

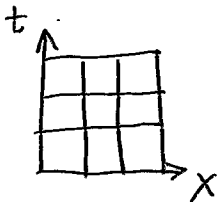
4.7 B-S.

$$\begin{cases} V_t = \mathcal{L}V \\ V(0, \tau) = B^L(\tau) \\ V(\infty, \tau) = B^R(\tau) \\ V(x, 0) = P(x) \end{cases}$$

$$\mathcal{L}V = \frac{\sigma^2 S^2}{R} V_{xx} - rV + rX V_x$$

$$\text{Put.} \begin{cases} B^L(\tau) = e^{-r\tau} \\ B^R(\tau) = 0 \\ P(x) = \max(1-x, 0) \end{cases}$$

Approximate: truncate



$$\begin{cases} [-\infty, \infty) \times [0, T] \rightarrow [0, X_{\max}] \times [0, T] \\ X_i = i \cdot \Delta X \\ \tau_n = n \cdot \Delta \tau \\ \Delta X = X_{\max} / N_X \\ \Delta \tau = T / N_\tau \end{cases}$$

$$\begin{cases} V_i(\tau) \doteq V(X_i, \tau) \\ V_i^n \doteq V(X_i, \tau_n) \end{cases}$$

$$V_{cl,i} = \mathcal{L}_h V_i$$

$$\mathcal{L}_h V_i = \frac{\sigma_i^2 X_i^2}{R} \delta_x^+ \delta_x^- V_i - V_i r_i + V_i X_i \delta_x^0 V_i + \text{B.C.'s} \quad i=1, 2, \dots, N_X-1$$

$$\begin{aligned} \text{or } \mathcal{L}_n V_i &= \frac{\sigma_i^2 X_i^2}{R} \cdot \frac{V_{i+1} - 2V_i + V_{i-1}}{\Delta X^2} - r_i V_i + r_i V_i \cdot \frac{V_{i+1} - V_{i-1}}{2\Delta X} \\ &= \left(\frac{\sigma_i^2 X_i^2}{R\Delta X^2} + \frac{r_i X_i}{2\Delta X} \right) V_{i+1} + \left(V_i - \frac{2\sigma_i^2 X_i^2}{R\Delta X^2} \right) V_i + \left(\frac{\sigma_i^2 X_i^2}{R\Delta X^2} - \frac{r_i X_i}{2\Delta X} \right) V_{i-1} \end{aligned}$$

$$\left\{ \begin{array}{l} \mathcal{L}_n V_i = a_i V_{i+1} + b_i V_i + c_i V_{i-1} \quad i=1, 2, \dots, N_x-1 \\ V_0 = B^L(\tau) \\ V_{N_x} = B^R(\tau) \end{array} \right\} \rightarrow \text{ODEs}$$

$$\text{let } \vec{V} = \begin{bmatrix} V_1 \\ V_2 \\ \vdots \\ V_{N_x-1} \end{bmatrix}$$

$$\text{Then, } \frac{d\vec{V}}{d\tau} = \begin{bmatrix} b_1 & c_1 & & & \\ a_2 & b_2 & c_2 & & \\ & a_3 & b_3 & c_3 & \\ & & & & c_{N_x-2} \\ & & & a_{N_x-1} & b_{N_x} \end{bmatrix} \vec{V} + \begin{bmatrix} a_1 B^L(\tau) \\ 0 \\ 0 \\ \vdots \\ c_{N_x-1} B^R(\tau) \end{bmatrix}$$

$$= A\vec{V} + \vec{g}(\tau)$$

$$A = \begin{bmatrix} b_1 & c_1 & & & \\ a_2 & b_2 & c_2 & & \\ & a_3 & b_3 & c_3 & \\ & & & & c_{N_x-2} \\ & & & a_{N_x-1} & b_{N_x-1} \end{bmatrix} \quad \vec{g}(\tau) = \begin{bmatrix} a_1 B^L(\tau) \\ 0 \\ 0 \\ \vdots \\ c_{N_x-1} B^R(\tau) \end{bmatrix}$$

