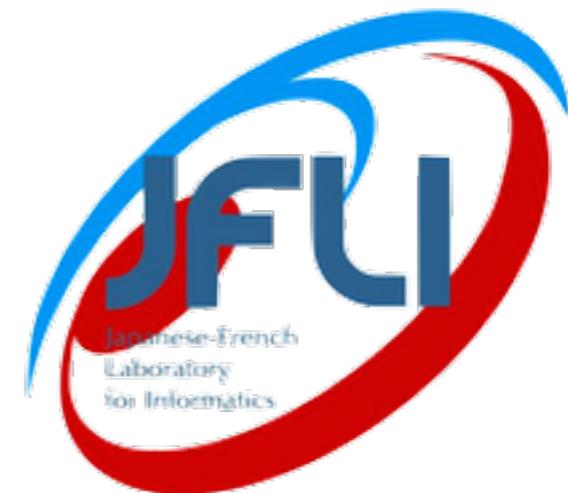


Relational Differential Dynamic Logic

Methods and Tools for Distributed Hybrid Systems
Amsterdam, 26/08/19

Jérémie Dubut
National Institute of Informatics
Japanese-French Laboratory of Informatics



Collaborations

Joint work with:

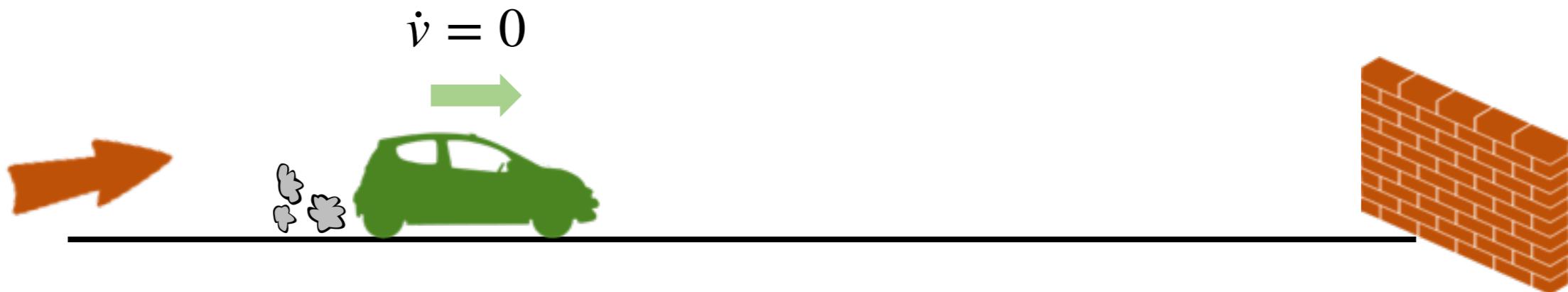
- from Tokyo: Ichiro Hasuo, Akihisa Yamada, David Sprunger, and Shin-ya Katsumata
- from France: Juraj Kolčák

Initiated by discussions with Kenji Kamijo, Yoshiyuki Shinya, and Takamasa Suetomi from Mazda Motor Corporation

Sources:

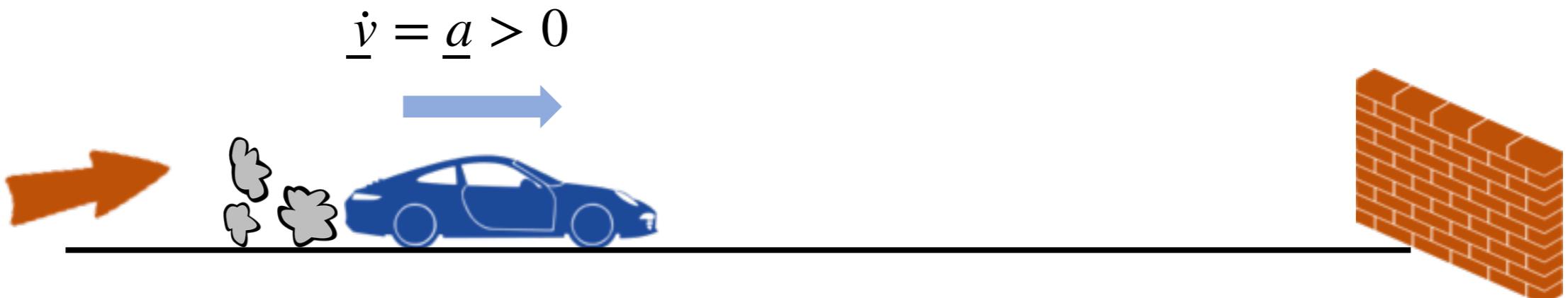
- J. Kolčák, I. Hasuo, J. Dubut, S. Katsumata, D. Sprunger, A. Yamada, Relational Differential Dynamic Logic. Preprint arXiv:1903.00153.
- some implementation on GitHub

Basic example: simplified ISO26262



In which cases will the vehicle crash hard?

Monotonicity property?



When $\bar{a} < \underline{a}$ which vehicle crash harder?

Elementary proof

Consider the following easy dynamics:

$$\underline{\dot{x}} = \underline{v}$$
$$\underline{\dot{v}} = \underline{a}$$



$$\dot{\bar{x}} = \bar{v}$$
$$\dot{\bar{v}} = \bar{a}$$



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$$\dot{\bar{x}} = \bar{v}$$
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Solving the equations:

$$\underline{v} = \underline{a} \cdot \underline{t} + \underline{v}_0$$

$$\bar{v} = \bar{a} \cdot \bar{t} + \bar{v}_0$$

$$\underline{x} = \frac{\underline{a}}{2} \cdot \underline{t}^2 + \underline{v}_0 \cdot \underline{t}$$

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Solving the equations:

$$\underline{v} = \underline{a} \cdot \underline{t} + \underline{v}_0 \quad \bar{v} = \bar{a} \cdot \bar{t} + \bar{v}_0$$

$$\underline{x} = \frac{\underline{a}}{2} \cdot \underline{t}^2 + \underline{v}_0 \cdot \underline{t} \quad \bar{x} = \frac{\bar{a}}{2} \cdot \bar{t}^2 + \bar{v}_0 \cdot \bar{t}$$

The time at which the vehicles reach the position x is:

$$\underline{t}(x) = \frac{\sqrt{\underline{v}_0^2 + 2\underline{a}x} - \underline{v}_0}{\underline{a}} \quad \bar{t}(x) = \frac{\sqrt{\bar{v}_0^2 + 2\bar{a}x} - \bar{v}_0}{\bar{a}}$$

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If $\underline{a} \leq \bar{a}$ and
 $\underline{v}_0 \leq \bar{v}_0$ then
 $\underline{v}(x) \leq \bar{v}(x)$
and the blue car
crashes harder!

