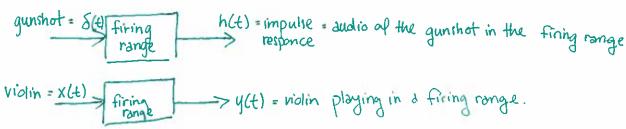
Sig Sys - Problem set 6 Ver 2

Impulse - a infinantly tall rectangle w/ area A (I for unit impulses) that happens almost instantly Impulse response - the result after you send the impulse through the system.



In order to find how the system changes the violin, you need to convolve the violin signal with the impulse response.

$$x(t) * h = y(t)$$

As you are convolving the 2, you are picking x(G) and then multiplying and shifting based on your impulse response. The result is y(t) which is how the violin sounds in the shooting range.

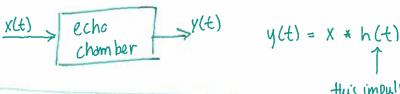
2.  $y(t) = \frac{1}{2}x(t-1) + \frac{1}{4}x(t-10)$  — Hus is a model for an echo chamber

It moves the original
Signal to t=1 w/
It moves another copy
to t=10 w/on amplitude
of \frac{1}{2}

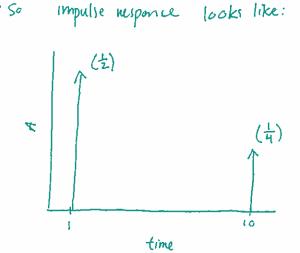
of \frac{1}{4}

It's an echo chamber because the signal "bances" with decreasing amplitudes as it loses energy from the bounces.

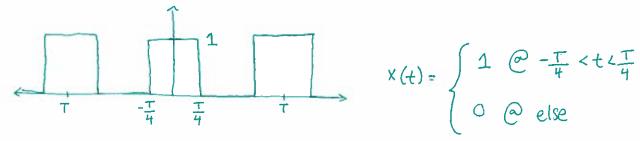
I've know that when we convolue with an impulse we multiply by the area of the impulse.



this impulse response must shift 2 copies of the signal and decrease the amplitudes



3. Find the Fourier series nep for this square wave



Integrate both sides 
$$w/\frac{1}{T}\int_{-T/2}^{T/2} x(t)e^{-J\frac{2\pi}{T}mt} = \frac{1}{T}\int_{-T/2}^{T/2} \sum_{k=\infty}^{\infty} C_k e^{J\frac{2\pi}{T}(k-m)} t$$

Switch the Summation

Summation & the integral probably bic adding is communative

(mis like aparticular coefficient) =  $\sum_{k=0}^{\infty} \frac{Ck}{l} \int_{-T/2}^{T/2} e^{j\frac{2\pi}{L}(k-m)t} dt$ 

$$= \sum_{k=-\infty}^{\infty} \frac{C_k}{1} \int_{-T/2}^{T/2} e^{j\frac{2\pi}{T}(k-m)t} dt$$

& So when you do out the moth you find that for any k that i) not = m the value is a so the summation value only depends on when k=m

$$\frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} X(t) e^{-\frac{1}{2} \frac{\pi}{T} mt} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2} \frac{\pi}{T}} (0) t}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} = \frac{C \ln \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} dt \qquad e^{\frac{1}{2}} dt \qquad e^{\frac{1}{2}} dt \qquad e^{\frac{1}{2}} dt}{T} dt \qquad e^{\frac{1}{2}} dt \qquad e^{\frac{1$$

= Cm = equation for the mth Fourier series coefficient

Finding on expression for CK

$$C_{k} = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} X(t) e^{-\frac{1}{2}} dt \qquad X(t) = \begin{cases} 1 & e^{-\frac{T}{4}} < t < \frac{T}{4} \\ 0 & e \text{ disc} \end{cases}$$

$$C_{k} = \frac{1}{T} \int_{-\frac{T}{4}}^{-\frac{T}{4}} 0 \cdot e^{-\frac{1}{2}} dt + \int_{-\frac{T}{4}}^{\frac{T}{4}} 1 \cdot e^{-\frac{1}{2}} dt + \int_{+\frac{T}{4}}^{\frac{T}{2}} 0 \cdot e^{-\frac{1}{2}} dt + \int_{+\frac{T}{4}}^{\frac{T}{2}} 0 \cdot e^{-\frac{1}{2}} dt + \int_{+\frac{T}{4}}^{\frac{T}{2}} (1 \cdot e^{-\frac{1}{2}} + e^{-\frac{1}{2}} + e^{-\frac{1}{2}} + e^{-\frac{1}{2}} + e^{-\frac{1}{2}} + e^{-\frac{1}{2}} dt + e$$

