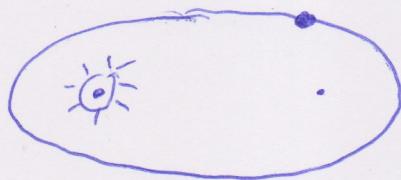


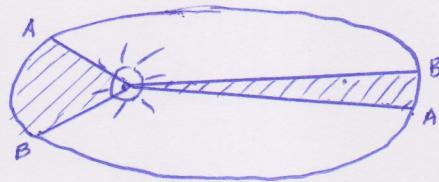
5. Two-Body Problem

Kepler's Laws - made by observation of Mars
(early 17th century)

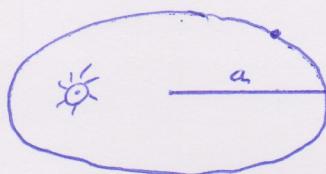
1. The orbit of each planet is an ellipse with the sun at one focus.



2. The line which joins a planet to the sun sweeps over equal areas in equal intervals of time



3. The squares of the periods of the planets are proportional to the cubes of the semi-major axes of their orbits.

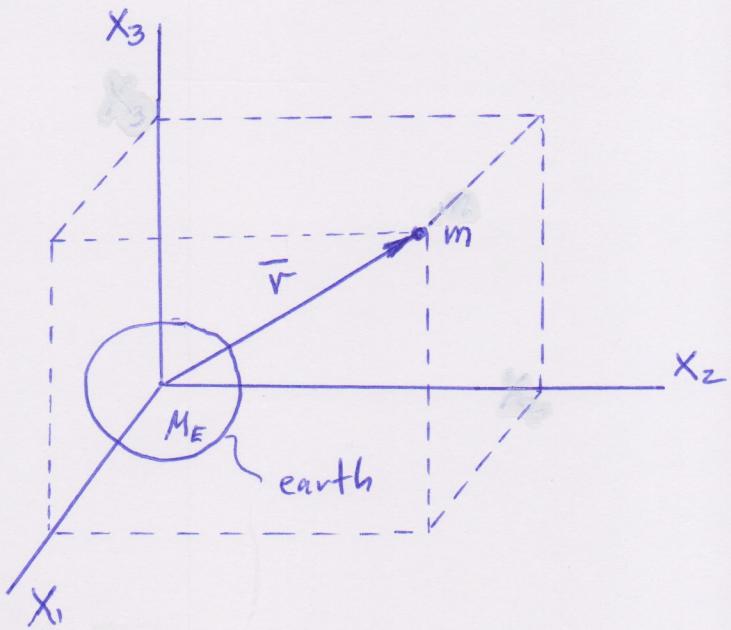


a = semi-major axis

T = period (time to complete one orbit)

$$T^2 \sim a^3$$

Satellite in orbit About Earth (neglecting atmospheric effects)



$$M_E = 5.974 \times 10^{24} \text{ kg}$$

$$\mu_E = GM_E \text{ (gravitational parameter)} \\ = 3.986 \times 10^5 \text{ km}^3/\text{sec}^2$$

Equatorial radius 6378 km

Polar radius 6357 km

Average radius 6368 km

In absence of all other bodies ($R_j = 0, j = 3, 4, 5, \dots$)
(4.9) reduces to

$$\frac{d^2\bar{r}}{dt^2} = -\frac{\mu}{r^3} \bar{r} \quad (5.1)$$

Will show that the angular momentum of the orbiting satellite remains constant

From cross product

$$\bar{r} \times \frac{d^2\bar{r}}{dt^2} = -\frac{\mu}{r^3} (\bar{r} \times \bar{r}) = 0 \quad (5.2)$$

Using

$$\frac{d}{dt} \left(\vec{r} \times \frac{d\vec{r}}{dt} \right) = \cancel{\frac{d\vec{r}}{dt} \times \frac{d\vec{r}}{dt}}^0 + \vec{r} \times \frac{d^2\vec{r}}{dt^2}$$

Therefore (5.2) becomes

$$\frac{d}{dt} \left(\vec{r} \times \frac{d\vec{r}}{dt} \right) = 0 \quad (5.3)$$

Integrate

$$\vec{r} \times \frac{d\vec{r}}{dt} = \vec{h} = \text{const.} \quad (\text{angular momentum}) \quad (5.4)$$

Note: \vec{h} is perpendicular to both \vec{r} & $\frac{d\vec{r}}{dt}$

Will now show that satellite trajectory is contained in a plane which passes through the origin.