$$\frac{d\vec{v}}{d\tau} = A\vec{v} + \vec{g} = \vec{F} - (\vec{v}, \tau)$$

$$\vec{v}^{n+1} = \vec{v}^n + \frac{\Delta\tau}{2} \{ (A\vec{v} + \vec{J})^{n+1} + (A\vec{v} + \vec{g})^n \}$$

$$(I - \frac{\Delta\tau}{2}A) \vec{v}^{n+1} = \vec{v}^n + \frac{\Delta\tau}{2}A^n \vec{v}^n + \frac{\Delta\tau}{2}(\vec{g}^{n+1} + \vec{g}^n)$$

$$M_1 \vec{v}^{n+1} = \vec{RHS} \Rightarrow \text{right-hand-side vector}$$

$$M_1 = \text{diag} \left(-\frac{\Delta\tau}{2}a, -\frac{\Delta\tau}{2}b, -\frac{\Delta\tau}{2}c \right)$$

↓
diagonal

Algorithm:

1. set \vec{v}^n to payoff ($v_j^0 = P(x_j)$)
2. for $n=0$ to $N_\tau-1$
3. Compute the $\vec{RHS} = (I + \frac{\Delta\tau}{2}A) \vec{v} + \frac{\Delta\tau}{2}(\vec{g}^{n+1} + \vec{g}^n)$
4. solve $(I - \frac{\Delta\tau}{2}A) \vec{v}^{n+1} = \vec{RHS}$
5. next n \hookrightarrow tri-diag \rightarrow store only the 3 vectors.
6. output results. Don't store v^n as $[*][*]$, only store $v[*]$

In step 3.

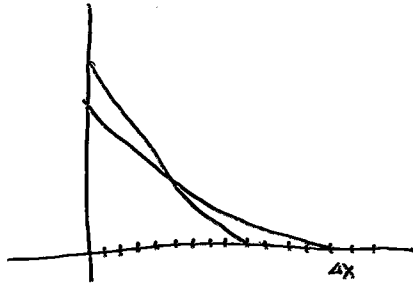
$$RHS_i = v_i + \frac{\Delta\tau}{2} (a_i v_{i-1} + b_i v_i + c_i v_{i+1}) \quad i = 2, \dots, N_x - 1 (N_x - 2)$$

$$RHS_1 = v_1 + \frac{\Delta\tau}{2} (b_1 v_1 + c_1 v_2)$$

⋮

$$RHS_{N_x-1} = v_{N_x-1} + \frac{\Delta\tau}{2} (a_{N_x-1} v_{N_x-2} + b_{N_x-1} v_{N_x-1}) \quad \cancel{c_{N_x-1} v_{N_x}}$$

results $V_j^{N_x-1}$



Example:
$$\begin{cases} u_t = u_{xx} & x \in [0, L] \\ u(0, t) = 0 \\ u(L, t) = 0 \\ u(x, 0) = u_0(x) \end{cases}$$

$$x_i = i \cdot \Delta x \quad \Delta x = \frac{L}{N_x}$$

$$\tau_n = n \cdot \Delta \tau$$

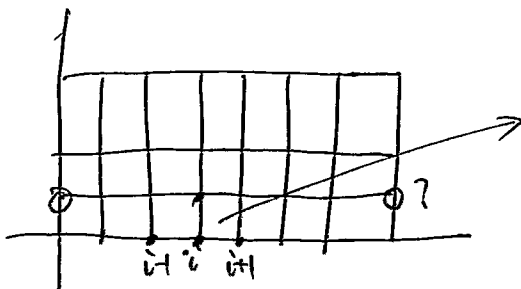
$$u(x_i, t) \approx \bar{u}_i(\tau)$$

$$\begin{cases} u_t = \frac{du}{dt} = \delta^+ \delta^- \bar{u}_i = \frac{u_{i+1} - 2u_i + u_{i-1}}{\Delta x^2} & i=1, 2, \dots, N_x-1 \\ \bar{u}_0(\tau) = 0 \\ \bar{u}_{N_x}(\tau) = 0 \end{cases}$$

+ forward euler

$$u_i^{n+1} = u_i^n + \Delta \tau (u_{i+1}^n - 2u_i^n + u_{i-1}^n) / \Delta x^2$$

$$= \frac{\Delta \tau}{\Delta x^2} u_{i+1}^n + (1 - \frac{2\Delta \tau}{\Delta x^2}) u_i^n + \frac{\Delta \tau}{\Delta x^2} u_{i-1}^n$$



steric
molecule

boundary condition.

