

Notes on the project ParametricAcousticArray

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Contents

1	List of Symbols	2
2	Theoretical Model and Analytical Solutions	2
2.1	KZK Equation	2
2.1.1	Rectangular profile	2
2.2	Spectral Solutions	3
2.3	Coupled PDEs	3
2.3.1	Rectangular Profile	3
2.4	Boundary Conditions	3
2.4.1	Circular Profile	3
2.4.2	Rectangular Profile	4
2.5	Sound Power of Waves	4
3	Numerical Methods	4
3.1	Discretization	4
3.1.1	Rectangular Profile	4
3.2	Matrix Form	6
3.2.1	Rectangular Profile	6
4	Test Data	7
4.1	Circular Profile	7
4.1.1	Kamakura 1989	7

1 List of Symbols

Symbols	Descriptions
a	radius of the source
c_0	isentropic speed of sound at ambient values of pressure and density
f_1, f_2	$f_1 = N_1 f_b, f_2 = N_2 f_b$, primary frequencies
f_b	$= \gcd f_2 - f_1$, basic frequency
f_m	$= (f_1 + f_2)/2$, mean primary frequency
k	$= 2\pi f/c_0$ wavenumber
k_m	$= (k_1 + k_2)/2$ wavenumber of the mean primary frequency
l_D	$= (\beta k \epsilon)^{-1}$, shock formation distance of a plane wave
L_x, L_y	the size of the rectangular transducer in x, y direction, and $L_y \geq L_x$
N_m	$= (N_1 + N_2)/2$
p	pressure
p_0	ambient pressure
\bar{p}	$= (p - p_0)/P_0$, acoustic pressure normalized to P_0
P_0	$\rho_0 c_0 v_0$, acoustic pressure peak amplitude on the source
R_D	$= ka^2/2 = \pi a^2/\lambda$, Rayleigh distance of primary frequency
R_{Dm}	$= k_m a^2/2 = \pi a^2/\lambda_m$, Rayleigh distance of mean primary frequency
t	time
\mathbf{u}	$= (u_x, u_y) = \mathbf{r}/L_y = (x, y)/L_y$
v_0	characteristic velocity peak amplitude on the source
x, y, z	dimensional coordinates, z along the direction of propagation
\mathbf{r}	(x, y)
r	$ \mathbf{r} $
α	absorption coefficient, in Neper per meter
β	parameter of nonlinearity
ϵ	$= u_0/c_0$, Mach number
σ	$= z/R_D$, nondimensional z
σ_D	$= l_D/R_D$
σ_{Dm}	$= l_{Dm}/R_{Dm}$
χ	$= x/R_{Dm}$
ψ	$= y/R_{Dm}$
γ	$= L_y/L_x$, the aspect ratio
$\nabla_{x,y}^2$	$= \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$
∇_u^2	$\frac{\partial^2}{\partial \mathbf{u}^2} = \frac{1}{u} \frac{\partial}{\partial u} \left(u \frac{\partial}{\partial u} \right)$
$\nabla_{\chi,\psi}^2$	$= \frac{\partial^2}{\partial \chi^2} + \frac{\partial^2}{\partial \psi^2}$

2 Theoretical Model and Analytical Solutions

2.1 KZK Equation

Nonlinear propagation of finite amplitude sound waves in dissipative and homogenous fluid is described by the following KZK equation:

$$\frac{\partial^2 \bar{p}}{\partial \sigma \partial \tau} = \alpha R_D \frac{\partial^3 \bar{p}}{\partial \tau^3} + \frac{1}{4} \nabla_{\perp}^2 \bar{p} + \frac{1}{2\sigma_D} \frac{\partial^2 (\bar{p})^2}{\partial \tau^2} \quad (1)$$

If the source is bifrequency, i.e.