Take dot product of (5.4) with \bar{r}

$$\underbrace{\left(\bar{r} \times \frac{d\bar{r}}{dt} \right)}_{0} \cdot \bar{r} = \bar{h} \cdot \bar{r}$$

$$0 \text{ since } \bar{r} \perp \underbrace{\left(\bar{r} \times \frac{d\bar{r}}{dt} \right)}_{\perp \bar{r} \therefore \frac{d\bar{r}}{dt}}$$

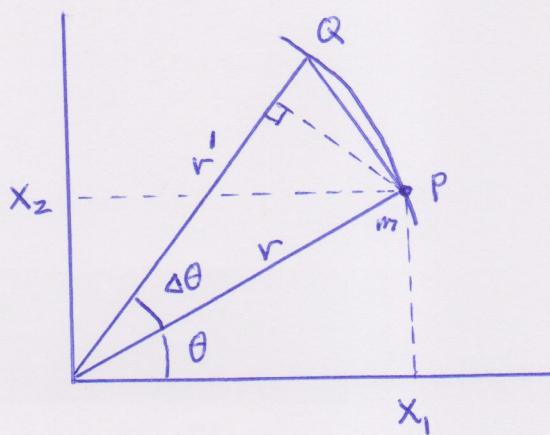
Therefore $\bar{h} \cdot \bar{r} = 0$

or $h_1 x_1 + h_2 x_2 + h_3 x_3 = 0 \quad (5.5)$

(equation of plane passing through the origin)

Proof of Kepler's 2nd Law

Suppose particle m is moving in the $x_1 x_2$ plane as shown.



Differential area swept out by radius vector as particle moves from P to Q is

$$\Delta A = \frac{1}{2} r' (r \sin \Delta \theta)$$

Time rate of change of this area is

$$\frac{\Delta A}{\Delta t} = \frac{1}{2} r' r \frac{\Delta \theta}{\Delta t} \cdot \frac{\sin \Delta \theta}{\Delta \theta}$$

Take limit as $\Delta t \rightarrow 0$

$$\frac{dA}{dt} = \frac{1}{2} r^2 \frac{d\theta}{dt} \quad (5.6)$$

From figure

$$\tan \theta = \frac{x_2}{x_1}$$

$$\theta = \tan^{-1} \frac{x_2}{x_1}$$

$$\frac{d\theta}{dt} = \frac{1}{1 + \left(\frac{x_2}{x_1}\right)^2} \frac{x_1 \dot{x}_2 - x_2 \dot{x}_1}{x_1^2}$$

$$\frac{d\theta}{dt} = \frac{\dot{x}_1 \dot{x}_2 - \dot{x}_2 \dot{x}_1}{x_1^2 + x_2^2}$$

Since $\bar{r} \times \frac{d\bar{v}}{dt} = \bar{h}$ (5.4)

$$\begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ x_1 & x_2 & 0 \\ \dot{x}_1 & \dot{x}_2 & 0 \end{vmatrix} = 0\hat{i} + 0\hat{j} + (\dot{x}_1 \dot{x}_2 - \dot{x}_2 \dot{x}_1)\hat{k}$$

$$= h_3 \hat{k} = h \hat{k} \quad (\text{since } h_1 = h_2 = 0)$$

$$\Rightarrow \dot{x}_1 \dot{x}_2 - \dot{x}_2 \dot{x}_1 = h \quad (\text{cancel } \hat{k})$$

Since $x_1^2 + x_2^2 = r^2$, get

$$\frac{d\theta}{dt} = \frac{\dot{x}_1 \dot{x}_2 - \dot{x}_2 \dot{x}_1}{x_1^2 + x_2^2} = \frac{h}{r^2} \quad (5.7)$$

Substitute (5.7) into (5.6)

$$\frac{dA}{dt} = \frac{1}{2} r^2 \frac{h}{r^2}$$

or

$$\boxed{\frac{dA}{dt} = \frac{1}{2} h = \text{constant}} \quad (\text{Kepler's 2nd law}) \quad (5.8)$$

Proof that the trajectory of the particle is a conic section (Kepler's 1st Law)

From (5.1)

$$\frac{d^2\bar{r}}{dt^2} = -\frac{\mu}{r^3} \bar{r}$$

$$\frac{d^2\bar{r}}{dt^2} \times \bar{h} = -\frac{\mu}{r^3} \bar{r} \times \bar{h} = -\mu \bar{r} \times (\bar{r} \times \frac{d\bar{r}}{dt})$$

