



THEORETICAL PHYSICS

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# Quantum Field Theory

## SI2410

### Fall 2014

## Preparation Questions

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The whole group

September 16, 2014

### 1 SEMINAR 1: FUNCTIONAL INTEGRALS AND INTRODUCTION TO RENORMALIZATION

#### 1.1 WHAT IS THE RELATIONSHIP BETWEEN THE FUNCTIONAL INTEGRAL FORMALISM AND THE N-POINT CORRELATION FUNCTION?

General generating functional is given by

$$Z[J] = \int D\phi \exp \left[ i \int d^4x [L + J(x)\phi(x)] \right] \quad (1.1)$$

where  $L$  is the Lagrangian density and  $J(x)\phi(x)$  is a source term.

The n-point correlation function is then given by

$$\langle 0|T\phi(x_1)\dots\phi(x_n)|0\rangle = \frac{1}{Z_0} \prod_{i=1}^n \left( -i \frac{\delta}{\delta J(x_i)} Z[J] \right) \quad (1.2)$$

where  $Z_0 = Z[J=0]$ . This yields

$$\langle 0|T\phi(x_1)\dots\phi(x_n)|0\rangle = \frac{\int D\phi \phi(x_1)\dots\phi(x_n) \exp \left[ i \int d^4x L \right]}{\int D\phi \exp \left[ i \int d^4x L \right]} \quad (1.3)$$

## 1.2 GIVEN A LAGRANGIAN DENSITY, HOW CAN THE FEYNMAN RULES OF A THEORY BE COMPUTED USING FUNCTIONAL INTEGRAL FORMALISM?

The Feynman rules can be computed by Taylor expanding

$$\exp iS = \exp \left[ i \int d^4x L_0 + L_{int} \right] \quad (1.4)$$

where  $S$  is the action,  $L_0$  is the free Lagrangian density and  $L_{int}$  is the interaction Lagrangian density. In the case of  $\lambda\phi^4$ -theory we have the Lagrangian

$$L = L_0 - \frac{\lambda}{4!} \phi^4 \quad (1.5)$$

and by Taylor expansion

$$\exp iS = \exp \left[ i \int d^4x L_0 - \frac{\lambda}{4!} \phi^4 \right] = \exp \left[ i \int d^4x L_0 \right] \left( 1 - i \int d^4x \frac{\lambda}{4!} \phi^4 + \dots \right) \quad (1.6)$$

where  $L_0$  gives the propagators and all terms of order higher than two yields interaction. In this case we read off the vertex factor in momentum space as

$$-i\lambda(2\pi)^4 \delta^4(\sum p) \quad (1.7)$$

## 1.3 WHAT COMPLICATIONS ARISE WHEN USING THE FUNCTIONAL INTEGRAL FORMALISM TO QUANTIZE THE ELECTROMAGNETIC FIELD? HOW IS IT SOLVED?

In QED the functional integral,

$$\int DA e^{iS[A]} \quad (1.8)$$

is badly divergent as a finite number of terms will have  $S[A] = 0$  due to gauge invariance,

$$A_\mu \rightarrow A_\mu + \frac{1}{e} \partial_\mu \alpha(x) \quad (1.9)$$

and thus  $A_\mu = \frac{1}{e} \partial_\mu \alpha(x)$  is equal to  $A_\mu = 0$ . Through the Faddeev-Popov procedure we introduce some gauge fixing function  $G(A) = 0$ , this would be  $G(A) = \partial_\mu A^\mu$  in the Lorentz gauge. By inserting  $\delta G(A)$  into our functional integral by using,

$$1 = \int D\alpha \delta(G(A^\alpha)) \det \left[ \frac{\delta G}{\delta \alpha} \right] \quad (1.10)$$

and a general  $G$  given by,  $G(A) = \partial^\mu A_\mu - \omega(x)$  where  $\omega$  is some scalar function. This gives as the photon propagator as,

$$\tilde{D}_F^{\mu\nu}(k) = \frac{-i}{k^2 + i\epsilon} \left( g^{\mu\nu} - (1 - \xi) \frac{k^\mu k^\nu}{k^2} \right) \quad (1.11)$$

where the choice of  $\xi$  gives our gauge.

