Finite Element solution of the fiber/matrix interface crack problem: convergence properties and mode mixity of the Virtual Crack Closure Technique

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Abstract

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1. Introduction

Bi-material interfaces represent the basic load transfer mechanism at the heart of Fiber Reinforced Polymer Composite (FRPC) materials. They are present at the macroscale, in the form of adhesive joints; at the mesoscale, as interfaces between layers with different orientations; at the microscale, as fiber-matrix interfaces. Bi-material interfaces have for long attracted the attention of researchers in Fracture Mechanics [1, 2], due to their hidden complexity.

The problem was first addressed in the 1950's by Williams [3], who derived through a linear elastic asymptotic analysis the stress distribution around an *open* crack (i.e. with crack faces nowhere in contact for any size of the crack) between two infinite half-planes of dissimilar materials and found the existence of a strong oscillatory behavior in the stress singularity at the crack tip of the form

$$r^{-\frac{1}{2}} \sin\left(\varepsilon \log r\right) \quad \text{with} \quad \varepsilon = \frac{1}{2\pi} \log\left(\frac{1-\beta}{1+\beta}\right); \tag{1}$$

in which β is one of the two parameters introduced by Dundurs [4] to characterize bi-material interfaces:

$$\beta = \frac{\mu_2 (\kappa_1 - 1) - \mu_1 (\kappa_2 - 1)}{\mu_2 (\kappa_1 + 1) + \mu_1 (\kappa_2 + 1)} \tag{2}$$

where $\kappa = 3 - 4\nu$ in plane strain and $\kappa = \frac{3-4\nu}{1+\nu}$ in plane stress, μ is the shear modulus, ν Poisson's coefficient, and indexes $1, 2$ refer to the two bulk materials joined at the interface. Defining a as the length of the crack, it was found that the size of the oscillatory region is in the order of $10^{-6}a$ [5]. Given the oscillatory behaviour of the crack tip singularity of the stress field of Eq. 1, the definition of Stress Intensity Factor (SIF) $\lim_{r\to 0} \sqrt{2\pi r}\sigma$ ceases to be valid as it returns logarithmically infinite terms [1]. Furthermore, it implies that the Mode mixity problem at the crack tip is ill-posed.

It was furthermore observed, always in the context of Linear Elastic Fracture Mechanics (LEFM), that an interpenetration zone exists close to the crack tip [6, 7] with a length in the order of 10^{-4} [6]. Following conclusions firstly proposed in[7], the presence of a *contact zone* in the crack tip neighborhood, of a length to be determined from the solution of the elastic problem, was introduced in [8] and shown to provide a physically consistent solution to the straight bi-material interface crack problem.

The curved bi-material interface crack, more often refered to as the fiber-matrix interface crack (or debond) due to its relevance in FRPCs, was first treated by England [9] and by Perlman and Sih [10], who provided the analytical solution of stress and displacement fields for a circular inclusion with respectively a single debond and an arbitrary number of debonds. Building on their work, Toya [11] particularized the solution and provided the expression of the Energy Release Rate (ERR) at the crack tip. The same problems exposed previously for the *open* straight bi-material were shown to exist also for the *open* fiber-matrix interface crack: the presence of strong oscillations in the crack tip singularity and crack face interpenetration after a critical initial flaw size.[1]

[1] For the fiber-matrix interface crack, flaw size is measured in terms of the angle $\Delta\theta$ sub-

In order to treat cases more complex than the single partially debonded fiber in an infinite matrix of [9, 10, 11], numerical studies followed. In the 1990's, París and collaborators [12] developed a Boundary Element Method (BEM) with the use of discontinuous singular elements at the crack tip and the Virtual Crack Closure Integral (VCCI) [13] for the evaluation of the Energy Release Rate (ERR). They validated their results [12] with respect to Toya's analytical solution [11] and analyzed the effect of BEM interface discretization on the stress field in the neighborhood of the crack tip [14]. Following Comninou's work on the straight crack [8], they furthermore recognized the importance of contact to retrieve a physical solution avoiding interpenetration [12] and studied the effect of the contact zone on debond ERR [15]. Their algorithm was then applied to investigate the fiber-matrix interface crack under different geometrical configurations and mechanical loadings [16, 17, 18, 19, 20, 21, 22].

Recently the Finite Element Method (FEM) was also applied to the solution of the fiber-matrix interface crack problem [23, 24, 25], in conjunction with the Virtual Crack Closure Technique (VCCT) [26, 27] for the evaluation of the ERR at the crack tip. In [23], the authors validated their model with respect to the BEM results of [12], but no analysis of the effect of the discretization in the crack tip neighborhood comparable to [14] was proposed. Thanks to the interest in evaluating the ERR of interlaminar delamination, different studies exist in the literature on the effect of mesh discretization on Mode I and Mode II ERR of the bi-material interface crack when evaluated with the VCCT in the context of the FEM [28, 29, 30]. However, no comparable analysis can be found in the literature on the application of the VCCT to the fiber-matrix interface crack (circular bi-material interface crack) problem in the context of a linear elastic FEM solution. It is this gap that the present work aims to address. We first present the FEM formulation of the problem, together with the main geometrical characteristics, material properties, boundary conditions and loading. We then propose a vectorial formulation of the VCCT and express

tended by half of the arc-crack, i.e. $a = 2\Delta\theta$.

the Mode I and Mode II ERR in terms of the FEM natural variables. With this tool, we derive an analytical estimate of the ERR convergence and compare it with numerical results.

2. FEM formulation of the fiber-matrix interface crack problem

In order to investigate the fiber-matrix interface crack problem, a 2-dimensional model of a single fiber inserted in a rectangular matrix element is considered (see Figure 1). Total element length and height are respectively $2L$ and L, where L is determined by the fiber radius R_f and the fiber volume fraction V_f by

$$L = \frac{R_f}{2}\sqrt{\frac{\pi}{V_f}}. \tag{3}$$

The fiber radius R_f is assumed to be equal to 1 μm. This choice is not dictated by physical considerations but for simplicity. It is thus useful to remark that, in a linear elastic solution as the one considered in the present work, the ERR is proportional to the geometrical dimensions of the model and, consequently, recalculation of the ERR for fibers of any size requires a simple multiplication.

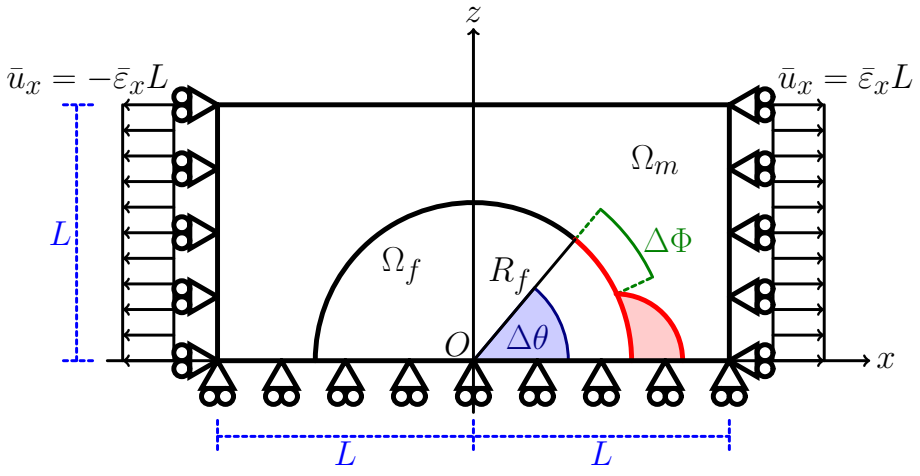


Figure 1: Schematic of the model with its main parameters.

As shown in Fig. 1, the debond is placed symmetrically with respect to the x axis and its size is characterized by the angle $\Delta\theta$ (which makes the full debond size equal to $2\Delta\theta$ and the full crack length equal to $R_f 2\Delta\theta$). A region $\Delta\Phi$ of variable size appears at the crack tip for large debond sizes (at least $\geq 60° - 80°$), in which the crack faces are in contact with each other and free to slide. Frictionless contact is thus considered between the two crack faces to allow free sliding and avoid interpenetration. Symmetry with respect to the x axis is applied on the lower boundary while the upper surface is left free. Kinematic coupling on the x-displacement is applied along the left and right sides of the model in the form of a constant x-displacement $\pm\bar{\varepsilon}_x L$, which corresponds to transverse strain $\bar{\varepsilon}_x$ equal to 1%.

Table 1: Summary of the mechanical properties of fiber and matrix. E stands for Young's modulus, μ for shear modulus and ν for Poisson's ratio.

Material	$E\,[GPa]$	$\mu\,[GPa]$	$\nu\,[-]$
Glass fiber	70.0	29.2	0.2
Epoxy	3.5	1.25	0.4

The model problem is solved with the Finite Element Method (FEM) within the Abaqus environment, a commercial FEM software [31]. The model is meshed with second order, 2D, plane strain triangular (CPE6) and rectangular (CPE8) elements. A regular mesh of quadrilateral elements with almost unitary aspect ratio is used at the crack tip. The angular size δ of an element in the crack tip neighborhood represents the main parameter of the numerical analysis. The crack faces are modeled as element-based surfaces and a small-sliding contact pair interaction with no friction is imposed between them. The Mode I, Mode II and total Energy Release Rates (ERRs) (respectively referred to as G_I, G_{II} and G_{TOT}) are evaluated using the VCCT [27], implemented in a in-house Python routine. A glass fiber-epoxy system is considered in the present work, and it is assumed that their response lies always in the linear elastic domain. The material properties of glass fiber and epoxy are reported in Table 1.

3. Vectorial formulation of the Virtual Crack Closure Technique (VCCT)

In order to express the VCCT formulation of the ERR in terms of FEM variables, we need to introduce a few rotation matrices in order to represent the discretized representation (FE mesh) of a crack along a circular interface. The position of the crack tip is characterized by the angular size of the crack (see Sec. 2 and Fig. 1 for reference) and the rotation corresponding to the crack tip reference frame is represented by the matrix $\underline{\underline{R}}_{\Delta\theta}$ defined as

$$\underline{\underline{R}}_{\Delta\theta} = \begin{bmatrix} \cos(\Delta\theta) & \sin(\Delta\theta) \\ -\sin(\Delta\theta) & \cos(\Delta\theta) \end{bmatrix}. \tag{4}$$

