

# Measuring precise radial velocities and cross-correlation function line-profile variations using a Skew Normal density<sup>★</sup>

U. Simola<sup>1</sup> <sup>★★</sup>, X. Dumusque<sup>2</sup> <sup>★★★</sup>, and Jessi Cisewski-Kehe<sup>3</sup>

<sup>1</sup> Department of Mathematics and Statistics, University of Helsinki, Helsinki, Finland

<sup>2</sup> Observatoire de Genève, Université de Genève, 51 ch. des Maillettes, CH-1290 Versoix, Switzerland

<sup>3</sup> Department of Statistics and Data Science, Yale University, New Haven, CT, USA

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## ABSTRACT

**Context.** Stellar activity is one of the primary limitations to the detection of low-mass exoplanets using the radial-velocity (RV) technique. Stellar activity can be probed by measuring time dependent variations in the shape of the cross-correlation function (CCF), often estimated using different parameters of the modelled CCF. Therefore to estimate the parameters of the CCF with high precision is essential to de-correlate the signal of an exoplanet from spurious RV signals caused by stellar activity.

**Aims.** We propose to estimate the parameters of the CCF by fitting a Skew Normal (SN) density which, unlike the commonly employed Normal density, includes a skewness parameter to capture the asymmetry of the CCF induced by stellar activity and also the natural asymmetry induced by convective blueshift.

**Methods.** The performances of the proposed method are compared to the commonly employed Normal density using both simulations and real observations, with different levels of activity and signal-to-noise ratio (SNR) levels.

**Results.** When considering real observations, the correlation between the RV and the asymmetry of the CCF and the correlation between the RV and the width of the CCF are stronger when using the parameters estimated with the SN rather than the ones obtained with the commonly employed Normal density. In particular the strongest correlations have been obtained when using the mean of the SN as an estimate for the RV. This suggests that the asymmetry of the CCF and the width of the CCF estimated using a SN density may be more sensitive to stellar activity, which can be helpful when estimating stellar rotational periods and in general when characterizing stellar activity signals. Using the proposed SN approach, the uncertainties estimated on the RV defined as the median of the SN are on average 10% smaller than the uncertainties calculated on the mean of the Normal. The uncertainties estimated on the asymmetry parameter of the SN are on average 15% smaller than the uncertainties measured on the Bisector Inverse Slope Span (BIS SPAN), which is the commonly used parameter to evaluate the asymmetry of the CCF.

**Conclusions.** The SN density shape appears to be useful for characterizing the CCF since the correlations used to probe stellar activity are stronger and the uncertainties associated of the estimates for RV and the asymmetry of the CCF are both smaller than those derived from other commonly employed strategies.

**Key words.** techniques: radial velocities – planetary systems – stars: activity – methods: data analysis

## 1. Introduction

When working with radial-velocities data (RVs), one of the main limitations to the detection of low-mass exoplanets is no longer the precision of the instruments used, but the different sources of variability induced by the stars (e.g. Feng et al. 2017; Dumusque et al. 2017; Rajpaul et al. 2015; Robertson et al. 2014). Stellar oscillations, granulation phenomena, and stellar activity can all induce apparent RV signals that are above the meter-per-second ( $\text{m s}^{-1}$ ) precision (e.g. Saar & Donahue 1997; Queloz et al. 2001; Desort et al. 2007; Dumusque et al. 2011; Dumusque 2016) reached by the best high-resolution spectrographs (HARPS, HARPS-N, Mayor et al. 2003; Cosentino et al. 2012). It is therefore mandatory to better understand stellar signals and to develop methods to correct for them, if in the near future we want to detect or confirm an Earth-twin planet using the RV technique. This is even more true now that instruments

like the Echelle SPectrograph for Rocky Exoplanet and Stable Spectroscopic Observations (ESPRESSO) (Pepe et al. 2014) and the EXtreme PREcision Spectrometer (EXPRES) (Fischer et al. 2016) should reach the precision and stability to detect such signals. However, if solutions are not found to mitigate the impact of stellar activity, the detection or confirmation of potential Earth-twins will be extremely challenging and false detections could plague the field.

One of the most challenging stellar signals to characterize and to correct for are the signals induced by stellar activity. Stellar activity is responsible for creating magnetic regions on the surface of stars, and those regions change locally the temperature and the convection, which can induce spurious RVs variations (e.g. Meunier et al. 2010; Dumusque et al. 2014; Borgniet et al. 2015). In theory, it should be easy to differentiate between the pure Doppler-shift induced by a planet, which shifts the entire stellar spectrum, and stellar activity, which modifies the shape of spectral lines and by doing so creates a spurious shift of the stellar spectrum (Saar & Donahue 1997; Hatzes 2002; Kurster et al. 2003; Lindegren & Dravins 2003; Desort et al. 2007; Lagrange et al. 2010; Meunier et al. 2010; Dumusque et al. 2014). However, on quiet GKM dwarfs, the main targets for precise

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<sup>★★</sup> e-mail: umberto.simola@helsinki.fi

<sup>★★★</sup> Branco Weiss Fellow–Society in Science (url: <http://www.society-in-science.org>)

RVs measurements, stellar activity can induce signals of a few  $\text{m s}^{-1}$ . This corresponds physically to variations smaller than 1/100th of a pixel on the detector, making changes in the shape of the spectral lines challenging to detect. In order to measure such tiny variations, a common approach is to average the information of all the lines in the spectrum by cross correlating the stellar spectrum with a synthetic or an observed stellar template (Baranne et al. 1996; Pepe et al. 2002; Anglada-Escudé & Butler 2012). The result of this operation gives us the cross-correlation function (CCF). To measure the Doppler-shift between different spectra, and therefore to retrieve the RVs of a star as a function of time, the variations of the CCF barycenter are calculated. The barycenter is generally estimated by the mean of a Normal density shape fit to the CCF. Variations in line shape between different spectra, which indicate the presence of signals induced by stellar activity, are measured by analyzing different parameters of models fit to the CCF. Usually the width of the CCF is estimated using the full-width half-maximum (FWHM) of the fitted Normal density and its asymmetry by calculating the CCF bisector and measuring the bisector inverse slope span (BIS SPAN, Queloz et al. 2001).

If a spurious RV signal is induced by activity, generally a strong correlation will be observed between the RV and chromospheric activity indicators like  $\log(R'_{HK})$  or H- $\alpha$  (Boisse et al. 2009; Dumusque et al. 2012; Robertson et al. 2014), but also between the RV and the FWHM of the CCF or its BIS SPAN (Queloz et al. 2001; Boisse et al. 2009; Queloz et al. 2009; Dumusque 2016). Therefore a common strategy when fitting a Keplerian signal to a set of estimated RVs in search of a planet is to include linear terms in the model to account for activity, such as the  $\log(R'_{HK})$ , the FWHM, and the BIS SPAN (Dumusque et al. 2017; Feng et al. 2017). It is also common to add a Gaussian process to the model to account for the correlated noise induced by stellar activity. The hyperparameters of the Gaussian process can be trained on different activity indicators (Haywood et al. 2014; Rajpaul et al. 2015) or directly on the RVs (Faria et al. 2016). It is therefore essential for mitigating stellar activity to obtain activity indicators that are the most correlated with the RVs but also for which we can obtain the best precision.

Several indicators have been developed that can be more sensitive to line asymmetry than the BIS SPAN. In Boisse et al. (2011), the authors developed  $V_{span}$ , which is the difference between the RV measured by fitting a Normal density to the upper and the lower parts of the CCF. This CCF asymmetry parameter is shown to be more sensitive than the BIS SPAN at low signal-to-noise ratio (SNR). Figueira et al. (2013) studied the use of new indicators, BIS-, BIS+, bi-Gauss and  $V_{asy}$ . The authors were able to show that when using bi-Gauss, the amplitude in asymmetry is 30% larger than when using BIS SPAN, therefore allowing the detection of lower levels of activity. They also demonstrated that  $V_{asy}$  seems to be a better indicator of line asymmetry at high SNR, as its correlation with RV is significantly stronger than any other correlation between the previously proposed asymmetry indicators and RV.

In all the methods described above, except bi-Gauss, the RV and the FWHM are derived using a Normal density fitted to the CCF, and the asymmetry is estimated using a separate approach. In this paper we propose to use a Skew Normal (SN) density to estimate with a single fit of the CCF the RV, the FWHM and the asymmetry of the CCF, as this function includes a skewness parameter (Azzalini 1985).

The paper is organized as follows. In Sec. 2 we introduce the SN density, describe its applicability for modelling the CCF and study how the SN parameters relate to the mean of the Normal

density, the FWHM, and the BIS SPAN of the CCF. In Sec. 3 we propose a linear model to correct for stellar activity signals in RVs, which extends the linear models previously proposed for this purpose (e.g. Dumusque et al. 2017; Feng et al. 2017). In Sec. 4 the performance of the SN fit to the CCF is investigated using simulations coming from the Spot Oscillation And Planet 2.0 code (SOAP 2.0, Dumusque et al. 2014), followed by an analysis of real observations in Sec. 5. In Sec. 6 error bars are computed for the different estimated CCF parameters. Finally the discussion of the results and the conclusions are presented in Secs. 7 and 8, respectively.