Using the vector identity

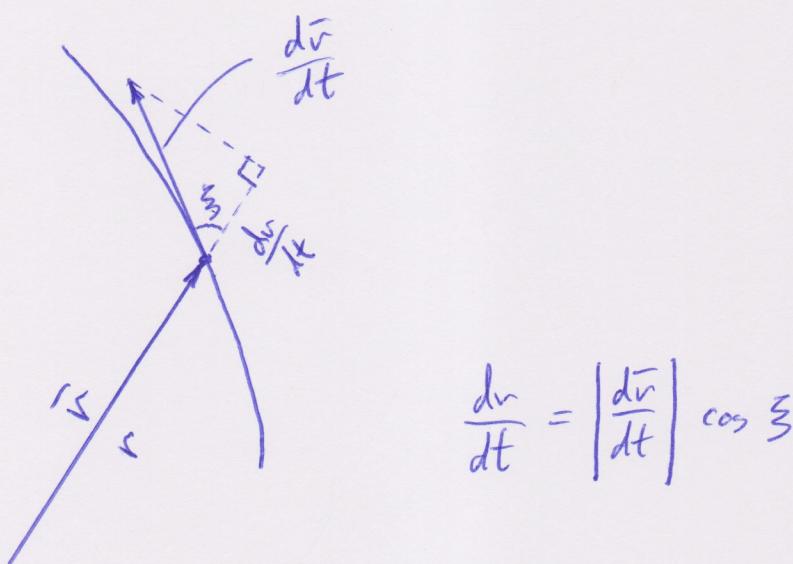
$$\bar{A} \times (\bar{B} \times \bar{C}) = (\bar{A} \cdot \bar{C}) \bar{B} - (\bar{A} \cdot \bar{B}) \bar{C}$$

$$\bar{r} \times (\bar{r} \times \frac{d\bar{r}}{dt}) = (\bar{r} \cdot \frac{d\bar{r}}{dt}) \bar{r} - \underbrace{(\bar{r} \cdot \bar{r})}_{r^2} \frac{d\bar{r}}{dt}$$

$$\frac{d^2\bar{r}}{dt^2} \times \bar{h} = \frac{\mu}{r^3} \left[r^2 \frac{d\bar{r}}{dt} - (\bar{r} \cdot \frac{d\bar{r}}{dt}) \bar{r} \right]$$

Note: $\frac{dr}{dt}$ is component of $\frac{d\bar{v}}{dt}$ in \bar{v} direction.

$$\frac{dr}{dt} \neq \left| \frac{d\bar{v}}{dt} \right|$$



$$\bar{r} \cdot \frac{d\bar{v}}{dt} = r \left| \frac{d\bar{v}}{dt} \right| \cos \xi = r \frac{dr}{dt}$$

Returning to previous equation

$$\begin{aligned} \frac{d^2\bar{v}}{dt^2} \times \bar{r} &= \frac{\mu}{r^3} \left[r^2 \frac{d\bar{v}}{dt} - r \frac{dr}{dt} \bar{r} \right] \\ &= \mu \frac{r \frac{d\bar{v}}{dt} - \frac{dr}{dt} \bar{r}}{r^2} = \mu \frac{d}{dt} \left(\frac{\bar{r}}{r} \right) \end{aligned}$$

Integrate

$$\frac{d\bar{v}}{dt} \times \bar{h} = \mu \left[\frac{\bar{r}}{r} + \bar{e} \right] \quad (5.9)$$

\bar{e} = constant of integration (eccentricity vector)
(lies in plane of trajectory)

Take dot product of (5.9) with \bar{r} and use
the identity $\bar{A} \cdot (\bar{B} \times \bar{C}) = (\bar{A} \times \bar{B}) \cdot \bar{C}$

$$\begin{aligned} \bar{r} \cdot \left(\frac{d\bar{v}}{dt} \times \bar{h} \right) &= \left(\bar{r} \times \frac{d\bar{v}}{dt} \right) \cdot \bar{h} = \bar{h} \cdot \bar{h} = h^2 \\ &= \bar{r} \cdot \mu \left[\frac{\bar{r}}{r} + \bar{e} \right] = \mu \left[\bar{r} \cdot \frac{\bar{r}}{r} + \bar{r} \cdot \bar{e} \right] \\ &= \mu [r + r e \cos \theta] \end{aligned}$$

where θ is the angle between \bar{r} and \bar{e} and is
called the true anomaly. In the literature,
the true anomaly is often denoted by f .

Solve for r

$$r = \frac{h^2/\mu}{1 + e \cos \theta}$$

Define the parameter p

$$p = \frac{h^2}{\mu} \quad (5.10)$$

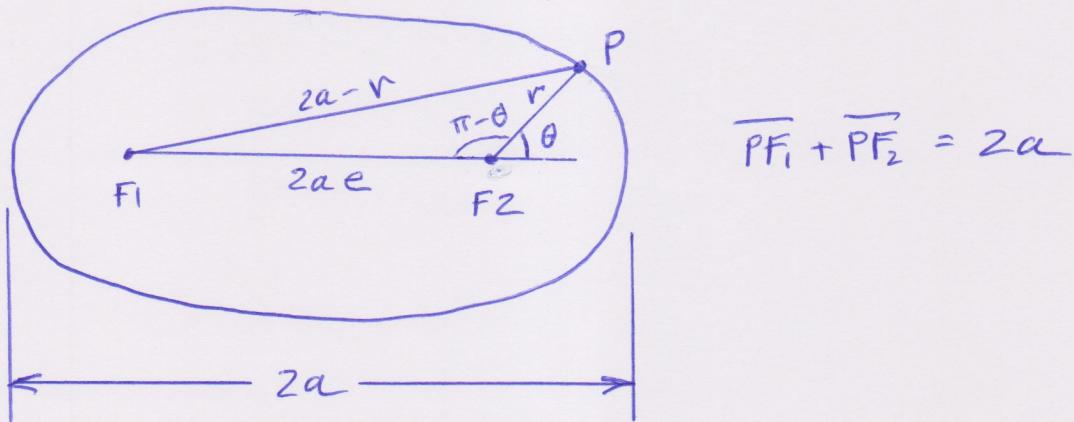
$$r = \frac{p}{1 + e \cos \theta} \quad (5.11)$$