Nodes belonging to the elements sharing the crack tip are involved in the VCCT estimation of the ERR and it is assumed that, given a sufficiently fine discretization, they are aligned with the crack propagation direction defined at the crack tip. However small the elements in the crack tip neighborhood, a misalignment always exists with respect to the assumed crack propagation direction. This is measured by the matrices $\underline{\underline{P}}_{\delta}(p)$, defined as

$$\underline{\underline{P}}_{\delta}(p) = \begin{bmatrix} \cos\left(\left(1+\frac{1-p}{m}\right)\delta\right) & \sin\left(\left(1+\frac{1-p}{m}\right)\delta\right) \\ -\sin\left(\left(1+\frac{1-p}{m}\right)\delta\right) & \cos\left(\left(1+\frac{1-p}{m}\right)\delta\right) \end{bmatrix} \tag{5}$$

and $\underline{\underline{Q}}_{\delta}(q)$, equal to

$$\underline{\underline{Q}}_{\delta}(q) = \begin{bmatrix} \cos\left(\frac{q-1}{m}\delta\right) & \sin\left(\frac{q-1}{m}\delta\right) \\ -\sin\left(\frac{q-1}{m}\delta\right) & \cos\left(\frac{q-1}{m}\delta\right) \end{bmatrix}, \tag{6}$$

respectively for the free and bonded nodes involved in the VCCT estimation. In Eqs. 5 and 6, δ is the angular size of an element in the crack tip neighborhood (see Sec. 2 and Fig. 1), m is the order of the element shape functions and p, q are indices referring to the nodes belonging respectively to free and bonded elements sharing the crack tip. Introducing the permutation matrix

$$\underline{\underline{P}}_{\pi} = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}, \tag{7}$$

it is possible to express the derivatives of rotation matrices $\underline{\underline{R}}_{\Delta\theta}$, $\underline{\underline{P}}_{\delta}$ and $\underline{\underline{Q}}_{\delta}$ with respect to their argument:

$$\frac{\partial \underline{\underline{R}}_{\Delta\theta}}{\partial \Delta\theta} = \underline{\underline{D}} \cdot \underline{\underline{R}}_{\Delta\theta}, \quad \frac{\partial \underline{\underline{P}}_{\delta}}{\partial \delta} = \left(1 + \frac{1-p}{m}\right) \underline{\underline{D}} \cdot \underline{\underline{P}}_{\delta}, \quad \frac{\partial \underline{\underline{Q}}_{\delta}}{\partial \delta} = \frac{q-1}{m} \underline{\underline{D}} \cdot \underline{\underline{Q}}_{\delta}. \quad (8)$$

By means of Eqs. 5 and 6, we can express the crack tip forces $\underline{F}_{xy} = \begin{bmatrix} F_x \\ F_y \end{bmatrix}$ and crack displacements $\underline{u}_{xy} = \begin{bmatrix} u_x \\ u_y \end{bmatrix}$ in the crack tip reference frame (where the tangential direction θ correspond to the direction of crack propagation) while taking into account the misalignment to the finite discretization as

$$\underline{F}_{r\theta} = \underline{\underline{Q}}_{\delta} \underline{\underline{R}}_{\Delta\theta} \underline{F}_{xy} \qquad \underline{u}_{r\theta} = \underline{\underline{P}}_{\delta} \underline{\underline{R}}_{\Delta\theta} \underline{u}_{xy} \quad (9)$$

where $\underline{F}_{r\theta} = \begin{bmatrix} F_r \\ F_\theta \end{bmatrix}$ and $\underline{u}_{r\theta} = \begin{bmatrix} u_r \\ u_\theta \end{bmatrix}$.

The crack tip forces can be expressed as a function of the crack opening displacement as

$$\underline{F}_{xy} = \underline{\underline{K}}_{xy} \underline{u}_{xy} + \widetilde{\underline{F}}_{xy}, \quad (10)$$

where $\underline{\underline{K}}_{xy}$ is in general a full matrix of the form $\underline{\underline{K}}_{xy} = \begin{bmatrix} K_{xx} K_{xy} \\ K_{yx} K_{yy} \end{bmatrix}$ and $\widetilde{\underline{F}}_{xy}$ represents the effect of the rest of the FE solution through the remaining nodes of the elements attached to the crack tip. As such, the term $\widetilde{\underline{F}}_{xy}$ can be expressed as a linear combination of the solution vector $\underline{u}_N$ of nodal displacements of the form $\widetilde{\underline{\underline{K}}}_N \underline{u}_N$. Equation 10 thus become

$$\underline{F}_{xy} = \underline{\underline{K}}_{xy} \underline{u}_{xy} + \widetilde{\underline{\underline{K}}}_N \underline{u}_N. \quad (11)$$

An exemplifying derivation of the relationships expressed in Equations 10 and 11 can be found in Appendix A. It is worthwhile to observe that another author [32] proposed a relationship of the form $\underline{F}_{xy} = \underline{\underline{K}}_{xy} \underline{u}_{xy}$. However, in [32],

this relationship is assumed *a priori* and manipulated to propose a revised version of the VCCT, based on the assumption that the matrix $\underline{\underline{K}}_{xy}$ should be diagonal to provide physically-consistent fracture mode partitioning. On the other hand, in the present work we derive the relationships of Eqs. 10 and 11 from the formulation of the Finite Element Method. According to our derivation, it seems correct that the matrix $\underline{\underline{K}}_{xy}$ should not in general be diagonal in order to take into account Poisson's effect. In fact, a positive crack opening displacement would cause a transverse displacement in the neighborhood of the crack tip. Given that material properties are different on the two sides of a bi-material interface, a net shear would be applied to the crack tip which would correspond to a net contribution to the crack tip force related to crack shear displacement. The analytical derivations presented in Appendix A confirm these physical considerations.

Based upon the work of Raju [33], we introduce the matrix $\underline{\underline{T}}_{pq}$ to represent the weights needed in the VCCT to account for the use of singular elements. As already done previously, indices p and q refer to nodes placed respectively on the free (crack face) and bonded side of the crack tip. Nodes are enumerated so that the crack tip has always index 1, i.e. the higher the index the further the node is from the crack tip. Matrix $\underline{\underline{T}}_{pq}$ has always a size of $d \times d$ where d is the number of geometrical dimensions of the system. An element $\underline{\underline{T}}_{pq}(i,j)$ with $i,j = 1,\ldots,d$ represents the weight to be assigned to the product of component i of the displacement extracted at node p with component j of the force extracted at node q. The expression of $\underline{\underline{T}}_{pq}$ for quadrilateral elements with or without singularity is reported in Appendix B. Notice that, given m is the order of the element shape functions, the element side has $m+1$ nodes and this represents the upper limit of indices p and q.

By using matrix $\underline{\underline{T}}_{pq}$, it is possible to express the total ERR G evaluated with the VCCT as

$$G_{TOT} = \frac{1}{2R_f\delta} \sum_{p=1}^{m+1} \sum_{q=1}^{m+1} Tr\left(\underline{u}^T_{r\theta,p} \underline{\underline{T}}^T_{pq} \underline{F}_{r\theta,q}\right). \tag{12}$$

Introducing the vector $\underline{G} = \begin{bmatrix} G_I \\ G_{II} \end{bmatrix}$ of fracture mode ERRs, Mode I and Mode II ERR evaluated with the VCCT can be expressed as

$$\underline{G} = \frac{1}{2R_f\delta} \sum_{p=1}^{m+1} \sum_{q=1}^{m+1} Diag\left(\underline{F}_{r\theta,q} \underline{u}^T_{r\theta,p} \underline{\underline{T}}^T_{pq}\right), \tag{13}$$

where $Diag\,()$ is the function that extracts as a column vector the diagonal of the matrix provided as argument. Substituting Equations 9 and 11 in Equations 12 and 13, we can express the Mode I, Mode II and total Energy Release Rate as a function of the crack displacements and the FE solution (mode details in **??**) as

$$\begin{aligned} G_{TOT} =& \frac{1}{2R_f\delta} \sum_{p=1}^{m+1} \sum_{q=1}^{m+1} Tr\left(\underline{\underline{Q}}_\delta \underline{\underline{R}}_{\Delta\theta} \underline{\underline{K}}_{xy,q} \underline{u}_{xy,q} \underline{u}^T_{xy,p} \underline{\underline{R}}^T_{\Delta\theta} \underline{\underline{P}}^T_\delta \underline{\underline{T}}^T_{pq}\right) + \\ &+ \frac{1}{2R_f\delta} \sum_{p=1}^{m+1} \sum_{q=1}^{m+1} Tr\left(\underline{\underline{Q}}_\delta \underline{\underline{R}}_{\Delta\theta} \underline{\widetilde{F}}_{xy,q} \underline{u}^T_{xy,p} \underline{\underline{R}}^T_{\Delta\theta} \underline{\underline{P}}^T_\delta \underline{\underline{T}}^T_{pq}\right) \end{aligned} \tag{14}$$

and

$$\begin{aligned} \underline{G} = \begin{bmatrix} G_I \\ G_{II} \end{bmatrix} =& \frac{1}{2R_f\delta} \sum_{p=1}^{m+1} \sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_\delta \underline{\underline{R}}_{\Delta\theta} \underline{\underline{K}}_{xy,q} \underline{u}_{xy,q} \underline{u}^T_{xy,p} \underline{\underline{R}}^T_{\Delta\theta} \underline{\underline{P}}^T_\delta \underline{\underline{T}}^T_{pq}\right) + \\ &+ \frac{1}{2R_f\delta} \sum_{p=1}^{m+1} \sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_\delta \underline{\underline{R}}_{\Delta\theta} \underline{\widetilde{\underline{K}}}_{N,q} \underline{u}_N \underline{u}^T_{xy,p} \underline{\underline{R}}^T_{\Delta\theta} \underline{\underline{P}}^T_\delta \underline{\underline{T}}^T_{pq}\right) \end{aligned} \tag{15}$$

4. Rotational invariance of G_{TOT}

Recalling Equation 14 and observing that matrix $\underline{\underline{T}}_{pq}$ is always equal to the identity matrix pre-multiplied by a suitable real constant (see Eq. B.1 in Appendix B), the total Energy Release Rate can be rewritten as

$$G_{TOT} = \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Tr\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\left(\underline{\underline{K}}_{xy,q}\underline{u}_{xy,q} + \underline{\widetilde{F}}_{xy,q}\right)\underline{u}^T_{xy,p}\underline{\underline{T}}^T_{pq}\underline{\underline{R}}^T_{\Delta\theta}\underline{\underline{P}}^T_{\delta}\right) =$$
$$= \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Tr\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\underline{F}_{xy,q}\underline{u}^T_{xy,p}\underline{\underline{T}}^T_{pq}\underline{\underline{R}}^T_{\Delta\theta}\underline{\underline{P}}^T_{\delta}\right), \tag{16}$$

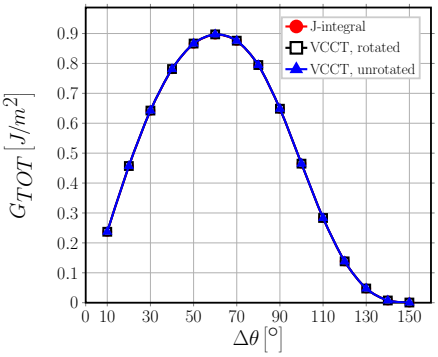
where $\underline{F}_{xy}$ and $\underline{u}_{xy}$ are the vectors of respectively the crack tip forces and crack displacements in the global $(x-y)$ reference frame. Given that $\underline{\underline{Q}}_{\delta}$, $\underline{\underline{P}}_{\delta}$ and $\underline{\underline{R}}_{\Delta\theta}$ all represent a linear transformation (a rigid rotation in particular), the invariance of the trace to linear transformations ensures that

$$\begin{aligned} G_{TOT} &= \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Tr\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\underline{F}_{xy,q}\underline{u}^T_{xy,p}\underline{\underline{T}}^T_{pq}\underline{\underline{R}}^T_{\Delta\theta}\underline{\underline{P}}^T_{\delta}\right) = \\ &= \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Tr\left(\underline{F}_{xy,q}\underline{u}^T_{xy,p}\underline{\underline{T}}^T_{pq}\right). \end{aligned} \tag{17}$$