## 2. The Skew Normal distribution

The Skew Normal (SN) distribution is a class of probability distributions which includes the Normal distribution as a special case (Azzalini 1985). The SN distribution has, in addition to a location and a scale parameter analogous to the Normal distribution's mean and standard deviation, a third parameter which describes the skewness (i.e. the asymmetry) of the distribution. Considering a random variable  $Y \in \mathbb{R}$  (where  $\mathbb{R}$  is the real line) which follows a SN distribution with location parameter  $\xi \in \mathbb{R}$ , scale parameter  $\omega \in \mathbb{R}^+$  (i.e., the positive real line), and skewness parameter  $\alpha \in \mathbb{R}$ , its density at some value  $y \in Y$  can be written as

$$SN(y; \xi, \omega, \alpha) = \frac{2}{\omega} \phi\left(\frac{y - \xi}{\omega}\right) \Phi\left(\frac{\alpha(y - \xi)}{\omega}\right), \quad (1)$$

where  $\phi$  and  $\Phi$  are, respectively, the density function and the distribution function of a standard Normal distribution<sup>1</sup>. The skewness parameter  $\alpha$  quantifies the asymmetry of the SN. Examples of SN densities under different skewness parameter values and the same location and scale parameters ( $\xi = 0$  and  $\omega = 1$ ) are displayed in Fig. 1. A usual Normal distribution is the special case of the SN distribution when the skewness parameter  $\alpha$  is equal to zero<sup>2</sup>. For reasons related to the interpretation of the parameters in Eq. 1 and computational issues with estimating  $\alpha$  near 0, a different parametrization is used in this work, which is referred to as the *centered parametrization* (CP). This CP is much closer to the parametrization of a Normal distribution, as it uses a mean parameter  $\mu$ , a variance parameter  $\sigma^2$  and a skewness parameter  $\gamma$ . In order to define the CP, we need to express the CP parameters ( $\mu, \sigma^2, \gamma$ ) as a function of ( $\xi, \omega^2, \alpha$ ). This can be done using the following relations:

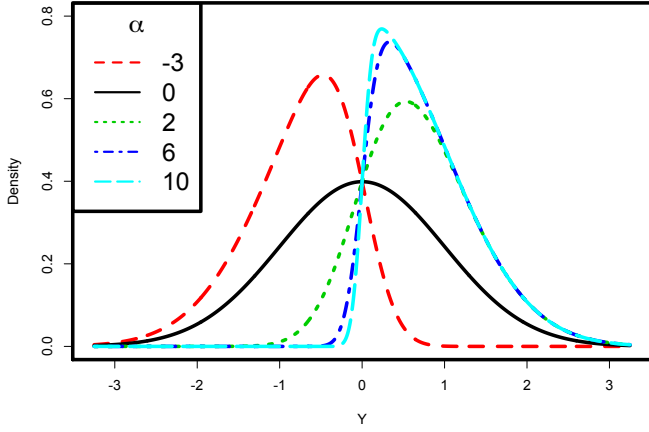
$$\mu = \xi + \omega\beta, \quad \sigma^2 = \omega^2(1 - \beta^2), \quad \gamma = \frac{1}{2}(4 - \pi)\beta^3(1 - \beta^2)^{-3/2}, \quad (2)$$

where  $\beta = \sqrt{\frac{2}{\pi}} \left( \frac{\alpha}{\sqrt{1 + \alpha^2}} \right)$  (e.g. Arellano & Azzalini 2010).

By using Eq. 2, the new set of parameters ( $\mu, \sigma^2, \gamma$ ) provides a clearer interpretation of the behavior of the SN distribution. For the  $\alpha$  values used in Fig. 1, the corresponding values of ( $\mu, \sigma^2, \gamma$ ) are displayed in Table 1. In particular,  $\mu$  and  $\sigma^2$  are the actual mean and variance of the distribution, rather than simply a location and scale parameter, and  $\gamma$  provides a measure of the

<sup>1</sup> A standard Normal distribution is a Normal distribution with a mean of 0 and a standard deviation of 1.

<sup>2</sup> This can be seen from Eq. 1. If  $\alpha = 0$  then  $\Phi\left(\frac{\alpha(y - \xi)}{\omega}\right) = \Phi(0) = 0.5$  and therefore  $SN(y; \xi, \omega, 0) = \frac{1}{\omega} \phi\left(\frac{y - \xi}{\omega}\right)$  which is the density of a Normal distribution. Note that  $\Phi(0) = 0.5$  because  $\Phi(0)$  is the probability that a standard Normal random variable is less than or equal than 0.



**Fig. 1.** Density function of a random variable  $Y$  following the SN distribution  $SN(\xi, \omega, \alpha)$  with location parameter  $\xi = 0$ , scale parameter  $\omega = 1$  and different values of the skewness parameter  $\alpha$  indicated by different colors and line types. Note that the solid black line has an  $\alpha = 0$  making it a Normal distribution.

skewness of the SN. Along with the mean of the SN, we consider the median of the distribution as a measure of its barycenter. See Table 1 for the medians of the SN densities displayed in Fig. 1.

**Table 1.** CP values ( $\mu, \sigma^2, \gamma$ ) along with the median corresponding to the  $\alpha$  values shown in Fig. 1, with location parameter  $\xi = 0$  and scale parameter  $\omega = 1$ . Values are rounded to three decimal places.

$\alpha$	$\mu$	$\sigma^2$	$\gamma$	Median
-3	-0.757	0.427	-0.667	-0.672
0	0.000	1.000	0.000	0.000
2	0.714	0.491	0.454	0.655
6	0.787	0.381	0.891	0.674
10	0.794	0.370	0.956	0.674

Further details about the parametrization from Eq. 1, called the *Direct Parametrization* (DP), the CP, and general statistical properties of the SN can be found in Azzalini & Capitanio (2014).

### 2.1. Fitting the Skew Normal density to the CCF

To fit the CCF using a SN density shape, we use a least-squares algorithm and the following model:

$$f_{CCF}(x_i) = C - A \times SN(x_i; \mu, \sigma^2, \gamma), \quad i = 1, \dots, n \quad (3)$$

where  $C$  is an unknown offset for the continuum of the CCF,  $A$  is the unknown amplitude of the CCF, sometimes referred to as the CCF contrast, and  $\mu, \sigma^2$  and  $\gamma$  are, respectively, the mean, the variance, and the skewness of the SN as defined above. The values  $x_1, \dots, x_n$  are the different values of the x-axis of the CCF, generally in velocity units (e.g.  $\text{m s}^{-1}$ ).

When using a Normal density shape model for the CCF, the estimated mean is used as the estimated RV and the FWHM<sup>3</sup> is used to quantify the width of the CCF. Because the Normal density is symmetric, the skewness is not defined and therefore a separate approach is necessary to estimate the skewness of the CCF. An estimated skewness parameter is generally obtained

by calculating the BIS SPAN of the CCF (see Sect. 1, and e.g. Queloz et al. 2001).

With the proposed SN approach, we propose two estimators of the RV: the mean and median of the SN model fit (referred to as SN mean RV and SN median RV, respectively), and present advantages and limitations for both of these choices in Sec. 5 and Sec. 6. The width of the SN, SN FWHM, is defined in the same way as for the Normal density<sup>4</sup>, and finally the skewness of the CCF is estimated by the  $\gamma$  parameter.

To evaluate the strength of the correlation between the estimated RVs and the different stellar activity indicators, we calculated the Pearson correlation coefficient,  $R$ , which in its general form is defined as:

$$R(x, y) = \frac{\text{cov}(x, y)}{\sigma(x)\sigma(y)}, \quad (4)$$

where  $x$  and  $y$  are two quantitative variables,  $\text{cov}(x, y)$  indicates the covariance between  $x$  and  $y$ , and  $\sigma(x)$  and  $\sigma(y)$  represent their standard deviations. A  $p$ -value for the statistical test with null hypothesis  $H_0 : R = 0$  is also generally provided.

### 3. Radial Velocity correction for stellar activity

Exoplanets only produce a pure RV signal. On the contrary, stellar activity, and, in particular, the presence of active regions on the stellar photosphere, do not produce blueshifts or redshifts of the entire stellar spectrum but can create spurious RV signals by modifying the shape of spectral lines. To track these variations in the shape of the spectral lines, a common approach consists in using the FWHM, the BIS SPAN, or other indicators such as those introduced in Boisse et al. (2011) or Figueira et al. (2013), which provide information on the width and asymmetry of the CCF. A strong correlation between the estimated RVs and one or more of these parameters provides an indication that stellar activity signals may be affecting the measurements.

When fitting for planetary signals in RV data, it is common to include linear dependencies with the BIS SPAN and the FWHM to take into account the signal induced by stellar activity (e.g. Dumusque et al. 2017; Feng et al. 2017). We propose to add additional parameters in the model to correct for stellar activity: (i) an amplitude parameter  $A$  of the CCF (referred to as the CCF contrast) and (ii) an interaction term for  $\gamma$  and SN FWHM (or the BIS SPAN and the FWHM in the Normal case). The stellar activity correction we propose can therefore be written as:

$$RV_{\text{activity}} = \beta_0 + \beta_1 A + \beta_2 \gamma + \beta_3 \text{SN FWHM} + \beta_4 (\gamma \text{SN FWHM}) + \epsilon, \quad (5)$$

where  $\beta_0$  is the intercept and  $\epsilon$  is the random error with mean equal to 0 and covariance matrix equal to  $\sigma^2 I$  ( $I$  defined as the identity matrix). The contrast parameter  $A$  accounts for the presence of a spot on the stellar surface, which produces a change in the amplitude of the CCF, in addition to changes in asymmetry or width (see e.g. Fig. 2 in Dumusque et al. 2014). The benefits of including a variable that quantifies the interaction between  $\gamma$  and SN FWHM (or BIS SPAN and FWHM) will be better understood through the results of the examples presented in Sec. 4. This interaction term can account for possible interactions between SN FWHM (or FWHM) and  $\gamma$  (or BIS SPAN), meaning that each variables' association with the response,  $RV_{\text{activity}}$ , depends also on the other variable.

