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► $\nabla \mathbf{u} = \begin{bmatrix} \frac{\partial u_x}{\partial x}, \frac{\partial u_x}{\partial y}, \frac{\partial u_x}{\partial z} \\ \frac{\partial u_y}{\partial x}, \frac{\partial u_y}{\partial y}, \frac{\partial u_y}{\partial z} \\ \frac{\partial u_z}{\partial x}, \frac{\partial u_z}{\partial y}, \frac{\partial u_z}{\partial z} \end{bmatrix}$

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 - Simplificaciones 1D y 2D.
 - Otras coordenadas.

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▶ Ecuación de conducción del calor: $\frac{\partial T}{\partial t} = \nabla \cdot (k \nabla T)$

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- **E**cuación de Poisson: $\nabla \cdot \mathbf{E} = \nabla \cdot (-\nabla \phi) = \rho_e$
- Conservación de momento en flujo incompresible (Navier–Stokes):

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla)\mathbf{u} - \nu \nabla^2 \mathbf{u} = F$$
$$\nabla \cdot \mathbf{u} = 0$$

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Partimos de la ecuación general en 3D:

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$$T|_{\Gamma_D} = T(\mathbf{x}^{\Gamma_D})$$

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Las condiciones de frontera deben cumplir que:

$$\Gamma_D \cap \Gamma_N = \emptyset$$

$$\Gamma_D \cup \Gamma_N = \Gamma$$

Ejemplo: Ecuación del calor 1D.

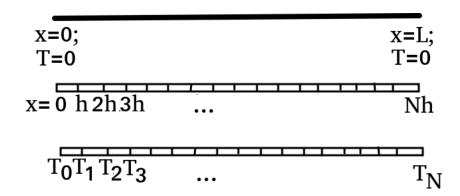
$$\begin{cases} \frac{\partial T(x,t)}{\partial t} = \frac{\partial}{\partial x} (k \frac{\partial T}{\partial x}) \\ T(x,0) = \sin(\frac{\pi x}{L}) \\ T(0,t) = T(L,t) = 0.0 \end{cases}$$

Ejemplo: Ecuación del calor 1D.

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Discretización espacial.

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- Discretización espacial.
- Aproximación de los operadores diferenciales.

Derivadas numéricas

ightharpoonup Derivada segunda centrada con error de orden h^2 .

$$f(x+h) = f(x) + f'(x)h + f''(x)\frac{h^2}{2} + f'''(x)\frac{h^3}{6} + f''''(x)\frac{h^4}{24} + \dots$$
$$f(x-h) = f(x) - f'(x)h + f''(x)\frac{h^2}{2} - f'''(x)\frac{h^3}{6} + f''''(x)\frac{h^4}{24} - \dots$$

Sumando las ecuaciones, obtenemos:

$$f(x+h) + f(x-h) = 2f(x) + f''(x)h^2 + f''''(x)\frac{h^4}{12} + \dots$$

$$f''(x) = \frac{f(x+h) - 2f(x) + f(x-h)}{h^2} - f''''(\zeta)\frac{h^2}{12}$$

$$f''(x) \approx \frac{f(x+h) - 2f(x) + f(x-h)}{h^2}; \quad \epsilon(O(h^2))$$

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$$\begin{bmatrix} \frac{\partial T_0}{\partial t} \\ \frac{\partial T_1}{\partial t} \\ \frac{\partial T_2}{\partial t} \\ \dots \\ \frac{\partial T_{N-1}}{\partial t} \\ \frac{\partial T_N}{\partial t} \end{bmatrix} = \frac{k}{h^2} \begin{bmatrix} 1 & 0 & 0 & 0 \dots & 0 \\ 1 & -2 & 1 & 0 \dots & 0 \\ 0 & 1 & -2 & 1 \dots & 0 \\ \dots & \dots & \dots & \dots & \dots \\ 0 & \dots & 1 & -2 & 1 \\ 0 & 0 & 0 & 0 \dots & 1 \end{bmatrix} \begin{bmatrix} 0 \\ T_1 \\ T_2 \\ \dots \\ T_{N-1} \\ 0 \end{bmatrix}$$

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Si definimos: $T^D = [T_0, T_1, T_2...T_N]^T$.

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Si definimos: $T^D = [T_0, T_1, T_2...T_N]^T$. Tendremos:

$$\frac{\partial T^D}{\partial t} = \frac{k}{h^2} \mathbf{S}^d T^D$$

Ejemplo: Ecuación del calor 1D.

$$\begin{cases} \frac{\partial T(x,t)}{\partial t} = k \frac{\partial^2 T(x,t)}{\partial x^2} \\ T(0,t) = T(L,t) = 0.0 \\ T(x,0) = \sin(\frac{\pi x}{L}) \end{cases}$$

- Discretización espacial.
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- Discretización espacial.
- Aproximación de los operadores diferenciales.
- Stencil de Diferencias finitas.
- Otros métodos: Elementos finitos, volúmenes finitos, elementos de borde, descomposición espectral...