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Differential dynamic logic in a nutshell

- A Hoare-triples-style syntax to formalise properties of hybrid system
- A sequent calculus to implement proofs of those properties
- A tool: KeYmaeraX

Ref: A. Platzer's group <http://symbolaris.com>

$$\Gamma \vdash [\alpha] P$$

Γ, P : sets of first order formulae of real arithmetic

α : hybrid program

$$\alpha ::= ?P \mid \alpha; \alpha \mid \dot{x} = f(x) \& Q \mid \alpha^\star \mid x := e \mid \dots$$

Invariants

$$\frac{\Gamma \vdash \mathbf{Inv} \quad \mathbf{Inv} \vdash [\alpha] \mathbf{Inv} \quad \mathbf{Inv} \vdash P}{\Gamma \vdash [\alpha] P} (\mathbf{Inv})$$

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To prove the statement $\Gamma \vdash [\alpha] P$

Invariants

It is enough to find an invariant such that:

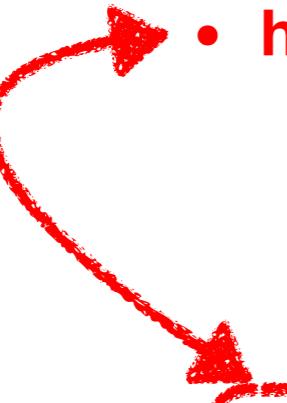
$$\frac{\Gamma \vdash \text{Inv} \quad \text{Inv} \vdash [\alpha] \text{ Inv} \quad \text{Inv} \vdash P}{\Gamma \vdash [\alpha] P} (\text{Inv})$$

To prove the statement $\Gamma \vdash [\alpha] P$

Invariants

It is enough to find an invariant such that:

- holds initially


$$\frac{\Gamma \vdash \text{Inv} \quad \text{Inv} \vdash [\alpha] \text{ Inv} \quad \text{Inv} \vdash P}{\Gamma \vdash [\alpha] P} (\text{Inv})$$

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Invariants

It is enough to find an invariant which:

- holds initially
- implies the post-condition

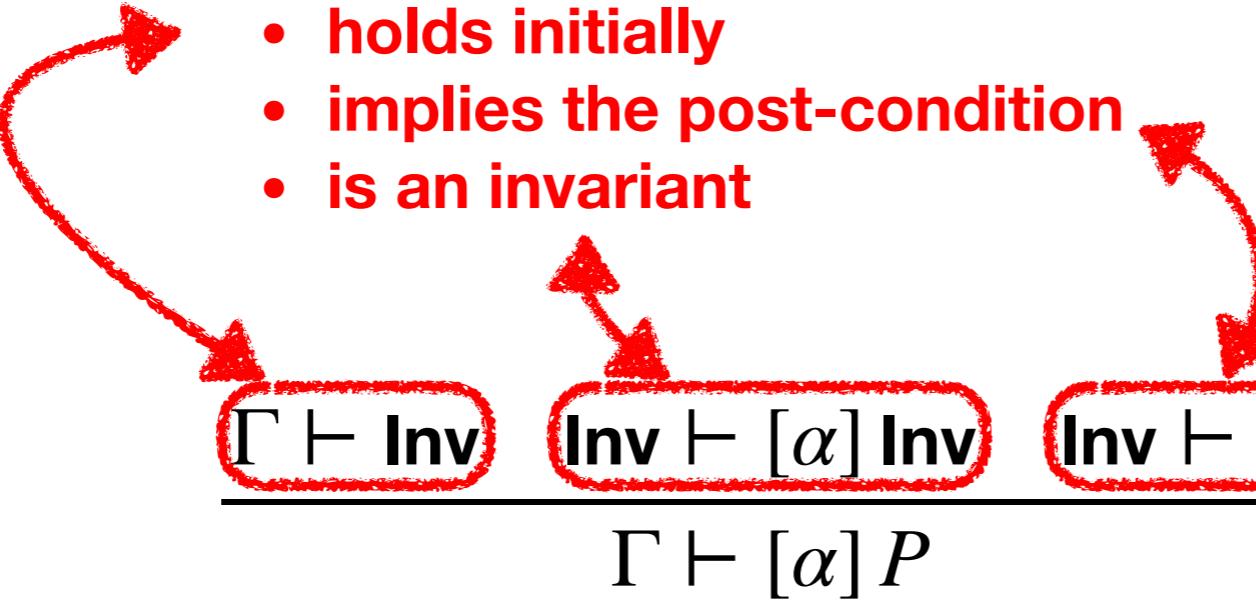
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To prove the statement $\Gamma \vdash [\alpha] P$

Invariants

It is enough to find an invariant which:

- holds initially
- implies the post-condition
- is an invariant

$$\frac{\Gamma \vdash \text{Inv} \quad \text{Inv} \vdash [\alpha] \text{ Inv} \quad \text{Inv} \vdash P}{\Gamma \vdash [\alpha] P} (\text{Inv})$$


To prove the statement $\Gamma \vdash [\alpha] P$

Loop invariants

$$\frac{\Gamma \vdash \mathbf{Inv} \quad \mathbf{Inv} \vdash [\alpha] \mathbf{Inv} \quad \mathbf{Inv} \vdash P}{\Gamma \vdash [\alpha^\star] P} (\mathbf{LI})$$

Differential invariants?

$$\dot{\mathbf{x}} = \mathbf{e} \ \& \ Q \quad \simeq \quad (\mathbf{?}Q; \mathbf{x} := \mathbf{x} + dt \cdot \mathbf{e})^{\star}; \mathbf{?}Q$$

$$\frac{\Gamma, Q \vdash \mathbf{Inv} \quad \mathbf{Inv}, Q \vdash \mathbf{Inv}(\mathbf{x} \leftarrow \mathbf{x} + dt \cdot \mathbf{e}) \quad \mathbf{Inv} \vdash P}{\Gamma \vdash [\dot{\mathbf{x}} = \mathbf{e} \ \& \ Q]P} \quad (\mathbf{dtI})$$

Assume that $P = \mathbf{Inv} \equiv f = 0$. We want something to ensure:

$$f(\omega) = 0 \Rightarrow f(\omega + dt \cdot \mathbf{e}(\omega)) = 0$$

It is enough to require that f is constant along the dynamics, that is, if ψ is a solution of $\dot{\mathbf{x}} = \mathbf{e}$, then $K : t \mapsto f(\psi(t))$ is constant, that is, its derivative is zero.

$$\dot{K}(t) = \sum_{x \in \mathbf{X}} \frac{\partial f}{\partial x}(\psi(t)) \cdot \dot{\psi}(t) = \sum_{x \in \mathbf{X}} \frac{\partial f}{\partial x}(\psi(t)) \cdot \mathbf{e}_x(\psi(t))$$

So it is enough that the function $\mathcal{L}_{\mathbf{e}} f = \sum_{x \in \mathbf{X}} \frac{\partial f}{\partial x} \cdot \mathbf{e}_x$ to be zero along the dynamics.