<sup>4</sup> Note that SN FWHM does not correspond to the width of the SN density at half maximum like in the Normal case.

<sup>3</sup> FWHM =  $2\sqrt{2\ln 2}\sigma$  with standard deviation  $\sigma$

The proposed model is analyzed using statistical tests on the parameters  $\beta_0, \beta_1, \beta_2, \beta_3$  and  $\beta_4$  where the null hypothesis is  $H_0 : \beta_i = 0$ , for  $i = 0, \dots, 4$ . The significance level for the tests are set at 0.05. The coefficient of determination,  $R^2$ , is used to assess how well the proposed linear combination of variables accounts for the variability of  $RV_{\text{activity}}$ .

The proposed function defined in Eq. 5 is the result of statistical and astronomical considerations. In particular, we considered issues with correlations between the model covariates, which can lead to issues because the matrix needed to calculate the estimates could be singular; this problem is known in statistics as *multicollinearity* and more discussion can be found in Belsley (1991). In the analysis of real data presented in this work, we never observed a correlation coefficient exceeding 0.66 between the asymmetry and width parameters and therefore the problem of multicollinearity appeared to be mitigated. We also investigated the statistical significant of the interactions between  $A$  and the width of the CCF, and  $A$  and the asymmetry of the CCF; those two interactions were not statistically significant.

#### 4. Simulation Study

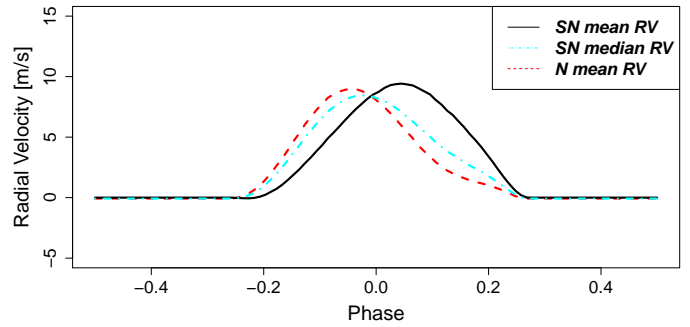
In order to evaluate the performance of the proposed SN approach for modelling the CCF and the benefit of using the proposed correction for stellar activity (see Eq. 5), we begin by considering a simulation study using spectra generated from the Spot Oscillation And Planet 2.0 code (SOAP 2.0, Dumusque et al. 2014).

For a given configuration of spots and faculae on the stellar surface, SOAP 2.0 outputs simulated CCF as a function of rotational phase. The code also returns the RV and the FWHM estimated using a Normal model for the CCF, and the BIS SPAN obtained by calculating the bisector of the CCF. SOAP 2.0 gives noiseless CCFs affected by stellar activity, which are used to compare the performances of a SN and Normal models of the CCF.

For the simulations discussed below, a star similar to the Sun was modelled with a solar disc of one solar radius seen equator-on, and a stellar rotational period set to 25.0 days. The stellar effective temperature is set to 5778 K, and a quadratic limb-darkening relation with linear and quadratic coefficients 0.29 and 0.34, respectively (Oshagh et al. 2013; Claret & Bloemen 2011). In order to make the result of the simulations more comparable to real data obtained with the HARPS spectrograph discussed in Sect. 5, the SOAP 2.0 CCFs were generated with a width of  $40 \text{ km s}^{-1}$  and considering initial spectra with a spectral resolution of  $R=115'000$ .

##### 4.1. Faculae

To see the impact of a facula on the parameters of different models of the CCF, we simulated the effect of an equatorial facula of size 3% relative to the visible stellar hemisphere. The facula is face-on when the phase equals to 0. Note that a 3% faculae is relatively large for the Sun; at maximum activity, large faculae generally have a size of 1% (e.g. Borgniet et al. 2015). In Fig. 2, we compare the barycentric variation of the CCF as measured when fitting a Normal density and using its mean (N mean RV), and when fitting a SN density and taking its mean (SN mean RV) or median (SN median RV). We see that all the different estimates of the CCF barycenter present a signal of similar amplitude, however the signal obtained with SN mean RV is notably different from the two others with a maximum amplitude happening at a different phase.



**Fig. 2.** RV estimates for N mean RV (red dashed line), SN mean RV (black line), and SN median RV (cyan dotted-dashed line). In this case, the CCFs were generated using SOAP 2.0 with an equatorial 3% facula on the simulated Sun. The star does one full rotation between phase -0.5 and 0.5, with the facula being seen face-on for phase 0. The variations observed in SN mean RV are notably different from the variations measured in SN median RV and N mean RV.

Correlations between the different RV estimates and the different CCF asymmetry or width estimates are displayed in Fig. 3. The correlation between  $\gamma$  and SN mean RV, and  $\gamma$  and SN median RV are stronger than the correlation between BIS SPAN and RV, with Pearson correlation coefficient values of  $R=0.46$ ,  $-0.67$  and  $-0.09$ , respectively. There is a stronger correlation between SN FWHM and SN mean RV ( $R = 0.98$ ) than between FWHM and N mean RV ( $R = 0.84$ ), but the correlation between SN FWHM and SN median RV is the weakest ( $R = 0.50$ ). This first analysis shows that in the case of a facula, using some parameters from the SN can lead to stronger correlations than the ones obtained by the usual Normal parameters. Therefore, the SN parameters may better probe stellar activity. We investigate this feature further in the following two sections, where we consider simulated data with a single spot and a spot plus a planet, and in Sec 5 with real observations.

The RV variations displayed in Fig. 2 are caused only by stellar activity, in this case a facula, so we applied the activity correction proposed in Eq. 5 to evaluate the performance of the model in this setting. The results of this correction are displayed in Fig. 4 and the statistical tests on the coefficients involved in Eq. 5 are summarized in Table 2. The proposed correction for stellar activity is able to account for the majority of the activity signal created by a facula, with a  $R^2$  larger than 0.99. In addition, the rms of the different estimates of the RV reduces from about  $3 \text{ m s}^{-1}$  before correcting for stellar activity to values below  $0.15 \text{ m s}^{-1}$  after correcting for stellar activity. We see a slightly smaller rms after correction when using the SN parameters compared to the Normal parameters. When comparing the correction proposed in Eq. 5 with what is generally used (i.e. a linear combination of only the asymmetry and width parameters), we see that the proposed correction is able to reduce the rms of the RV residuals by a factor of 2. Looking at the significance of the coefficients in Table 2, we observe that all the SN or Normal parameters are relevant for the correction.

##### 4.2. Spot

Next we consider the effects on the CCF model parameters due to the presence of an equatorial spot of size 1% relative to the visible stellar hemisphere. The spot is face-on when the phase equals to 0. Note that this is a large spot for the Sun, as large



**Fig. 3.** (left) Correlations between the different asymmetry parameters and their corresponding RV estimates in the case of an equatorial 3% facula on the simulated Sun. (right) Correlations between the different width parameters and their corresponding RV estimates for the same facula. In the presence of a facula, both the shape and the width of the CCF change as the star rotates, producing statistically significant correlations.

**Table 2.** P-values for the estimated coefficients from the model in Eq. 5 for correcting stellar activity induced by an equatorial 3% facula on the simulated Sun. All the parameters corresponding to the Normal or SN variables are statistically significant. The estimated  $R^2$  show that the proposed correction for stellar activity explains the vast majority of the spurious variability present in the different RV estimates.

Parameter	N mean RV	SN mean RV	SN median RV
$\beta_0$	0.033	0.00020	0.61
$\beta_1$	$2.22e-16$	$2.22e-16$	$2.22e-16$
$\beta_2$	0.0034	$2.22e-16$	$2.22e-16$
$\beta_3$	0.00016	$1.091e-6$	$9.75e-7$
$\beta_4$	$2.22e-16$	$2.22e-16$	$2.22e-16$
$R^2$	0.9978	0.9985	0.9981

induced by this simulated spot. In contrast to the case of the facula, all the different estimates of the CCF barycenter for the spot have the same shape in variation; however, the amplitude for SN mean RV is slightly smaller.

Fig. 6 shows the correlations between the asymmetry parameters and the different estimates of the CCF barycenter (i.e. SN mean RV, SN median RV and N mean RV). The correlation between  $\gamma$  and SN median RV is the strongest ( $R = -0.93$ ), followed by the BIS SPAN - N mean RV correlation ( $R = -0.86$ ) and the  $\gamma$  - SN mean RV correlation ( $R = -0.85$ ). The correlations between the width and the CCF barycenter are more circular and no substantial correlation is observed. Similar to the case with a facula, some parameters of the SN gives stronger correlations compared to the Normal parameters.

As before, the original RV estimates are corrected using Eq. 5. The results of this correction are displayed in Fig. 7 and the statistical tests on the coefficients involved in Eq. 5 are summarized in Table 3. As with the facula, the proposed correction is

spots are generally on the order of 0.1% (e.g. Borgniet et al. 2015). In Fig. 5, we show the barycentric variations of the CCF





**Fig. 4.** (top) The spurious estimated RVs (black dots) caused by a facula in the simulated data using a Normal and a SN model, the estimated RVs using Eq. 5 (red crosses), and the estimated RVs using the usual correction for stellar activity (green triangles), based on  $RV_{\text{activity}} = \beta_0 + \beta_1 \gamma + \beta_2 \text{SN FWHM}$  for the SN fit and on  $RV_{\text{activity}} = \beta_0 + \beta_1 \text{BIS SPAN} + \beta_2 \text{FWHM}$  for the Normal fit. (bottom) The residuals from the model fit using Eq. 5 (red crosses) and the residuals from the usual correction (green triangles). The std. are also reported in the legend, and the residuals have a smaller systematic component when using the proposed model of Eq. 5 compared to the usual model. The tests of statistical significance on the parameters are presented in Table 2.

able, in the case of the spot, to capture almost entirely the stellar activity signal when considering the SN or Normal parameters, with  $R^2$  values for the linear combination larger than 0.99.