$$f_i = N_i f_b, \quad N_i \in \mathbb{N}^+, \quad i = 1, 2 \quad (2)$$

2.1.1 Rectangular profile

The Eq. (1) is written as

$$\frac{\partial^2 \bar{p}}{\partial \sigma \partial \tau} = \alpha R_D \frac{\partial^3 \bar{p}}{\partial \tau^3} + \frac{1}{\pi \gamma} \nabla_u^2 \bar{p} + \frac{1}{2\sigma_D} \frac{\partial^2 (\bar{p})^2}{\partial \tau^2} \quad (3)$$

The transformations are:

$$\sigma_m = z/R_{Dm}, \quad T = (1 + \sigma_m) \bar{p}, \quad \tau_m = \tau - \frac{\pi \gamma}{4} \frac{\xi^2}{N_m(1 + \sigma_m)}, \quad (\chi, \psi) = \frac{(u_x, u_y)}{1 + \sigma_m} \quad (4)$$

Then, the transformed KZK equation is obtained by Eq. (3)

$$\frac{\partial^2 T}{\partial \sigma_m \partial \tau_m} = \frac{\alpha_m}{N_m^2} R_{Dm} \frac{\partial^3 T}{\partial \tau_m^3} + \frac{N_m}{\pi \gamma (1 + \sigma_m)^2} \nabla_{\chi, \psi}^2 T + \frac{1}{2N_m \sigma_{Dm} (1 + \sigma_m)} \frac{\partial^2 (T^2)}{\partial \tau_m^2} \quad (5)$$

2.2 Spectral Solutions

A solution of Eq. (3) is sought in the form of a Fourier series

$$\bar{p} = \sum_{n=1}^{\infty} (g_n \sin n\tau + h_n \cos n\tau) = \sum_{n=1}^{\infty} \bar{p}_n \quad (6)$$

where g_n and h_n are functions of spatial coordinates.

The solution of Eq. (5) is

$$T = \sum_{n=1}^{\infty} [g_n(\chi, \psi, \sigma_m) \sin n\tau_m + h_n(\chi, \psi, \sigma_m) \cos n\tau_m] = \sum_{n=1}^{\infty} T_n \quad (7)$$

2.3 Coupled PDEs

We thus obtain the following set of coupled partial differential equations for g_n and h_n .

2.3.1 Rectangular Profile

$$\begin{aligned} \frac{\partial g_n}{\partial \sigma} = & -\alpha_n R_{Dm} g_n + \frac{1}{\pi \gamma n (1 + \sigma_m)^2} \nabla_{\chi, \psi}^2 h_n + \frac{n}{2N_m \sigma_{Dm} (1 + \sigma_m)} \\ & \times \left[\frac{1}{2} \sum_{m=1}^{n-1} (g_m g_{n-m} - h_m h_{n-m}) - \sum_{m=n+1}^{\infty} (g_{m-n} g_m + h_{m-n} h_m) \right] \end{aligned} \quad (8)$$

$$\begin{aligned} \frac{\partial h_n}{\partial \sigma} = & -\alpha_n R_{Dm} h_n - \frac{1}{\pi \gamma n (1 + \sigma_m)^2} \nabla_{\chi, \psi}^2 g_n + \frac{n}{2N_m \sigma_{Dm} (1 + \sigma_m)} \\ & \times \left[\sum_{m=1}^{n-1} g_m h_{n-m} + \sum_{m=n+1}^{\infty} (h_{m-n} g_m - g_{m-n} h_m) \right] \end{aligned} \quad (9)$$

2.4 Boundary Conditions

The boundary condition for g_n and h_n at $\sigma = \sigma_m = 0$ is

$$\begin{aligned} g_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} c \sin n\tau_m d\tau_m \\ h_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} c \cos n\tau_m d\tau_m \end{aligned}$$

2.4.1 Circular Profile

The boundary condition at $\sigma = \sigma_m = 0$ is

$$\bar{p}(\sigma_m = 0, \xi, \tau) = c(\xi, \tau) \quad (10)$$

Suppose it is

$$\bar{p}(\sigma_m = 0, \xi, \tau) = c_1(\xi) \sin N_1 \tau + c_2(\xi) \sin N_2 \tau \quad (11)$$

