödinger equation in 1D

For the time-dependent Schrödinger equation

$$i\hbarrac{\partial}{\partial t}|\psi(t,x)
angle=\hat{H}|\psi(t,x)
angle$$

can be decomposed by the method of separation of variables when the potential is time-independent V=V(x), that is

 $\hat{H}|\psi(x)
angle=E|\psi(x)
angle,$ 

where the hamiltonian is given by

$$\hat{H} = -rac{\hbar^2}{2m}
abla^2 + V(x).$$

### Problem

The goal of this project is to solve time-independent Schrödinger equation in 1-dimension called "Eigen Probelm", that is

$$oxed{-rac{\hbar^2}{2m}rac{d^2}{dx^2}\psi(x)+V(x)\psi(x)=E\psi(x),}$$

#### Discrete Schrödinger equation

The time-independent Schrödinger equation is

$$-rac{\hbar^2}{2m}rac{d^2}{dx^2}|\psi
angle + V|\psi
angle = E|\psi
angle,$$

where is the column vector

$$|\psi
angle=\psi_i,\quad i=0,1,\dots,n-1,$$

potential V is a diagonal matrix

$$V=V_i\,\delta_{i,j},\quad i,j=0,1,\dots,n-1,$$

and the 1-dimension laplacian operator  $abla^2$  approximate to a n imes n matrix, since

$$rac{d^2f(x)}{dx^2}pproxrac{f(x-\delta_x)-2f(x)+f(x+\delta_x)}{\delta_x^2},$$

we define

$$rac{d^2}{dx^2} f_i = rac{1}{\delta_x} (f_{i-1} - 2f_i + f_{i+1}) \,, \quad i = 0, 1, \dots, n-1,$$

where  $f_{-1}=f_n=0$ , plugging we have

$$-rac{\hbar^2}{2m}rac{d^2}{dx^2}\psi_i + V_i\,\delta_{i,j}\psi_i = E\psi_i,$$

or

$$-\frac{\hbar^2}{2m}\frac{1}{\delta x^2}\begin{pmatrix} -2 & 1 & 0 & 0 & \cdots & 0 & 0 & 0 \\ 1 & -2 & 1 & 0 & \cdots & 0 & 0 & 0 & 0 \\ 0 & 1 & -2 & 1 & \cdots & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -2 & \cdots & 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \cdots & -2 & 1 & 0 \\ 0 & 0 & 0 & 0 & \cdots & 0 & 1 & -2 \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-3} \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \\ \psi_{n-2} \\ \psi_{n-1} \end{pmatrix} + \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_$$

So, now it becomes solving this eigen problem for a large system

$$\begin{pmatrix} \frac{\hbar^2}{\delta_x^2 m} + V_0 & -\frac{\hbar^2}{2\delta_x^2 m} & 0 & 0 & \cdots & 0 & 0 & 0 \\ -\frac{\hbar^2}{2\delta_x^2 m} & \frac{\hbar^2}{\delta_x^2 m} + V_1 & -\frac{\hbar^2}{2\delta_x^2 m} & 0 & \cdots & 0 & 0 & 0 \\ 0 & -\frac{\hbar^2}{2\delta_x^2 m} & \frac{\hbar^2}{\delta_x^2 m} + V_2 & -\frac{\hbar^2}{2\delta_x^2 m} & \cdots & 0 & 0 & 0 \\ 0 & 0 & -\frac{\hbar^2}{2\delta_x^2 m} & \frac{\hbar^2}{\delta_x^2 m} + V_3 & \cdots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \cdots & \frac{\hbar^2}{\delta_x^2 m} + V_{n-3} & -\frac{\hbar^2}{2\delta_x^2 m} & 0 \\ 0 & 0 & 0 & 0 & \cdots & -\frac{\hbar^2}{2\delta_x^2 m} & \frac{\hbar^2}{\delta_x^2 m} + V_{n-2} & -\frac{\hbar^2}{2\delta_x^2 m} \\ 0 & 0 & 0 & 0 & \cdots & 0 & -\frac{\hbar^2}{2\delta_x^2 m} & \frac{\hbar^2}{\delta_x^2 m} + V_{n-1} \end{pmatrix}$$

# Algorithm

## import package

import numpy as np
import matplotlib as mpl
import matplotlib.pyplot as plt
from scipy.linalg import eigh\_tridiagonal

## setting plotting default

Matplotlib: how to change default style

How to change the font size on a matplotlib plot

```
Im []:
    mpl.rcParams['figure.dpi'] = 400
    mpl.rcParams['lines.linewidth'] = 0.5
    mpl.rcParams['lines.color'] = 'red'
    mpl.rcParams['lines.color'] = 'red'
    mpl.rcParams['text.usetex'] = True
    mpl.rcParams['font.family'] = 'Times New Roman'
    mpl.rcParams['font.size'] = 8
    mpl.rcParams['axes.prop_cycle'] = mpl.cycler('color', 'bgrcmyk')
    mpl.rcParams['axes.linewidth'] = 0.2
    plt.rcParams['mathtext.fontset'] = "cm"
    plt.rc('axes', labelsize=8)
    plt.rc('xtick', labelsize=5)
    plt.rc('ytick', labelsize=5)
```