In Fig. 7, we see that the proposed activity correction is able to reduce the signal of a spot from a raw RV rms larger than  $4.80 \text{ m s}^{-1}$  down to a rms of  $0.38 \text{ m s}^{-1}$ , for any of the different RV estimates. When comparing the activity correction proposed in this paper with what is commonly used, i.e. only a linear combination of the width and asymmetry of the CCF, we see that our solution is capable of reducing the RV residual rms by a factor of 3.5, which is even more than the factor 2 found in the case of the facula. The Normal or SN parameters involved in Eq. 5 are all statistically significant to explain the activity signal as seen in Table 3, except the width of the CCF when the RVs are derived with the mean or median of the SN density.

#### 4.3. Spot and planet

The final simulation includes a planetary signal influencing the CCF along with the 1% spot modelled previously (see Sec. 4.2). The purpose of this example is to check if we are able to disentangle these two different sources of variations when using the parameters derived using a Normal or a SN model for the CCF. In this scenario the planet is injected with a semi-amplitude of  $10 \text{ m s}^{-1}$  with no eccentricity and with a period corresponding to one-third of the stellar rotational period, i.e. one-third of 25 days.

Fig. 8 shows the variations observed in the CCF barycenter parameters. As in the case of the spot, all RV estimates show similar variations, with SN mean RV showing a slightly smaller amplitude.



**Fig. 5.** RV estimates for N mean RV (red dashed line), SN mean RV (black line) or SN median RV (cyan dotted-dashed line) using CCFs generated from SOAP 2.0 with an equatorial 1% spot on the simulated Sun. The star does one full rotation between phase -0.5 and 0.5, with the spot being seen face-on at phase 0. The SN mean RV seems to have the smallest spurious variations caused by the spot.

**Table 3.** P-values for the different coefficients used in Eq. 5 for the correction of stellar activity induced by an equatorial 1% spot on the simulated Sun. All the parameters corresponding to the Normal or SN parameters are statistically significant to explain the spurious RV variations caused by this spot, except for the width of the CCF when using SN mean RV or SN median RV as RV estimates. The estimated  $R^2$  show that the proposed correction for stellar activity explains the vast majority of the spurious variability seen in the different RV estimates.

Parameter	N mean RV	SN mean RV	SN median RV
$\beta_0$	0.4975	0.21	0.21
$\beta_1$	$2e-16$	$2e-16$	$2e-16$
$\beta_2$	$2e-16$	$2e-16$	$2e-16$
$\beta_3$	0.017	0.13	0.11
$\beta_4$	$2e-16$	$2e-16$	$2e-16$
$R^2$	0.9959	0.9936	0.9952

The correlation between the different CCF parameters are displayed in Fig. 9. The correlations are weaker than in the case of the spot due to the planet inducing changes in RV without affecting the width or the asymmetry of the CCF. However, the order of the strength of the correlations between the CCF asymmetry parameters and RV are comparable with the ones obtained for the spot-only model:  $\gamma$ -SN median RV has the strongest correlation ( $R = -0.84$ ), followed by the correlation between BIS SPAN-N mean RV ( $R = -0.78$ ) and finally by the correlation between  $\gamma$ -SN mean RV ( $R = -0.75$ ). The patterns observed in the width-RV phase space correlations in Fig. 9 follow a circle, similar to the spot-only model; no substantial correlation is observed between those two parameters.

In order to correct the estimated RVs from the spurious variation caused by the spot, the proposed model for correcting the activity is added to a planetary signal model that takes into account the RV variations caused by a planet. The observed RVs can therefore be modelled as a combination of the activity and the planetary signals:

$$RV = RV_{\text{activity}} + RV_{\text{planet}}, \quad (6)$$

where  $RV_{\text{activity}}$  can be found in Eq. 5, and  $RV_{\text{planet}}$ , in the case with no eccentricity, can be modelled by the following sinusoidal function:

$$RV_{\text{exoplanet}} = K \sin\left(\frac{2\pi}{P}(t - t_0)\right), \quad (7)$$

**Table 4.** P-values for the different coefficients used in Eq. 5 for the correction of stellar activity induced by an equatorial 1% spot on the simulated Sun, and a planet with a period of one-third the rotational period of the star and a semi-amplitude of  $10 \text{ m s}^{-1}$ . All the parameters corresponding to the Normal or SN variables are statistically significant, except the width of the CCF. Note that since nonlinear least squares was required, the residual standard error rather than the  $R^2$  is displayed as a reference.

Parameter	N mean RV	SN mean RV	SN median RV
$\beta_0$	0.00063	$2e-16$	$1.42e-09$
$\beta_1$	$2e-16$	$2e-16$	$2e-16$
$\beta_2$	$2e-16$	$2e-16$	$2e-16$
$\beta_3$	0.067	0.40	0.38
$\beta_4$	$2e-16$	$2e-16$	$2e-16$
$K$	$2e-16$	$2e-16$	$2e-16$
$P$	$2e-16$	$2e-16$	$2e-16$
$t_0$	$2e-16$	$2e-16$	$2e-16$
Residuals	$0.71 \text{ m s}^{-1}$	$0.66 \text{ m s}^{-1}$	$0.70 \text{ m s}^{-1}$

with amplitude  $K$ , orbital period  $P$ , and an epoch at the periastris  $t_0$ . The previous three unknown parameters define the planetary orbit.

The proposed model from Eq. 6 was fitted to the RV data and the results of the estimated model are summarized in Table 4. Except for the width parameters with coefficient  $\beta_3$ , all the other Normal or SN parameters are significantly useful to explain the RV variation induced by a spot plus a planet. We also observe that the RV residuals, once corrected for stellar activity and the presence of the planet, are comparable in terms of rms for all the three different RV estimates, with SN mean RV giving a slightly smaller value.

## 5. Real data application

In this Section we present the analysis conducted on the star Alpha Centauri B (Alpha Cen B), comparing the result of fitting a CCF using the SN density defined in Sec. 2.1 with the commonly employed approach based on fitting a Normal density for estimating the RV and retrieving the BIS SPAN for evaluating the asymmetry of the CCF. Four other stars have been analyzed with the proposed method and details can be found in Appendix A. For all the stars considered in the presented work, only CCFs that were derived from spectra that had at least a SNR of 10 at 550 nm were selected.

### 5.1. Comparison for Alpha Cen B of the different CCF parameters derived with the Normal and the Skew Normal

A total of 1808 CCFs that were derived from the spectra of Alpha Centauri B taken in 2010 by the HARPS spectrograph have been analyzed. Note that more observations were carried out that year, however only the data that was not significantly affected by contamination from Alpha Centauri A was used (see Dumusque et al. 2012). The selected observations represent probably, among all RV data existing, the best sampled and most precise RV data set showing strong solar-like activity signal (Thompson et al. 2017; Dumusque et al. 2012).

First, the correlation between  $\gamma$  and BIS SPAN is evaluated. In the left panel of Fig. 10, we see that the relationship between  $\gamma$  and the BIS SPAN is linear, with a slope equal to 720



**Fig. 6.** (left) Correlations between the different asymmetry parameters and their corresponding RV estimates in the case of an equatorial 1% spot on the simulated Sun. (right) Correlations between the different width parameters and their corresponding RV estimates for the same spot. In the presence of a spot, both the shape and the width of the CCF change as the star rotates. However, only the asymmetry produces a statistically significant correlation with the different RV estimates. The width parameters and their corresponding RV estimates present weak correlations and, in general, much weaker correlations compared to the results obtained when an equatorial 3% facula is present on the simulated Sun.

and a strong Pearson correlation coefficient of  $R = 0.95$ . This strong correlation suggests that both  $\gamma$  and BIS SPAN are measuring similar asymmetries for the CCFs. This strong correlation allows as well to convert the dimensionless  $\gamma$  parameters into  $\text{m s}^{-1}$  using the slope of the correlation, in this case  $720 \text{ m s}^{-1}$ .

The right plot of Fig. 10 displays the comparison between the RVs estimated using the SN density and the Normal density. The amplitude of the activity signal is slightly stronger for the SN mean RV (in the top-right plot the black circles of the SN mean RV tend to show more variability), while the signal measured using N mean RV or SN median RV are comparable. This behavior is similar to the one observed for the facula simulated with SOAP 2.0 in see Sec. 4.1, suggesting that the activity signal could be due to faculae present on Alpha Centauri B. Another argument that suggests the presence of faculae, is the strong positive measured correlation between  $\gamma$  and SN mean RV and SN FWHM and SN mean RV, as displayed in Fig. 11.

In Sec. 4.2, we saw that a spot induces a negative correlation between  $\gamma$  and SN mean RV, while weak correlations are measured between SN FWHM and SN mean RV. Dumusque (2014) had previously suggested that the activity of Alpha Centauri B could be due to faculae.