General equation for conic section in polar coordinates
 (Proof of Kepler's 1st Law)

e is the eccentricity of the conic section

Geometric Considerations

A. Ellipse $e < 1$



$$\overrightarrow{PF_1} + \overrightarrow{PF_2} = 2a$$

$$(2a - r)^2 = r^2 + (2ae)^2 - 2(2ae)(r) \cos(\pi - \theta)$$

$$4a^2 - 4ar + r^2 = r^2 + 4a^2e^2 - 4aer \cos(\pi - \theta)$$

$\underbrace{-\cos \theta}$

$$a - r = ae^2 + er \cos \theta$$

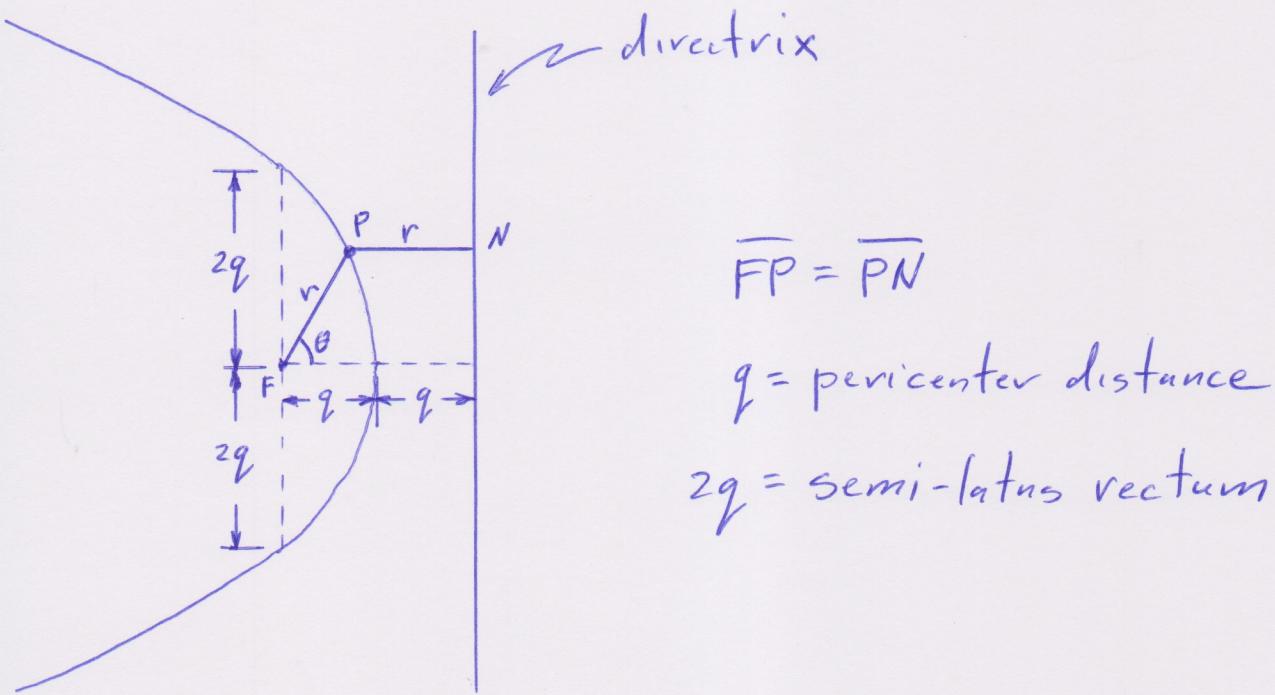
$$r = \frac{a(1-e^2)}{1+e \cos \theta}$$

(5.12)

Circle is special case of ellipse with $e = 0$

$$r = a$$

B. Parabola $e=1$

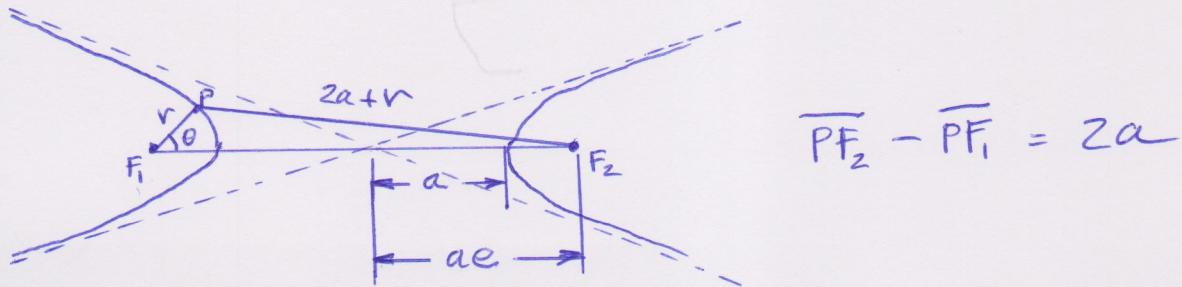


$$r + r \cos \theta = 2q$$

$$r = \frac{2q}{1 + \cos \theta}$$

(5.13)