Differential invariants

$$\frac{\Gamma, Q \vdash f = 0 \quad \Gamma \vdash [\dot{x} = e \& Q] \mathcal{L}_e f = 0}{\Gamma \vdash [\dot{x} = e \& Q] f = 0} \text{(DI)}$$

Monotonicity property?

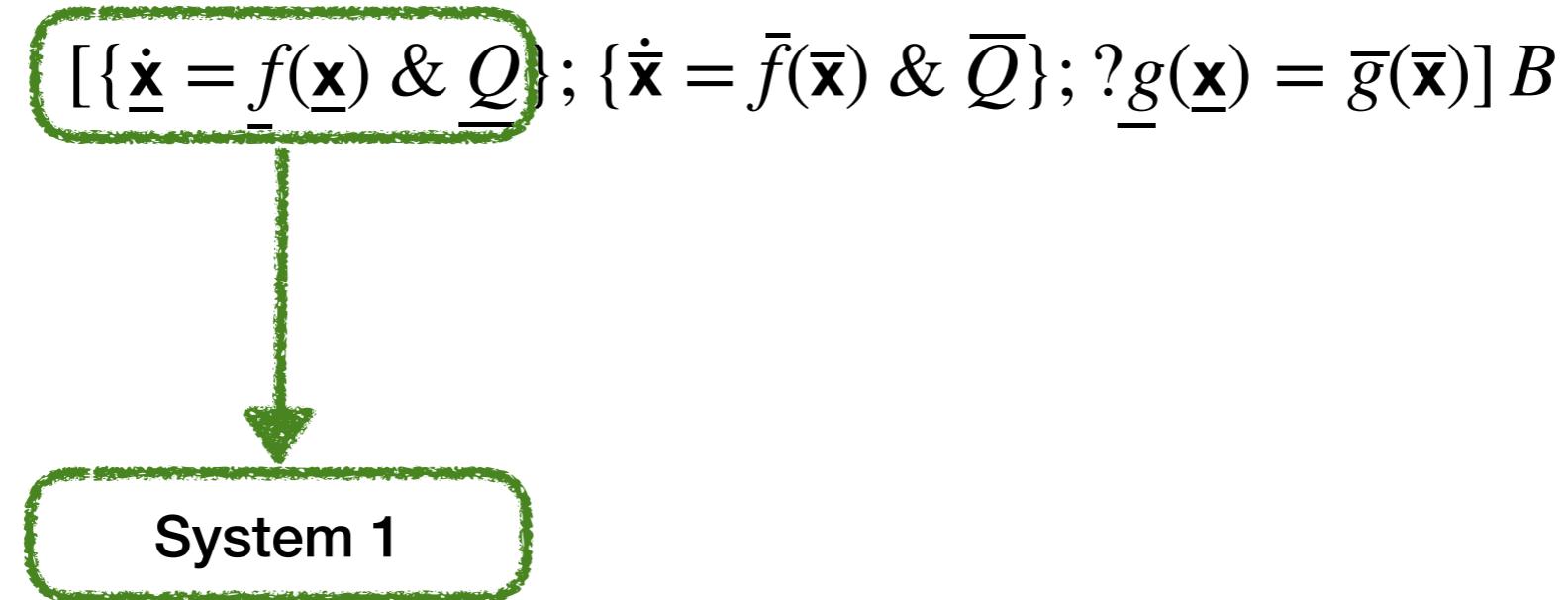


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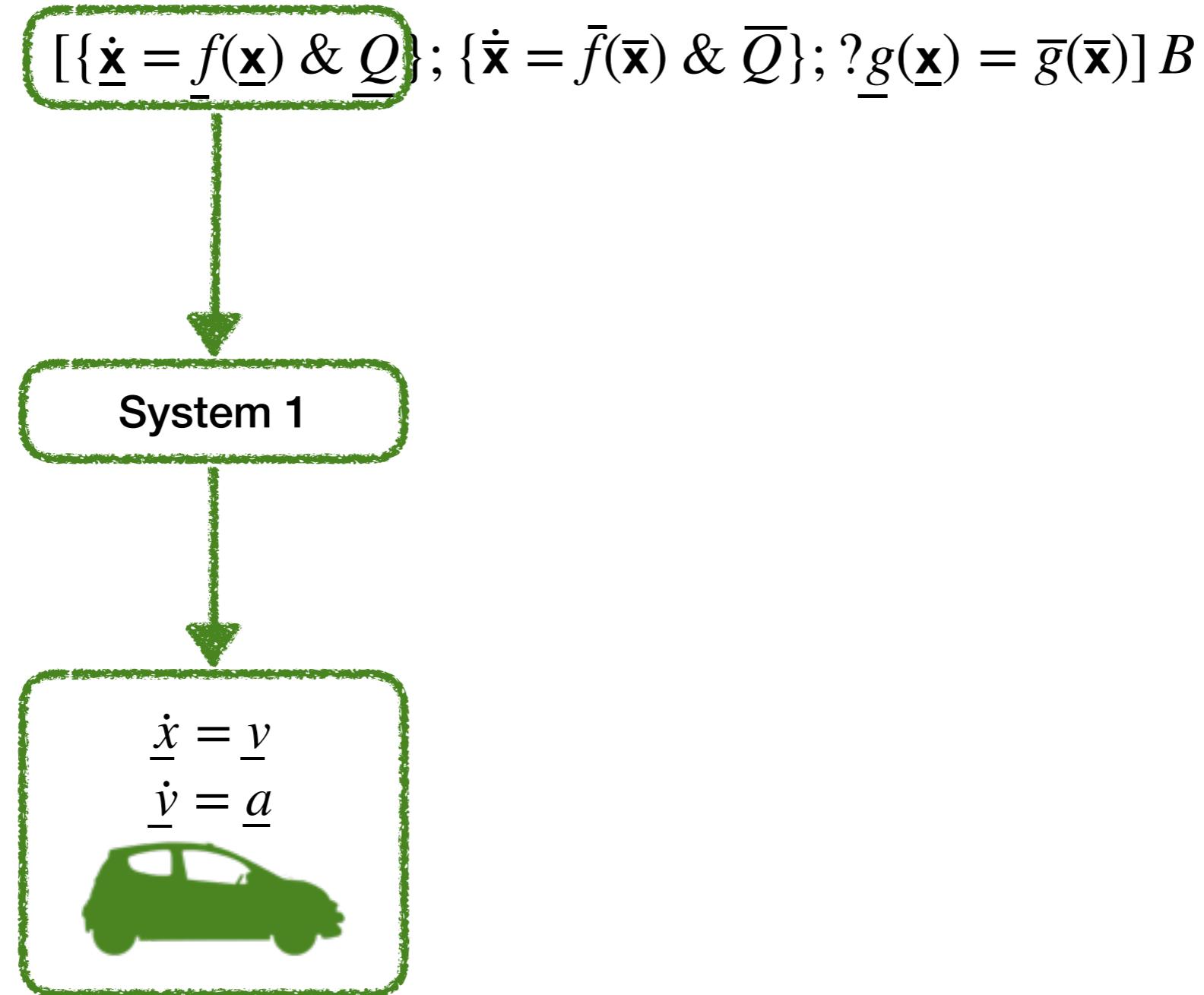
Relational formulae

$$[\{\dot{\underline{\mathbf{x}}} = \underline{f}(\underline{\mathbf{x}}) \ \& \ \underline{Q}\}; \{\dot{\bar{\mathbf{x}}} = \bar{f}(\bar{\mathbf{x}}) \ \& \ \bar{Q}\}; ?\underline{g}(\underline{\mathbf{x}}) = \bar{g}(\bar{\mathbf{x}})] B$$

