# Higgs boson production at the Large Hadron Collider: accurate theoretical predictions at higher orders in QCD

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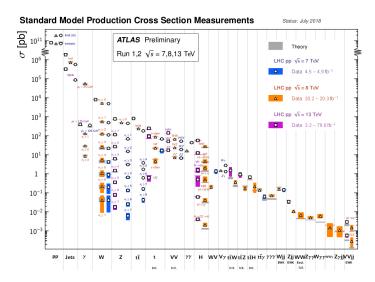


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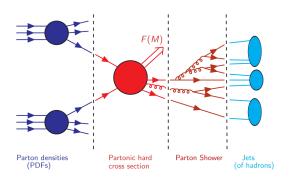
### Outline

- QCD and collider physics
  - QCD Factorization
  - Partonic cross section and perturbative QCD
- All order perturbative resummation
  - Higher order radiative corrections
  - Resummation of large logarithmic corrections
- Precise and fast predictions for Higgs boson physics
  - Higgs production at the LHC
  - HTurbo numerical code
  - Preliminary results & Conclusions

#### LHC physics



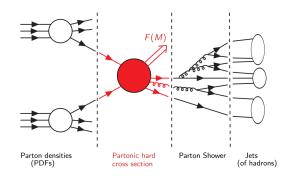
#### Factorization theorem



Compute hadronic cross sections is a hard problem  $\longrightarrow$  QCD Factorization

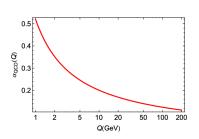
$$\sigma^{\mathrm{F}}(p_1, p_2) = \int_0^1 dx_1 dx_2 \ f_{\alpha}(x_1, \mu_F^2) * f_{\beta}(x_2, \mu_F^2) * \hat{\sigma}_{\alpha\beta}^{\mathrm{F}}(x_1 p_1, x_2 p_2, \alpha_s(\mu_R^2), \mu_F^2)$$

#### Partonic cross section



- Parton densities (PDFs)  $f_{\alpha}(x_i, \mu_F^2)$ : non perturbative but universal
- Partonic cross section  $\hat{\sigma}_{\alpha\beta}^{F}$ : process dependent but computable as perturbative series in  $\alpha_s$

#### QCD coupling



 Running coupling given by Renormalization Group Equation (RGE)

$$lpha_s(\mu) = rac{1}{eta_0 \log \left(rac{\mu^2}{\Lambda_{ ext{QCD}}^2}
ight)}$$

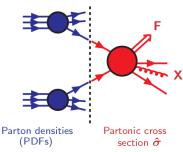
- $\beta_0$  LO of the  $\beta$  function, is > 0
- $\Lambda_{\rm QCD}$ , parameter that defines value of the coupling at large scales

QCD is weakly coupled for  $\mu >> \Lambda_{\rm QCD} \longrightarrow$  asymptotically free

Perturbative Quantum Chromodynamics (pQCD)

#### Perturbative QCD

- Born cross section is the leading-order (LO) term of the perturbative series
- $\sigma^{(1)}, \sigma^{(2)}, \sigma^{(3)}$  are the NLO, NNLO, N3LO corrections



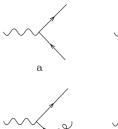
$$\hat{\sigma} = \sigma^{\text{Born}} \Big( 1 + \alpha_s \sigma^{(1)} + \alpha_s^2 \sigma^{(2)} + \alpha_s^3 \sigma^{(3)} + ... \Big)$$

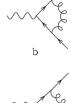
Lower order predictions strongly depend on the auxiliary and unphysical renormalization and factorization scales  $\longrightarrow$  Need higher order corrections to increase theoretical accuracy!

All order perturbative resummation

#### Higher order corrections

- Calculation of higher order corrections is not an easy task due to infrared (IR) soft and collinear singularities
- Final state singularities cancel by combining real and virtual contributions
- Initial state collinear singularities factorized inside the PDFs

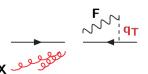




 $q_{\perp}$  resummation

Study the differential  $q_{\perp}$  distribution

$$h_1(p_1) + h_2(p_2) \longrightarrow F(M, \mathbf{q}_{\perp}) + X$$



$$\int_0^{Q_\perp^2} dq_\perp^2 \frac{d\hat{\sigma}}{dq_\perp^2} \sim c_0 + \alpha_s (c_{12}L^2 + c_{11}L + c_{10}) + ..., \quad \text{where} \quad L = \ln(q_\perp/M^2)$$

