

HTurbo: Fast predictions for Higgs production at the LHC

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① QCD in a nutshell

- Factorization in QCD
- Partonic cross section
- Perturbative QCD

② Dealing with divergences

- Higher order corrections
- Resummation in QCD

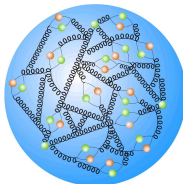
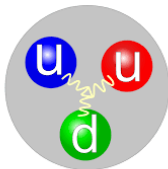
③ HTurbo

- Higgs production at the LHC: HRes and HqT
- HTurbo: Fast predictions for Higgs production
- Results & Conclusions

QCD in a nutshell

QCD in a nutshell

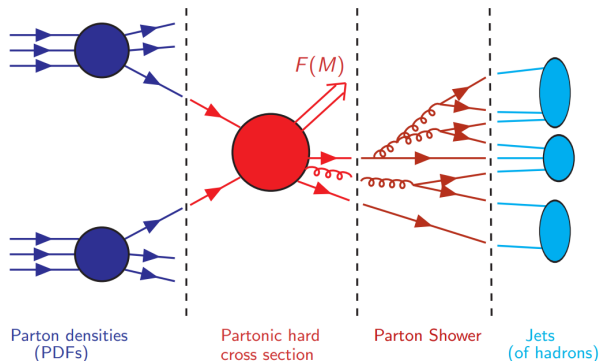
Address nuclear structure



- Hadrons made of partonic objects \longrightarrow non perturbative physics
- Smash two hadrons to explore internal structure \longrightarrow LHC physics

QCD in a nutshell

Factorization theorem

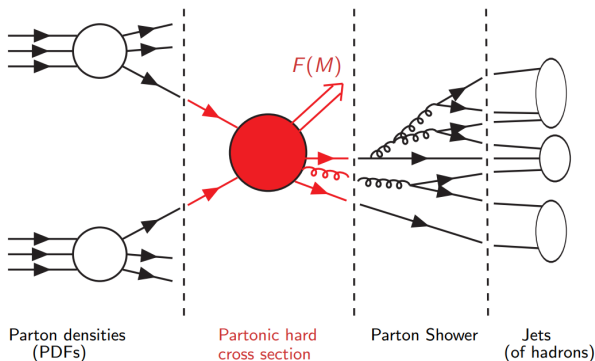


Factorize the problem \longrightarrow Convolute the PDFs with the partonic $\hat{\sigma}_{ij}$

$$\sigma = \int_0^1 dx_1 dx_2 f_\alpha(x_1, \mu_F) * f_\beta(x_2, \mu_F) * \hat{\sigma}_{\alpha\beta}(\alpha_s(\mu_R), \mu_F)$$

QCD in a nutshell

Factorization theorem



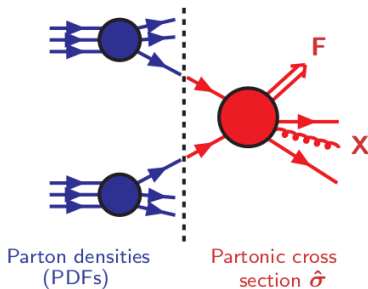
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Partonic cross section and pQCD

Perturbative QCD

- Factorization theorem helps us to understand short range interactions
- Write cross section as a perturbative series



$$\hat{\sigma} = \sigma^{\text{Born}} \left(1 + \frac{\alpha_s}{2\pi} \sigma^{(1)} + \left(\frac{\alpha_s}{2\pi} \right)^2 \sigma^{(2)} + \left(\frac{\alpha_s}{2\pi} \right)^3 \sigma^{(3)} + \dots \right)$$

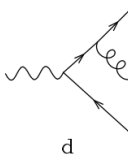
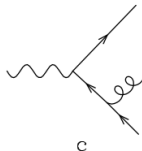
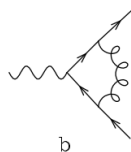
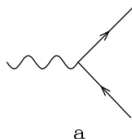
Leading order predictions can strongly depend on the renormalization and factorization scales → **Need higher order corrections**

Dealing with divergences

Perturbative QCD

Higher order corrections

- ① Higher order corrections are **not an easy task** due to **infrared (IR) singularities**
- ② Soft and collinear singularities
- ③ Impossible direct use of numerical techniques

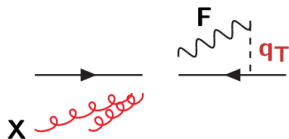


Resummation in QCD

Resumming large logs

Produce a final state from hadron collisions and study the q_\perp distribution

$$h_1(p_1) + h_2(p_2) \longrightarrow F(M, \mathbf{q}_\perp) + X$$



Truncated fixed order predictions \rightarrow **divergent $\ln^m(M^2/q_\perp^2)$ appear.**
Then the q_\perp distribution need to be evaluated by replacing the partonic cross section as follows