C. Hyperbola $e > 1$



$$(2ar + r)^2 = r^2 + (2ae)^2 - 2(r)(2ae) \cos\theta$$

$$4a^2 + 4ar + r^2 = r^2 + 4a^2e^2 - 4are \cos\theta$$

$$a + r = ae^2 - re \cos\theta$$

$$r = \frac{a(e^2 - 1)}{1 + e \cos\theta}$$

(5.14)

Summary

<u>Conic Section</u>	<u>e</u>	<u>P</u>
circle	0	a
ellipse	<1	$a(1-e^2)$
parabola	1	$2q$
hyperbola	>1	$a(e^2-1)$

Proof of Kepler's 3rd Law $T^2 \sim a^3$

From (5.10) $P = h^2/\mu$

$$h = \sqrt{\mu P} = \sqrt{\mu a(1-e^2)} \quad (\text{for ellipse}) \quad (5.15)$$

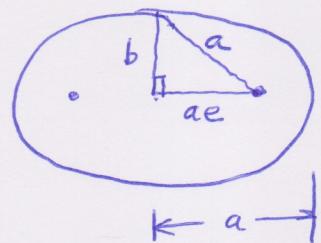
From (5.8)

$$\frac{dA}{dt} = \frac{1}{2} h$$

Area of ellipse is

$$A = \pi a b$$

~~area of ellipse~~



$$b^2 + a^2 e^2 = a^2$$

$$b^2 = a^2(1 - e^2)$$

$$A = \pi a^2 \sqrt{1 - e^2} \quad (5.16)$$

Integrate (5.8) $\frac{dA}{dt} = \frac{1}{2} h$

$$\int_0^A dA = \frac{1}{2} h \int_0^T dt$$

T = period of revolution (orbital period)

$$A = \frac{1}{2} h T \quad (5.17)$$

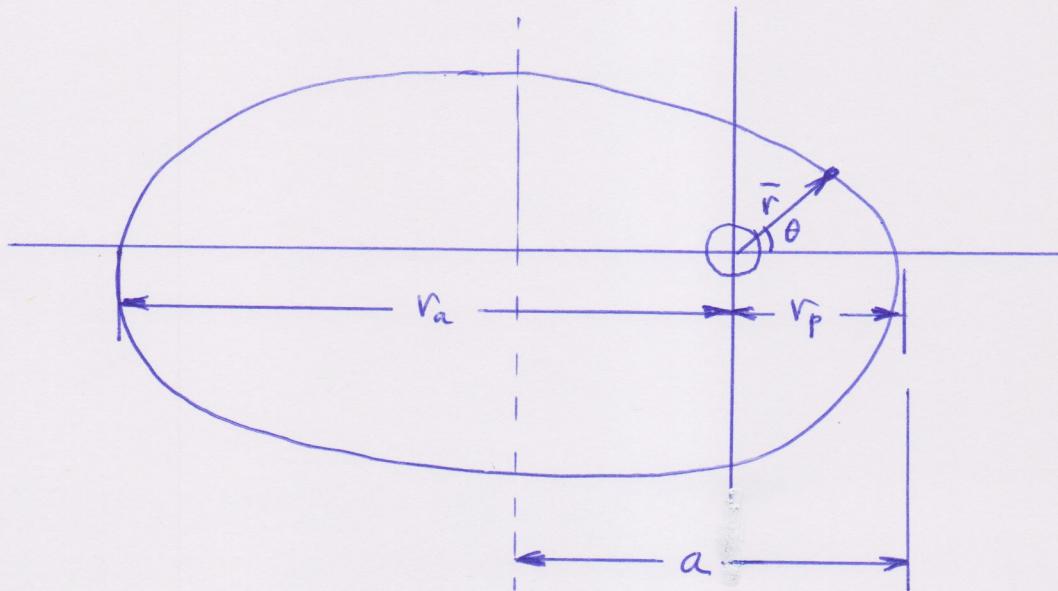
Equate (5.16) to (5.17) & use (5.15)

$$\pi a^2 \sqrt{1-e^2} = \frac{1}{2} \underbrace{\sqrt{\mu a(1-e^2)}}_h T$$

$$T^2 = \frac{4\pi^2}{\mu} a^3$$

Kepler's 3rd Law (5.18)