$$C_{k} = \frac{1}{T} \left( \frac{1}{-J^{\frac{2\pi}{T}}k} e^{-J^{\frac{2\pi}{T}}k^{\frac{\pi}{T}}} - \frac{1}{-J^{\frac{2\pi}{T}}k} e^{+J^{\frac{\pi}{T}}k^{\frac{\pi}{T}}} \right)$$

$$Sin(\theta) = \frac{1}{2J}e^{-J\theta}$$

$$C_{k} = \frac{1 \cdot \pi}{T \cdot 3\pi J k} e^{-JJ^{\frac{\pi}{T}}k} e^{-JJ^{\frac{\pi}{T}}k} e^{-J\theta}$$

$$C_{k} = \frac{1 \cdot \pi}{2J^{\frac{\pi}{T}}k} e^{-JJ^{\frac{\pi}{T}}k} e^{-J\theta}$$

$$C_{k} = \frac{1 \cdot \pi}{2J^{\frac{\pi}{T}}k} e^{-J\theta} e^{-J\theta}$$

$$C_{k} = \frac{1 \cdot \pi}{2J^{\frac{\pi}{T}}k} e^{-J\theta}$$

$$C_{k} = \frac{1}{-\pi k} \left( \frac{1}{2j} e^{j\theta} - \frac{1}{2j} e^{-j\theta} \right)$$

$$C_{k} = \frac{1}{-\pi k} \sin \left( -\frac{\pi}{2k} \right)$$

$$C_{k} = \frac{1}{-\pi k} \sin \left( -\frac{\pi}{2k} \right) \text{ if add}$$

- c) In the three rappoximations of the square wave Inotice that as I increase the # of Fairier coefficients the square were becomes from defined. However it does appear that the approximate signal comnot jump from 0 to 1 immediately, that there must be some slope to make it a continuous function and that no matter what the Gibbs phenomenon ( ) will always be present. Even though the spikes will never go away, higher approximations of the wore still result in a more better approximation because the Gibbs spike lose area as they get smaller.
- 4. Suppose X(t) is a periodic signal w/ fundamental period T, and has a Fourier series representation w/ coefficient Cx. Consider anw signal y(t)= x(t-T.) where |T.12T such that y(t) is a delayed version of x(t). Find the Crs of y(t) in terms of Ck

$$C_{k} = \frac{1}{T} \int_{-T/2}^{T/2} x(t-T_{1})e^{-\frac{2\pi}{T}Jkt} dt \qquad \Upsilon = t-T_{1}$$

$$\frac{d\Upsilon}{dt} = 1 \qquad d\Upsilon = 1 dt$$

$$C_{k} = \frac{1}{T} \int_{-T/2}^{T/2-T_{1}} x(\Upsilon)e^{-\frac{2\pi}{T}Jk(\Upsilon+T_{1})} d\Upsilon$$

4 cont

# Because shifting the window of integration doesnot change the Ck for a pendiction signal  $Ck_{min} \frac{1}{T} \int_{-T/2}^{T/2} X(\tau) e^{-J\frac{2\pi}{T}k(\tau+T_i)} d\tau \qquad \text{that's just a phase shift}$   $Ck_{min} \frac{1}{T} \int_{-T/2}^{T/2} X(\tau) e^{-J\frac{2\pi}{T}k\tau} d\tau \qquad \text{constant}$   $Ck_{min} = \frac{1}{e^{J\frac{2\pi}{T}kT_i}} \cdot \frac{1}{T} \int_{-T/2}^{T/2} X(\tau) e^{-J\frac{2\pi}{T}k\tau} d\tau \qquad \text{constant}$   $Ck_{min} = \frac{1}{e^{J\frac{2\pi}{T}kT_i}} \cdot \frac{1}{T} \int_{-T/2}^{T/2} X(\tau) e^{-J\frac{2\pi}{T}k\tau} d\tau \qquad \text{constant}$  = Ck  $Ck_{min} = \frac{1}{e^{J\frac{2\pi}{T}kT_i}} \cdot Ck$ 

 $\bigcirc$ 