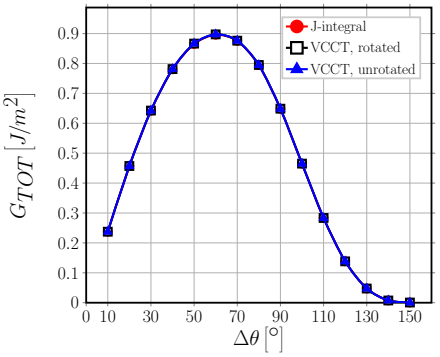
As G_{TOT} was defined according to Equation 12 and given that $Tr\,(AB) = Tr\,(BA)$, it holds that

$$G_{TOT} = \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1}\underline{u}^T_{r\theta,p}\underline{\underline{T}}^T_{pq}\underline{F}_{r\theta,q} = \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Tr\left(\underline{F}_{r\theta,q}\underline{u}^T_{r\theta,p}\underline{\underline{T}}^T_{pq}\right) =$$
$$= \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Tr\left(\underline{F}_{xy,q}\underline{u}^T_{xy,p}\underline{\underline{T}}^T_{pq}\right) = \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1}\underline{u}^T_{xy,p}\underline{\underline{T}}^T_{pq}\underline{F}_{xy,q} \tag{18}$$

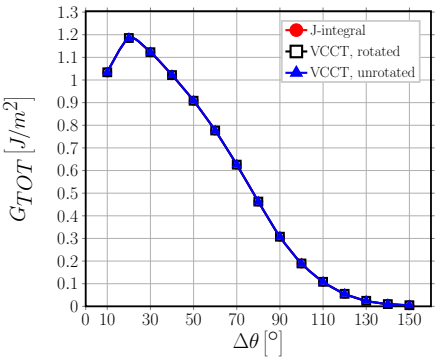
which shows that the total Energy Release Rate is invariant to rigid rotations and can be calculated equivalently with forces and displacements expressed in the local crack tip reference frame or the global reference frame. The analytical result is confirmed by the numerical solution of the fiber-matrix interface crack with different element orders and model fiber volume fractions, as shown in Figure 2.



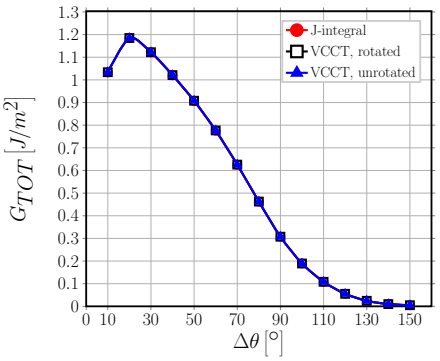
(a) $V_f = 0.1\%$, 1^{st} order elements, $\delta = 0.05^\circ$.



(b) $V_f = 0.1\%$, 2^{nd} order elements, $\delta = 0.05^\circ$.



(c) $V_f = 40\%$, 1^{st} order elements, $\delta = 0.05^\circ$.



(d) $V_f = 40\%$, 2^{nd} order elements, $\delta = 0.05^\circ$.

Figure 2: Numerical invariance of the total Energy Release Rate: G_{TOT} computed with the VCCT with rotated forces and displacements (label *rotated*), with the VCCT with forces and displacements in the global reference frame (label *unrotated*) and with J-integral method (label *J-integral*).

The result of Equation 18 has also a physical implication: given that the stress and displacement fields in the vicinity of the crack tip are the same, two cracks with different crack propagation directions are energetically equivalent with respect to the total Energy Release Rate. For two such cracks, given that laws of the type $G_{TOT} \geq G_c$ govern crack propagation, if G_c did not depend on mode ratio, crack orientation would not at all affect its growth.

5. Convergence analysis

5.1. Analytical considerations

Substituting Equations 8 in the derivative of Equation 13, we can investigate the dependency of Mode I and Mode II ERR with respect to the size δ of an element in the crack tip neighborhood through

$$\begin{aligned}\frac{\partial \underline{\underline{G}}}{\partial \delta} =& -\frac{1}{2R_f\delta^2}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\underline{\underline{K}}_{xy}\underline{u}_{xy}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) - \frac{1}{2R_f\delta^2}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\widetilde{\underline{\underline{K}}}_{N}\underline{u}_{N}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) + \\ &+\frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\underline{\underline{K}}_{xy}\underline{u}_{xy}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{D}}^T\underline{\underline{T}}_{pq}^T\right) + \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\widetilde{\underline{\underline{K}}}_{N}\underline{u}_{N}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{D}}^T\underline{\underline{T}}_{pq}^T\right) + \\ &+\frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{DQ}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\underline{\underline{K}}_{xy}\underline{u}_{xy}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) + \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{DQ}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\widetilde{\underline{\underline{K}}}_{N}\underline{u}_{N}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) + \\ &+\frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\underline{\underline{K}}_{xy}\frac{\partial \underline{u}_{xy}}{\partial\delta}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) + \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\widetilde{\underline{\underline{K}}}_{N}\frac{\partial \underline{u}_{N}}{\partial\delta}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) + \\ &+\frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\underline{\underline{K}}_{xy}\underline{u}_{xy}\frac{\partial \underline{u}_{xy}^T}{\partial\delta}\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) + \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\widetilde{\underline{\underline{K}}}_{N}\underline{u}_{N}\frac{\partial \underline{u}_{xy}^T}{\partial\delta}\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right);\end{aligned} \tag{19}$$

which, after refactoring, provides

$$\begin{aligned}\frac{\partial \underline{\underline{G}}}{\partial \delta} =& \frac{1}{\delta}\underline{\underline{G}} + \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\left(\underline{\underline{K}}_{xy}\underline{u}_{xy} + \widetilde{\underline{\underline{K}}}_{N}\underline{u}_{N}\right)\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{D}}^T\underline{\underline{T}}_{pq}^T\right) + \\ &+\frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{DQ}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\left(\underline{\underline{K}}_{xy}\underline{u}_{xy} + \widetilde{\underline{\underline{K}}}_{N}\underline{u}_{N}\right)\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) + \\ &+\frac{1}{R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\underline{\underline{K}}_{xy}\frac{\partial \underline{u}_{xy}}{\partial\delta}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) + \frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\widetilde{\underline{\underline{K}}}_{N}\frac{\partial \underline{u}_{N}}{\partial\delta}\underline{u}_{xy}^T\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right) + \\ &+\frac{1}{2R_f\delta}\sum_{p=1}^{m+1}\sum_{q=1}^{m+1} Diag\left(\underline{\underline{Q}}_{\delta}\underline{\underline{R}}_{\Delta\theta}\widetilde{\underline{\underline{K}}}_{N}\underline{u}_{N}\frac{\partial \underline{u}_{xy}^T}{\partial\delta}\underline{\underline{R}}_{\Delta\theta}^T\underline{\underline{P}}_{\delta}^T\underline{\underline{T}}_{pq}^T\right).\end{aligned} \tag{20}$$

Following the asymptotic analysis of [3, 1], in the case of an *open crack* the displacement in the crack tip neighborhood will have a functional form of the type

$$u\left(\delta\right) \sim \sqrt{\delta}\left(\sin,\cos\right)\left(\epsilon\log\delta\right) \quad \text{with} \quad \epsilon = \frac{1}{2\pi}\log\left(\frac{1-\beta}{1+\beta}\right) \tag{21}$$

and β is Dundurs' parameter introduced in Section 1. Application of Equation 21 to the terms on the right hand side of Eq. 20 provides:

$$\underline{u}_{xy}, \underline{u}_N \sim u\left(\delta\right) \sim \sqrt{\delta}\left(\sin, \cos\right)\left(\epsilon \log \delta\right) \xrightarrow{\delta \to 0} 0; \tag{22}$$

$$\underline{u}_{xy}\underline{u}_{xy}^T, \underline{u}_N\underline{u}_{xy}^T \sim u^2\left(\delta\right) \sim \delta\left(\sin^2, \cos^2, \sin\cdot\cos\right)\left(\epsilon \log \delta\right) \xrightarrow{\delta \to 0} 0; \tag{23}$$

$$\frac{\partial \underline{u}_{xy}}{\partial \delta}\underline{u}_{xy}^T, \frac{\partial \underline{u}_N}{\partial \delta}\underline{u}_{xy}^T \sim -\frac{1}{2}\left(\sin^2, \cos^2, \sin\cdot\cos\right)\left(\epsilon \log \delta\right) + \left(-\sin^2, \cos^2, \pm \sin\cdot\cos\right)\left(\epsilon \log \delta\right) \xrightarrow{\delta \to 0} finite; \tag{24}$$

$$\underline{G} \sim \frac{1}{\delta}\underline{u}_{xy}\underline{u}_{xy}^T \sim \frac{1}{\delta}u^2\left(\delta\right) \sim \left(\sin^2, \cos^2, \sin\cdot\cos\right)\left(\epsilon \log \delta\right) \xrightarrow{\delta \to 0} finite. \tag{25}$$

In Equations 22-25, the multiplication by a trigonometric function of the type $\left(\sin, \cos, \sin^2, \cos^2, \sin\cdot\cos\right)$ prevents the divergence of the asymptote. Recalling Eqs. 5 and 6, in the limit of $\delta \to 0$ the rotation matrices become equal to the identity matrix:

$$\underline{\underline{P}}_\delta, \underline{\underline{Q}}_\delta \xrightarrow{\delta \to 0} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}. \tag{26}$$

Applying the results of Equations 22-26 to Eq. 20, it can be shown that the derivative of $\underline{G}$ can be split in a factor that goes to 0 in the limit of $\delta \to 0$ and in a factor independent of δ:

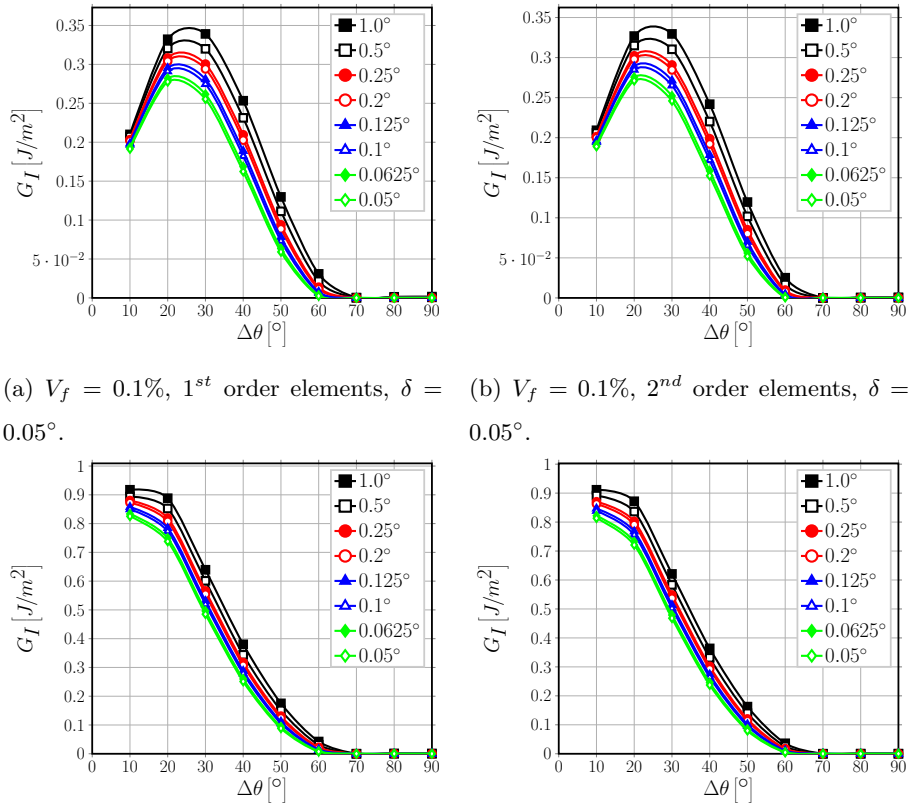
$$\lim_{\delta \to 0} \frac{\partial \underline{G}}{\partial \delta} \sim \frac{1}{\delta}\left(\cancelto{0}{\underline{F}\left(\delta\right)} + \underline{C}\right). \tag{27}$$

Thus, asymptotically, the Mode I and Mode II Energy Release Rate behave like the logarithm of the angular size δ of the elements in the crack tip neighborhood:

$$\lim_{\delta \to 0} \frac{\partial \underline{G}}{\partial \delta} \sim \frac{1}{\delta} \quad \xrightarrow{\int d\delta} \quad \lim_{\delta \to 0} \underline{G} \sim \underline{A}\log(\delta) + \underline{B}. \tag{28}$$

5.2. Numerical results

Evaluations of the Mode I, Mode II and total Energy Release Rate using the VCCT applied to the FE solution of the fiber-matrix interface crack in the single fiber model of Sec. 2 are reported respectively in Fig. 3, Fig. 4 and Fig. 5.



