

空气动力学 HW 3

3.8

解:

proof:

uniform flow

$$\psi = V_{\infty} y$$

$$u = \frac{\partial \psi}{\partial y} = V_{\infty}$$

$$v = -\frac{\partial \psi}{\partial x} = 0$$

$$\nabla \cdot \vec{V} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \text{ incompressible}$$

$$\nabla \times \vec{V} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ u & v & w \end{vmatrix}$$

$$= \vec{i} \cdot \left(\frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \right) + \vec{j} \cdot \left(\frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right)$$

$$+ \vec{k} \cdot \left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$

$$= 0\vec{i} + 0\vec{j} + 0\vec{k}, \text{ irrotational}$$

3.10

解:

proof:

$$\phi = V_{\infty} x = V_{\infty} r \cos \theta$$

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2} = \frac{\partial^2 V_{\infty} x}{\partial x^2} + 0 + 0$$

$$= \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \phi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} + \frac{\partial^2 \phi}{\partial z^2}$$

$$= \frac{1}{r} \frac{\partial}{\partial r} (r V_{\infty} \cos \theta) + \frac{1}{r^2} \cdot V_{\infty} r \cdot (-\cos \theta) + 0$$

$$= 0, \text{ 满足 Laplace 方程}$$

$$\psi = V_{\infty} y = V_{\infty} r \sin \theta$$

$$\nabla^2 \psi = \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = 0 + \frac{\partial^2 V_{\infty} y}{\partial y^2} + 0$$

$$= \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \psi}{\partial \theta^2} + \frac{\partial^2 \psi}{\partial z^2}$$

$$= \frac{1}{r} \frac{\partial}{\partial r} (r V_{\infty} \sin \theta) + \frac{1}{r^2} V_{\infty} r \cdot (-\sin \theta) + 0$$

$$= 0, \text{ 满足 Laplace 方程}$$

3.14

$$\text{解: Doublet flow: } \psi = \frac{-k}{2\pi} \cdot \frac{\sin \theta}{r}$$

$$V_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = \frac{1}{r} \cdot \frac{-k}{2\pi} \cdot \frac{\cos \theta}{r} = -\frac{k \cos \theta}{2\pi r^2}$$

$$\Rightarrow \phi = \frac{k \sin \theta}{2\pi r}$$

3.15

解:

无升力圆柱绕流

$$\psi = V_{\infty} r \sin \theta \left(1 - \frac{R^2}{r^2} \right)$$

$$V_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = \frac{1}{r} V_{\infty} r \left(1 - \frac{R^2}{r^2} \right) \cdot \cos \theta$$

$$= V_{\infty} \left(1 - \frac{R^2}{r^2} \right) \cos \theta$$

$$V_{\theta} = -\frac{\partial \psi}{\partial r} = (-1) \cdot V_{\infty} \sin \theta \cdot \left(1 + \frac{R^2}{r^2} \right)$$

$$C_p = 1 - \frac{V^2}{V_{\infty}^2} = 1 - \frac{V_r^2 + V_{\theta}^2}{V_{\infty}^2}$$

$$= 1 - \frac{V_{\infty}^2 \cos^2 \theta \cdot \left(1 - \frac{R^2}{r^2} \right)^2 + V_{\infty}^2 \sin^2 \theta \left(1 + \frac{R^2}{r^2} \right)^2}{V_{\infty}^2}$$

$$= 1 - \left[1 + \frac{R^4}{r^4} + 2 \frac{R^2}{r^2} (\sin^2 \theta - \cos^2 \theta) \right]$$

$$= 2 \frac{R^2}{r^2} (\cos^2 \theta - \sin^2 \theta) - \frac{R^4}{r^4} \quad \text{ANS}$$

on the surface

$$C_p|_{R=r} = 2(\cos^2 \theta - \sin^2 \theta) - 1$$

$$= 2\cos^2 \theta - 2\sin^2 \theta - 1$$

$$= 1 - 4\sin^2 \theta$$

3.17

解:

有升力圆柱绕流

$$\psi = V_{\infty} r \sin \theta \left(1 - \frac{R^2}{r^2}\right) + \frac{\Gamma}{2\pi} \ln \frac{r}{R}$$

$$V_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = \frac{1}{r} \cdot V_{\infty} r \cos \theta \left(1 - \frac{R^2}{r^2}\right) \\ = V_{\infty} \cos \theta \cdot \left(1 - \frac{R^2}{r^2}\right)$$

$$V_{\theta} = -\frac{\partial \psi}{\partial r} = (-1) \cdot \left[V_{\infty} \sin \theta \left(1 + \frac{R^2}{r^2}\right) + \frac{\Gamma}{2\pi} \cdot \frac{1}{r} \right]$$

$$\therefore \frac{V_r}{V_{\infty}} = \cos \theta \left(1 - \frac{R^2}{r^2}\right), \text{ independent of } V_{\infty};$$

$$\frac{V_{\theta}}{V_{\infty}} = (-1) \left[\sin \theta \cdot \left(1 + \frac{R^2}{r^2}\right) + \frac{\Gamma}{2\pi r} \cdot \frac{1}{V_{\infty}} \right],$$

it's dependent of V_{∞} .