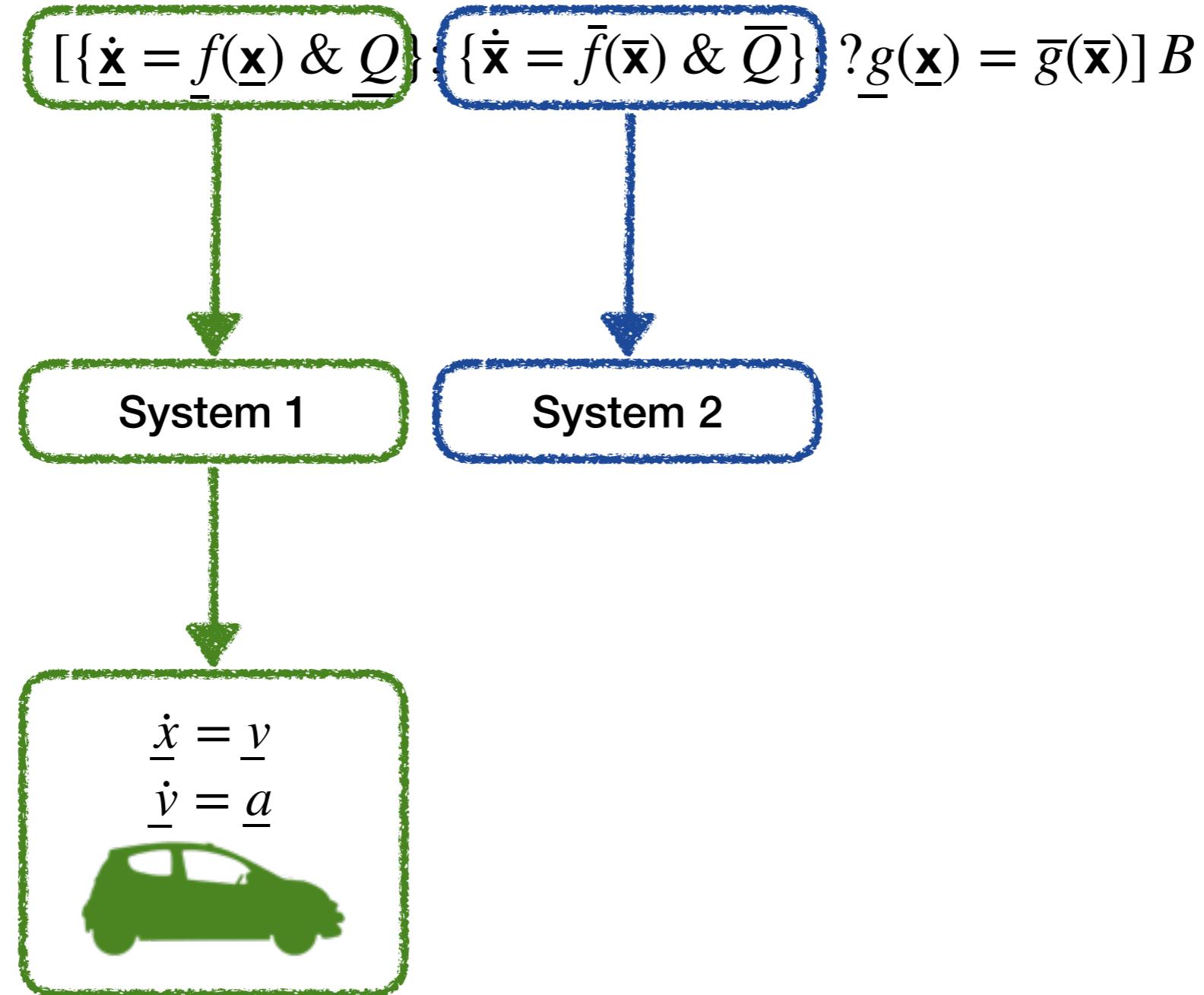
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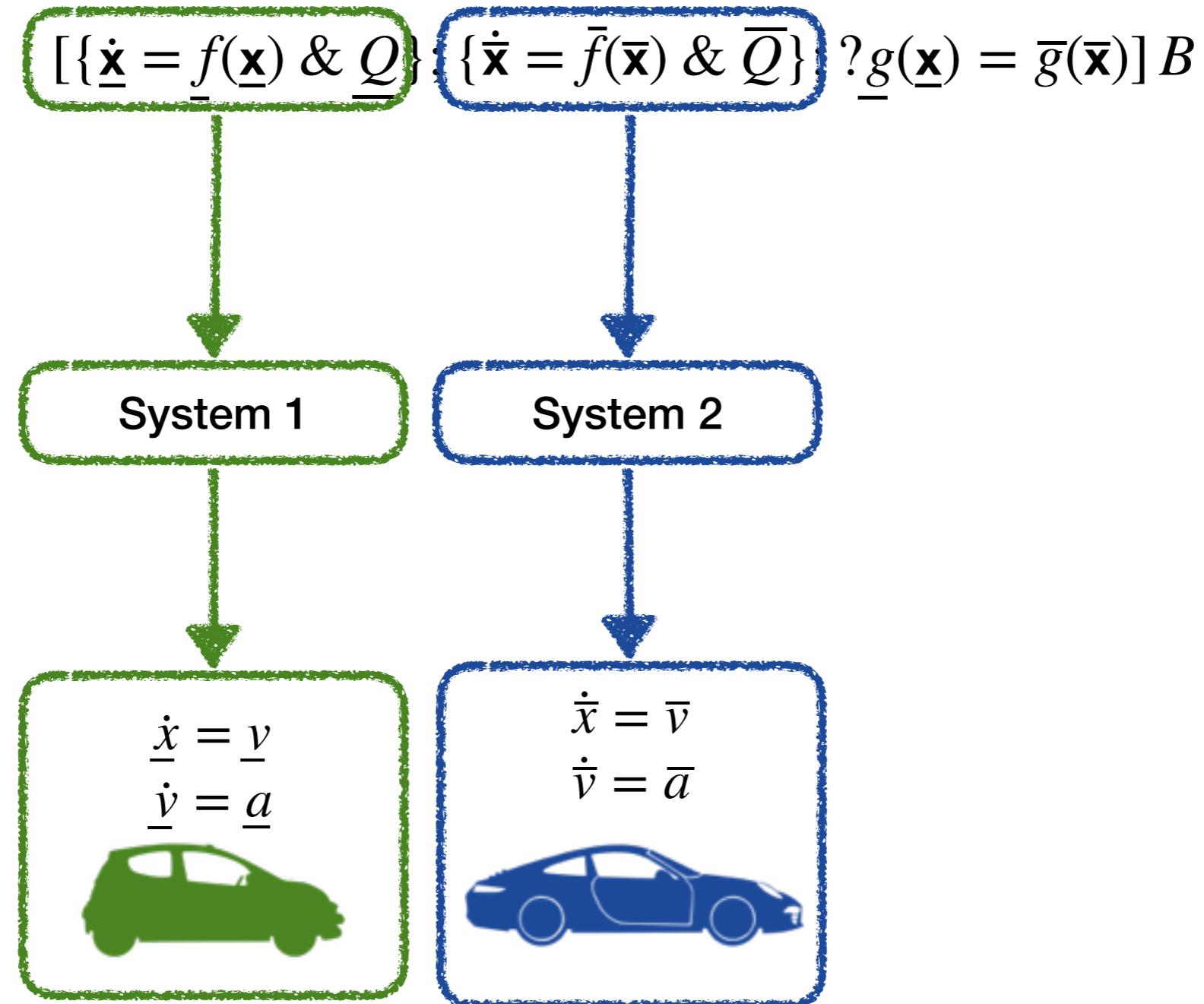
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