(a) $V_f = 0.1\%$, 1^{st} order elements, $\delta = 0.05°$.

(b) $V_f = 0.1\%$, 2^{nd} order elements, $\delta = 0.05°$.

(c) $V_f = 40\%$, 1^{st} order elements, $\delta = 0.05°$.

(d) $V_f = 40\%$, 2^{nd} order elements, $\delta = 0.05°$.

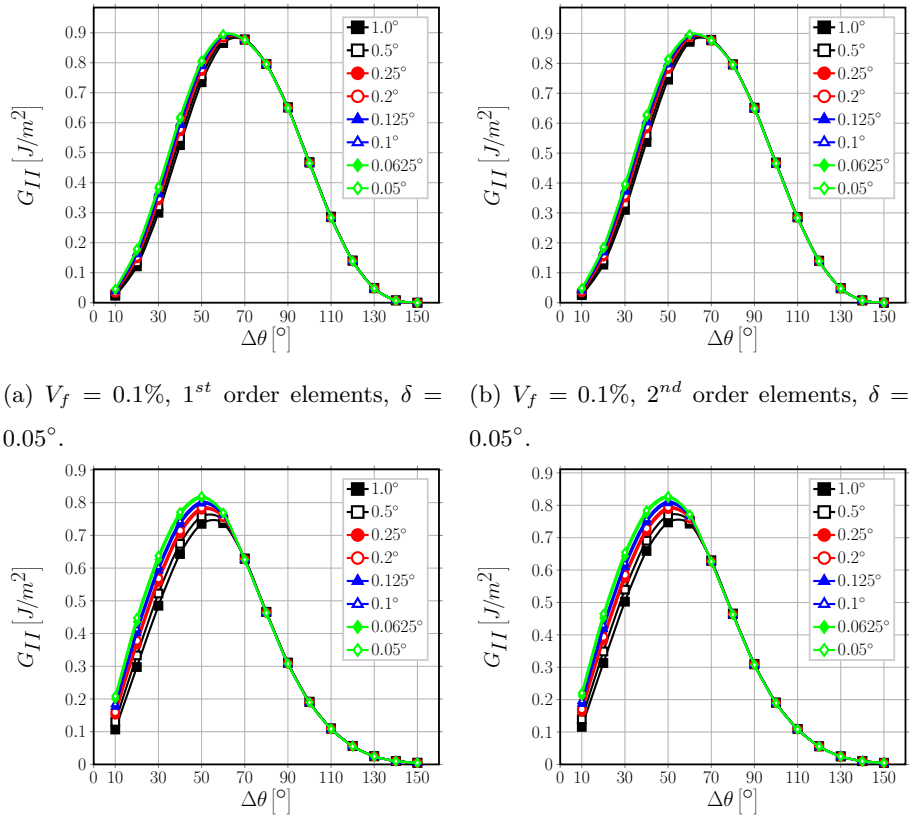
Figure 3: Effect of the size δ of an element at the crack tip on Mode I ERR.

Results for Mode I ERR in Fig. 3 show clearly the transition from the *open* crack regime, where Mode I ERR is different from zero, to the *closed* crack regime of the debond, where $G_I = 0$. Looking at Fig. 3, the crack is *open* for $\Delta\theta \leq 60°$ and it is *closed*, i.e. a contact zone is present, for $\Delta\theta \geq 70°$. As expected from the analysis of the previous section, and given that Mode I ERR

is different from zero only in the *open* crack regime, a significant dependence on the element size δ can be observed in Fig. 3 when using both 1^{st} and 2^{nd} order elements and with both an effectively infinite ($V_f = 0.1\%$) and finite size ($V_f = 40\%$) matrix. At first sight, it is immediate to see from Fig. 3 that a decrease in δ leads to a decrease in G_I. However, two further effects can be observed due to the refinement of the mesh at the crack tip, i.e. the decrease of the element size δ. First, the occurrence of the peak G_I is shifted to lower angles for very low volume fractions: it occurs at $\Delta\theta = 30^\circ$ with $\delta = 1.0^\circ, 0.5^\circ$ and at $\Delta\theta = 20^\circ$ with $\delta \leq 0.25^\circ$ for both 1^{st} and 2^{nd} order elements and $V_f = 0.1\%$. Second, the apperance of the contact zone, i.e. the switch to the *closed* crack regime, is anticipated to smaller debonds: it occurs at $\Delta\theta = 70^\circ$ with $\delta \geq 0.2^\circ$ and at $\Delta\theta = 60^\circ$ with $\delta < 0.2^\circ$ for both 1^{st} and 2^{nd} order elements and both $V_f = 0.1\%$ and $V_f = 40\%$.

Observing Figure 4, it possible to notice the existence of two distinct regimes in the behavior of G_{II} with respect to the element size δ. For $\Delta\theta \leq 60^\circ$ G_{II} depends on the value of δ, while $\Delta\theta \geq 70^\circ$ it is effectively independent of the element size at the crack tip for both 1^{st} and 2^{nd} order elements and both an effectively infinite ($V_f = 0.1\%$) and finite size ($V_f = 40\%$) matrix. Comparing the value of $\Delta\theta$ at which the change from the δ-dependency regime to the δ-independency regime occurs for G_{II} with Mode I ERR in Fig. 3, it is possible to observe that the δ-dependency regime change of Mode II ERR coincides with the onset of the contact zone, i.e. the transition from *open* crack regime to the *closed* crack regime. The result confirms the analytical considerations of the previous section: for an *open* crack both Mode I and Mode II ERR depend on the element size δ at the crack tip.

Further observation of Figure 4 reveals that, in the *open* crack regime, decreasing the element size δ causes an increase of Mode II ERR. Similarly to Mode I ERR, a shift of the peak G_{II} can also observed for $V_f = 0.1\%$: the maximum value of G_{II} occurs at $\Delta\theta = 70^\circ$ for $\delta > 0.25^\circ$ for 1^{st} order elements and for $\delta > 0.5^\circ$ for 2^{nd} order elements, while it is shifted to $\Delta\theta = 60^\circ$ for



(a) $V_f = 0.1\%$, 1^{st} order elements, $\delta = 0.05°$.

(b) $V_f = 0.1\%$, 2^{nd} order elements, $\delta = 0.05°$.

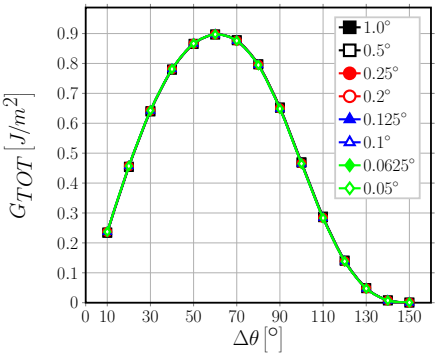
(c) $V_f = 40\%$, 1^{st} order elements, $\delta = 0.05°$.

(d) $V_f = 40\%$, 2^{nd} order elements, $\delta = 0.05°$.

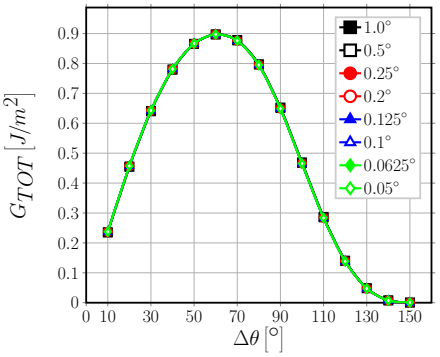
Figure 4: Effect of the size δ of an element at the crack tip on Mode II ERR.

$\delta \leq 0.25°$ for 1^{st} order elements and for $\delta \leq 0.5°$ for 2^{nd} order elements.

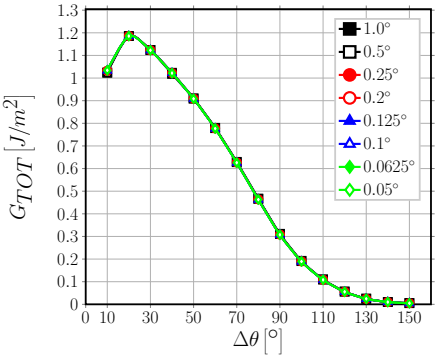
Analysis of the total ERR in Figure 5 leads to an observation that was not predicted by the considerations of the previous section: G_{TOT} is effectively independent of the element size δ in both the *open* and the *closed* crack regimes. Given that $G_{II} = G_{TOT}$ for the *closed* crack, it explains the independency of G_{II} from δ after the onset of the contact zone.



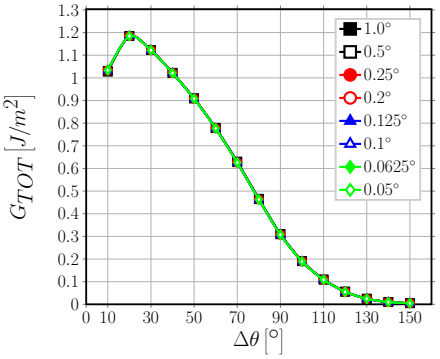
(a) $V_f = 0.1\%$, 1^{st} order elements, $\delta = 0.05°$.



(b) $V_f = 0.1\%$, 2^{nd} order elements, $\delta = 0.05°$.



(c) $V_f = 40\%$, 1^{st} order elements, $\delta = 0.05°$.



(d) $V_f = 40\%$, 2^{nd} order elements, $\delta = 0.05°$.

Figure 5: Effect of the size δ of an element at the crack tip on total ERR.