4.9

$$u_t = u_{xx}, \quad x \in (-\infty, \infty)$$

$$u_j^{n+1} = u_j^n + \frac{\Delta t}{\Delta x^2} (u_{j+1}^n - 2u_j^n + u_{j-1}^n)$$

$$\lambda = \frac{\Delta t}{\Delta x^2}$$

$$u_j^{n+1} = \lambda u_{j+1}^n + (1-2\lambda) u_j^n + \lambda u_{j-1}^n$$

Example:

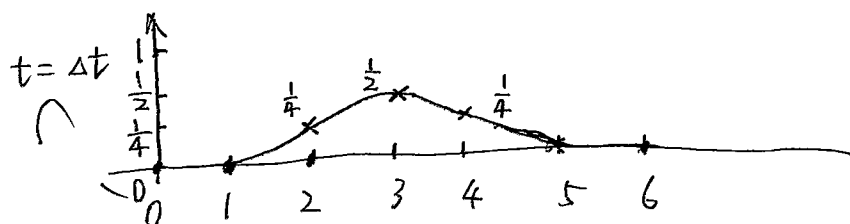
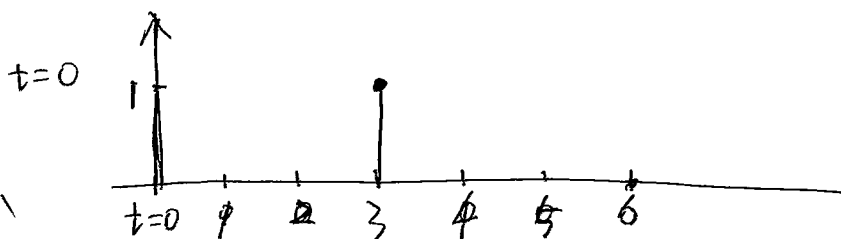
$$\text{IVP. } \begin{cases} u_t = u_{xx} & x \in (0, 6) \\ u(x, 0) = \begin{cases} 1 & x=3 \\ 0 & x \neq 3 \end{cases} \\ u(0, t) = 0 \\ u(6, t) = 0 \end{cases}$$

$$\Delta x = 1, \quad N_x = 6$$

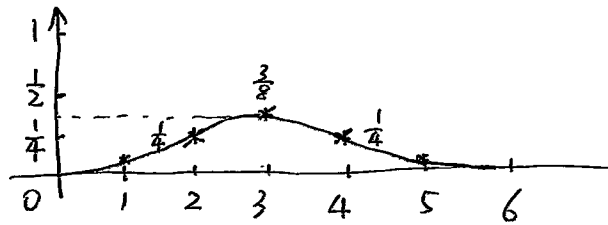
$$\Delta t = \frac{1}{4},$$

then

$$\begin{aligned} u_j^{n+1} &= \frac{1}{4} u_{j+1}^n + \frac{1}{2} u_j^n + \frac{1}{4} u_{j-1}^n \\ &= \frac{1}{4} (u_{j+1}^n + 2u_j^n + u_{j-1}^n) \end{aligned}$$