2.4.2 Rectangular Profile

The boundary condition at $\sigma = \sigma_m = 0$ is

$$\bar{p}(\sigma_m = 0, \mathbf{u}, \tau) = c(\mathbf{u}, \tau) \quad (12)$$

Suppose it is

$$\bar{p}(\mathbf{u}, \sigma_m = 0, \tau) = c_1(\mathbf{u}) \sin N_1 \tau + c_2(\mathbf{u}) \sin N_2 \tau \quad (13)$$

$$\begin{aligned} &= c_1(\mathbf{u}) \left[\cos \left(\frac{\pi \gamma N_1}{4 N_m} \xi^2 \right) \sin N_1 \tau_m + \sin \left(\frac{\pi \gamma N_1}{4 N_m} \xi^2 \right) \cos N_1 \tau_m \right] \\ &+ c_2(\mathbf{u}) \left[\cos \left(\frac{\pi \gamma N_1}{4 N_m} \xi^2 \right) \sin N_2 \tau_m + \sin \left(\frac{\pi \gamma N_1}{4 N_m} \xi^2 \right) \cos N_2 \tau_m \right] \end{aligned} \quad (14)$$

Then

$$g_{N_i}(\chi, \psi, \sigma_m = 0) = c_i(\chi, \psi) \cos \left[\frac{\pi \gamma N_i}{4 N_m} (\chi^2 + \psi^2) \right], \quad i = 1, 2 \quad (15)$$

$$h_{N_i}(\chi, \psi, \sigma_m = 0) = c_i(\chi, \psi) \sin \left[\frac{\pi \gamma N_i}{4 N_m} (\chi^2 + \psi^2) \right], \quad i = 1, 2 \quad (16)$$

$$g_n(\chi, \psi, \sigma_m = 0) = h_n(\chi, \psi, \sigma_m = 0) = 0, \quad n \neq N_1, N_2 \quad (17)$$

2.5 Sound Power of Waves

The overall power of a wave is found by integrating the intensity across the entire field. Within the parabolic approximation the linear plane wave impedance relation is valid [1], and the nondimensional power of a wave is given by (Eq. (2.6) in [2])

$$\bar{\mathcal{P}}(\sigma_m) = \frac{1}{2\pi} \int_0^{2\pi} \int_0^\infty \bar{p}^2(\xi, \sigma_m, \tau) \xi \, d\xi \, d\tau$$

The nondimensional power of an individual harmonic component is given by (Eq (2.7) in [2])

$$\bar{\mathcal{P}}_n(\sigma_m) = \frac{1}{2\pi} \int_0^{2\pi} \int_0^\infty \bar{p}_n^2(\xi, \sigma_m, \tau) \xi \, d\xi \, d\tau \quad (18)$$

Transform the coordinates into (ζ, σ_m) , Eq. (18) can be written into

$$\begin{aligned} \bar{\mathcal{P}}_n(\sigma_m) &= \frac{1}{2\pi} \int_0^{2\pi} \int_0^\infty [g_n(\zeta, \sigma_m) \sin n\tau_m + h_n(\zeta, \sigma_m) \cos n\tau_m]^2 \, d\zeta \, d\tau \\ &= \frac{1}{2\pi} \int_0^{2\pi} \int_0^\infty [(g_n(\zeta, \sigma_m) \cos \varphi_n + h_n(\zeta, \sigma_m) \sin \varphi_n) \sin n\tau \\ &\quad + (-g_n(\zeta, \sigma_m) \sin \varphi_n + h_n(\zeta, \sigma_m) \cos \varphi_n) \cos n\tau]^2 \, \zeta \, d\zeta \, d\tau \end{aligned} \quad (19)$$

$$\begin{aligned} &= \frac{1}{2} \int_0^\infty \left[(g_n(\zeta, \sigma_m) \cos \varphi_n + h_n(\zeta, \sigma_m) \sin \varphi_n)^2 \right. \\ &\quad \left. + (-g_n(\zeta, \sigma_m) \sin \varphi_n + h_n(\zeta, \sigma_m) \cos \varphi_n)^2 \right] \zeta \, d\zeta \end{aligned} \quad (20)$$