Satellite in Elliptic Orbit



$$r = \frac{a(1-e^2)}{1+e \cos \theta}$$

$$r_p = r(\theta=0) = \frac{a(1-e^2)}{1+e} = a(1-e) \quad (5.19)$$

$$r_a = r(\theta=\pi) = \frac{a(1-e^2)}{1-e} = a(1+e) \quad (5.20)$$

Divide (5.19) by (5.20)

$$\frac{v_p}{v_a} = \frac{1-e}{1+e} \quad (5.21)$$

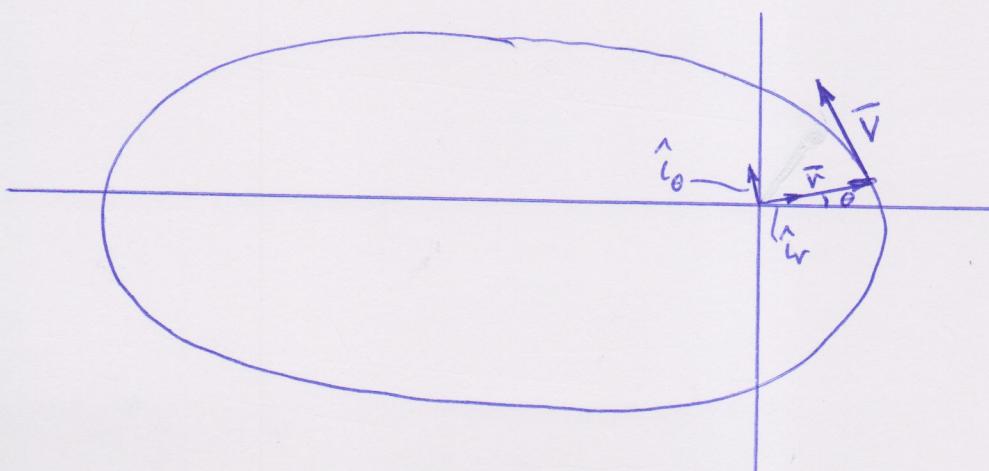
Solve for e

$$e = \frac{v_a - v_p}{v_a + v_p} \quad (5.22)$$

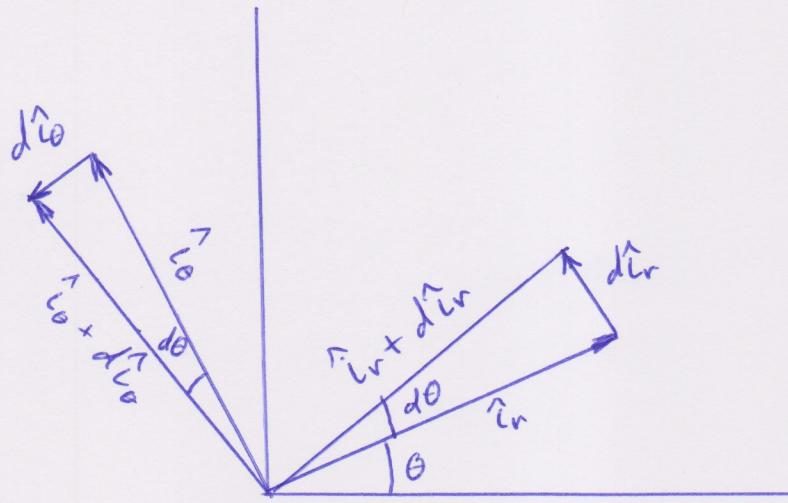
Using (5.12) and (5.19) can write

$$r = \frac{v_p(1+e)}{1+e \cos \theta} \quad (5.23)$$

Radial and Angular Velocity Components



Note: Unit vectors \hat{i}_r and \hat{i}_θ depend on θ



$$d\hat{e}_r = 1 \cdot d\theta \hat{e}_\theta \Rightarrow \frac{d\hat{e}_r}{d\theta} = \hat{e}_\theta$$

$$d\hat{e}_\theta = 1 \cdot d\theta (-\hat{e}_r) \Rightarrow \frac{d\hat{e}_\theta}{d\theta} = -\hat{e}_r$$

$$\bar{r} = r \hat{e}_r$$

$$\bar{V} = \frac{d\bar{r}}{dt} = \frac{dr}{dt} \hat{e}_r + r \frac{d\hat{e}_r}{d\theta} \frac{d\theta}{dt} = \frac{dr}{dt} \hat{e}_r + r \frac{d\theta}{dt} \hat{e}_\theta$$

$$V_r = \frac{dr}{dt} \quad (5.24)$$

$$V_\theta = r \frac{d\theta}{dt} \quad (5.25)$$

$$\text{Since } r = \frac{a(1-e^2)}{1+e\cos\theta}$$

$$\begin{aligned}\frac{dr}{dt} &= -\frac{a(1-e^2)}{(1+e\cos\theta)^2} (-e\sin\theta) \frac{d\theta}{dt} \\ &= \frac{a^2(1-e^2)^2 \frac{d\theta}{dt}}{(1+e\cos\theta)^2} \cdot \frac{e\sin\theta}{a(1-e^2)} \\ &= r^2 \frac{d\theta}{dt} \cdot \frac{e\sin\theta}{a(1-e^2)}\end{aligned}$$