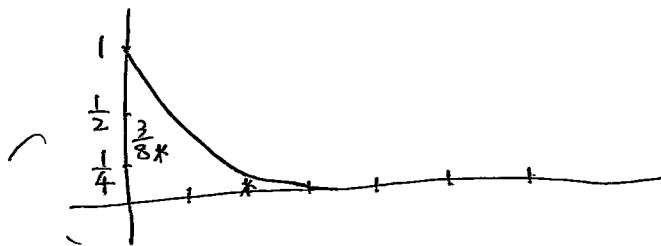
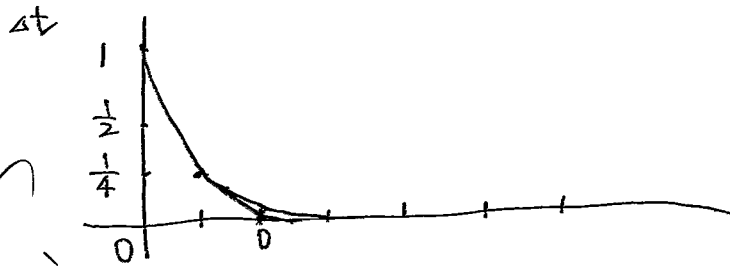
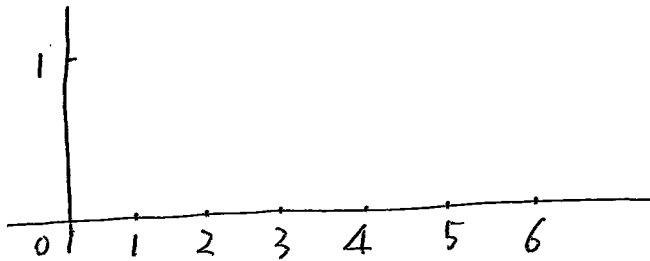


$$t = 2\Delta t$$



Example 2.

$$\begin{cases} u(x, 0) = 0 \\ u(0, t) = 1 \\ u(6, t) = 0 \end{cases}$$



Looking at $u_j^{n+1} = \lambda u_{j+1}^n + (1-2\lambda) u_j^n + \lambda u_{j-1}^n$

in matrix form.

$$\begin{aligned}
 u(0, t) &= \beta^L \\
 u(6, t) &= \beta^R \\
 \begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \\ u_5 \end{bmatrix}^{n+1} &= \begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \\ u_5 \end{bmatrix}^n \begin{bmatrix} \frac{1}{2} & \frac{1}{4} & & & \\ \frac{1}{4} & \frac{1}{2} & \frac{1}{4} & & \\ & & & & \\ & & & & \\ & & & & \end{bmatrix} + \begin{bmatrix} \frac{1}{4} \beta^L \\ 0 \\ 0 \\ \vdots \\ \frac{1}{4} \beta^R \end{bmatrix}
 \end{aligned}$$

~~the~~

$$\begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \\ u_5 \end{bmatrix}^{n+1} = \begin{bmatrix} \frac{1}{2} & \frac{1}{4} & & & \\ \frac{1}{4} & \frac{1}{2} & \frac{1}{4} & & \\ & \frac{1}{4} & \frac{1}{2} & \frac{1}{4} & \\ & & \frac{1}{4} & \frac{1}{2} & \frac{1}{4} \\ & & & \frac{1}{4} & \frac{1}{2} \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \\ u_5 \end{bmatrix}^n + \begin{bmatrix} \frac{1}{4} \beta^L \\ 0 \\ 0 \\ 0 \\ \frac{1}{4} \beta^R \end{bmatrix}$$

3.5 Some Basic Theory

1. has $\Delta x, \Delta t \rightarrow 0$

does $u_j^n \rightarrow u(x_j, t_n)$?

2. How fast does it converge?

3. Are there practical limits on $\Delta t \nmid \Delta x$?

Def. The local truncation error of a finite difference approx. is the amount by which a smooth enough soln. fails to satisfy the approximation.

EX: $u_t = u_{xx}$

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} = \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{\Delta x^2}$$

Substitute in u .

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} = \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{\Delta x^2} + \tau$$

$$\begin{aligned} \delta_t^+ u_j^n &= \frac{(u_j^n + \Delta t \cdot \partial_t u_j^n + \frac{\Delta t^2}{2} \partial_t^2 u_j^n + \dots) - u_j^n}{\Delta t} \\ &= \partial_t u_j^n + \frac{\Delta t}{2} \partial_t^2 u_j^n \end{aligned}$$