where

$$\varphi_n = \frac{n}{N_m} (1 + \sigma_m) \zeta^2$$

3 Numerical Methods

3.1 Discretization

3.1.1 Rectangular Profile

$$\sigma_i = i\Delta\sigma, \quad i = 0, 1, \dots, I-1, I, \quad \sigma_0 = 0, \quad \sigma_I = \sigma_{\max}, \quad \Delta\sigma = \frac{\sigma_{\max}}{I}$$

$$\chi_j = j\Delta\chi, \quad j = 0, 1, \dots, J-1, J, \quad \chi_0 = 0, \quad \chi_J = \chi_{\max}, \quad \Delta\chi = \frac{\chi_{\max}}{J}$$

$$\psi_k = k\Delta\psi, \quad k = 0, 1, \dots, K-1, K, \quad \psi_0 = 0, \quad \psi_K = \psi_{\max}, \quad \Delta\psi = \frac{\psi_{\max}}{K}$$

For $\partial g_n / \partial \sigma$ or $\partial h_n / \partial \sigma$, we adopt Backward Implicit Finite Difference

$$\left. \frac{\partial g_n(\chi, \psi, \sigma_m)}{\partial \sigma_m} \right|_{i,j,k} = \frac{g_n^{i,j,k} - g_n^{i-1,j,k}}{\Delta\sigma_m}, \quad i = 1, 2, \dots, I \quad (21)$$

where

$$\begin{aligned} g_n^{i,j,k} &\approx g_n(\chi_j, \psi_k, \sigma_{m,i}) \\ h_n^{i,j,k} &\approx h_n(\chi_j, \psi_k, \sigma_{m,i}) \end{aligned} \quad (22)$$

For $\nabla_{\chi,\psi}^2 g_n$ or $\nabla_{\chi,\psi}^2 h_n$, 4 cases are considered:

(1) $j = 1, 2, \dots, J-1$, and $k = 1, 2, \dots, K-1$

$$\nabla_{\chi,\psi}^2 \big|_{i,j,k} = \frac{g_n^{i,j-1,k} - 2g_n^{i,j,k} + g_n^{i,j+1,k}}{(\Delta\chi)^2} + \frac{g_n^{i,j,k-1} - 2g_n^{i,j,k} + g_n^{i,j,k+1}}{(\Delta\psi)^2} \quad (23)$$

(2) $j = 1, 2, \dots, J-1$, and $k = 0$

$$\nabla_{\chi,\psi}^2 \big|_{i,j,0} = \frac{g_n^{i,j-1,0} - 2g_n^{i,j,0} + g_n^{i,j+1,0}}{(\Delta\chi)^2} + \frac{-2g_n^{i,j,0} + 2g_n^{i,j,1}}{(\Delta\psi)^2} \quad (24)$$

(3) $j = 0$, and $k = 1, 2, \dots, K-1$

$$\nabla_{\chi,\psi}^2 \big|_{i,0,k} = \frac{-2g_n^{i,0,k} + 2g_n^{i,1,k}}{(\Delta\chi)^2} + \frac{g_n^{i,0,k-1} - 2g_n^{i,0,k} + g_n^{i,0,k+1}}{(\Delta\psi)^2} \quad (25)$$

(4) $j = 0$, and $k = 0$

$$\nabla_{\chi,\psi}^2 \big|_{i,0,0} = \frac{-2g_n^{i,0,0} + 2g_n^{i,1,0}}{(\Delta\chi)^2} + \frac{-2g_n^{i,0,0} + 2g_n^{i,0,1}}{(\Delta\psi)^2} \quad (26)$$

where the symmetric conditions are used

$$\begin{aligned} \left. \frac{\partial g_n}{\partial \chi} \right|_{\chi=0} &= \left. \frac{\partial h_n}{\partial \chi} \right|_{\chi=0} = 0, \quad n = 1, 2, \dots, N \\ \left. \frac{\partial g_n}{\partial \psi} \right|_{\psi=0} &= \left. \frac{\partial h_n}{\partial \psi} \right|_{\psi=0} = 0, \quad n = 1, 2, \dots, N \end{aligned} \quad (27)$$