$\alpha_{S}L^{2}$	$\alpha_{\mathcal{S}}L$	 $\mathcal{O}(lpha_{\mathcal{S}})$
$\alpha_S^2 L^4$	$\alpha_S^2 L^3$	 $\mathcal{O}(\alpha_S^2)$
$\alpha_S^n L^{2n}$	$\alpha_S^n L^{2n-1}$	 $\mathcal{O}(\alpha_S^n)$
dominant logs		 

Truncated fixed order predictions  $\rightarrow$  enhanced  $\alpha_s^n \ln^m(M^2/q_\perp^2)$  appear

 $q_{\perp}$  resummation

Separate partonic  $q_{\perp}$  distribution as follows

$$\frac{d\hat{\sigma}_{ab}}{dq_{\perp}^2} = \left[\frac{d\hat{\sigma}_{ab}^{(\rm res.)}}{dq_{\perp}^2}\right]_{\rm l.a.} + \left[\frac{d\hat{\sigma}_{ab}^{(\rm fin.)}}{dq_{\perp}^2}\right]_{\rm f.o.} \quad , \quad {\rm such \ that}$$

$$\begin{split} &\int_0^{q_\perp^2} dq_\perp^2 \frac{d\hat{\sigma}_{ab}^{(\mathrm{res.})}}{dq_\perp^2} \sim \sum \alpha_s^n \log^m \frac{M^2}{q_\perp^2} \quad \mathrm{for} \quad q_\perp \to 0 \\ &\lim_{q_\perp \to 0} \int_0^{q_\perp^2} dq_\perp^2 \frac{d\hat{\sigma}_{ab}^{(\mathrm{fin.})}}{dq_\perp^2} = 0 \end{split}$$

Resummed and finite components can be matched (LL+LO, NLL+NLO, NNLO+NNLL, ...) to have uniform accuracy in a wide range of  $q_{\perp}$ 

 $q_{\perp}$  resummation

Resummation holds in impact parameter space b

$$rac{d\hat{\sigma}_{\mathsf{a}\mathsf{b}}^{(\mathrm{res.})}}{dq_{\perp}^{2}} = rac{\mathsf{M}^{2}}{\hat{\mathsf{s}}} \int db \; rac{\mathsf{b}}{2} \; J_{0}(\mathit{b}q_{\perp}) \; \mathcal{W}_{\mathsf{a}\mathsf{b}}(\mathit{b}, \mathsf{M})$$

with  $\mathcal{W}_{ab}$  also expressed in Mellin space (with respect to  $z=M^2/\hat{s}$ )

$$\mathcal{W}_N(b, M) = \mathcal{H}_N(\alpha_s) \times \exp\{\mathcal{G}_N(\alpha_s, L)\}$$
 being  $L \equiv \log(M^2 b^2)$ 

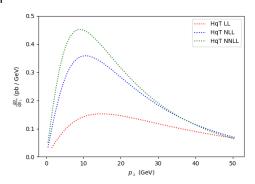
- Large logarithms exponentiated in the universal Sudakov form factor  $\mathcal{G}_N(\alpha_s, L)$
- Constant (b-independent) terms factorized in the process dependent hard factor  $\mathcal{H}_N(\alpha_s)$

Accurate predictions for Higgs boson at the LHC

# HqT and HRes

#### Predictions for Higgs $q_{\perp}$ distribution

- q⊥ resummation implemented in numerical codes HqT and HRes [Catani, de Florian, Ferrera, Grazzini, Tommasini]
- Higher order accuracy require high computation times
- Codes producing fast and accurate predictions are needed for precision era of the LHC



### **HTurbo**

#### Starting point DYTurbo

Numerical code **DYTurbo** [Camarda et al.] ref. at 1910.07049, fast and precise  $q_{\perp}$  resummation and several improvements for Drell-Yan  $(h_1h_2 \rightarrow V + X \rightarrow I^+I^- + X)$ 

- First goal: set up a numerical code for Higgs boson production starting from DYTurbo
- Set LO amplitude  $gg \rightarrow H$
- Set Sudakov and Hard coefficients for Higgs production
- Compare with HRes and HqT

Final goal: extend theoretical accuracy up to N<sup>3</sup>LL+N<sup>3</sup>LO

### HTurbo

#### Starting point DYTurbo

Both Sudakov factor  $\mathcal{G}_N$  and hard coefficient  $\mathcal{H}_N$  can be expanded as perturbative series in  $\alpha_s$ 

$$\mathcal{G}_{N}(\alpha_{s},L) = L g^{(1)}(\alpha_{s}L) + g^{(2)}(\alpha_{s}L) + \frac{\alpha_{s}}{\pi}g^{(3)}(\alpha_{s}L) + \dots$$
$$\mathcal{H}_{N}(\alpha_{s}) = 1 + \alpha_{s}\mathcal{H}^{(1)} + \alpha_{s}^{2}\mathcal{H}^{(2)} + \dots$$