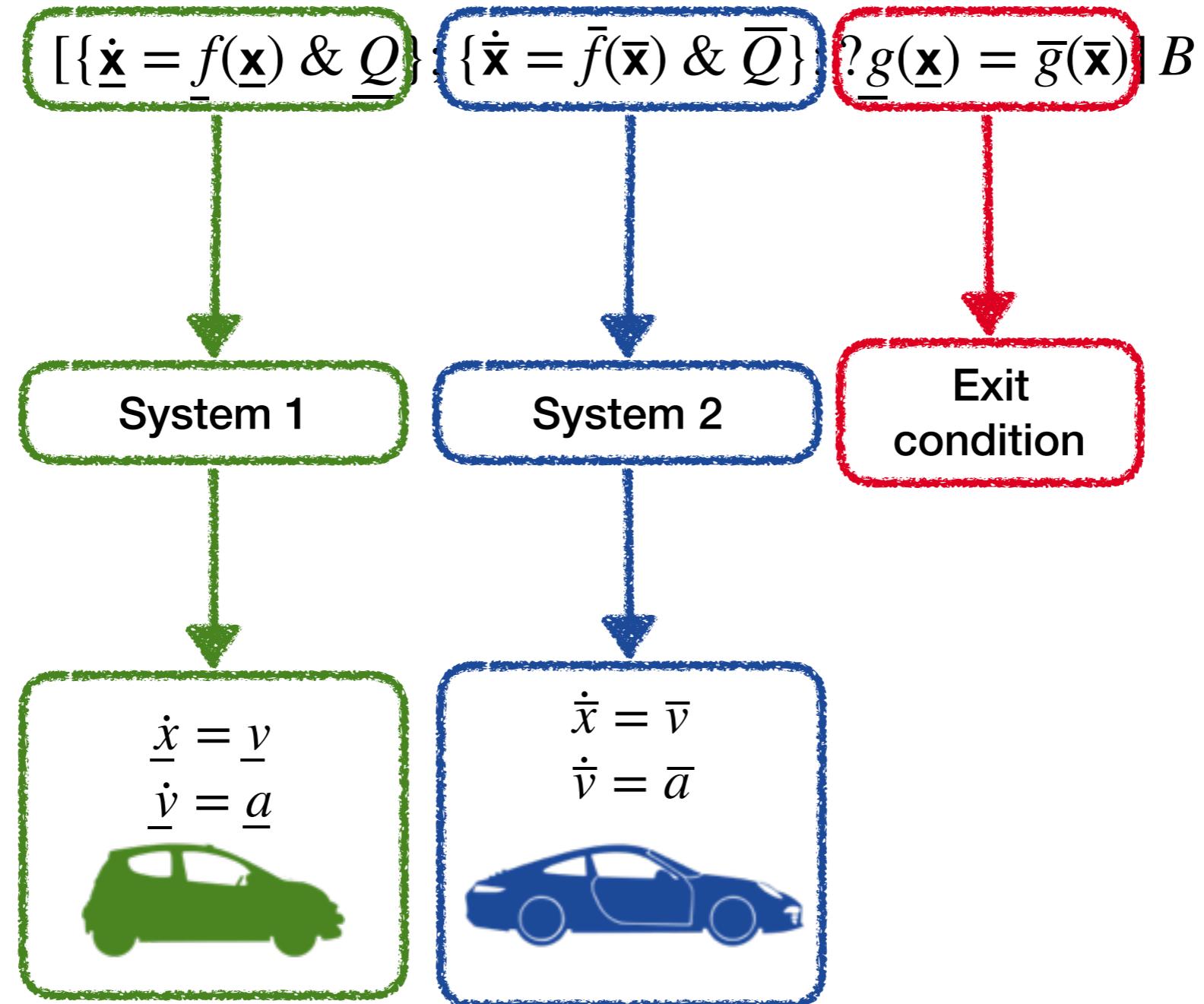
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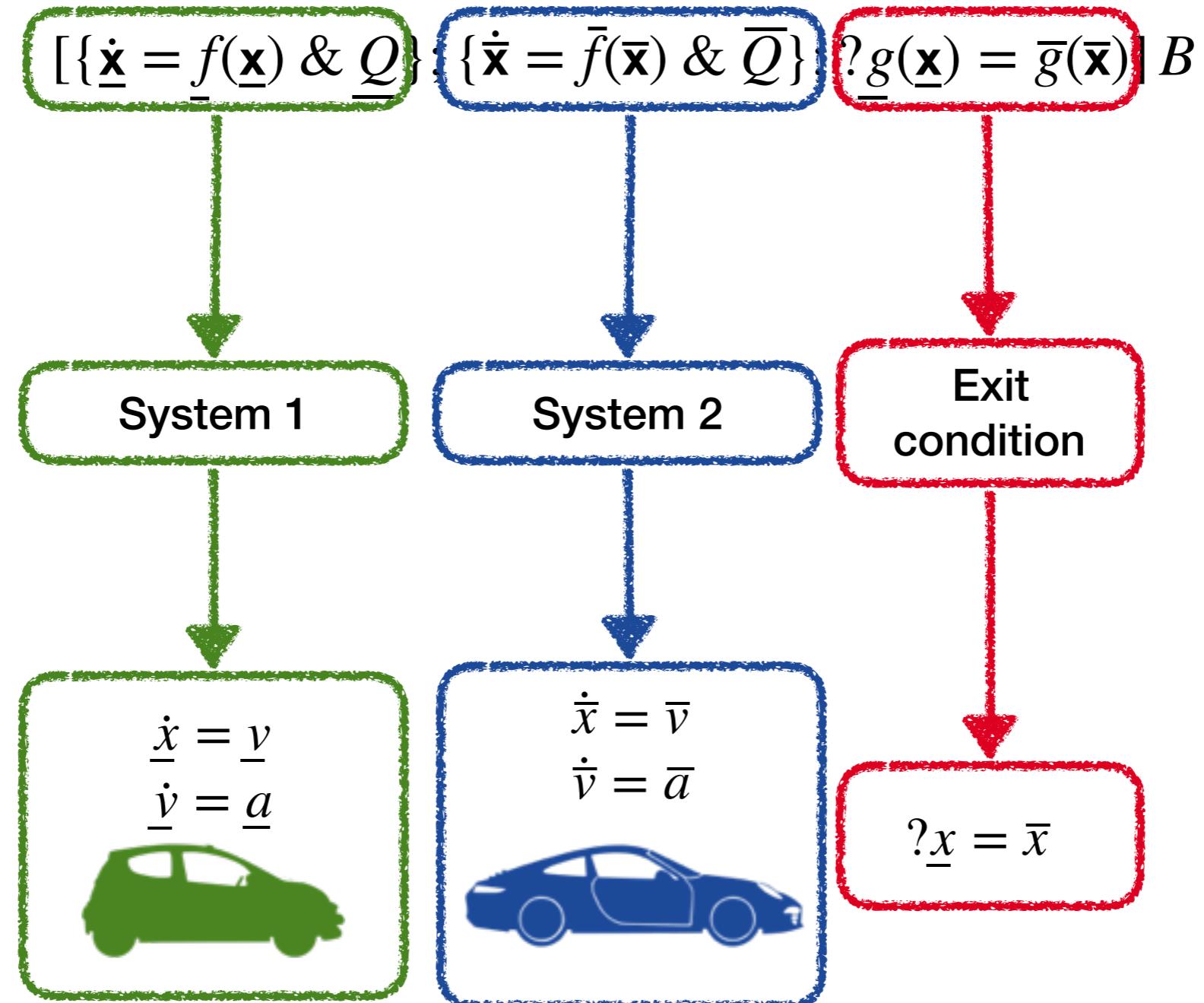
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