Similar to the analyses presented in Sec. 4, looking at Fig. 11 we compare the correlations between the asymmetry or the width parameters of the CCF and the RV. For this analysis, we also include the asymmetry parameters derived in Boisse et al. (2011),  $V_{span}$  and in Figueira et al. (2013), BIS-, BIS+, Bi Gauss and  $V_{asy}$ , as these authors found those asymmetry parameters more correlated to the RVs than BIS SPAN. It is clear in the case of Alpha Cen B, that the correlation found between  $\gamma$  and SN mean RV is the strongest. The Pearson correlation coefficient is  $R = 0.74$ , while it is at best  $R = 0.42$  for all the other asymmetry-RV correlations not derived using the SN density fit. When looking at the correlations between the width and the RV





**Fig. 7.** (top) The spurious estimated RVs (black dots) caused by a spot in the simulated data, the estimated RVs using Eq. 5 (red crosses) and the estimated RVs using the usual correction for stellar activity (green triangles), based on  $RV_{\text{activity}} = \beta_0 + \beta_1 \gamma + \beta_2 \text{SN FWHM}$  for the SN fit and on  $RV_{\text{activity}} = \beta_0 + \beta_1 \text{BIS SPAN} + \beta_2 \text{FWHM}$  for the normal fit. (bottom) The residuals from the model fit using Eq. 5 (red crosses) and the residuals from the usual correction (green triangles). The std. are also reported in the legend, and the residuals have a smaller systematic component when using the proposed model compared to the usual model. The tests of statistical significance on the parameters are presented in Table 3.

estimates for Alpha Cen B, the correlation is strongest for the SN parameters, with  $R = 0.82$  for SN FWHM-SN mean RV, while it is  $R = 0.70$  for FWHM and N mean RV.

Results illustrating the performance of the stellar activity correction proposed in Sec. 3 are displayed in Fig. 12. For Alpha Cen B, the RV estimated with SN mean RV has a std that is 35% larger than the std of the RV estimated with the N mean RV, and the std of SN median RV is 9% larger than that of the N mean RV. Even though we see these differences in the estimated RV, once we correct for stellar activity by using Eq. 5, the rms of the residuals are essentially the same for all three approaches. Although when studying the correlations between the different parameters we see that the parameters retrieved by using the SN density are more sensitive to stellar activity than the ones obtained with a Normal density fit, the proposed linear model for correcting from stellar activity is not able to perform better in the SN case than in the Normal case. When comparing

the new correction for stellar activity proposed in Sec. 3 with the usual correction, that uses only a linear combination of the width and the asymmetry of the CCF, we end up with very similar results. The new correction is however able to get RV rms that are  $6 \text{ cm s}^{-1}$  smaller than the usual correction.

When looking at the statistical tests of the different parameters used for correcting activity in Table 5, we see that the BIS SPAN (coefficient  $\beta_2$ ) does not statistically help in explaining stellar activity when using the parameters derived from the Normal density fit. However all the other parameters in the Normal and SN cases are statistically significant for modelling stellar activity. By analyzing the values of the coefficient of determination,  $R^2$ , we see that the model for the SN mean RV is able to capture the highest percentage of variability in the estimated RV. This is not a surprising result since for the three different RV estimates, we arrive to the same RV residual rms after correction



**Fig. 8.** RV estimates for N mean RV (red dashed line), SN mean RV (black line) or SN median RV (cyan dotted-dashed line). In this case, the CCFs have been generated using SOAP 2.0, considering an equatorial 1% spot on the simulated Sun in addition to a planet with a period of one-third of the rotational period of the star and with an amplitude of  $10 \text{ m s}^{-1}$ . The star does one full rotation between phase -0.5 and 0.5, with the spot being seen face-on at phase 0.

**Table 5.** P-values for the different coefficients used in Eq. 5 for the correction from stellar activity in Alpha Cen B data. All the variables corresponding to the Normal or SN parameters are statistically significant, except for the asymmetry of the CCF when using the BIS SPAN with the N mean RV. The evaluation of the  $R^2$  shows that the proposed linear combination better explains variations in RVs due to the stellar activity when the RVs are derived using SN mean RV.

Parameter	N mean RV	SN mean RV	SN median RV
$\beta_0$	0.49	0.90	0.027
$\beta_1$	$2.22e-16$	$2.22e-16$	$2.22e-16$
$\beta_2$	0.33	$2.22e-16$	$1.23e-11$
$\beta_3$	$2.22e-16$	$2.22e-16$	$2.22e-16$
$\beta_4$	$2.22e-16$	$2.22e-16$	$2.22e-16$
$R^2$	0.57	0.78	0.66

for activity, but before correction, SN mean RV shows the largest RV rms (see Fig 12).

### 5.2. Comparison for HD192310, HD10700, HD215152 and Corot-7 of the different CCF parameters derived with the Normal and the Skew Normal

In the previous section we evaluated the improvement obtained by the SN parameters compared to the Normal parameters and the BIS SPAN. We carryout similar analyses for four other main sequence stars: HD192310, HD10700, HD215152 and Corot-7. The same correlations plots and activity-after-correction plots displayed in the previous section for Alpha Cen B can be found for those new four stars in Appendix A.

The correlations between the parameters of these additional stars are similar to those obtained for Alpha Cen B. The correlation between  $\gamma$  and SN mean RV is the strongest among all the asymmetry-RV correlations. Concerning the correlations between the width parameters and the estimated RV, the strongest correlation often is between SN FWHM and SN mean RV. However, there is one exception in the case of HD10700 where the Pearson correlation coefficient between FWHM and N mean RV is equal to  $R = 0.53$ , while it is  $R = 0.42$  between SN FWHM and SN mean RV.

Except for the special case discussed above for HD10700, the analyses of those four stars, in addition to the analyses on Alpha Cen B, shows that the parameters derived when using a SN density are more sensitive to activity. Therefore using the SN parameters, and in particular defining RV as SN mean RV, can lead to better detection of stellar activity over the Normal parameters. This is the case for the evaluation of the asymmetry-RV correlations for Alpha Cen B, HD10700, HD215152, HD192310 and Corot-7, and the width-RV correlation for Alpha Cen B, HD215152, HD192310 and Corot-7 (see Appendix A).

In terms of correction for stellar activity, in the case of Alpha Cen B we see that although the RV rms is larger when measuring the RVs using SN mean RV (compared to the RVs obtained using N mean RV), once corrected for activity using the new model proposed in Sec. 3, both RVs estimates provides similar residuals. We also observe slight differences in the rms between the new proposed correction and the usual correction, which uses only a linear combination of the width and the asymmetry of the CCF. For Corot-7 the new correction is able to provide RV residual rms  $30 \text{ cm s}^{-1}$  smaller than the one obtained with the usual correction.

### 5.3. Detection limits when using the estimated RVs from the Normal or the Skew Normal models

In the previous section, we saw that the estimated RV measured when considering a SN or a Normal density resulted in different amplitudes, especially when using the SN mean RV. However, once corrected for stellar activity using the linear combination presented in Eq. 5, as shown in the bottom plots of Fig 12, the rms of the residuals are essentially the same for all three approaches. In this section, we investigate the ability of the three different RV estimators' (N mean RV, SN mean RV, and SN median RV) to detect planetary signals among stellar activity. To carryout this test, the minimum detected amplitude of a planetary signal is estimated at different orbital periods when considering data affected by stellar activity.

In order to obtain CCFs affected by realistic stellar activity signals, the CCFs from Alpha Cen B used previously were considered. To simulate a planetary signal, the CCFs were blue- or red-shifted with the desired amplitude, period, and phase. Several RV data sets with the same stellar signal, but different planetary signals were injected using parameters corresponding to the following grid:

- period of 3, 5, 7, 9, 11, 15, 20, 25, and 30 days,
- amplitude from  $0.5$  to  $3 \text{ m s}^{-1}$  by steps of  $0.05 \text{ m s}^{-1}$ ,
- 10 different phases, evenly sampled between  $0$  and  $2\pi$ .

For each of the 4500 simulations considered, the activity correction presented in Eq. 5 was applied, and signals in the residuals were investigated using a Generalized Lomb-Scargle periodogram (Lomb 1976; Scargle 1982; Zechmeister & Kürster 2009). If a signal with a p-value<sup>5</sup> smaller than 1% had a period compatible with the injected planetary period within an error budget of 20%, the signal was considered significant and the corresponding planet considered detected. Finally, for each period considered, the minimum amplitude at which at least 50% of the planets with different phases were detected was found. The results are shown in Fig. 13. All three of the different RV estimators give similar detection limits. Though more investigations should be considered, these results suggest that any of the RV estimators can be used when searching for a planetary signal in RV data contaminated by stellar activity.

<sup>5</sup> The p-values were estimated using a bootstrap procedure.



**Fig. 9.** Evaluation of the correlation between the RVs and the asymmetry parameters of the simulated data with a 1% spot and an injected planetary signal. The shape of the CCF changes as the spot moves, producing statistically significant correlations only between the estimated RVs and the asymmetry parameter. The correlations between the estimated RVs and the width parameter of the CCF are weaker than the case with only a spot.

Although the detection limits remain relatively constant for the periods  $\leq 15$  days, there is a notable increase in the detection limit at 20 and 30 days. This is likely due to the periods of the simulated planets, 20 days and 30 days, are close to the first harmonic or to the rotational period of the star (36.7 days, Dumusque et al. 2012) and therefore close to the semi-periodicity of the stellar activity signal. In such a case, the activity correction absorbs part of the planetary signal and a stronger planetary signal is needed at these periods in order for them to be detectable.

