## Notes on Bosonic String Amplitudes

## Jian-ming Zheng

Department of Physics, Tsinghua University, Beijing 100084, China

This is a short note on the basics of bosonic strings and the calculation of their amplitudes.

## 1 The world-sheet action and string spectrum

String theory describes an 1-dimensional extended object. The 1-dimensional string sweeps a 2-dimensional surface  $\Sigma$ . The string itself is parametrized by the coordinate  $\sigma$ . Along with the proper time  $\tau$  of the trajectory, they parametrize the world-sheet  $\Sigma$ . The world-sheet  $\Sigma$  is embedded in the D-dimensional spacetime, with induced metric  $h_{ab} = G_{\mu\nu}\partial_a X^{\mu}\partial_b X^{\nu}$ . A simple and natural guess of the action is just the area of the world-sheet with tension  $T = \frac{1}{2\pi\alpha'}$ :

$$S_{NG} = -\frac{1}{2\pi\alpha'} \int_{\Sigma} d\tau d\sigma \sqrt{-\det h_{ab}}.$$
 (1)

This action is called Nambu-Goto action and has the following symmetry:

- 1. Spacetime Poincare symmetry:  $X^{\prime\mu}(\tau,\sigma) = \Lambda^{\mu}_{\nu}X^{\nu}(\tau,\sigma) + a^{\mu}$ ;
- 2. Diffeomorphism:  $\tau' = \tau'(\tau, \sigma), \sigma' = \sigma'(\tau, \sigma)$ .

A more convenient action for quantization is the Polyakov action, which involves a dynamical metric  $\gamma_{ab}$ :

$$S_P = -\frac{1}{4\pi\alpha'} \int_{\Sigma} d\tau d\sigma \sqrt{-\gamma} \gamma^{ab} \partial_a X^{\mu} \partial_b X_{\mu}. \tag{2}$$

Here we have taken the spacetime metric to be flat. This action classically agrees with the Nambu-Goto action, obtained by varying the metric:

$$\delta_{\gamma} S_{P} = -\frac{1}{4\pi\alpha'} \int_{\Sigma} d\tau d\sigma \sqrt{-\gamma} \delta \gamma^{ab} (h_{ab} - \frac{1}{2} h_{cd} \gamma_{ab} \gamma^{cd}). \tag{3}$$

The solution implies  $h_{ab}\sqrt{-h} = \gamma_{ab}\sqrt{-\gamma}$ . This can be used to eliminate  $\gamma$  in the action, and gives the NG action.

The Polyakov action has the following symmetries:

1. Spacetime Poincare symmetry;

2. Diffeomorphism:

$$X^{\prime\mu}(\tau^{\prime},\sigma^{\prime}) = X^{\mu}(\tau,\sigma); \quad \frac{\partial \sigma^{\prime c}}{\partial \sigma^{a}} \frac{\partial \sigma^{\prime d}}{\partial \sigma^{b}} \gamma^{\prime}_{cd}(\tau^{\prime},\sigma^{\prime}) = \gamma_{ab}(\tau,\sigma). \tag{4}$$

;

3. Weyl transformation:

$$X^{\prime\mu}(\tau,\sigma) = X^{\mu}(\tau,\sigma); \quad \gamma_{ab}^{\prime}(\tau,\sigma) = e^{2\omega(\tau,\sigma)}\gamma_{ab}(\tau,\sigma). \tag{5}$$

We view the latter two types of invariance as world-sheet gauge symmetries.

## 1.1 String spectrum