Had previously found that

$$\frac{dA}{dt} = \frac{1}{2} r^2 \frac{d\theta}{dt} = \frac{1}{2} h = \text{const.} \Rightarrow h = r^2 \frac{d\theta}{dt}$$

Therefore

$$\boxed{V_r = \frac{dr}{dt} = \frac{he\sin\theta}{a(1-e^2)}} \quad (5.26)$$

$$V_\theta = r \frac{d\theta}{dt} = \frac{h}{r}$$

$$V_\theta = r \frac{d\theta}{dt} = \frac{h(1+e \cos \theta)}{a(1-e^2)}$$
(5.27)

At $\theta = 0$ (perigee)

$$V_r = 0 \quad V_\theta = \frac{h(1+e)}{a(1-e^2)} = \frac{h}{a(1-e)} = \frac{h}{r_p} \Rightarrow V_p = \frac{h}{r_p}$$

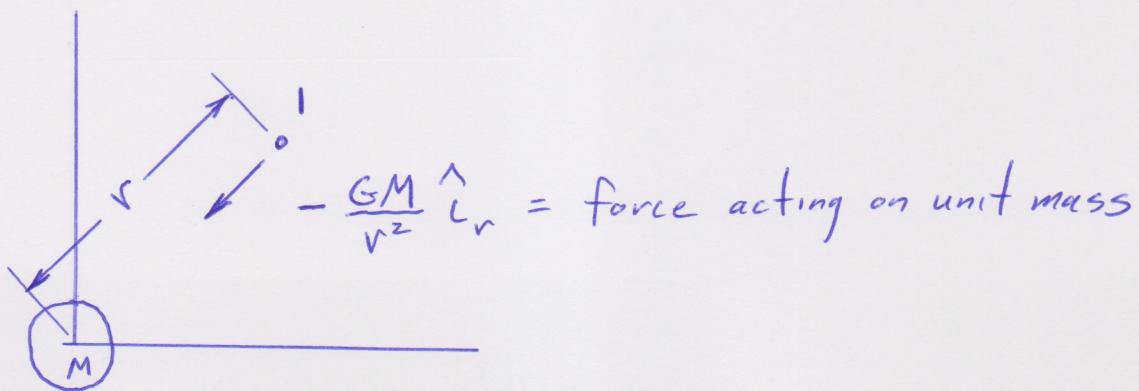
At $\theta = \pi$ (apogee)

$$V_r = 0 \quad V_\theta = \frac{h(1-e)}{a(1-e^2)} = \frac{h}{a(1+e)} = \frac{h}{r_a} \Rightarrow V_a = \frac{h}{r_a}$$

Since $h = \text{constant}$

$$V_p V_p = V_a V_a$$
(5.28)

Potential Energy



To move the unit mass from $r=R$ to $r=\infty$, apply force

$$+ \frac{GM}{r^2} \hat{e}_r$$

Work done by the applied force is

$$W = \int_R^\infty \frac{GM}{r^2} dr = -\frac{GM}{r} \Big|_R^\infty = 0 - \left(-\frac{GM}{R}\right) = \frac{GM}{R}$$

Potential energy of system is increased by $\frac{GM}{R}$

Set potential energy $V=0$ at ∞

$V(R) = -\frac{GM}{R}$

(5.29) (per unit mass of the satellite)

The Energy Integral

From equation of relative motion (5.1)

$$\frac{d^2\bar{r}}{dt^2} = -\frac{\mu}{r^3} \bar{r}$$

Dot with $\frac{d\bar{r}}{dt}$

$$\frac{d\bar{r}}{dt} \cdot \frac{d^2\bar{r}}{dt^2} = -\frac{\mu}{r^3} \bar{r} \cdot \frac{d\bar{r}}{dt}$$

Note that

$$\frac{d\bar{r}}{dt} \cdot \frac{d^2\bar{r}}{dt^2} = \frac{1}{2} \frac{d}{dt} \left(\frac{d\bar{r}}{dt} \cdot \frac{d\bar{r}}{dt} \right) = \frac{d}{dt} \left(\frac{V^2}{2} \right)$$

$$\bar{r} \cdot \frac{d\bar{r}}{dt} = V \frac{dV}{dt}$$

Therefore

$$\frac{d}{dt} \left(\frac{V^2}{2} \right) = -\frac{\mu}{r^2} \frac{dV}{dt}$$

Integrate

$$\frac{V^2}{2} = \frac{\mu}{r} + \varepsilon \quad \varepsilon = \text{constant of integration}$$

or

$$\underbrace{\frac{V^2}{2} - \frac{\mu}{r}}_{\substack{\text{specific KE} \\ \text{PE}}} = \varepsilon \quad (5.30)$$

$T + V = \varepsilon = \text{const.}$

Eq. (5.39) shows that the sum of the kinetic and potential energy of the orbiting satellite remains constant.