$$\begin{aligned} \Delta x^2 \delta_x^+ \delta_x^- u_j^n &= u_j^n + \Delta x \cdot \partial_x u_j^n + \frac{\Delta x^2}{2} \partial_x^2 u_j^n + \frac{\Delta x^3}{3!} \partial_x^3 u_j^n + \frac{\Delta x^4}{4!} \partial_x^4 u_j^n(\xi, t_n) \\ &\quad - 2u_j^n \\ &\quad + u_j^n - \Delta x \partial_x u_j^n + \frac{\Delta x^2}{2} \partial_x^2 u_j^n - \frac{\Delta x^3}{3!} \partial_x^3 u_j^n + \frac{\Delta x^4}{4!} \partial_x^4 u(\xi, t_n) \end{aligned}$$

$$= \Delta x^2 \partial_t^2 u_j^n + \frac{2}{4!} \Delta x^4 \partial_t^4 u(\xi, t_n)$$

so

$$\frac{\partial}{\partial t} u_j^n + \frac{\Delta t}{2} \partial_t^2 u(x_j, u) = \partial_x^2 u_j^n + \frac{2}{4!} \Delta x^2 \partial_x^4 u(\xi, t_n) + \tau$$

$$\Rightarrow \tau_j^n = \underbrace{\left(\frac{\partial}{\partial t} - \frac{\partial^2}{\partial x^2} \right) u_j^n}_0 + \frac{\Delta t}{2} \cdot \frac{\partial^2 u(x_j, u)}{\partial t^2} - \frac{2}{4!} \Delta x^2 \partial_x^4 u(\xi, t_n)$$

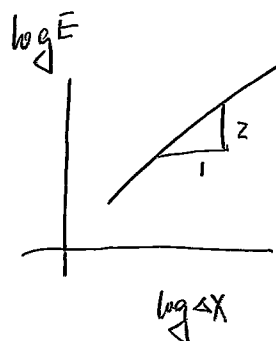
$$\tau_j^n = O(\Delta t, \Delta x^2)$$

Def. A FDE is consistent if $\tau = O(\Delta t^p, \Delta x^q)$

with $p, q \geq 1$

Thm: A necessary condition for convergence is consistency.

4.12



$$K^* X = S^*$$

$$\Delta t = O(\Delta t)$$

change strike price

$$V = \frac{V^*}{K^*}$$

we need stability.

we can write the FDE as

$$\vec{u}^{n+1} = Q \vec{u}^n + \vec{g}$$

for example.

$$\frac{du_j}{dt} = a_j u_{j-1} + b_j u_j + c_j u_{j+1}$$

for the heat equ.

$$\frac{du_j}{dt} = \frac{1}{\Delta x^2} (u_{j-1} - 2u_j + u_{j+1})$$

$$\text{or } \frac{d\vec{u}}{dt} = \frac{1}{\Delta x^2} A \vec{u} + \vec{g}, \quad A = \text{diag}(1, -2, 1)$$

if we use forward euler

$$\vec{u}^{n+1} = \vec{u}^n + \frac{\Delta t}{\Delta x^2} A \vec{u}^n + \Delta t \vec{g}^n$$

$$= (I + \frac{\Delta t}{\Delta x^2} A) \vec{u}^n + (\Delta t \vec{g}^n)$$

Q

if we use trapezoidal rule.

$$u^{n+1} = u^n + \frac{\Delta t}{2\Delta x^2} (A(u^{n+1} + u^n) + \frac{\Delta t}{2}(\vec{g}^{n+1} + \vec{g}^n))$$

$$(I - A \frac{\Delta t}{\Delta x^2}) u^{n+1} = (I + A \Delta t / \Delta x^2) \vec{u}^n + \frac{\Delta t}{2} (\vec{g}^{n+1} + \vec{g}^n)$$

$$\vec{u}^{n+1} = \underbrace{(I - (\frac{\Delta t}{2\Delta x^2})A)^{-1} (I + \frac{\Delta t}{2\Delta x^2}A)}_Q \vec{u}^n + \vec{\bar{g}}$$

suppose $\vec{\hat{g}} = 0$

$$\vec{u}^{n+1} = Q \vec{u}^n = Q^2 \vec{u}^{n-1} = \dots = Q^{n+1} \vec{u}^0$$

$$\vec{u}^n = Q^n \vec{u}^0 \quad n = \frac{T}{\Delta t}$$

In norm

$$\|\vec{u}^n\| = \|Q^n \vec{u}^0\| \leq \|Q^n\| * \|\vec{u}^0\|$$

Def: The FDE is stable if $\exists C_T$ independent of $\Delta t \neq \Delta x$.