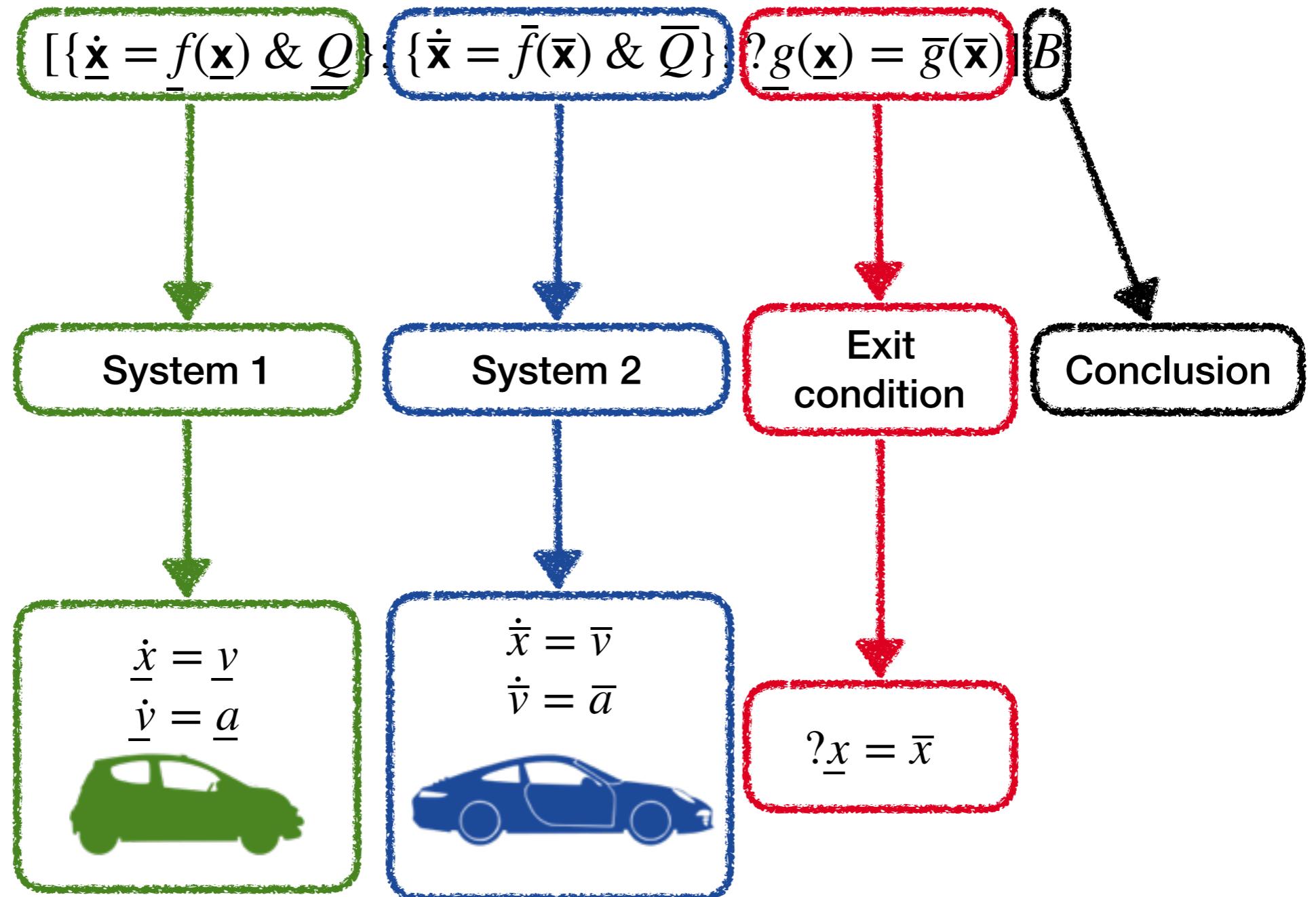
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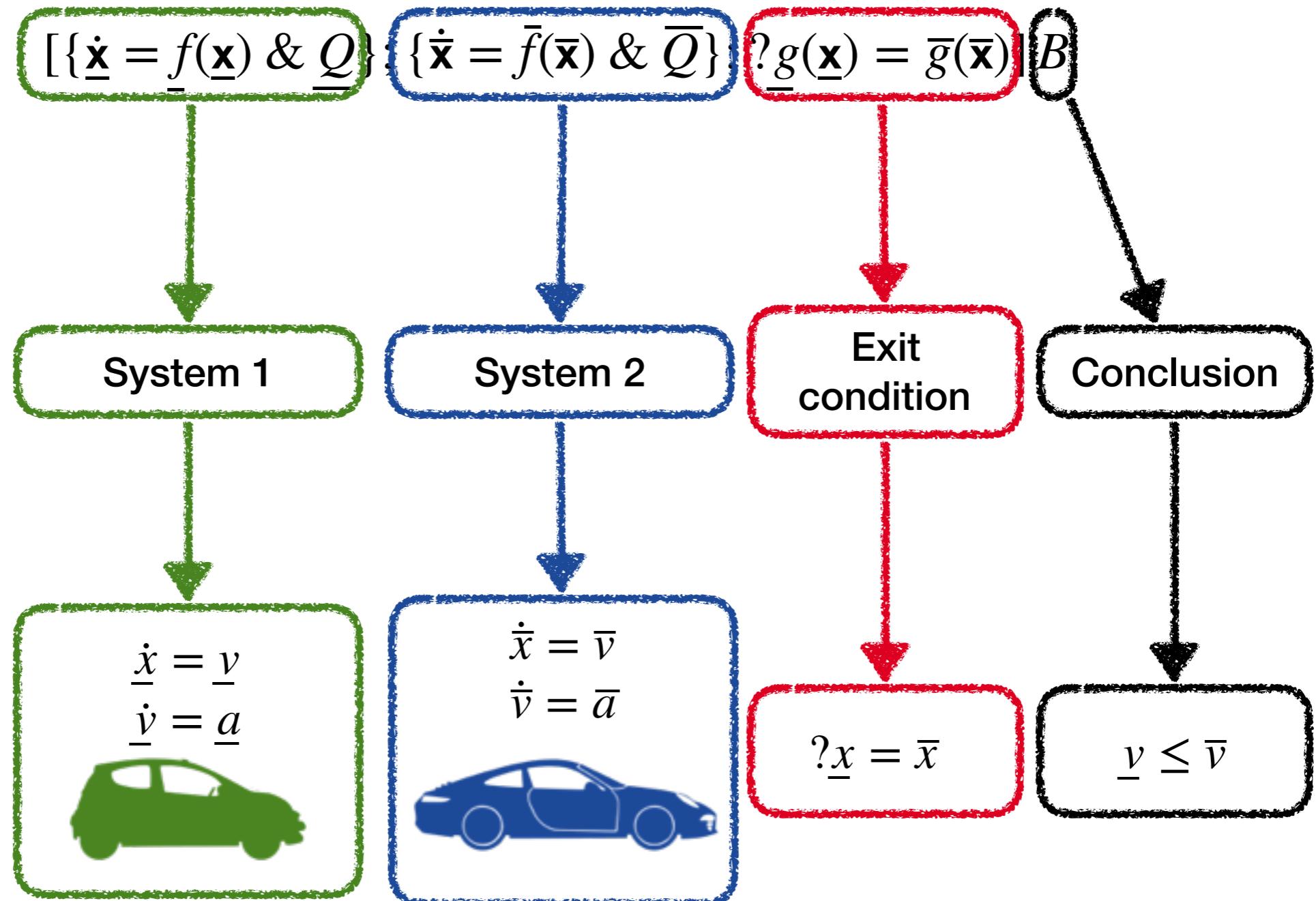
Relational formulae