For each new order implement a factor of  $\mathcal{G}_N$  and Hard  $\mathcal{H}_N$ 

$$LL(\sim \alpha_s^n L^{n+1}) : g^{(1)}, \hat{\sigma}^{(0)}$$

$$NLL(\sim \alpha_s^n L^n) : g^{(2)}, \mathcal{H}^{(1)}$$

$$NNLL(\sim \alpha_s^n L^{n-1}) : g^{(3)}, \mathcal{H}^{(2)}$$

Start by building predictions up to NNLO+NNLL, then add  $$N^3LO\!+\!N^3LL$$ 

# **HTurbo**

#### Code optimization

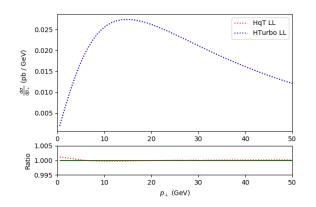
### Reimplementation of **HqT** and **HRes** for $q_T$ -resummation

- ullet C++ structure with **Fortran** interfaces o Multi-threading
- Optimization in the integration routines / integral transforms
  - Factorize boson and decay kinematics
  - Gauss-Legendre quadrature rules (1-dim.)
  - Vegas/Cuhre through Cuba (multi-dim.)

### Comparison HRes and HTurbo - speed performance

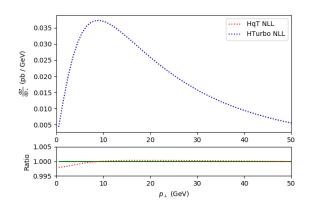
Predictions	HRes	HTurbo
resummed NNLL	10h	10'
combined NNLO+NNLL	20h	1h

#### Comparison HTurbo and HqT - LL



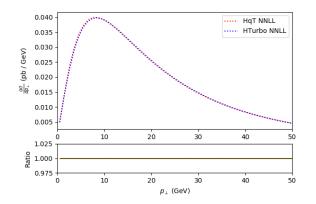
- HTurbo  $q_{\perp}$  distribution vs HRes and HqT at LL
- Excellent numerical agreement up to the 0.1% level

#### Comparison HTurbo and HqT - NLL



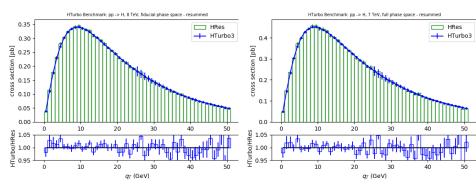
- HTurbo  $q_{\perp}$  distribution vs HRes and HqT at NLL
- ullet Excellent numerical agreement up to the 0.1% level

#### Comparison HTurbo and HqT - NNLL



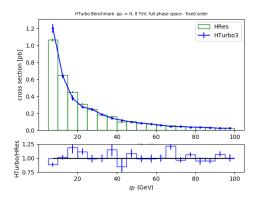
- HTurbo  $q_{\perp}$  distribution vs HRes and HqT at NNLL
- ullet Excellent numerical agreement up to the 0.1% level

#### Comparison HTurbo and HRes - resummed



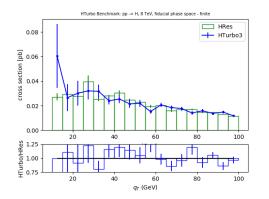
- Represent full (LHS) and fiducial (RHS) phase space √
- ullet Agreement up to resummed NNLL  $\longrightarrow$  ready for N<sup>3</sup>LL

#### Comparison HTurbo and HRes - fixed order



- Represent fiducial (RHS) phase space HRes vs HTurbo √
- Agreement up to fixed order NL0  $\longrightarrow$  ready for N<sup>3</sup>LL

#### Comparison HTurbo and HRes - finite



- Represent fiducial phase space HRes vs HTurbo √
- Agreement up to finite NNLL+NLO → ready for N<sup>3</sup>LL

# Summary & Conclusions

- Fast and accurate predictions are needed towards the precision era of the LHC
- ② Developing a novel numerical code, **HTurbo**, which implements  $q_{\perp}$  resummation for Higgs boson production
- 4 HTurbo is faster than any of the existing codes
- Outlook of thesis work:
  - Add N<sup>3</sup>LO+N<sup>3</sup>LL prediction
  - Perform phenomenological studies comparing with LHC data

### Thank you!



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