#### 1.4 WHAT ARE THE PROPERTIES OF GRASSMANN NUMBERS AND HOW ARE THEY USED TO QUANTIZE SPINOR FIELDS?

The Grassman numbers are anticommuting numbers with the following properties,

$$\begin{aligned}
 \theta\eta &= -\eta\theta \\
 \theta^2 &= 0 \\
 \int 1 d\theta &= 0 \\
 \int \theta d\theta &= 1 \\
 \int \frac{\partial f}{\partial \theta} d\theta &= 0
 \end{aligned}
 \tag{1.12}$$

To quantize spinor fields a grassman valued source field is introduced into  $Z[\eta, \bar{\eta}]$  such that,

$$Z[\eta, \bar{\eta}] = \int D\eta \int D\bar{\eta} e^{i \int d^4x [L + \bar{\eta}\psi + \bar{\psi}\eta]}.$$
(1.13)

A Grassman field is given by,

$$\eta(x) = \sum_i \eta_i \phi_i$$
(1.14)

where  $\eta_i$  is a grassman number and  $\phi_i$  is a field basis and in the equation above the base is 4-spinor Dirac fields.

#### 1.5 WHAT IS THE SUPERFICIAL DEGREE OF DIVERGENCE AND HOW CAN IT BE COMPUTED? USE QED AS AN EXAMPLE.

The superficial degree of divergence is for a Feynmann diagram's integral defined as,

$$D = (\text{power of momenta in numerator} - \text{power of momenta in denominator})$$
(1.15)

and tells us something about the divergence of a given diagram. In general d-dimensional QED this can be written as,

$$D = d + \frac{d-4}{2}V - \frac{d-2}{2}N_\gamma - \frac{d-1}{2}N_e$$
(1.16)

where V is the number of vertices and  $N_i$  the number of  $i$  = photon, electron external lines. D does not tell us everything about the divergence, three diagrams (c,f and g in Fig. 1.1) always diverge, D can only tell us if a given diagram diverge or not if the diagram does not contain any of those three.

#### 1.6 HOW DOES RENORMALIZED PERTURBATION THEORY RELATE THE BARE AND PHYSICAL MASSES?

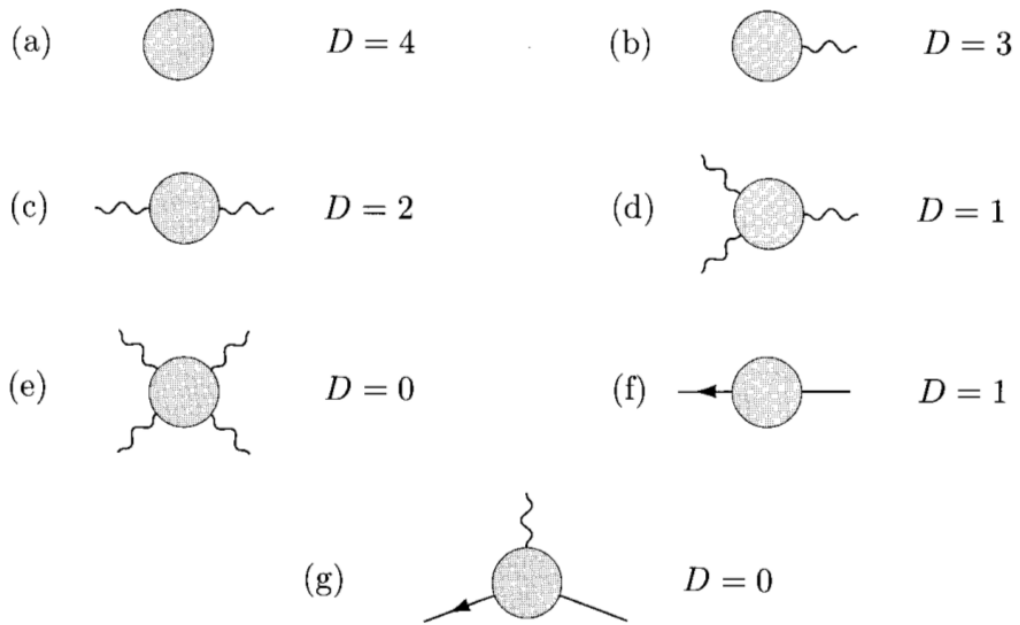


Figure 1.1: The seven QED amplitudes whose superficial degree of divergence ( $D$ ) is  $\geq 0$ .