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$$T_{(t+\Delta t)}^{D} = \left(\mathbf{I} + \frac{\Delta t k}{h^2} \mathbf{S}^d\right) T_{(t)}^{D}$$

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$$\begin{split} T_{(t+\Delta t)}^D - T_{(t)}^D &= \frac{\Delta t k}{h^2} \mathbf{S}^d T_{(t+\Delta t)}^D \\ T_{(t+\Delta t)}^D - \frac{\Delta t k}{h^2} \mathbf{S}^d T_{(t+\Delta t)}^D &= T_{(t)}^D \\ \left(\mathbf{I} - \frac{\Delta t k}{h^2} \mathbf{S}^d \right) T_{(t+\Delta t)}^D &= T_{(t)}^D \\ T_{(t+\Delta t)}^D &= \left(\mathbf{I} - \frac{\Delta t k}{h^2} \mathbf{S}^d \right)^{-1} T_{(t)}^D \end{split}$$

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- Crank–Nicolson:

$$\frac{T_{(t+\Delta t)} - T_{(t)}}{\Delta t} = \frac{k}{h^2} \mathbf{S}^d \left(\alpha T_{(t+\Delta t)} + (1-\alpha) T_{(t)} \right)$$



Ecuaciones en derivadas parciales: esquema temporal de Crank-Nicolson

$$\begin{split} \frac{T_{(t+\Delta t)} - T_{(t)}}{\Delta t} &= \frac{k}{h^2} \mathbf{S}^d \left(\frac{1}{2} T_{(t+\Delta t)} + \frac{1}{2} T_{(t)} \right) \\ T_{(t+\Delta t)} - T_{(t)} &= \frac{\Delta t k}{2h^2} \mathbf{S}^d T_{(t+\Delta t)} + \frac{\Delta t k}{2h^2} \mathbf{S}^d T_{(t)} \\ T_{(t+\Delta t)} - \frac{\Delta t k}{2h^2} \mathbf{S}^d T_{(t+\Delta t)} &= \frac{\Delta t k}{2h^2} \mathbf{S}^d T_{(t)} + T_{(t)} \\ \left(\mathbf{I} - \frac{\Delta t k}{2h^2} \mathbf{S}^d \right) T_{(t+\Delta t)} &= \left(\frac{\Delta t k}{2h^2} \mathbf{S}^d + \mathbf{I} \right) T_{(t)} \\ T_{(t+\Delta t)} &= \left(\mathbf{I} - \frac{\Delta t k}{2h^2} \mathbf{S}^d \right)^{-1} \left(\frac{\Delta t k}{2h^2} \mathbf{S}^d + \mathbf{I} \right) T_{(t)} \end{split}$$

$$\frac{\partial T^D}{\partial t} = \frac{k}{h^2} \mathbf{S}^d T^D$$

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- Crank–Nicolson: $\frac{T_{(t+\Delta t)}-T_{(t)}}{\Delta t} = \frac{k}{h^2} S^d \left(\alpha T_{(t+\Delta t)} + (1-\alpha)T_{(t)}\right)$
- Runge–Kutta y semi-implícitos

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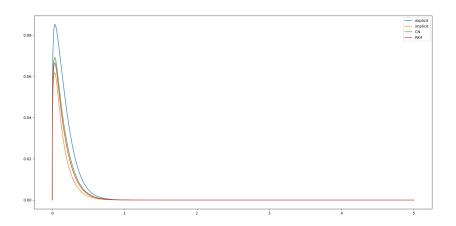
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- ► Runge–Kutta y semi-implícitos
- ► Métodos multipasos
- ► Implementaciones y costo.



Ecuaciones en derivadas parciales. Comparación de costo y estabilidad

	Explícito	Implícito	C-N	RK4
$\Delta t = 0.625 \text{ms} - h = 0.025$	Inestable	1.55	1.57	3.2
$\Delta t = 0.3125 \text{ms-}h = 0.025$	0.09	3.33	3.85	6.4
$\Delta t = 0.1 \text{ms} - h = 0.025$	0.24	9.03	9.91	20.3
$\Delta t = 0.1 \text{ms-}h = 0.01$	0.84	31.1	32.6	76

Errores



Ejemplo: Ecuación del calor 2D.

$$\begin{cases} \frac{\partial T(x,t)}{\partial t} = \nabla \cdot (k\nabla T(x,t)) \\ T(0,y,t) = 100.0 \\ T(L,y,t) = 0.0 \\ \frac{\partial T(x,0,t)}{\partial y} = 0 \\ \frac{\partial T(x,L,t)}{\partial y} = 0 \\ T(x,y,0) = 0.0 \end{cases}$$