Resummation in QCD

Resumming large logs

Replace partonic q_\perp distribution as follows

$$\frac{d\hat{\sigma}_{ab}}{dq_\perp^2} \rightarrow \left[\frac{d\hat{\sigma}_{ab}^{(\text{res.})}}{dq_\perp^2} \right]_{\text{l.a.}} + \left[\frac{d\hat{\sigma}_{ab}^{(\text{fin.})}}{dq_\perp^2} \right]_{\text{f.o.}}$$

such that

$$\int dq_\perp^2 \frac{d\hat{\sigma}_{ab}^{(\text{res.})}}{dq_\perp^2} \sim \sum \alpha_s^n \log^m \frac{M^2}{Q_\perp^2} \quad \text{for } q_\perp \rightarrow 0$$
$$\int dq_\perp^2 \frac{d\hat{\sigma}_{ab}^{(\text{fin.})}}{dq_\perp^2} \sim 0 \quad \text{for } q_\perp \rightarrow 0$$

Resummation in QCD

q_\perp subtraction

Resummation holds in impact parameter space b

$$\frac{d\hat{\sigma}_{ab}^{\text{res.}}}{dq_\perp^2} = \frac{M^2}{\hat{s}} \int db \frac{b}{2} J_0(bq_\perp) \mathcal{W}_{ab}(b, M)$$

Which is expressed in Mellin space (with respect to $z = M^2/\hat{s}$)

$$\mathcal{W}_N(b, M) = \mathcal{H}_N(\alpha_s) \times \exp\{\mathcal{G}_N(\alpha_s, L)\} \quad \text{being} \quad L \equiv \log(M^2 b^2)$$

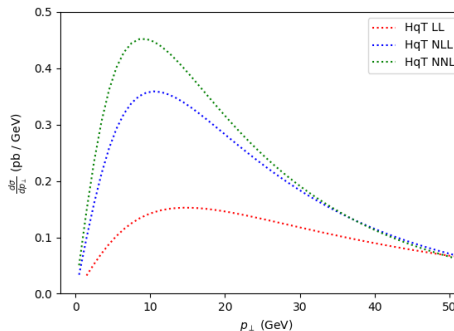
- Resummed effects exponentiated in a universal Sudakov factor
- Process-dependence factorized in the hard factor

HTurbo: Fast predictions for Higgs production

HqT and HRes

Predictions for Higgs q_{\perp} distribution

- HqT and HRes [de Florian, G.F., Grazzini, Tommasini] produce q_{\perp} distributions
- Higher order corrections require **high computation time**
- Codes producing fast predictions are needed for precision era of the LHC



HTurbo

Higgs distribution from Drell-Yan

Start from DYTurbo [Camarda et al.] ref. at [1910.07049](#) and [/dyturbo](#)

- 1 Matrix element $gg \rightarrow H$ as LO amplitude
- 2 Set Sudakov and Hard coefficients for Higgs production

$$\mathcal{G}(\alpha_s, L) = L g^{(1)}(\alpha_s L) + g^{(2)}(\alpha_s L) + \frac{\alpha_s}{\pi} g^{(3)}(\alpha_s L) + \dots$$

$$\mathcal{H}(\alpha_s) = 1 + \frac{\alpha_s}{\pi} \mathcal{H}^{(1)} + \left(\frac{\alpha_s}{\pi}\right)^2 \mathcal{H}^{(2)} + \dots$$

Build predictions up to NNLO

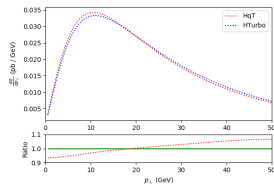
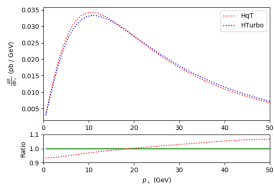
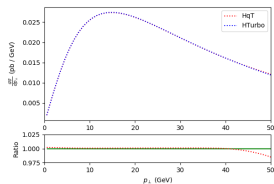
$$\text{LL}(\sim \alpha_s^n L^{n+1}) : g^{(1)}, \hat{\sigma}^{(0)}$$

$$\text{NLL}(\sim \alpha_s^n L^n) : g^{(2)}, \mathcal{H}^{(1)}$$

$$\text{NNLL}(\sim \alpha_s^n L^{n-1}) : g^{(3)}, \mathcal{H}^{(2)}$$

Results

Comparison HTurbo and HqT



- HTurbo produces qt distributions that match HRes and HqT
- Discrepancy at the 1/1000 level at NLL, at the 1/10 for NLL

Summary & Conclusions

- ① Fast predictions are required towards the precision era of the LHC
- ② HTurbo produces $q\bar{t}$ distributions that perfectly match HRes and HqT
- ③ Predictions by HTurbo are much faster than any of the existing codes
- ④ Next steps: Implement PDF evolution N3LO distributions

Thank you!



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