Discrete Eq. (8) and Eq. (9), we have

(1) $j = 1, 2, \dots, J-1$, and $k = 1, 2, \dots, K-1$

$$\begin{aligned} g_n^{i,j,k} - g_n^{i-1,j,k} &= -\Delta\sigma_m \alpha_n R_{\text{Dm}} g_n^{i,j,k} + \mathcal{G}_n^{i-1,j,k} + \frac{\Delta\sigma_m}{\pi\gamma n(1 + \sigma_{m,i})^2} \\ &\times \left[\frac{h_n^{i,j-1,k} - 2h_n^{i,j,k} + h_n^{i,j+1,k}}{(\Delta\chi)^2} + \frac{h_n^{i,j,k-1} - 2h_n^{i,j,k} + h_n^{i,j,k+1}}{(\Delta\psi)^2} \right] \end{aligned} \quad (28)$$

$$\begin{aligned} h_n^{i,j,k} - h_n^{i-1,j,k} &= -\Delta\sigma_m \alpha_n R_{\text{Dm}} h_n^{i,j,k} + \mathcal{H}_n^{i-1,j,k} - \frac{\Delta\sigma_m}{\pi\gamma n(1 + \sigma_{m,i})^2} \\ &\times \left[\frac{g_n^{i,j-1,k} - 2g_n^{i,j,k} + g_n^{i,j+1,k}}{(\Delta\chi)^2} + \frac{g_n^{i,j,k-1} - 2g_n^{i,j,k} + g_n^{i,j,k+1}}{(\Delta\psi)^2} \right] \end{aligned} \quad (29)$$

(2) $j = 1, 2, \dots, J-1$, and $k = 0$

$$g_n^{i,j,0} - g_n^{i-1,j,0} = -\Delta\sigma_m \alpha_n R_{\text{Dm}} g_n^{i,j,0} + \mathcal{G}_n^{i-1,j,0} + \frac{\Delta\sigma_m}{\pi\gamma n(1+\sigma_{m,i})^2} \times \left[\frac{h_n^{i,j-1,0} - 2h_n^{i,j,0} + h_n^{i,j+1,0}}{(\Delta\chi)^2} + \frac{-2h_n^{i,j,0} + 2h_n^{i,j,1}}{(\Delta\psi)^2} \right] \quad (30)$$

$$h_n^{i,j,0} - h_n^{i-1,j,0} = -\Delta\sigma_m \alpha_n R_{\text{Dm}} h_n^{i,j,0} + \mathcal{H}_n^{i-1,j,0} - \frac{\Delta\sigma_m}{\pi\gamma n(1+\sigma_{m,i})^2} \times \left[\frac{g_n^{i,j-1,0} - 2g_n^{i,j,0} + g_n^{i,j+1,0}}{(\Delta\chi)^2} + \frac{-2g_n^{i,j,0} + 2g_n^{i,j,1}}{(\Delta\psi)^2} \right] \quad (31)$$

(3) $j = 0$, and $k = 1, 2, \dots, K-1$

$$g_n^{i,0,k} - g_n^{i-1,0,k} = -\Delta\sigma_m \alpha_n R_{\text{Dm}} g_n^{i,0,k} + \mathcal{G}_n^{i-1,0,k} + \frac{\Delta\sigma_m}{\pi\gamma n(1+\sigma_{m,i})^2} \times \left[\frac{-2h_n^{i,0,k} + 2h_n^{i,1,k}}{(\Delta\chi)^2} + \frac{h_n^{i,0,k-1} - 2h_n^{i,0,k} + h_n^{i,0,k+1}}{(\Delta\psi)^2} \right] \quad (32)$$

$$h_n^{i,0,k} - h_n^{i-1,0,k} = -\Delta\sigma_m \alpha_n R_{\text{Dm}} h_n^{i,0,k} + \mathcal{H}_n^{i-1,0,k} - \frac{\Delta\sigma_m}{\pi\gamma n(1+\sigma_{m,i})^2} \times \left[\frac{-2g_n^{i,0,k} + 2g_n^{i,1,k}}{(\Delta\chi)^2} + \frac{g_n^{i,0,k-1} - 2g_n^{i,0,k} + g_n^{i,0,k+1}}{(\Delta\psi)^2} \right] \quad (33)$$