## 6. Estimation of standard errors for the CCF parameters

In this section, we investigate how the photon noise influences the CCF parameters derived either by a Normal density or a SN density fit. Because a CCF is obtained from a cross-correlation, each point of a CCF is correlated with the other points. Therefore, we cannot simply vary each point in the CCF by their re-

spective error bars and then recalculate the best SN or Normal density fit to see how the CCF noise influences the estimation of the parameters of interest (i.e., N mean RV, SN mean RV, SN median RV, FWHM, SN FWHM, BIS SPAN and  $\gamma$ ). Instead, we go to the individual spectra where the individual measurements can be considered independent for the other points. The standard error on each point of a spectrum is given by the photon noise, which follows a Poisson distribution and is therefore estimated by taking the square-root of the measured flux.

The following method was carried out in order to estimate the error bars on the different parameters derived from the CCF. We first modify the values of all the points in the spectrum given their respective error bars. To do so random Gaussian noise centered on the value of the point and with standard deviation the square-root of the flux was added across each spectrum. The CCF was calculated using this spectrum according to the method presented in Pepe et al. (2002), then fitted either a Normal or SN density to this CCF and recorded the different parameters. This



**Fig. 10.** (left) Correlation between  $\gamma$  and the BIS SPAN for Alpha Centauri B. The strong correlation suggests these two parameters are similarly measuring the asymmetry. (top right) RVs as function of Julian Day for Alpha Centauri B in 2010. The RVs are estimated using the mean of a Normal fitted to the CCF (red triangles), or the mean (black circles) or median (cyan crosses) of a SN density fitted to the CCF. (bottom right) Differences between the RVs estimated with the SN density and those from the Normal density.

process was repeated a hundred times in order to obtain a distribution for each CCF parameter. The standard deviations of the resulting distributions provide standard errors for each CCF parameters.

The standard errors were computed for each CCF parameter for all the HARPS measurements of HD215152, HD192310 and Corot-7. These three stars include measurements that cover the range of SNR measured at 550 nm (SNR550) from 10 to 500, which represent a very low SNR and the saturation limit of the HARPS detector, respectively. HD10700 and Alpha Cen B were not included because they have a large number of measurements, which would require a substantial computational effort. The variation of the noise for each CCF parameter as a function of SNR550 is displayed in Fig. 14. The top plots show the standard errors of the three different estimated RV's, the width, and the asymmetry estimates. Note that because BIS SPAN and  $\gamma$  do not have the same units, the estimated slopes of the correlation between those two parameters to transform  $\gamma$  in  $\text{m s}^{-1}$  were used (see Fig. 10 and Table A.1 for the value of the slope for each star). The bottom plots show the ratio between the standard errors measured when using the SN parameters and the Normal parameters. Values smaller or larger than one will imply that standard errors from the SN parameters are respectively more or less precise than the Normal parameters.

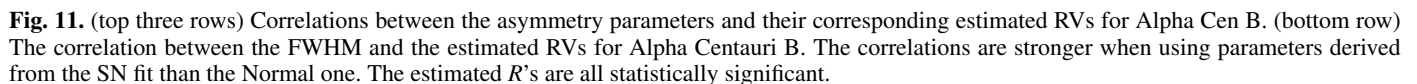
The different RV estimates all appear to follow a similar exponential decay. Considering each RV estimate individually, it appears that all the measurements follow the same curve even though the measurements are from three different stars. This suggests that for the stars considered, the precision in each RV estimate may be driven by the SNR of the analyzed spectra. This is not surprising as the three stars studied have very similar spectral types. They are all main sequence K-dwarfs.

When comparing the three different estimates for the RV, the SN mean RV presents standard errors that are 60% larger than the N mean RV. However, the SN median RV gives errors 10% more precise than N mean RV. The parameters describing the width of the CCF, FWHM and SN FWHM, have comparable standard errors. Finally, for the asymmetry parameters,  $\gamma$  has standard errors that are 15% more precise than BIS SPAN. In conclusion, when fitting a SN density to the CCF and using as RV estimate SN median RV, we are able to improve by 10% the precision on the estimated RV. using the SN density we are also able to improve by 15% the precision on the estimated asymmetry parameter of the CCF. Both these improvements do not lead to a worse precision on the estimated width of the CCF. SN mean RV should not be used to derive precise estimated RVs, except perhaps in specific conditions described below, as the precision on this parameter is 60% worse than the precision on the RVs derived from N mean RV.

## 7. Discussion

When fitting a SN density to the CCF, parameters that describe the RV (i.e. the CCF barycenter), the amplitude (sometimes called the CCF contrast), the width and the asymmetry of the CCF can all be estimated in a single model framework. For the estimation of the RV, we investigated the use of the mean and the median of the SN density. Then the width is derived using the variance of the SN density ( $\text{SN FWHM} = 2\sqrt{2\ln(2)\sigma^2}$ ) and the asymmetry by using  $\gamma$ , the naturally available skewness parameter of the SN density.

To evaluate the performance of the proposed SN framework, tests on both simulated and real data were carried out and compared to the commonly employed approach consisting of fitting a Normal density shape (for estimating the RV and FWHM) and



We also propose a new model to correct the estimated RV data for stellar activity signals. Generally, when fitting for plan-

The conclusion found for the SOAP 2.0 simulations that the correlation between the CCF asymmetry and the estimated RV, as well as between the CCF width and the estimated RV, is stronger for the SN parameters than the Normal ones also holds when using real data. However, the stronger asymmetry-RV correlation for the simulated data is between  $\gamma$  and SN median RV, while it is between  $\gamma$  and SN mean RV for the real data. There could be several reasons for this difference. It may be due to



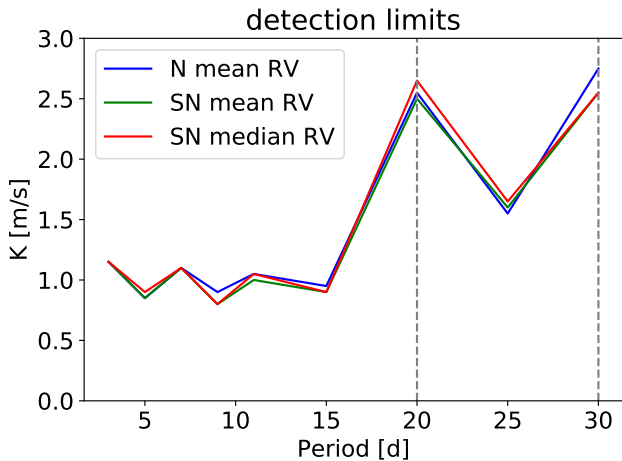


**Fig. 12.** (top) The RVs (black dots) for Alpha Centauri B estimated using a SN and a Normal fit. (bottom) The residuals from the model fit using Eq. 5 (New corr. std, black dots) and the residuals from the usual correction (Usual corr. std, blue triangles), based on  $RV_{\text{activity}} = \beta_0 + \beta_1\gamma + \beta_2\text{SN FWHM}$  for the SN fit and on  $RV_{\text{activity}} = \beta_0 + \beta_1\text{BIS SPAN} + \beta_2\text{FWHM}$  for the normal fit. The residuals have a smaller systematic component when using the proposed model of Eq. 5 (black dots) compared to the usual model (blue triangles).

the fact that SOAP 2.0 uses the spectrum of a spot as input to model the stellar activity induced by a facula. Because the temperature between a spot and a faculae is significantly different, their spectra should be different. Additionally, there are expected to be multiple active regions on a star at different locations in longitude and latitude, while the SOAP 2.0 data used in the simulation study included only a single active region. Nevertheless, in both the simulated and the real cases, the parameters derived from the SN are generally more sensitive to activity than the parameters derived from the Normal. Also, the correlation between the asymmetry and the SN mean RV is always stronger than the parametrization of the CCF presented in Boisse et al. (2011) and Figueira et al. (2013). Since if an apparent RV signal is induced by activity generally a strong correlation will be observed between the RV and the FWHM of the CCF or its BIS SPAN, to use the SN density to fit the CCF leads to a better understanding of the spurious variations in RV caused by stellar activity.

Considering the different RV estimates of the real data, the amplitude of stellar activity tends to be largest for the SN mean RV, followed by the SN median RV, and finally by N mean RV. Although there appears to be a larger difference between the SN mean RV and the SN median RV, the SN median RV is similar to the N mean RV. Hence, the mean of the SN density appears to be more sensitive to variation in the CCF shape than the median of the SN or the mean of the Normal.

Having an estimator of RV that is more sensitive to stellar activity, such as the SN mean RV, can help with better probing stellar rotational periods or better understanding the covariance of stellar signals when fitting a Gaussian Process to the RVs (e.g. Faria et al. 2016; Haywood et al. 2014). We saw in the previous section that the SN mean RV estimator is 60% noisier than the N mean RV estimator. This is not necessarily a negative aspect. Indeed, in the case where photon-noise is not dominant, stellar activity can be better characterized if its effect can be amplified



**Fig. 13.** Detection limits of planetary signals once the stellar activity signal is removed from the raw RVs using the model proposed in Eq. 5. *[[Jessi: For the black-and-white versions, it would be helpful to have different line types plotted]]* The different RV estimators have similar detection limits.

by measuring the SN mean RV, even if the uncertainties related to this parameter are increased by 60%.

Even if the different RV estimators generally have different amplitudes, once the new correction for stellar activity is applied, the residuals of the model have similar rms. When comparing the activity correction proposed in this paper with the usual correction that only uses a linear combination of the CCF asymmetry and width, for the simulations based on the presence of a facula or a spot the new proposed correction almost entirely explains the spurious variations in RV. However, when moving to real data, there is just a slight improvement by using the proposed correction function for stellar activity. Anyway, we encourage the community to use this new simple linear model if the goal is to reduce activity without considering the computationally expensive use of Gaussian Process. A further analysis of the p-value for each coefficient can be performed to see the ones that are not relevant and that can therefore be removed.