6. Conclusions & Outlook

Acknowledgements

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Table 2: Summary of linear regression results and main statistical tests for Mode I ERR

$\mathbf{V_f}$ [%]	Order	$\boldsymbol{\Delta\theta}$ [°]	$\mathbf{A}\left[\frac{\mathbf{J}}{\mathbf{m^2}}\right]$	$\mathbf{B}\left[\frac{\mathbf{J}}{\mathbf{m^2}}\right]$	$\mathbf{r}$ [−]	$\mathbf{r^2}$ [−]	$\mathbf{p(A)}$ [−]	$\mathbf{p(B)}$ [−]
0.1	1	10.0	0.0064	0.2113	0.9933	0.9866	7.48E-07	3.49E-14
		20.0	0.0183	0.3331	0.9996	0.9992	1.44E-10	2.40E-16
		30.0	0.0280	0.3392	1.0000	1.0000	2.25E-16	4.26E-21
		40.0	0.0304	0.2524	0.9997	0.9995	4.38E-11	7.94E-15
		50.0	0.0235	0.1278	0.9985	0.9970	8.61E-09	2.01E-11
		60.0	0.0094	0.0284	0.9854	0.9709	7.75E-06	6.14E-07
0.1	2	10.0	0.0069	0.2103	0.9962	0.9924	1.36E-07	1.03E-14
		20.0	0.0187	0.3277	0.9997	0.9994	7.85E-11	1.62E-16
		30.0	0.0280	0.3296	1.0000	1.0000	3.28E-16	7.29E-21
		40.0	0.0298	0.2408	0.9997	0.9995	4.82E-11	1.04E-14
		50.0	0.0225	0.1177	0.9984	0.9967	1.10E-08	3.27E-11
		60.0	0.0081	0.0228	0.9811	0.9626	1.66E-05	2.17E-06
40	1	10.0	0.0311	0.9196	0.9963	0.9927	1.03E-07	9.33E-15
		20.0	0.0501	0.8882	1.0000	0.9999	1.21E-13	2.33E-19
		30.0	0.0510	0.6374	0.9998	0.9996	1.66E-11	2.58E-16
		40.0	0.0419	0.3760	0.9988	0.9976	4.56E-09	5.25E-13
		50.0	0.0279	0.1713	0.9980	0.9961	2.22E-08	2.52E-11
		60.0	0.0108	0.0391	0.9901	0.9804	3.44E-06	9.46E-08
40	2	10.0	0.0336	0.9148	0.9988	0.9977	3.45E-09	5.09E-16
		20.0	0.0504	0.8719	1.0000	1.0000	3.70E-14	8.26E-20
		30.0	0.0506	0.6191	0.9999	0.9997	7.63E-12	1.35E-16
		40.0	0.0414	0.3608	0.9994	0.9989	4.95E-10	6.80E-14
		50.0	0.0269	0.1593	0.9982	0.9964	1.66E-08	2.31E-11
		60.0	0.0097	0.0329	0.9890	0.9781	4.96E-06	1.99E-07

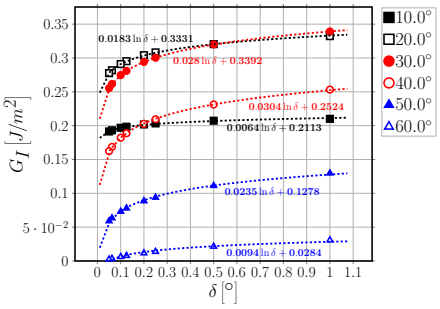
References

[1] M. Comninou, An overview of interface cracks, Engineering Fracture Mechanics 37 (1) (1990) 197–208. doi:10.1016/0013-7944(90)90343-f.

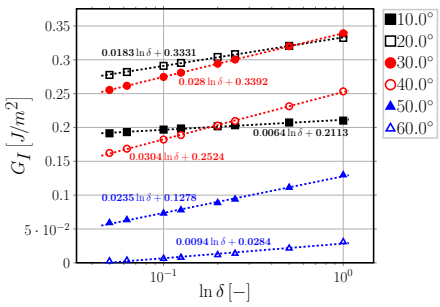
Table 3: Summary of linear regression results and main statistical tests for Mode II ERR

Order	$\mathbf{V_f}$ [%]	$\mathbf{\Delta\theta}$ [°]	$\mathbf{A}\left[\frac{\mathbf{J}}{\mathbf{m^2}}\right]$	$\mathbf{B}\left[\frac{\mathbf{J}}{\mathbf{m^2}}\right]$	**r** [−]	$\mathbf{r^2}$ [−]	**p(A)** [−]	**p(B)** [−]
0.1	1.0	10.0	-0.0076	0.0228	-0.9996	0.9991	2.09E-10	1.64E-11
		20.0	-0.0194	0.1211	-1.0000	1.0000	1.99E-15	2.02E-18
		30.0	-0.0290	0.3007	-0.9999	0.9998	4.12E-12	1.97E-16
		40.0	-0.0311	0.5270	-0.9995	0.9989	4.13E-10	1.05E-15
		50.0	-0.0240	0.7375	-0.9979	0.9958	2.32E-08	1.66E-15
		60.0	-0.0095	0.8685	-0.9835	0.9672	1.12E-05	1.22E-15
0.1	2.0	10.0	-0.0078	0.0249	-0.9996	0.9992	1.91E-10	1.06E-11
		20.0	-0.0196	0.1272	-1.0000	1.0000	3.48E-15	2.78E-18
		30.0	-0.0288	0.3108	-0.9999	0.9998	1.45E-12	5.47E-17
		40.0	-0.0305	0.5387	-0.9995	0.9990	3.32E-10	6.55E-16
		50.0	-0.0229	0.7478	-0.9979	0.9959	2.17E-08	1.09E-15
		60.0	-0.0082	0.8744	-0.9806	0.9615	1.81E-05	8.26E-16
40.0	1.0	10.0	-0.0344	0.1055	-0.9997	0.9995	3.82E-11	2.73E-12
		20.0	-0.0500	0.2977	-1.0000	0.9999	4.22E-14	5.66E-17
		30.0	-0.0505	0.4866	-0.9999	0.9997	6.44E-12	4.82E-16
		40.0	-0.0420	0.6454	-0.9996	0.9991	2.12E-10	9.66E-16
		50.0	-0.0275	0.7386	-0.9985	0.9971	9.01E-09	1.44E-15
		60.0	-0.0099	0.7402	-0.9926	0.9853	1.41E-06	5.13E-16
40.0	2.0	10.0	-0.0353	0.1145	-0.9998	0.9995	2.92E-11	1.50E-12
		20.0	-0.0504	0.3130	-1.0000	0.9999	4.00E-14	4.17E-17
		30.0	-0.0502	0.5039	-0.9999	0.9998	2.87E-12	1.69E-16
		40.0	-0.0410	0.6615	-0.9996	0.9992	2.02E-10	6.89E-16
		50.0	-0.0263	0.7502	-0.9987	0.9973	6.87E-09	7.76E-16
		60.0	-0.0090	0.7458	-0.9921	0.9842	1.79E-06	3.37E-16

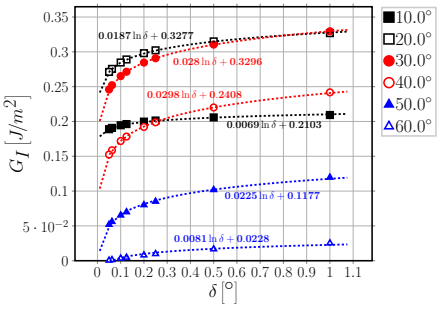
[2] D. Hills, J. Barber, Interface cracks, International Journal of Mechanical Sciences 35 (1) (1993) 27–37. doi:10.1016/0020-7403(93)90062-y.



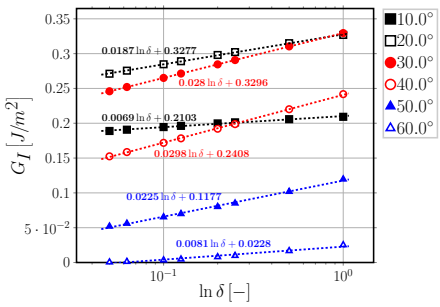
(a) 1^{st} order elements, $\delta = 0.05°$.



(b) 1^{st} order elements, $\delta = 0.05°$.



(c) 2^{nd} order elements, $\delta = 0.05°$.



(d) 2^{nd} order elements, $\delta = 0.05°$.

Figure 6: Logarithmic dependence on δ of Mode I ERR: interpolation of numerical results for $V_f = 0.1\%$.

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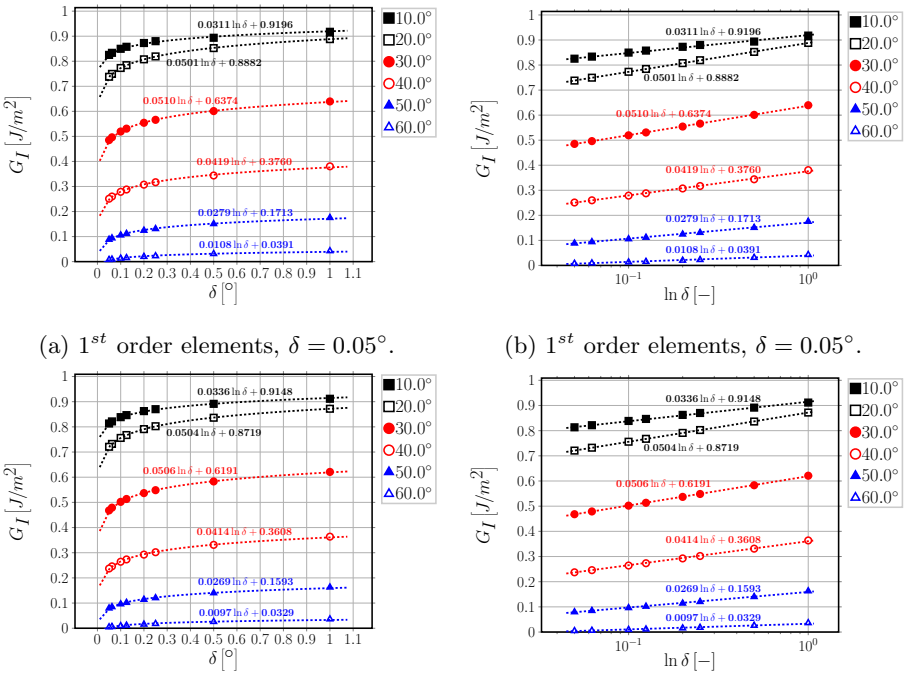

(a) 1^{st} order elements, $\delta = 0.05°$.

(b) 1^{st} order elements, $\delta = 0.05°$.

(c) 2^{nd} order elements, $\delta = 0.05°$.

(d) 2^{nd} order elements, $\delta = 0.05°$.

Figure 7: Logarithmic dependence on δ of Mode I ERR: interpolation of numerical results for $V_f = 40\%$.

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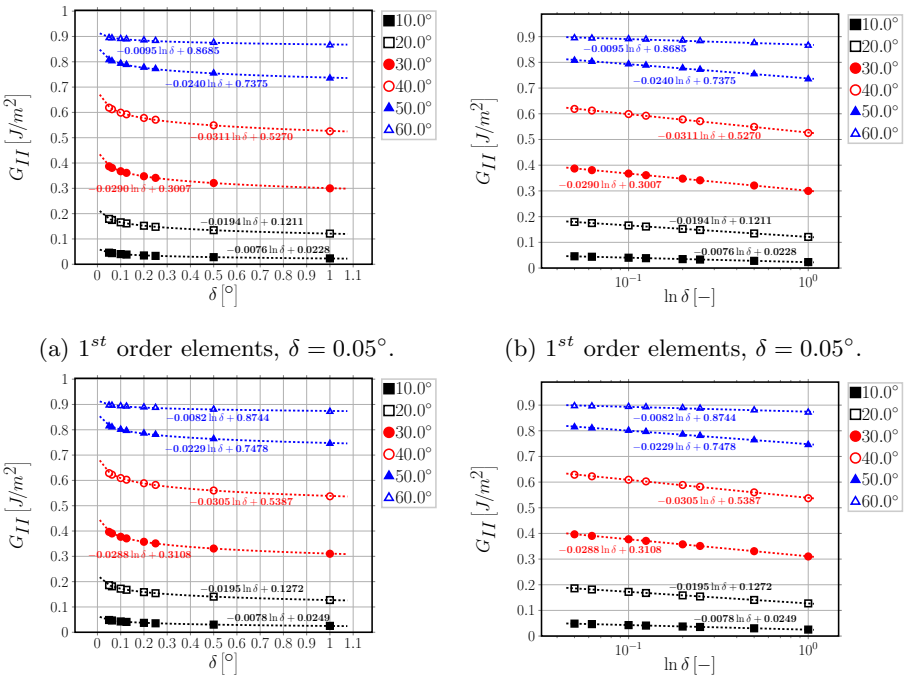

(a) 1^{st} order elements, $\delta = 0.05°$.