So the streamlines shape will change.

3.18

解:

$$L' = 6 \text{ N/m}$$

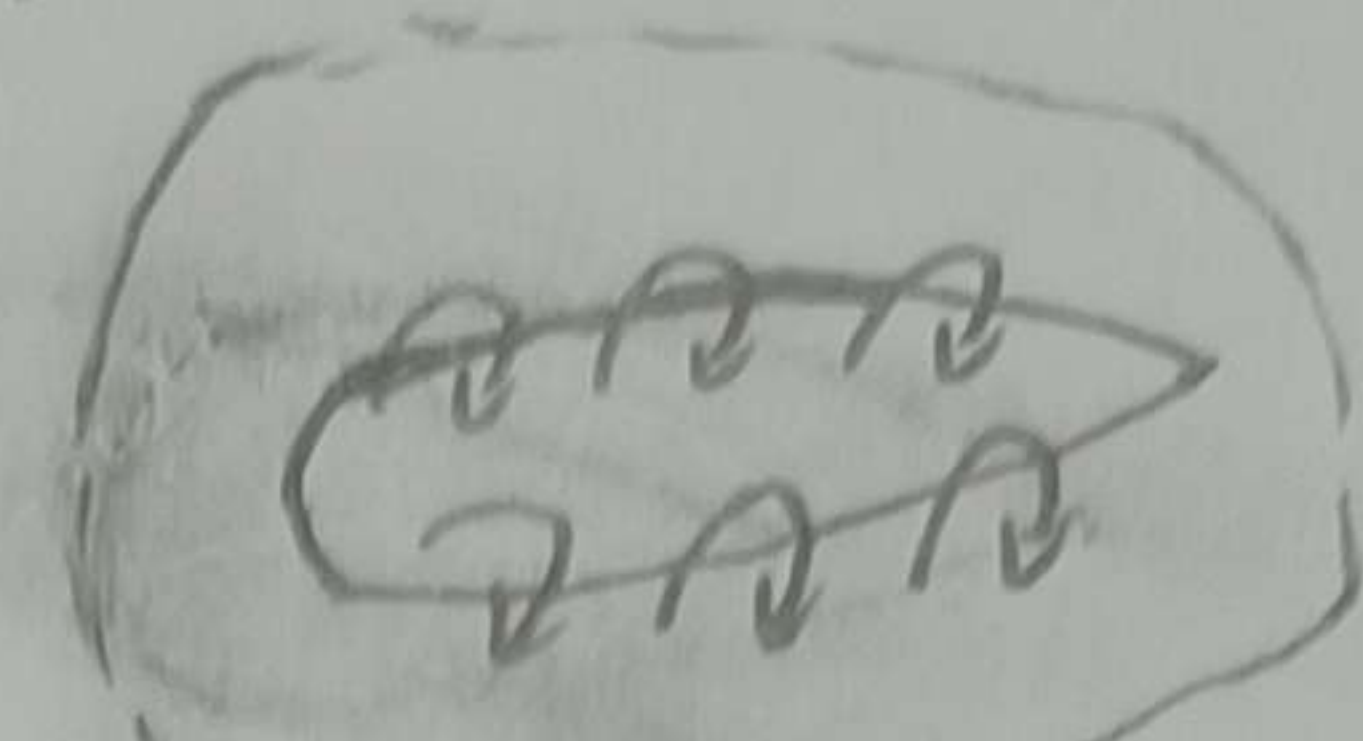
$$V_{\infty} = 30 \text{ m/s}$$

$$L' = \rho_{\infty} V_{\infty} \Gamma$$

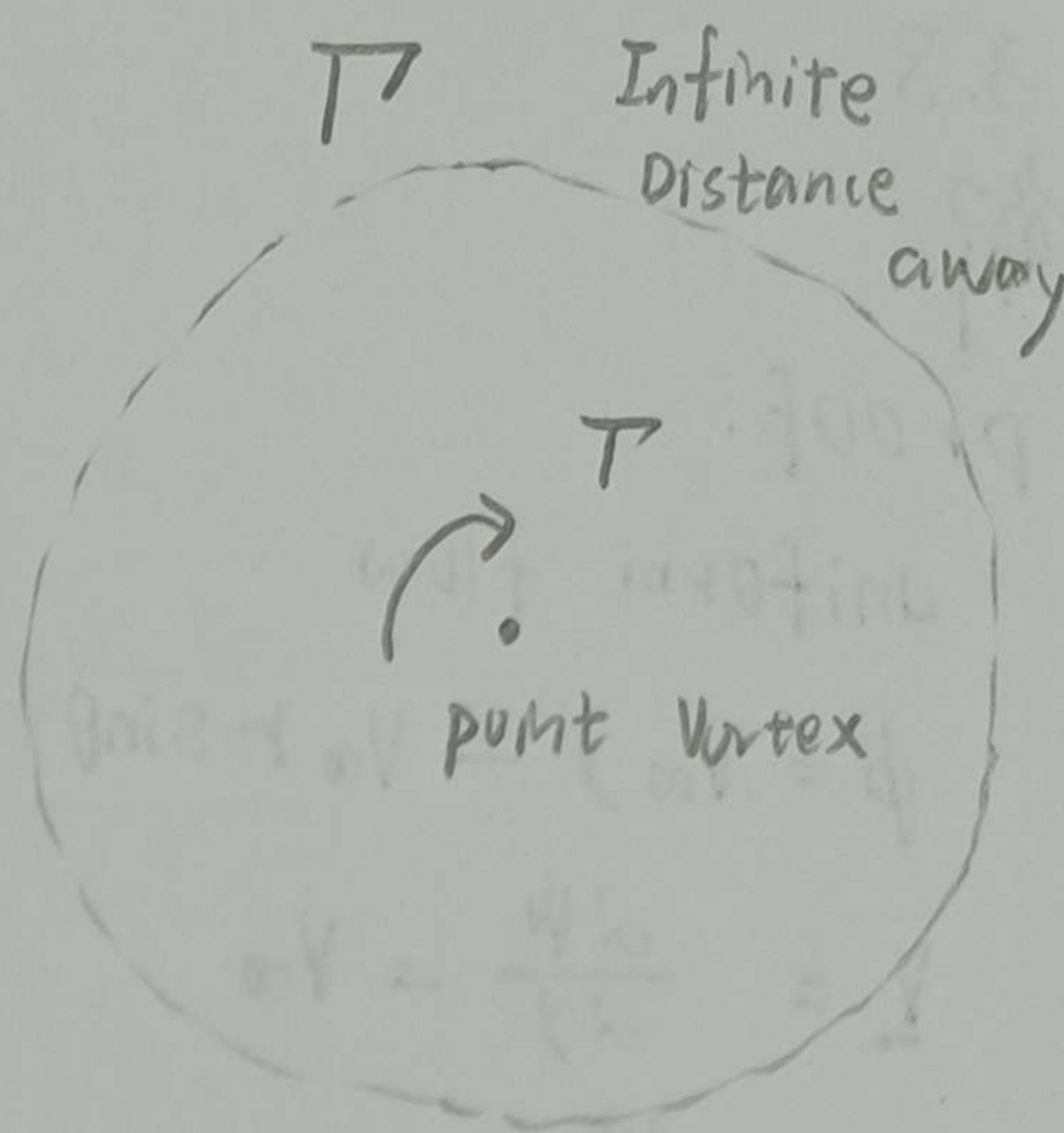
$$\therefore \Gamma = \frac{L'}{\rho_{\infty} V_{\infty}} = \frac{6}{1.23 \times 30} = 0.1626 \text{ m}^2/\text{s}$$

3.20

解:

 Γ 

Distributed Singularities

 \Rightarrow 

point vortex

airfoil 上的 flow 可合成为点源或点涡的分布,

涡强度的叠加是为总环量 Γ , Γ 的值在所有包含 airfoil 的闭合曲线上相同,airfoil 成为纸面上一个 speck, 分布的点涡形成一个强度为 Γ 的较大涡,

这与圆柱体的单点涡等价,

airfoil 上的升力与圆柱体相同, 即有 $L' = \rho_{\infty} V_{\infty} \Gamma$

3.21

解:

Since the volume flow rates between the streamlines are the same.

$$\Rightarrow \psi_3 - \psi_2 = \psi_2 - \psi_1$$

$$\psi_3 = 2\psi_2 - 0$$

$$\psi = V_\infty r \sin\theta - \frac{K}{2\pi} \cdot \frac{\sin\theta}{r} = V_\infty r \sin\theta \left(1 - \frac{R^2}{r^2}\right)$$

for ψ_2 passes through $(\frac{6}{5}R, \frac{\pi}{2})$

$$\psi_2 = V_\infty \cdot \frac{6}{5}R \left(1 - \frac{25}{36}\right) = V_\infty R \cdot \frac{11}{30}$$

$$\therefore \psi_3 = \frac{11}{15} V_\infty R = V_\infty r \sin\frac{\pi}{2} \left(1 - \frac{R^2}{r^2}\right)$$

$$r^2 - \frac{11}{15}R \cdot r - R^2 = 0$$

$$r = \frac{\frac{11}{15}R \pm \sqrt{\left(\frac{11}{15}R\right)^2 + 4R^2}}{2} = \frac{\frac{11}{15}R \pm \sqrt{\frac{1021}{225}}R}{2}$$

$$\therefore r = 1.432R \quad (\text{舍去负值})$$