A test was carried out to see if some RV estimates were better at finding planets in RV data affected by observed stellar signals. When using the new correction proposed in this paper to mitigate the effect from stellar activity and when assessing planetary detection based on a periodogram analysis, we find very similar detection limits for the three RV estimators presented in this paper. *[[Umberto: trying to understand why we get better results for P=20 days when using the usual correction]]* Therefore it seems that any of the RV estimate can be used to search for planetary signals.

Finally, we investigated the precision of each of the SN and Normal parameters including the BIS SPAN. It turns out that the SN mean RV should not be used to get precise RVs as the standard errors on this parameter is 60% greater than the N mean RV. However, the SN median RV is 10% more precise than the N mean RV. Regarding the asymmetry estimates, we observe that  $\gamma$  has a precision 15% better than the BIS SPAN.

## 8. Conclusion

When searching for low-mass exoplanets using the RV technique, it is necessary to retrieve precise estimates of the RV and

also to account for variations induced by stellar activity in order to avoid false detections. Stellar activity such as spots and faculae can lead to shape variations in the spectra features, which then results in shape variations of the CCF. The correlations between the width or asymmetry of the CCF and the estimated RV are commonly used as a way to detect if the RVs are affected by stellar activity signals. Because the presence of real planets would result in only a shift in the CCF (not a change on its shape), strong correlations between the shape features of the CCF and the estimated RVs suggest that stellar activity may be present.

In this paper, a new approach for measuring shape changes in the CCF is proposed using the SN density to estimate the RV and the shape of the CCF. This new method is compared to a commonly used method based on a Normal density fit to the CCF. The mean of the Normal density is used as the estimated RV and the FWHM estimates the width of the CCF. Because the Normal density does not have any skewness, another method is necessary to estimate the asymmetry of the CCF; for this goal the BIS SPAN is often employed. In addition, the proposed SN approach is compared to other parameterizations of the CCF asymmetry that have been shown to be sensitive to activity signals (Boisse et al. 2011; Figueira et al. 2013).

In the different tests carried out for this work, the SN parameters performed at least as well (and often times better) as the parameters from the Normal approach and the BIS SPAN. The SN parameters  $\gamma$ , SN FWHM and SN mean RV always show stronger correlation than the correlations between any of the parameters derived by the Normal and the BIS SPAN, or the different asymmetry parametrizations presented in Boisse et al. (2011) and Figueira et al. (2013). This suggests the SN parameters may be better at probing stellar activity signals than the other methods. In addition, the uncertainties measured on SN median RV and  $\gamma$  are respectively 10% and 15% smaller than the corresponding uncertainties on N mean RV and BIS SPAN.

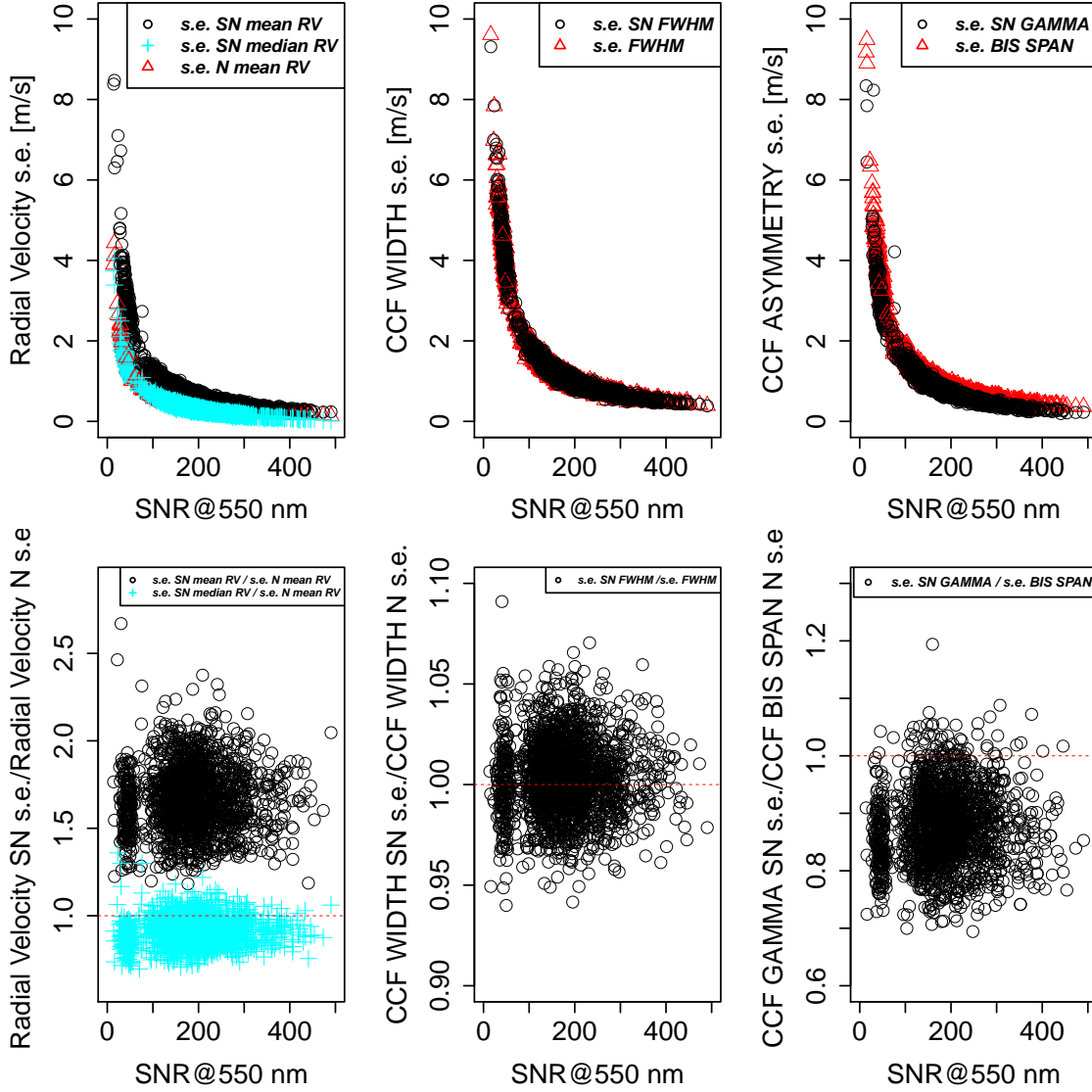
Because of the advantages of using the proposed SN approach over the commonly employed approach based on the Normal density fit to the CCF and the BIS SPAN or the asymmetry parameters described in Boisse et al. (2011) and Figueira et al. (2013), to fit a SN density to the CCF may be more useful for detecting stellar activity than the previously proposed parametrizations. Correlations between  $\gamma$  and the SN mean RV, and between the width and the SN mean RV can be used to probe stellar activity signals in RV data, and the SN median RV can be used to estimate RV.

## 9. Acknowledgements

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## Appendix A: Appendix

In this Appendix, a similar analysis as the one presented in Sec. 5 is discussed for four main-sequence stars: HD192310 (K2V),



**Fig. 14.** Bootstrap analyses on the stars HD215152, HD192310 and Corot-7. (top) Comparison between the standard errors from the bootstrap analysis of the estimated RVs, FWHM, and asymmetry parameters using the SN fit and the common strategy (Normal fit and BIS SPAN). (bottom) Ratio between the standard errors retrieved on the parameters derived from the common strategy and the corresponding standard errors retrieved on the parameters derived from the SN fit. When using SN mean RV (black circles), the standard errors are in average 60% larger than the standard errors of N mean RV (red triangles). However, the standard errors for SN median RV (cyan crosses) are on average 10% smaller than the standard errors coming from the N mean RV. The use of the asymmetry SN parameter  $\gamma$  leads to standard errors in average 15% smaller than the standard errors related to the BIS SPAN. Note that for the asymmetry, the error in BIS SPAN is in  $\text{m s}^{-1}$ . To be able to compare the errors in  $\gamma$  and BIS SPAN, we multiplied the error in  $\gamma$  by the slope of the correlation between  $\gamma$  and BIS SPAN.

HD10700 (G8V), HD215152 (K3V) and finally Corot-7 (K0V). [\[\[Xavier: All those data are public and were found on the ESO archive. We have an acknowledgement for that\]\]](#) [\[\[Jessi: Are there papers that can/should be cited here for the datasets used?\]\]](#)

Table A.1 summarizes the results obtained by the SN and Normal densities. These results are consistent with those from the analysis of Alpha Centauri B. The correlation between  $\gamma$  and SN mean RV is stronger than the correlation between the BIS SPAN and N mean RV or between the asymmetry parameters derived in Boisse et al. (2009) and Figueira et al. (2013) and N mean RV for all the considered stars. The correlation between SN FWHM and SN mean RV is stronger than the correlation between FWHM and N mean RV for three of the four stars. Also for all these stars the originally estimated RVs were corrected

from spurious variations in RVs caused by stellar activity using Eq. 5. Fig. A.2–A.6 show the resulting corrected RVs. Once corrected from stellar activity, the Normal and SN residuals are comparable for the stars HD192310 and HD10700. However, the rms of the residuals for HD215152 are  $0.05 \text{ m s}^{-1}$  lower for the SN model than the Normal model, and the rms of the residuals for Corot-7 are  $0.34 \text{ m s}^{-1}$  lower. HD215152 and Corot-7 have in average a lower SNR at 550 nm than HD10700 and HD192310. The average SNR at 550 nm for the stars HD10700, HD192310, HD215152 and Corot-7 are respectively 273, 207, 141 and 44.