Relational formulae



Relational formulae



Relational invariant

$$\frac{\Gamma, \underline{Q}, \overline{Q} \vdash \mathbf{Inv} \quad \mathbf{Inv} \vdash [\underline{\delta}; \overline{\delta}; ?E] \mathbf{Inv} \quad \mathbf{Inv}, E \vdash B}{\Gamma \vdash [\underline{\delta}; \overline{\delta}; ?E] B} (\mathbf{RI})$$

$$\underline{\delta} \equiv \dot{\underline{x}} = \underline{f}(\underline{x}) \ \& \ \underline{Q}$$

$$\overline{\delta} \equiv \dot{\bar{x}} = \bar{f}(\bar{x}) \ \& \ \overline{Q}$$

$$E \equiv \underline{g}(\underline{x}) = \overline{g}(\bar{x})$$

Relational invariant, for ISO26262

We know that:

$$\underline{v}(x) = \sqrt{\underline{v}_0^2 + 2\underline{a}x}$$

$$\bar{v}(x) = \sqrt{\bar{v}_0^2 + 2\bar{a}x}$$

So:

$$\frac{\underline{v}(x)^2 - \underline{v}_0^2}{2\underline{a}} = x = \frac{\bar{v}(x)^2 - \bar{v}_0^2}{2\bar{a}}$$

And then:

$$R \equiv \bar{a}(\underline{v}^2 - \underline{v}_0^2) = \underline{a}(\bar{v}^2 - \bar{v}_0^2)$$

is a relational invariant.

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is a relational invariant.

That is, one has to prove the following statements:

- $\underline{a} \leq \bar{a}, \underline{v} \geq \underline{v}_0, \bar{v} \geq \bar{v}_0, R \vdash \bar{v} \geq \underline{v}$
- $\underline{v} = \underline{v}_0, \bar{v} = \bar{v}_0 \vdash R$
- $R \vdash [\{\dot{x} = \underline{v}, \dot{v} = \underline{a}\}; \{\dot{x} = \bar{v}, \dot{v} = \bar{a}\}; ?_x = \bar{x}] R$

The invariant implies
the property
(easy proof)

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The invariant holds initially
(easy proof)

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We know that:

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So:

$$\frac{\underline{v}(x)^2 - \underline{v}_0^2}{2\underline{a}} = x = \frac{\bar{v}(x)^2 - \bar{v}_0^2}{2\bar{a}}$$

And then:

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The invariant is
preserved by the
dynamics
(easy proof?)

Problems to tackle

We want a method that:

- does not require the solutions of the differential equations in any way
- transforms the two dynamics into one unique, so that we can use known methods from differential invariants.

Two systems into one?

What we have now:

- one system on $\underline{x}, \underline{v}$
- one system on \bar{x}, \bar{v}

that take different times to arrive at a particular position.

What we want:

- one system on $\underline{x}, \underline{v}, \bar{x}, \bar{v}$

such that the positions are synchronized, that is, at all time t :

$$\underline{x}(t) = \bar{x}(t)$$

Reparametrisation of dynamics

Crucial idea: reparametrise the time of \bar{x}, \bar{v}

Time stretch function: derivable function $k : \mathbb{R} \rightarrow \mathbb{R}$ with

$$\dot{k} > 0$$

Reparamatrised dynamics: let $\dot{x} = f(x)$ be a differential equation.
It reparametrisation by k is $\dot{x} = \dot{k}(t) \cdot f(x)$.

x is a solution of $\dot{x} = f(x)$ iff
 $x \circ k$ is a solution of $\dot{x} = \dot{k}(t) \cdot f(x)$.

Is it OK?

Fix some initial condition x_0 . Then

$$\{x(t) \mid x \text{ sol. of } \dot{x} = f(x) \text{ at } x_0\} = \{x(t) \mid x \text{ sol. of } \dot{x} = k(t) \cdot f(x) \text{ at } x_0\}.$$

Why is it OK then?