(b) 1^{st} order elements, $\delta = 0.05°$.

(c) 2^{nd} order elements, $\delta = 0.05°$.

(d) 2^{nd} order elements, $\delta = 0.05°$.

Figure 8: Logarithmic dependence on δ of Mode II ERR: interpolation of numerical results for $V_f = 0.1\%$.

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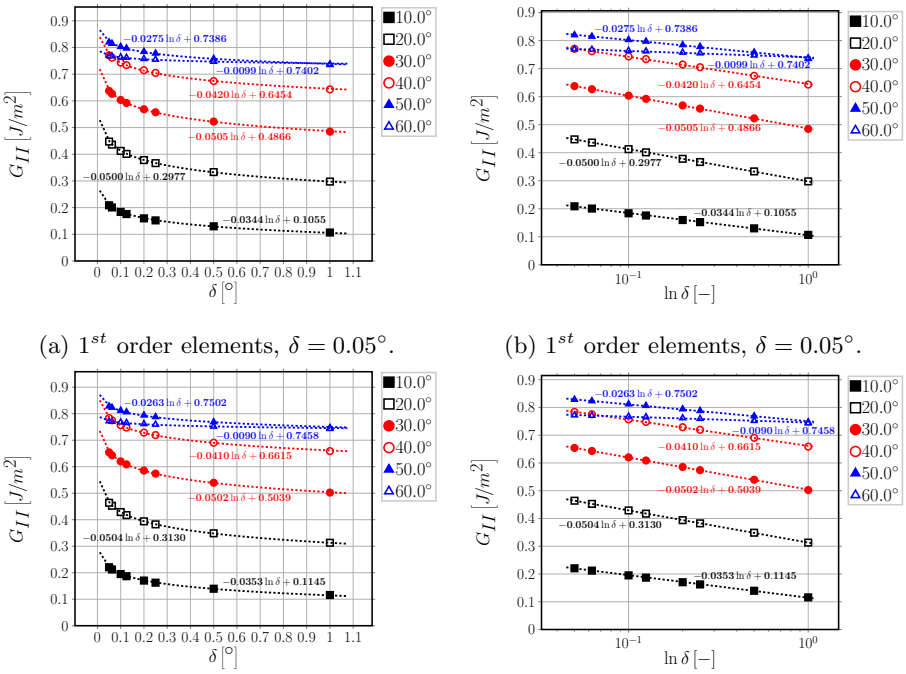

(a) 1^{st} order elements, $\delta = 0.05°$.

(b) 1^{st} order elements, $\delta = 0.05°$.

(c) 2^{nd} order elements, $\delta = 0.05°$.

(d) 2^{nd} order elements, $\delta = 0.05°$.

Figure 9: Logarithmic dependence on δ of Mode II ERR: interpolation of numerical results for $V_f = 40\%$.

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Appendix A. Derivation of the relationship between crack tip forces and displacements for first order quadrilateral elements

Appendix A.1. Foundational relations

In the isoparametric formulation of the Finite Element Method, the element Jacobian J and its inverse J^{-1} can be expressed in general as

$$\underline{\underline{J}} = \left[\underline{e}_\xi | \underline{e}_\eta\right] = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial x}{\partial \eta} \\ \frac{\partial y}{\partial \xi} & \frac{\partial y}{\partial \eta} \end{bmatrix} \qquad \underline{\underline{J}}^{-1} = \left[\underline{e}^x | \underline{e}^y\right] = \begin{bmatrix} \frac{\partial \xi}{\partial x} & \frac{\partial \xi}{\partial y} \\ \frac{\partial \eta}{\partial x} & \frac{\partial \eta}{\partial y} \end{bmatrix} \tag{A.1}$$

where $\{e_\xi, e_\eta\}$ and $\{e^x, e^y\}$ are respectively the covariant and contravariant basis vectors of the mapping between global $\{x, y\}$ and local element $\{\xi, \eta\}$ coordinates:

$$\underline{e}_\xi = \begin{bmatrix} \frac{\partial x}{\partial \xi} \\ \frac{\partial y}{\partial \xi} \end{bmatrix} \quad \underline{e}_\eta = \begin{bmatrix} \frac{\partial x}{\partial \eta} \\ \frac{\partial y}{\partial \eta} \end{bmatrix}, \tag{A.2}$$

$$\underline{e}_x = \begin{bmatrix} \frac{\partial \xi}{\partial x} \\ \frac{\partial \eta}{\partial x} \end{bmatrix} \quad \underline{e}_y = \begin{bmatrix} \frac{\partial \xi}{\partial y} \\ \frac{\partial \eta}{\partial y} \end{bmatrix}. \tag{A.3}$$

Denoting by d the number of geometrical dimensions of the problem ($d = 2$ in the present work) and by $\underline{p}$ the $d \times 1$ position vector in global coordinates, we can formally introduce the $3\,(d-1) \times d$ matrix operator of partial differentiation $\underline{\underline{\widetilde{B}}}$ such that

$$\underline{\varepsilon}\left(\underline{p}\right) = \underline{\underline{\widetilde{B}}} \cdot \underline{u}\left(\underline{p}\right), \tag{A.4}$$

where $\underline{u}$ and $\underline{\varepsilon}$ are respectively the $d\times1$ displacement vector and the $3\left(d-1\right)\times 1$ strain vector in Voigt notation. Denoting by n the number of nodes of a generic element ($n = s \times m$ where s represents the number of sides of the element and m the order of the shape functions), we can furthermore introduce the $d \times d \cdot n$ matrix $\underline{\underline{\mathrm{N}}}$ of shape functions such that

$$\underline{u} = \underline{\underline{N}} \cdot \underline{u}_N, \tag{A.5}$$

where $\underline{u}_N$ is the $d\cdot n\times1$ vector of element nodal variables. Having introduced $\underline{\underline{\widetilde{B}}}$ and $\underline{\underline{N}}$ in Equations A.4 and A.5 respectively, it is possible to define the $3\left(d-1\right) \times d \cdot n$ matrix $\underline{\underline{B}}$ of derivatives (with respect to global coordinates) of shape functions as

$$\underline{\underline{B}} = \underline{\underline{\widetilde{B}}} \cdot \underline{\underline{N}}. \tag{A.6}$$

We introduce the linear elastic material behavior in the form of the $3\left(d-1\right)\times 3\left(d-1\right)$ rigidity matrix $\underline{\underline{D}}$ such that

$$\underline{\sigma} = \underline{\underline{D}} \cdot \underline{\varepsilon}, \tag{A.7}$$

where $\underline{\sigma}$ the $3\left(d-1\right)\times1$ stress vector in Voigt notation. It is finally possible to define the $n \times n$ element stiffness matrix $\underline{\underline{k_e}}$ as

$$\underline{\underline{k_e}} = \int_{V_e(x,y)} \left(\underline{\underline{B}}^T \underline{\underline{D}} \cdot \underline{\underline{B}}\right) dV_e\left(x,\ldots,y\right) = \int_{V_e(\xi,\eta)} \left(\underline{\underline{B}}^T \underline{\underline{D}} \cdot \underline{\underline{B}}\right) \sqrt{g} dV_e\left(\xi,\ldots,\eta\right), \tag{A.8}$$

where $g = det\left(\underline{\underline{J}}^T \underline{\underline{J}}\right)$ and V_e is the element volume. Given that isoparametric elements are always defined between -1 and 1 in each dimension, Equation A.8 can simplified to

$$\underline{\underline{k_e}} = \int_{-1}^{1} \cdots \int_{-1}^{1} \left(\underline{\underline{B}}^T \underline{\underline{D}} \cdot \underline{\underline{B}}\right) \sqrt{g} d\xi,\ldots,d\eta, \tag{A.9}$$

which is amenable to numerical integration by means of a Gaussian quadrature of the form

$$\underline{\underline{k_e}} \approx \sum_{i=1}^{N} \cdots \sum_{j=1}^{N} w_i \dots w_j \left(\underline{\underline{B}}^T (\xi_i, \dots, \eta_j) \cdot \underline{\underline{D}} \cdot \underline{\underline{B}} (\xi_i, \dots, \eta_j) \sqrt{g} \right), \tag{A.10}$$

where $(\xi_i, \dots, \eta_j)$ are the coordinates of the N Gaussian quadrature points. The element stiffness matrix as evaluated in Eq. A.10 is in general a full symmetric (in the case of linear elasticity) matrix of the form

$$k_e = \begin{bmatrix} k_{e|11} & k_{e|12} & k_{e|13} & k_{e|14} & k_{e|15} & k_{e|16} & k_{e|17} & k_{e|18} \\ k_{e|12} & k_{e|22} & k_{e|23} & k_{e|24} & k_{e|25} & k_{e|26} & k_{e|27} & k_{e|28} \\ k_{e|13} & k_{e|23} & k_{e|33} & k_{e|34} & k_{e|35} & k_{e|36} & k_{e|37} & k_{e|38} \\ k_{e|14} & k_{e|24} & k_{e|34} & k_{e|44} & k_{e|45} & k_{e|46} & k_{e|47} & k_{e|48} \\ k_{e|15} & k_{e|25} & k_{e|35} & k_{e|45} & k_{e|55} & k_{e|56} & k_{e|57} & k_{e|58} \\ k_{e|16} & k_{e|26} & k_{e|36} & k_{e|46} & k_{e|56} & k_{e|66} & k_{e|67} & k_{e|68} \\ k_{e|17} & k_{e|27} & k_{e|37} & k_{e|47} & k_{e|57} & k_{e|67} & k_{e|77} & k_{e|78} \\ k_{e|18} & k_{e|28} & k_{e|38} & k_{e|48} & k_{e|58} & k_{e|68} & k_{e|78} & k_{e|88} \end{bmatrix}. \tag{A.11}$$

Appendix A.2. Calculation of displacements and reaction forces

$$u_x = u_{x,M} - u_{x,F} \qquad u_y = u_{y,M} - u_{y,F} \tag{A.12}$$

$$u_r = \cos(\Delta\theta)\, u_x + \sin(\Delta\theta)\, u_y \qquad u_\theta = -\sin(\Delta\theta)\, u_x + \cos(\Delta\theta)\, u_y \tag{A.13}$$

$$F_r = \cos(\Delta\theta)\, F_{x,CT} + \sin(\Delta\theta)\, F_{y,CT} \qquad F_\theta = -\sin(\Delta\theta)\, F_{x,CT} + \cos(\Delta\theta)\, F_{y,CT} \tag{A.14}$$