**Table A.1.** Notable correlations between the asymmetry or the FWHM parameters and the RVs for four stars: HD192310, HD10700, HD215152 and Corot 7. The complete results of the analyses of the correlations for the four stars are presented in Fig. A.1–A.7.

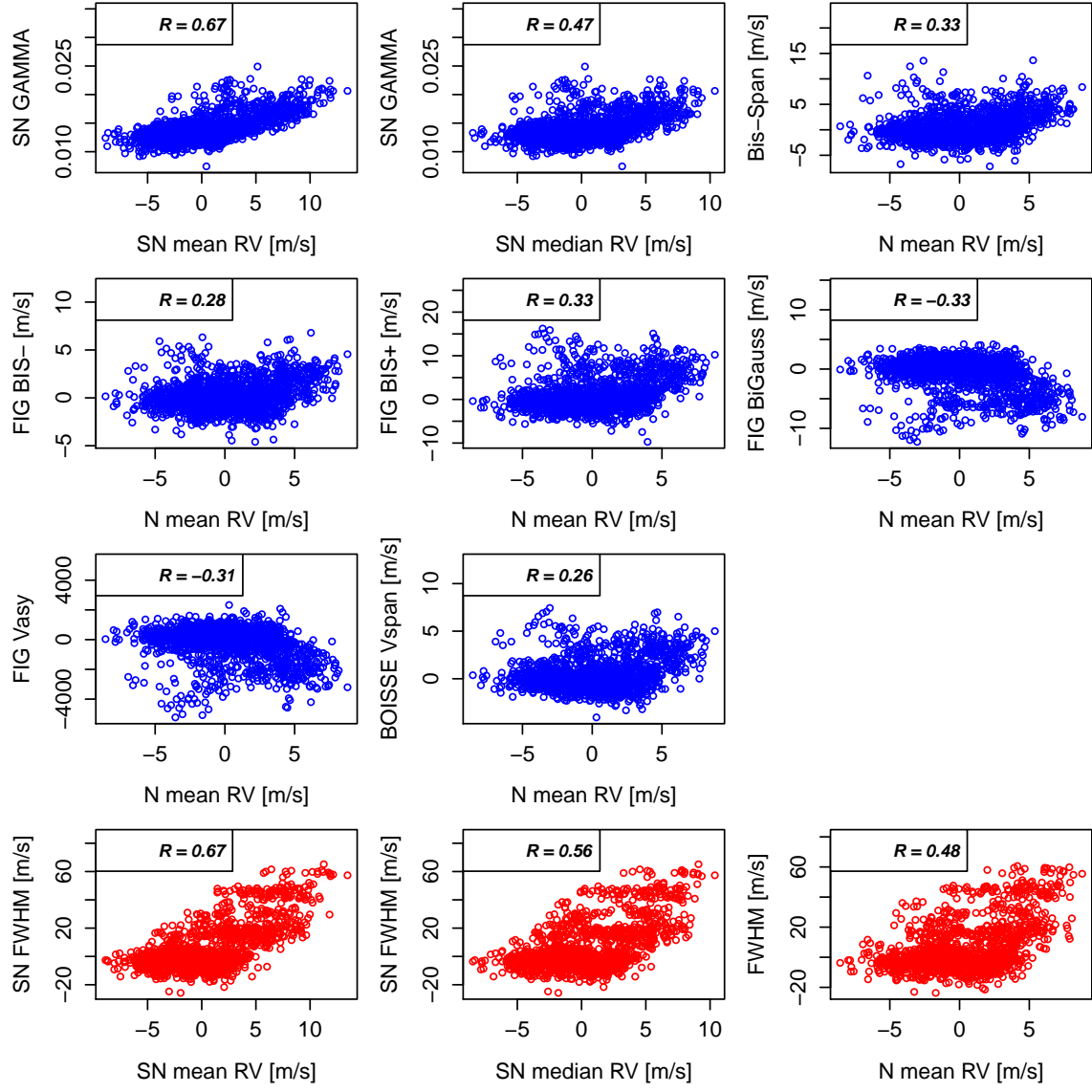
Star	# CCFs	R(SN $\gamma$ , Bis-Span)	slope(SN $\gamma$ , Bis-Span)	R(SN $\gamma$ , SN mean RV)
HD192310	1577	0.888	786	0.669(0.64; 0.695)
HD10700	7928	0.78	604	0.322(0.302; 0.342)
HD215152	273	0.763	794	0.571(0.485; 0.646)
Corot 7	173	0.814	607	0.561(0.450; 0.656)

Star	R(Bis-Span, N mean RV)	R(FIG BiGaussian, N mean RV)	R(SN FWHM, SN mean RV)	R(FWHM, N mean RV)
HD192310	0.329(0.285; 0.373)	−0.333(−0.376; −0.289)	0.666(0.637; 0.692)	0.476(0.4367; 0.514)
HD10700	−0.073(−0.095; −0.0051)	0.127(0.105; 0.148)	0.421(0.403; 0.439)	0.529(0.513; 0.545)
HD215152	−0.067(−0.184; 0.052)	0.269(0.155; 0.376)	0.210(0.094; 0.321)	−0.138(−0.253; −0.020)
Corot 7	0.092(−0.058; 0.238)	−0.335(−0.228; −0.082)	−0.709(0.626; 0.776)	0.595(0.489; 0.683)

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**Fig. A.1.** (top three rows) Correlations between the asymmetry parameters and their corresponding RVs for HD192310. (bottom row) Correlations between the FWHM and the estimated RVs. The correlations are consistently stronger when using parameters derived from the SN than the Normal. The estimated  $R$  are all statistically significant.





**Fig. A.2.** (top) The RVs (black dots) for HD192310 estimated using a SN and a Normal fit. (bottom) The residuals from the model fit using Eq. 5 (New corr. std—black dots) and the residuals from the usual correction (Usual corr. std—blue triangles), based on  $RV_{\text{activity}} = \beta_0 + \beta_1\gamma + \beta_2\text{SN FWHM}$  for the SN fit and on  $RV_{\text{activity}} = \beta_0 + \beta_1\text{BIS SPAN} + \beta_2\text{FWHM}$  for the normal fit. The residuals for both the proposed correction from stellar activity are comparable.



**Fig. A.3.** (top three rows) Correlations between the asymmetry parameters and their corresponding RVs for HD10700. (bottom row) Correlations between the FWHM and the RVs for HD10700. The correlations are consistently stronger when using SN mean RV compared to N mean RV for the asymmetry parameters; however, the correlation between the FWHM and the N mean RV, only for this quiet star, is stronger than the analogous correlations with the estimated SN RVs. The estimated  $R$  are statistically significant, except for the correlation between FIG BIS and RV ( $p$ -values=0.36).



**Fig. A.4.** (top) The RVs (black dots) for HD10700 estimated using a SN and a Normal fit. (bottom) The residuals from the model fit using Eq. 5 (New corr. std–black dots) and the residuals from the usual correction (Usual corr. std–blue triangles), based on  $RV_{\text{activity}} = \beta_0 + \beta_1 \gamma + \beta_2 \text{SN FWHM}$  for the SN fit and on  $RV_{\text{activity}} = \beta_0 + \beta_1 \text{BIS SPAN} + \beta_2 \text{FWHM}$  for the normal fit. The residuals for both the proposed correction from stellar activity are comparable.





**Fig. A.6.** (top) The RVs (black dots) for HD215152 estimated using a SN and a Normal fit. (bottom) The residuals from the model fit using Eq. 5 (New corr. std–black dots) and the residuals from the usual correction (Usual corr. std–blue triangles), based on  $RV_{\text{activity}} = \beta_0 + \beta_1\gamma + \beta_2\text{SN FWHM}$  for the SN fit and on  $RV_{\text{activity}} = \beta_0 + \beta_1\text{BIS SPAN} + \beta_2\text{FWHM}$  for the normal fit. The residuals for both the proposed correction from stellar activity are comparable.





**Fig. A.7.** (top three rows) Correlations between the asymmetry parameters and their corresponding RVs for Corot 7. (bottom row) Correlations between the FWHM and the RVs for Corot 7. The correlations are consistently stronger when using parameters derived from the SN than the Normal. The p-values associated with each  $R$  are not statistically significant for the correlation between N mean RV and BIS SPAN (p-values=0.23) and the correlation between N mean RV and FIG BIS- (p-values=0.11).



**Fig. A.8.** (top) The RVs (black dots) for Corot 7 estimated using a SN and a Normal fit. (bottom) The residuals from the model fit using Eq. 5 (New corr. std–black dots) and the residuals from the usual correction (Usual corr. std–blue triangles), based on  $RV_{\text{activity}} = \beta_0 + \beta_1 \gamma + \beta_2 \text{SN FWHM}$  for the SN fit and on  $RV_{\text{activity}} = \beta_0 + \beta_1 \text{BIS SPAN} + \beta_2 \text{FWHM}$  for the normal fit. The residuals have a smaller systematic component when using the proposed model of Eq. 5 (black dots) compared to the usual model (blue triangles). Moreover, once corrected for stellar activity using Eq. 5, the remaining std from the SN models are  $0.334 \text{ m s}^{-1}$  smaller than the remaining std of the Normal model.