(4) $j = 0$, and $k = 0$

$$g_n^{i,0,0} - g_n^{i-1,0,0} = -\Delta\sigma_m \alpha_n R_{\text{Dm}} g_n^{i,0,0} + \mathcal{G}_n^{i-1,0,0} + \frac{\Delta\sigma_m}{\pi\gamma n(1+\sigma_{m,i})^2} \times \left[\frac{-2h_n^{i,0,0} + 2h_n^{i,1,0}}{(\Delta\chi)^2} + \frac{-2h_n^{i,0,0} + 2h_n^{i,0,1}}{(\Delta\psi)^2} \right] \quad (34)$$

$$h_n^{i,0,0} - h_n^{i-1,0,0} = -\Delta\sigma_m \alpha_n R_{\text{Dm}} h_n^{i,0,0} + \mathcal{H}_n^{i-1,0,0} - \frac{\Delta\sigma_m}{\pi\gamma n(1+\sigma_{m,i})^2} \times \left[\frac{-2g_n^{i,0,0} + 2g_n^{i,1,0}}{(\Delta\chi)^2} + \frac{-2g_n^{i,0,0} + 2g_n^{i,0,1}}{(\Delta\psi)^2} \right] \quad (35)$$

where

$$\mathcal{G}_n^{i,j,k} = \frac{n\Delta\sigma_m}{2N_m\sigma_{\text{Dm}}(1+\sigma_{m,i})} \left[\frac{1}{2} \sum_{m=1}^{n-1} (g_m^{i,j,k} g_{n-m}^{i,j,k} - h_m^{i,j,k} h_{n-m}^{i,j,k}) - \sum_{m=n+1}^N (g_m^{i,j,k} g_{m-n}^{i,j,k} + h_m^{i,j,k} h_{m-n}^{i,j,k}) \right] \quad (36)$$

$$\mathcal{H}_n^{i,j,k} = \frac{n\Delta\sigma_m}{2N_m\sigma_{\text{Dm}}(1+\sigma_{m,i})} \left[\sum_{m=1}^{n-1} g_m^{i,j,k} h_{n-m}^{i,j,k} + \sum_{m=n+1}^N (g_m^{i,j,k} h_{m-n}^{i,j,k} - h_m^{i,j,k} g_{m-n}^{i,j,k}) \right] \quad (37)$$

3.2 Matrix Form

3.2.1 Rectangular Profile

$$\begin{bmatrix} \mathbf{E}_n & -\mathbf{A}_n^i \\ \mathbf{A}_n^i & \mathbf{E}_n \end{bmatrix} \begin{bmatrix} \mathbf{G}_n^i \\ \mathbf{H}_n^i \end{bmatrix} = \begin{bmatrix} \mathbf{G}_n^{i-1} \\ \mathbf{H}_n^{i-1} \end{bmatrix} + \begin{bmatrix} \mathcal{G}_n^{i-1} \\ \mathcal{H}_n^{i-1} \end{bmatrix} \quad (38)$$

where

$$\mathbf{E}_n = (1 + \Delta\sigma_m \alpha_n R_{\text{Dm}}) \mathbf{I}_{JK \times JK}$$

$$\mathbf{G}_n^i = \begin{bmatrix} \mathbf{g}_n^{i,0,\cdot} \\ \mathbf{g}_n^{i,1,\cdot} \\ \vdots \\ \mathbf{g}_n^{i,J-1,\cdot} \end{bmatrix}, \quad \mathbf{g}_n^{i,j,\cdot} = \begin{bmatrix} \mathbf{g}_n^{i,j,0} \\ \mathbf{g}_n^{i,j,1} \\ \vdots \\ \mathbf{g}_n^{i,j,K-1} \end{bmatrix}, \quad \mathbf{H}_n^i = \begin{bmatrix} \mathbf{h}_n^{i,0,\cdot} \\ \mathbf{h}_n^{i,1,\cdot} \\ \vdots \\ \mathbf{h}_n^{i,J-1,\cdot} \end{bmatrix}, \quad \mathbf{h}_n^{i,j,\cdot} = \begin{bmatrix} \mathbf{h}_n^{i,j,0} \\ \mathbf{h}_n^{i,j,1} \\ \vdots \\ \mathbf{h}_n^{i,j,K-1} \end{bmatrix}$$