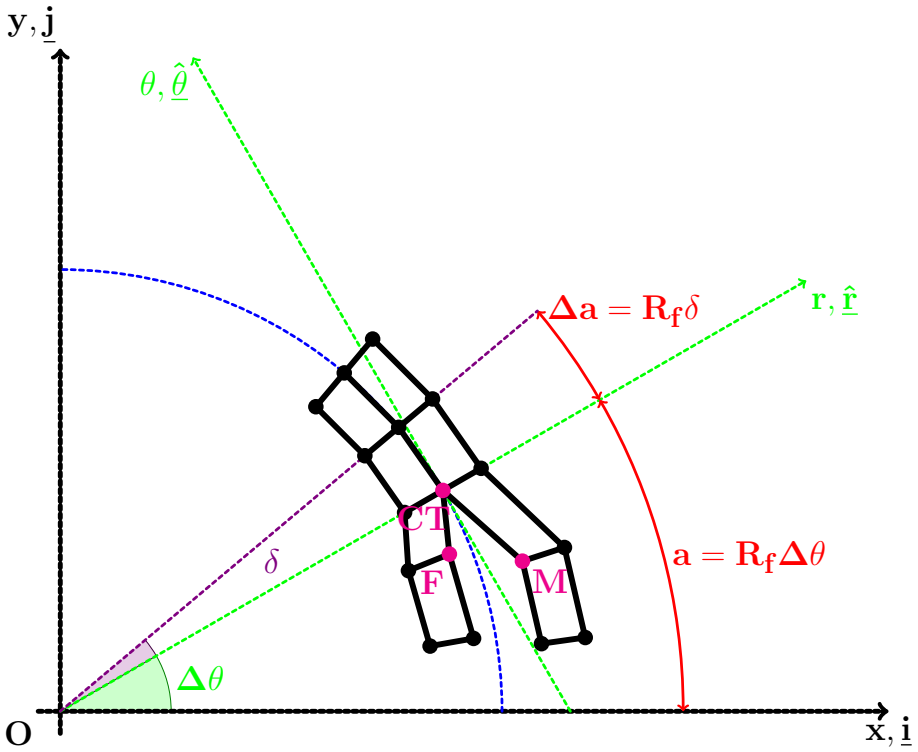

Figure A.10: Schematic representation of the discretized crack tip geometry for 1^{st} order quadrilateral elements.

$$
\begin{cases}
\left(k_{e,M|11}+k_{e,M|33}\right)u_{x,MCT}+\left(k_{e,M|12}+k_{e,M|34}\right)u_{y,MCT}+ \\
+k_{e,M|13}u_{x,M}+k_{e,M|14}u_{y,M}+\left(k_{M|17}+k_{M|35}\right)u_{N,MC|7}+\left(k_{M|18}+k_{M|36}\right)u_{N,MC|8}+ \\
+\sum_{i=5}^{6}k_{M|1i}u_{N,MC|i}+\sum_{i=7}^{8}k_{M|3i}u_{N,MB|i}+k_{M|31}u_{x,NCOI}+k_{M|32}u_{y,NCOI}=0 \\
\\
\left(k_{e,M|21}+k_{e,M|43}\right)u_{x,MCT}+\left(k_{e,M|22}+k_{e,M|44}\right)u_{y,MCT}+ \\
+k_{e,M|23}u_{x,M}+k_{e,M|24}u_{y,M}+\left(k_{M|27}+k_{M|45}\right)u_{N,MC|7}+\left(k_{M|28}+k_{M|46}\right)u_{N,MC|8}+ \\
+\sum_{i=5}^{6}k_{M|2i}u_{N,MC|i}+\sum_{i=7}^{8}k_{M|4i}u_{N,MB|i}+k_{M|41}u_{x,NCOI}+k_{M|42}u_{y,NCOI}=0 \\
\\
\left(k_{e,F|77}+k_{e,F|55}\right)u_{x,FCT}+\left(k_{e,F|78}+k_{e,F|56}\right)u_{y,FCT}+ \\
+k_{e,F|75}u_{x,F}+k_{e,F|76}u_{y,F}+\left(k_{F|71}+k_{F|53}\right)u_{N,FC|1}+\left(k_{F|72}+k_{F|54}\right)u_{N,FC|2}+ \\
+\sum_{i=2}^{3}k_{F|7i}u_{N,FC|i}+\sum_{i=1}^{2}k_{F|5i}u_{N,FB|i}+k_{F|57}u_{x,NCOI}+k_{F|58}u_{y,NCOI}=0 \\
\\
\left(k_{e,F|87}+k_{e,F|65}\right)u_{x,FCT}+\left(k_{e,F|88}+k_{e,F|66}\right)u_{y,FCT}+ \\
+k_{e,F|85}u_{x,F}+k_{e,F|86}u_{y,F}+\left(k_{F|81}+k_{F|63}\right)u_{N,FC|1}+\left(k_{F|82}+k_{F|64}\right)u_{N,FC|2}+ \\
+\sum_{i=2}^{3}k_{F|8i}u_{N,FC|i}+\sum_{i=1}^{2}k_{F|6i}u_{N,FB|i}+k_{F|67}u_{x,NCOI}+k_{F|68}u_{y,NCOI}=0 \\
\\
u_{x,FCT}-u_{x,MCT}-u_{x,DUMMY}=0
\end{cases}
$$

$$
\begin{cases}
\left(k_{e,M|11}+k_{e,M|33}\right)u_{x,MCT}+\left(k_{e,M|12}+k_{e,M|34}\right)u_{y,MCT}+ \\
+k_{e,M|13}u_{x,M}+k_{e,M|14}u_{y,M}+\left(k_{M|17}+k_{M|35}\right)u_{N,MC|7}+\left(k_{M|18}+k_{M|36}\right)u_{N,MC|8}+ \\
+\sum_{i=5}^{6}k_{M|1i}u_{N,MC|i}+\sum_{i=7}^{8}k_{M|3i}u_{N,MB|i}+k_{M|31}u_{x,NCOI}+k_{M|32}u_{y,NCOI}=0 \\
\\
\left(k_{e,M|21}+k_{e,M|43}\right)u_{x,MCT}+\left(k_{e,M|22}+k_{e,M|44}\right)u_{y,MCT}+ \\
+k_{e,M|23}u_{x,M}+k_{e,M|24}u_{y,M}+\left(k_{M|27}+k_{M|45}\right)u_{N,MC|7}+\left(k_{M|28}+k_{M|46}\right)u_{N,MC|8}+ \\
+\sum_{i=5}^{6}k_{M|2i}u_{N,MC|i}+\sum_{i=7}^{8}k_{M|4i}u_{N,MB|i}+k_{M|41}u_{x,NCOI}+k_{M|42}u_{y,NCOI}=0 \\
\\
\left(k_{e,F|77}+k_{e,F|55}\right)u_{x,MCT}+\left(k_{e,F|78}+k_{e,F|56}\right)u_{y,MCT}+ \\
+k_{e,F|75}u_{x,F}+k_{e,F|76}u_{y,F}+\left(k_{F|71}+k_{F|53}\right)u_{N,FC|1}+\left(k_{F|72}+k_{F|54}\right)u_{N,FC|2}+ \\
+\sum_{i=2}^{3}k_{F|7i}u_{N,FC|i}+\sum_{i=1}^{2}k_{F|5i}u_{N,FB|i}+k_{F|57}u_{x,NCOI}+k_{F|58}u_{y,NCOI}=0 \\
\\
\left(k_{e,F|87}+k_{e,F|65}\right)u_{x,MCT}+\left(k_{e,F|88}+k_{e,F|66}\right)u_{y,MCT}+ \\
+k_{e,F|85}u_{x,F}+k_{e,F|86}u_{y,F}+\left(k_{F|81}+k_{F|63}\right)u_{N,FC|1}+\left(k_{F|82}+k_{F|64}\right)u_{N,FC|2}+ \\
+\sum_{i=2}^{3}k_{F|8i}u_{N,FC|i}+\sum_{i=1}^{2}k_{F|6i}u_{N,FB|i}+k_{F|67}u_{x,NCOI}+k_{F|68}u_{y,NCOI}=0 \\
\\
u_{x,FCT}=u_{x,MCT} \\
u_{y,FCT}=u_{y,MCT} \\
\\
R_{x,DUMMY}=R_{x,FCT}=-R_{x,MCT}=F_{x,CT} \\
R_{y,DUMMY}=R_{y,FCT}=-R_{y,MCT}=F_{y,CT}
\end{cases}
\tag{A.16}
$$

$$
\begin{cases}
\left(k_{e,M|11}+k_{e,M|33}+k_{e,F|77}+k_{e,F|55}\right)u_{x,MCT}+\left(k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}\right)u_{y,MCT}+ \\
+k_{e,M|13}u_{x,M}+k_{e,M|14}u_{y,M}+k_{e,F|75}u_{x,F}+k_{e,F|76}u_{y,F}+ \\
+\left(k_{M|31}+k_{F|57}\right)u_{x,NCOI}+\left(k_{M|32}+k_{F|58}\right)u_{y,NCOI}+ \\
+\left(k_{M|17}+k_{M|35}\right)u_{N,MC|7}+\left(k_{M|18}+k_{M|36}\right)u_{N,MC|8}+\left(k_{F|71}+k_{F|53}\right)u_{N,FC|1}+\left(k_{F|72}+k_{F|54}\right)u_{N,FC|2}+ \\
+\sum_{i=5}^{6}k_{M|1i}u_{N,MC|i}+\sum_{i=7}^{8}k_{M|3i}u_{N,MB|i}+\sum_{i=2}^{3}k_{F|7i}u_{N,FC|i}+\sum_{i=1}^{2}k_{F|5i}u_{N,FB|i}=0 \\
\\
\left(k_{e,M|21}+k_{e,M|43}+k_{e,F|87}+k_{e,F|65}\right)u_{x,MCT}+\left(k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}\right)u_{y,MCT}+ \\
+k_{e,M|23}u_{x,M}+k_{e,M|24}u_{y,M}+k_{e,F|85}u_{x,F}+k_{e,F|86}u_{y,F}+ \\
+\left(k_{M|41}+k_{F|67}\right)u_{x,NCOI}+\left(k_{M|42}+k_{F|68}\right)u_{y,NCOI}+ \\
+\left(k_{M|27}+k_{M|45}\right)u_{N,MC|7}+\left(k_{M|28}+k_{M|46}\right)u_{N,MC|8}+\left(k_{F|81}+k_{F|63}\right)u_{N,FC|1}+\left(k_{F|82}+k_{F|64}\right)u_{N,FC|2}+ \\
+\sum_{i=2}^{3}k_{F|8i}u_{N,FC|i}+\sum_{i=1}^{2}k_{F|6i}u_{N,FB|i}+\sum_{i=5}^{6}k_{M|2i}u_{N,MC|i}+\sum_{i=7}^{8}k_{M|4i}u_{N,MB|i}=0 \\
\\
u_{x,FCT}=u_{x,MCT} \\
u_{y,FCT}=u_{y,MCT} \\
\\
R_{x,DUMMY}=R_{x,FCT}=-R_{x,MCT}=F_{x,CT} \\
R_{y,DUMMY}=R_{y,FCT}=-R_{y,MCT}=F_{y,CT}
\end{cases}
\tag{A.17}
$$