But may be depend on time T. \Rightarrow

$$\|Q^n\| \leq C_T \quad \forall n.$$

THM: A necessary condition for convergence is stability.

The PBE itself satisfies a simulator property. Assume

Dirichlet B.C's $\beta^L = \beta^R = 0$ for heat eqn.

$$\begin{cases} u_t = u_{xx} \\ u(R,t) = u(L,t) = 0 \\ u(x,t) = u(x) \end{cases} \quad x \in [0, L]$$

$$\int_0^L u_t u = \int_0^L u \cdot u_{xx}$$

$$\int_0^L u_t u \, dx = \int_0^L u u_{xx} \, dx$$

$$\frac{d}{dt} \int_0^L \frac{u^2}{2} \, dx = \int_0^L u u_{xx} \, dx$$

call $\|u\|_2 \equiv \sqrt{\int_0^L u^2 \, dx}$

so, $\frac{1}{2} \cdot \frac{d}{dt} \|u\|_2^2 = \underbrace{\int_0^L u \cdot u_{xx} \, dx}_{\substack{\text{Boundary} \\ \text{Conditions}}} = \underbrace{u \cdot u_x \Big|_0^L}_{\text{this} = 0} - \underbrace{\int_0^L (u_x)^2 \, dx}_{\|u_x\|^2}$

$$\frac{d}{dt} \|u\|_2^2 = -2 \|u_x\|^2 \leq 0$$

$$\Rightarrow \|u\|_2^2 \leq \|u_0\|_2^2$$

$$\Rightarrow \|u\|_2 \leq \|u_0\|_2$$

Def: An initial value problem is well-posed if the solution satisfies.

$$\|u\| \leq C_T \cdot \|u_0\|$$

Note: Suppose $\begin{cases} v_t = v_{xx} \\ v(x, 0) = u_0(x) + \delta \end{cases}$

then $v_t - u_t = (v - u)_{xx} = (v - u)_t$

call $w = v - u$

$$\begin{cases} w_t = w_{xx} \\ w(x, 0) = \delta \end{cases}$$

then

$$\|w\| \leq \|\delta\|$$

$$\|v - u\| \leq \|\delta\|$$

Def: The approximation u_j^n converges to $u(x_j, t_n)$ if $t_n < T$

$$\|u_j^n - u(x_j, t_n)\| \rightarrow 0$$

$$\text{as } \Delta t, \Delta x \rightarrow 0$$

THM: (LAX-Richtmyer equivalence)

For a consistent approximation to a well-posed initial value problem, stability is necessary and sufficient for convergence.

stability + consistency \Leftrightarrow convergence.

$$\begin{array}{ccc} \uparrow & & \uparrow \\ \|Q^n\| \leq C_T & & \tau = O(\Delta t^p, \Delta x^q) \\ \downarrow & & p, q \geq 1 \\ \text{const.} & & \end{array}$$

For heat equ.

$$\|u\| \leq \|u_0\|$$

$$\Rightarrow \|u\| \leq \underbrace{\|Q^n\|}_{\leq 1} \|u_0\|$$

we'll require
 $\|Q^n\| \leq 1$

$$\|Q^n\| \leq \|Q\|^n$$

and therefore require $\|Q\| \leq 1$

The $\|A\|_2 = \sqrt{\rho(A^T, A)}$ matrix

$$\rho = |\lambda|_{\max}$$

if $A = A^T$,

then $\|A\|_2 = \rho(A)$

EX: heat equation. forward euler in time.

$$\vec{u}^{n+1} = \underbrace{(I + \frac{\Delta t}{\Delta x^2} A)}_Q \vec{u}^n$$

$$\begin{cases} Ax = \lambda x \\ Ix = x \end{cases}$$

$$A = \text{diagonal}(1, -2, 1)$$

$$\lambda_Q = I + \frac{\Delta t}{\Delta x^2} \lambda_A$$

$$\begin{bmatrix} -2 & 1 & \\ & -2 & 1 \\ & & -2 \end{bmatrix}$$