- we do not care about « at time t , the vehicle is at position x with speed v »
- we care about « at position x , the vehicle has speed v »

Which reparametrisation to choose?

Fix $\underline{v}_0, \bar{v}_0$ and note $(\underline{\psi}_x, \underline{\psi}_v)$ the solution of $\dot{\underline{x}} = \underline{v}, \dot{\underline{v}} = \underline{a}$ at $\underline{x} = 0, \underline{v} = \underline{v}_0$, $(\bar{\psi}_x, \bar{\psi}_v)$ the solution of $\dot{\bar{x}} = \bar{v}, \dot{\bar{v}} = \bar{a}$ at $\bar{x} = 0, \bar{v} = \bar{v}_0$.

We have a time stretch function $k : \mathbb{R} \longrightarrow \mathbb{R}$ such that for every x

$$k(\underline{t}(x)) = \bar{t}(x)$$

given by:

$$k(t) = \frac{\sqrt{at^2 + 2\underline{v}_0\bar{a}t + \bar{v}_0^2} - \bar{v}_0}{\bar{a}}$$

Which reparametrisation to choose?

Fix $\underline{v}_0, \bar{v}_0$ and note $(\underline{\psi}_x, \underline{\psi}_v)$ the solution of $\dot{\underline{x}} = \underline{v}, \dot{\underline{v}} = \underline{a}$ at $\underline{x} = 0, \underline{v} = \underline{v}_0$, $(\bar{\psi}_x, \bar{\psi}_v)$ the solution of $\dot{\bar{x}} = \bar{v}, \dot{\bar{v}} = \bar{a}$ at $\bar{x} = 0, \bar{v} = \bar{v}_0$.

We have a time stretch function $k : \mathbb{R} \longrightarrow \mathbb{R}$ such that for every x

$$k(\underline{t}(x)) = \bar{t}(x)$$

So we want to look at:

$$\dot{\underline{x}} = \underline{v}, \dot{\underline{v}} = \underline{a}, \dot{\bar{x}} = \dot{k}(t) \cdot \bar{v}, \dot{\bar{v}} = \dot{k}(t) \cdot \bar{a}$$

But we have:

$$\bar{\psi}_x(k(t)) = \underline{\psi}_x(t)$$

Then:

$$\dot{k}(t) \cdot \bar{\psi}_v(k(t)) = \underline{\psi}_v(t)$$

So we want to look at:

$$\dot{\underline{x}} = \underline{v}, \dot{\underline{v}} = \underline{a}, \dot{\bar{x}} = \frac{\underline{v}}{\bar{v}} \cdot \bar{v}, \dot{\bar{v}} = \frac{\underline{v}}{\bar{v}} \cdot \bar{a}$$

How to generalize to relational formulae?

In general, from

$$\dot{\underline{x}} = \underline{f}(\underline{x})$$

and

$$\dot{\bar{x}} = \bar{f}(\bar{x})$$

under the exit condition

$$\underline{g}(\underline{x}) = \bar{g}(\bar{x})$$

we consider:

$$\dot{\underline{x}} = \underline{f}(\underline{x}), \dot{\bar{x}} = \frac{\mathcal{L}_{\underline{f}} \underline{g}}{\mathcal{L}_{\bar{f}} \bar{g}} \cdot \bar{f}(\bar{x})$$

Synchronisation rule

$$\frac{\Gamma, \underline{Q}, \bar{Q} \vdash E \quad \Gamma \vdash [\underline{\delta}] \mathcal{L}_{\underline{f}} \underline{g} > 0 \quad \Gamma \vdash [\bar{\delta}] \mathcal{L}_{\bar{f}} \bar{g} > 0 \quad \Gamma \vdash [\delta] B}{\Gamma \vdash [\underline{\delta}; \bar{\delta}; ?E] B} (\mathbf{Syn})$$

$$\underline{\delta} \equiv \dot{\underline{x}} = \underline{f}(\underline{x}) \ \& \ \underline{Q}$$

$$\bar{\delta} \equiv \dot{\bar{x}} = \bar{f}(\bar{x}) \ \& \ \bar{Q}$$

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$$\bar{\delta} \equiv \dot{\underline{x}} = \underline{f}(\underline{x}), \dot{\bar{x}} = \frac{\mathcal{L}_{\underline{f}} \underline{g}}{\mathcal{L}_{\bar{f}} \bar{g}} \cdot \bar{f}(\bar{x}) \ \& \ \underline{Q} \wedge \bar{Q}$$

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- $\underline{a} \leq \bar{a}, \underline{v} \geq \underline{v}_0, \bar{v} \geq \bar{v}_0, R \vdash \bar{v} \geq \underline{v}$
- $\underline{v} = \underline{v}_0, \bar{v} = \bar{v}_0 \vdash R$
- $R \vdash [\{\dot{x} = \underline{v}, \dot{v} = \underline{a}\}; \{\dot{x} = \bar{v}, \dot{v} = \bar{a}\}; ?\underline{x} = \bar{x}] R$

The invariant is
preserved by the
dynamics
(easy proof?)

Relational invariant, for ISO26262

Let's prove the following statement with the (**Syn**) rule:

$$R \vdash [\{\dot{x} = \underline{v}, \dot{v} = \underline{a}\}; \{\dot{\bar{x}} = \bar{v}, \dot{\bar{v}} = \bar{a}\}; ?\underline{x} = \bar{x}] R$$

with $R \equiv \bar{a}(\underline{v}^2 - \underline{v}_0^2) = \underline{a}(\bar{v}^2 - \bar{v}_0^2)$, that is:

- $\underline{v} > 0, \underline{a} \geq 0 \vdash [\dot{x} = \underline{v}, \dot{v} = \underline{a}] \underline{v} > 0$ (easy with differential invariant)
- $\bar{v} > 0, \bar{a} \geq 0 \vdash [\dot{\bar{x}} = \bar{v}, \dot{\bar{v}} = \bar{a}] \bar{v} > 0$ (easy with differential invariant)
- $R \vdash [\{\dot{x} = \underline{v}, \dot{v} = \underline{a}, \dot{\bar{x}} = \frac{\underline{v}}{\bar{v}} \cdot \bar{v}, \dot{\bar{v}} = \frac{\underline{v}}{\bar{v}} \cdot \bar{a}\}] R$

using differential invariant rule:

$$R \vdash [\{\dot{x} = \underline{v}, \dot{v} = \underline{a}, \dot{\bar{x}} = \frac{\underline{v}}{\bar{v}} \cdot \bar{v}, \dot{\bar{v}} = \frac{\underline{v}}{\bar{v}} \cdot \bar{a}\}] 2\bar{a}\underline{v}\underline{a} = 2\underline{a}\bar{v}\frac{\underline{v}}{\bar{v}}\bar{a}$$

Summary

Guideline:

- start with two independent systems and compare them under some conditions
- synchronize them by reparametrising one of them using the (**Syn**) rule
- use usual invariant techniques from dL

Case studies:

- monotonicity properties
- abstraction