To evaluate Σ

$$\Sigma = \frac{V^2}{2} - \frac{\mu}{r} \quad (\text{constant for all } \theta)$$

Therefore at periape ($\theta=0$)

$$\Sigma = \frac{V_p^2}{2} - \frac{\mu}{r_p}$$

At periape, regardless of the type of conic trajectory, the velocity and radius vectors are orthogonal so that using $\vec{h} = \vec{r} \times \frac{d\vec{r}}{dt}$

$$h = r_p V_p$$

$$\text{Using } p = h^2/\mu \Rightarrow h = \sqrt{\mu p}$$

$$r_p V_p = \sqrt{\mu p} \Rightarrow V_p^2 = \frac{\mu p}{r_p^2}$$

Therefore

$$\Sigma = \frac{\mu p}{2r_p^2} - \frac{\mu}{r_p} = \frac{\mu}{r_p} \left(\frac{p}{2r_p} - 1 \right)$$

Since $r_p = \frac{P}{1+e \cos 0^\circ} = \frac{P}{1+e}$

$$\epsilon = \frac{\mu(1+e)}{P} \left(\frac{P(1+e)}{2P} - 1 \right) = \frac{\mu(1+e)(e-1)}{2P} = -\frac{\mu(1-e^2)}{2P}$$

$$\boxed{\epsilon = -\frac{\mu(1-e^2)}{2P}}$$

(5.31)

Valid for any
conic trajectory

<u>Conic Trajectory</u>	<u>e</u>	<u>P</u>	<u>E</u>
elliptic	<1	$a(1-e^2)$	$-\frac{\mu}{2a}$
parabolic	1	$2g$	0
hyperbolic	>1	$a(e^2-1)$	$+\frac{\mu}{2a}$

Note that the total energy for all conic trajectories is independent of eccentricity and depends only on the value of "a".

In particular, for a circular orbit of radius r_c 41

$$\frac{V_c^2}{2} - \frac{\mu}{r_c} = -\frac{\mu}{2r_c}$$

$$\frac{V_c^2}{2} = \frac{\mu}{2r_c}$$

$$V_c = \sqrt{\frac{\mu}{r_c}} \quad (5.32)$$

For a 185 mile high circular orbit about earth:

$$r_c = 6368 \text{ km} + (185 \text{ miles}) \left(\frac{1.61 \text{ km}}{1 \text{ mile}} \right) = 6670 \text{ km}$$

$$V_c = \sqrt{\frac{\mu}{r_c}} = \sqrt{\frac{3.986 \times 10^5 \frac{\text{km}^3}{\text{sec}^2}}{6670 \text{ km}}}$$

$$V_c = 7.730 \text{ km/sec} = 17,300 \text{ mph}$$

$$T = \frac{\text{circumference}}{\text{speed}} = \frac{2\pi \cdot 6670 \text{ km}}{7.730 \frac{\text{km}}{\text{sec}}} \cdot \frac{1 \text{ min}}{60 \text{ sec}}$$

$$T = 90 \text{ min}$$

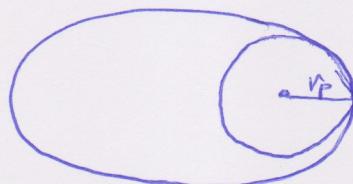
Since

$$V_p r_p = V_a r_a = h$$

$$V_p^2 = \frac{h^2}{r_p^2} = \frac{\mu r_p}{r_p^2} = \frac{\mu a(1-e^2)}{r_p a(1-e)} = \frac{\mu(1+e)}{r_p}$$

For a circular orbit at perigee height

$$V_{cp}^2 = \frac{\mu}{r_p}$$



Therefore

$$\frac{V_p^2}{V_{cp}^2} = 1+e$$

$$e = \frac{V_p^2}{V_{cp}^2} - 1 \quad (5.33)$$

Escape if $e \geq 1 \Rightarrow$

escape if $\boxed{\frac{V_p}{V_{cp}} \geq \sqrt{2}} \quad (5.34)$

At 185 miles above earth's surface

$$V_{\text{escape}} = \sqrt{2} V_{cp} = \sqrt{2} \left(7,730 \frac{\text{km}}{\text{sec}} \right) = 10,90 \frac{\text{km}}{\text{sec}}$$

$$= 24,900 \text{ mph}$$

Also, since

$$V_a^2 = \frac{h^2}{r_a^2} = \frac{\mu a (1-e^2)}{r_a a (1+e)} = \frac{\mu (1-e)}{r_a}$$

$$V_{ca}^2 = \frac{\mu}{r_a}$$

$$\frac{V_a^2}{V_{ca}^2} = 1 - e$$

$$e = 1 - \frac{V_a^2}{V_{ca}^2}$$

(5.35)