$$\begin{aligned}
\mathcal{G}_n^i &= \begin{bmatrix} \mathcal{G}_n^{i,0,\cdot} \\ \mathcal{G}_n^{i,1,\cdot} \\ \vdots \\ \mathcal{G}_n^{i,J-1,\cdot} \end{bmatrix}, \quad \mathcal{G}_n^{i,j,\cdot} = \begin{bmatrix} \mathcal{G}_n^{i,j,0} \\ \mathcal{G}_n^{i,j,1} \\ \vdots \\ \mathcal{G}_n^{i,j,K-1} \end{bmatrix}, \quad \mathcal{H}_n^i = \begin{bmatrix} \mathcal{H}_n^{i,0,\cdot} \\ \mathcal{H}_n^{i,1,\cdot} \\ \vdots \\ \mathcal{H}_n^{i,J-1,\cdot} \end{bmatrix}, \quad \mathcal{H}_n^{i,j,\cdot} = \begin{bmatrix} \mathcal{H}_n^{i,j,0} \\ \mathcal{H}_n^{i,j,1} \\ \vdots \\ \mathcal{H}_n^{i,j,K-1} \end{bmatrix} \\
\mathbf{A}_n^i &= \frac{\Delta\sigma_m}{\pi\gamma n(1+\sigma_{m,i})^2(\Delta\chi)^2} \mathbf{A}_{n,\chi} + \frac{\Delta\sigma_m}{\pi\gamma n(1+\sigma_{m,i})^2(\Delta\psi)^2} \mathbf{A}_{n,\psi} \\
\mathbf{A}_{n,\chi} &= \begin{bmatrix} \mathbf{B} & & \mathbf{0} \\ & \ddots & \\ \mathbf{0} & & \mathbf{B} \end{bmatrix}_{JK \times JK}, \quad \mathbf{B} = \begin{bmatrix} -2 & 2 & & & 0 \\ 1 & -2 & 1 & & \\ & & & \ddots & \\ & & & 1 & -2 & 1 \\ 0 & & & & 1 & -2 \end{bmatrix}_{K \times K} \\
\mathbf{A}_{n,\psi} &= \begin{bmatrix} -2\mathbf{I}_{K \times K} & 2\mathbf{I}_{K \times K} & & & \mathbf{0}_{K \times K} \\ \mathbf{I}_{K \times K} & -2\mathbf{I}_{K \times K} & \mathbf{I}_{K \times K} & & \\ & & & \ddots & \\ & & & \mathbf{I}_{K \times K} & -2\mathbf{I}_{K \times K} & \mathbf{I}_{K \times K} \\ \mathbf{0}_{K \times K} & & & & \mathbf{I}_{K \times K} & -2\mathbf{I}_{K \times K} \end{bmatrix}_{JK \times JK}
\end{aligned}$$

4 Test Data

4.1 Circular Profile

4.1.1 Kamakura 1989

Fig. 6

$$\begin{aligned}
\Delta\sigma &= 3 \times 10^{-4}, \quad \sigma_{\max} = 10, \quad \Delta\xi = 0.03, \quad \xi_{\max} = 6, \quad T = 23.7^\circ\text{C}, \quad h_r = 65.2\% \\
M &= 18, \quad f_1 = 25\text{kHz}, \quad f_2 = 30\text{kHz}, \quad P_0(25\text{kHz}) = 109.5\text{dB}, \quad P_0(30\text{kHz}) = 108.5\text{dB} \\
a &= 0.21\text{m}
\end{aligned}$$

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