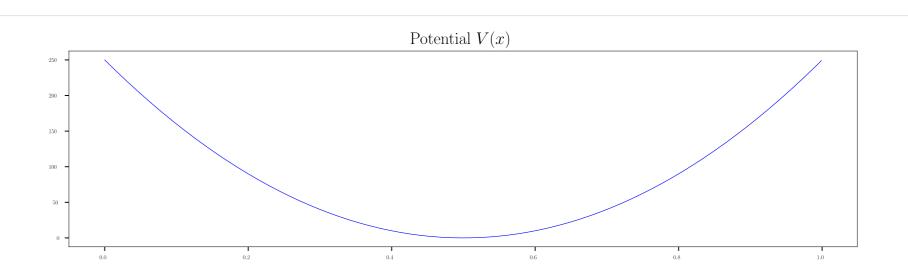
file:///Users/apple/Desktop/中正大學/三年級第一學期/量子物理/schrodinger\_equation/1D/Project\_output/1D\_TISE.html

seting the position space

D\_TISE  $x = [x_i, x_i + dx, \dots, x_f]$ 

#### define the potential function

In [ ]:
 f=plt.figure()
 def V(x):
 v = 1000 \*(x-0.5)\*\*2
 #v[x<0.5]=0
 return v
 plt.plot(x,V(x))
 plt.title("Potential \$V(x)\$")
 plt.show()
 f.savefig('assets/Potential.pdf')</pre>



### Solving the eigen problem

Since the discrete Hamiltonian is an tridiagonal matrix in the form

$$\hat{H} = egin{pmatrix} b_1 & c_1 & & & 0 \ a_2 & b_2 & c_2 & & & \ & a_3 & b_3 & \ddots & & \ & & \ddots & \ddots & c_{n-1} \ 0 & & & a_n & b_n \end{pmatrix}$$

we can use tridiagonal matrix algorithm to solve it, which only need  $\mathcal{O}(n)$  operations rather than  $\mathcal{O}(n^3)$  operations by Gaussian elimination.

Here using eigh\_tridiagonal(d,e) function, built up by scipy.

In [ ]:
 d = 1/dx\*\*2 + V(x)
 e = -1/(2\*dx\*\*2) \* np.ones(len(x)-1)
 w,v = eigh\_tridiagonal(d,e)

### Define the Eigenfunction $\psi_i$ and Eigenvalue $E_i$

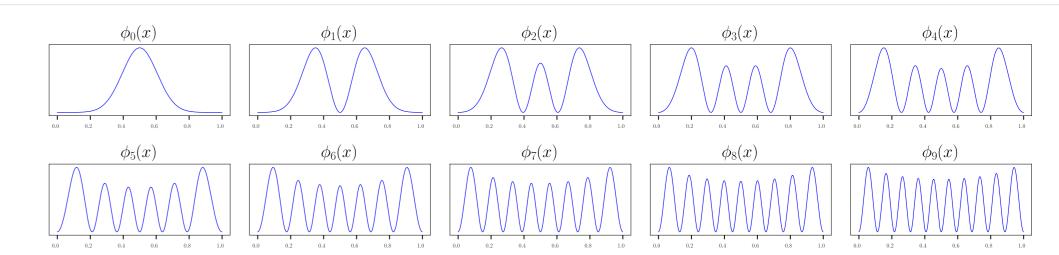
- ullet Eigenfunction  $\psi_i$  ,  $i=0,1,\ldots,\operatorname{length}\left(x
  ight)$
- ullet Eigenvalue  $E_{i}$ ,  $i=0,1,\ldots,\operatorname{length}\left(x
  ight)$

In [ ]: phi = v.T E = w

# Plot the Eigenfunction

$$\left|\phi_i(x)
ight|^2, \quad i=0,1,\ldots,9$$

In []:
 row = 2
 col = 5
 f=plt.figure()
 for i in range(row\*col):
 plt.subplot(row,col,i+1)
 plt.plot(x,phi[i]\*\*2)
 plt.title("\$\phi\_{"}+str(i)+"}(x)\$")
 plt.yticks([])
 f.tight\_layout()
 f.savefig('assets/Eigenfunction.pdf')



In []:
 E\_N = 100
 f, (a0,a1) = plt.subplots(1, 2, gridspec\_kw={'width\_ratios': [1, 10]})
 for i in range(E\_N):
 a0.axhline(y = E[i])
 al.plot(E[0:E\_N], ', ', markersize=2)
 a0.set\_yticks([])
 a0.set\_yticks([])
 al.set\_yticks([])
 al.set\_yticks([])
 al.set\_ytlabel('\$n\$')
 al.set\_ylabel('Energy')
 f.suptitle('Eigenvalue \$E\_n\$')
 f.savefig('assets/Eigenvalue.pdf')