$$
\begin{cases}
u_{y,MCT} = -\frac{k_{e,M|11}+k_{e,M|33}+k_{e,F|77}+k_{e,F|55}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} u_{x,MCT} + \\
-\frac{k_{e,M|13} u_{x,M}+k_{e,M|14} u_{y,M}+k_{e,F|75} u_{x,F}+k_{e,F|76} u_{y,F}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} + \\
-\frac{\left(k_{M|31}+k_{F|57}\right) u_{x,NCOI}+\left(k_{M|32}+k_{F|58}\right) u_{y,NCOI}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} + \\
-\frac{\left(k_{M|17}+k_{M|35}\right) u_{N,MC|7}+\left(k_{M|18}+k_{M|36}\right) u_{N,MC|8}+\left(k_{F|71}+k_{F|53}\right) u_{N,FC|1}+\left(k_{F|72}+k_{F|54}\right) u_{N,FC|2}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} + \\
-\frac{\sum_{i=5}^{6} k_{M|1i} u_{N,MC|i}+\sum_{i=7}^{8} k_{M|3i} u_{N,MB|i}+\sum_{i=2}^{3} k_{F|7i} u_{N,FC|i}+\sum_{i=1}^{2} k_{F|5i} u_{N,FB|i}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} \\
\\
\left[\left(k_{e,M|21}+k_{e,M|43}+k_{e,F|87}+k_{e,F|65}\right)+\frac{k_{e,M|11}+k_{e,M|33}+k_{e,F|77}+k_{e,F|55}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left(k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}\right)\right] u_{x,MCT} + \\
+\left(k_{e,M|23}-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} k_{e,M|13}\right) u_{x,M} + \\
+\left(k_{e,M|24}-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} k_{e,M|14}\right) u_{y,M} + \\
+\left(k_{e,F|85}-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} k_{e,M|75}\right) u_{x,F} + \\
+\left(k_{e,F|86}-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} k_{e,M|76}\right) u_{y,F} + \\
+\left[\left(k_{M|41}+k_{F|67}\right)-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left(k_{M|31}+k_{F|57}\right)\right] u_{x,NCOI} + \\
+\left[\left(k_{M|42}+k_{F|68}\right)-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left(k_{M|32}+k_{F|58}\right)\right] u_{y,NCOI} + \\
+\left(k_{M|27}+k_{M|45}\right) u_{N,MC|7}+\left(k_{M|28}+k_{M|46}\right) u_{N,MC|8}+\left(k_{F|81}+k_{F|63}\right) u_{N,FC|1}+\left(k_{F|82}+k_{F|64}\right) u_{N,FC|2} + \\
-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left[\left(k_{M|17}+k_{M|35}\right) u_{N,MC|7}+\left(k_{M|18}+k_{M|36}\right) u_{N,MC|8}\right] + \\
-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left[\left(k_{F|71}+k_{F|53}\right) u_{N,FC|1}+\left(k_{F|72}+k_{F|54}\right) u_{N,FC|2}\right] \\
+\sum_{i=2}^{3} k_{F|8i} u_{N,FC|i}+\sum_{i=1}^{2} k_{F|6i} u_{N,FB|i}+\sum_{i=5}^{6} k_{M|2i} u_{N,MC|i}+\sum_{i=7}^{8} k_{M|4i} u_{N,MB|i} + \\
-\frac{\sum_{i=5}^{6} k_{M|1i} u_{N,MC|i}+\sum_{i=7}^{8} k_{M|3i} u_{N,MB|i}+\sum_{i=2}^{3} k_{F|7i} u_{N,FC|i}+\sum_{i=1}^{2} k_{F|5i} u_{N,FB|i}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} = 0 \\
\\
u_{x,FCT} = u_{x,MCT} \\
u_{y,FCT} = u_{y,MCT} \\
\\
R_{x,DUMMY} = R_{x,FCT} = -R_{x,MCT} = F_{x,CT} \\
R_{y,DUMMY} = R_{y,FCT} = -R_{y,MCT} = F_{y,CT}
\end{cases}
\tag{A.18}
$$

$$
\begin{cases}
u_{y,MCT} = -\frac{k_{e,M|11}+k_{e,M|33}+k_{e,F|77}+k_{e,F|55}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} u_{x,MCT} + \\
-\frac{k_{e,M|13} u_{x,M}+k_{e,M|14} u_{y,M}+k_{e,F|75} u_{x,F}+k_{e,F|76} u_{y,F}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} + \\
-\frac{\left(k_{M|31}+k_{F|57}\right) u_{x,NCOI}+\left(k_{M|32}+k_{F|58}\right) u_{y,NCOI}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} + \\
-\frac{\left(k_{M|17}+k_{M|35}\right) u_{N,MC|7}+\left(k_{M|18}+k_{M|36}\right) u_{N,MC|8}+\left(k_{F|71}+k_{F|53}\right) u_{N,FC|1}+\left(k_{F|72}+k_{F|54}\right) u_{N,FC|2}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} + \\
-\frac{\sum_{i=5}^{6} k_{M|1i} u_{N,MC|i}+\sum_{i=7}^{8} k_{M|3i} u_{N,MB|i}+\sum_{i=2}^{3} k_{F|7i} u_{N,FC|i}+\sum_{i=1}^{2} k_{F|5i} u_{N,FB|i}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} \\
\\
\left[\left(k_{e,M|21}+k_{e,M|43}+k_{e,F|87}+k_{e,F|65}\right)+\frac{k_{e,M|11}+k_{e,M|33}+k_{e,F|77}+k_{e,F|55}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left(k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}\right)\right] u_{x,MC} \\
+\left(k_{e,M|23}-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} k_{e,M|13}\right) u_x + \\
+\left(k_{e,M|24}-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} k_{e,M|14}\right) u_y + \\
+\left(k_{e,M|23}+k_{e,F|85}-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left(k_{e,M|13}+k_{e,M|75}\right)\right) \cancelto{\approx 0}{u_{x,F}} + \\
+\left(k_{e,M|24}+k_{e,F|86}-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left(k_{e,M|14}+k_{e,M|76}\right)\right) \cancelto{\approx 0}{u_{y,F}} + \\
+\left[\left(k_{M|41}+k_{F|67}\right)-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left(k_{M|31}+k_{F|57}\right)\right] u_{x,NCOI} + \\
+\left[\left(k_{M|42}+k_{F|68}\right)-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left(k_{M|32}+k_{F|58}\right)\right] u_{y,NCOI} + \\
+\left(k_{M|27}+k_{M|45}\right) u_{N,MC|7}+\left(k_{M|28}+k_{M|46}\right) u_{N,MC|8}+\left(k_{F|81}+k_{F|63}\right) u_{N,FC|1}+\left(k_{F|82}+k_{F|64}\right) u_{N,FC|2} + \\
-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left[\left(k_{M|17}+k_{M|35}\right) u_{N,MC|7}+\left(k_{M|18}+k_{M|36}\right) u_{N,MC|8}\right] + \\
-\frac{k_{e,M|22}+k_{e,M|44}+k_{e,F|88}+k_{e,F|66}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}}\left[\left(k_{F|71}+k_{F|53}\right) u_{N,FC|1}+\left(k_{F|72}+k_{F|54}\right) u_{N,FC|2}\right] \\
+\sum_{i=2}^{3} k_{F|8i} u_{N,FC|i}+\sum_{i=1}^{2} k_{F|6i} u_{N,FB|i}+\sum_{i=5}^{6} k_{M|2i} u_{N,MC|i}+\sum_{i=7}^{8} k_{M|4i} u_{N,MB|i} + \\
-\frac{\sum_{i=5}^{6} k_{M|1i} u_{N,MC|i}+\sum_{i=7}^{8} k_{M|3i} u_{N,MB|i}+\sum_{i=2}^{3} k_{F|7i} u_{N,FC|i}+\sum_{i=1}^{2} k_{F|5i} u_{N,FB|i}}{k_{e,M|12}+k_{e,M|34}+k_{e,F|78}+k_{e,F|56}} = 0 \\
\\
u_{x,FCT} = u_{x,MCT} \\
u_{y,FCT} = u_{y,MCT} \\
\\
F_{x,CT} = R_{x,FCT} = \\
= \left(k_{e,F|77}+k_{e,F|55}\right) u_{x,FCT}+\left(k_{e,F|78}+k_{e,F|56}\right) u_{y,FCT} + \\
+k_{e,F|75} \cancelto{\approx 0}{u_{x,F}} + k_{e,F|76} \cancelto{\approx 0}{u_{y,F}} + \\
+\sum_{i=1}^{4} k_{e,F|7i} u_{N,FC|i}+\sum_{i=1,i\neq(5,6)}^{8} k_{e,F|5i} u_{N,FB|i} \\
F_{y,CT} = R_{y,FCT} = \\
= \left(k_{e,F|87}+k_{e,F|65}\right) u_{x,FCT}+\left(k_{e,F|88}+k_{e,F|66}\right) u_{y,FCT} + \\
+k_{e,F|85} \cancelto{\approx 0}{u_{x,F}} + k_{e,F|86} \cancelto{\approx 0}{u_{y,F}} + \\
+\sum_{i=1}^{4} k_{e,F|8i} u_{N,FC|i}+\sum_{i=1,i\neq(5,6)}^{8} k_{e,F|6i} u_{N,FB|i}
\end{cases}
\tag{A.19}
$$

$$\begin{cases} F_{x,CT} & = \left(k_{e,F|77} + k_{e,F|55}\right) u_{x,MCT} + \left(k_{e,F|78} + k_{e,F|56}\right) u_{y,MCT} + \\ & + \sum_{i=1}^{4} k_{e,F|7i} u_{N,FC|i} + \sum_{i=1,i\neq(5,6)}^{8} k_{e,F|5i} u_{N,FB|i} \\ F_{y,CT} & = \left(k_{e,F|87} + k_{e,F|65}\right) u_{x,MCT} + \left(k_{e,F|88} + k_{e,F|66}\right) u_{y,MCT} + \\ & + \sum_{i=1}^{4} k_{e,F|8i} u_{N,FC|i} + \sum_{i=1,i\neq(5,6)}^{8} k_{e,F|6i} u_{N,FB|i} \end{cases} \tag{A.20}$$

$$\begin{cases} F_{x,CT} & = K_{xx} u_x + K_{xy} u_y + \\ & + \sum_{i=1}^{4} K_{FC,x|i} u_{N,FC|i} + \sum_{i=1,i\neq(3,4,5,6)}^{8} K_{FB,x|i} u_{N,FB|i} + \\ & + \sum_{i=5}^{8} K_{FC,x|i} u_{N,MC|i} + \sum_{i=7}^{8} K_{MB,x|i} u_{N,FB|i} \\ F_{y,CT} & = K_{yx} u_x + K_{yy} u_y + \\ & + \sum_{i=1}^{4} K_{FC,y|i} u_{N,FC|i} + \sum_{i=1,i\neq(3,4,5,6)}^{8} K_{FB,y|i} u_{N,FB|i} + \\ & + \sum_{i=5}^{8} K_{FC,y|i} u_{N,MC|i} + \sum_{i=7}^{8} K_{MB,y|i} u_{N,FB|i} \end{cases} \tag{A.21}$$

$$\begin{cases} F_{x,CT} & = K_{xx} u_x + K_{xy} u_y + \widetilde{F}_x \\ F_{y,CT} & = K_{yx} u_x + K_{yy} u_y + \widetilde{F}_y \end{cases} \tag{A.22}$$

Appendix B. Expression of $\underline{\underline{T}}_{pq}$ for quadrilateral elements with or without singularity

The expression of $\underline{\underline{T}}_{pq}$ for quadrilateral elements with or without singularity is

$$
\underline{\underline{T}}_{pq} = \begin{cases} \underline{\underline{I}} \; for \; p = q < 2 \\ \underline{\underline{0}} \; otherwise \end{cases} \quad \text{for } 1^{st} \text{ order quadrilateral elements}
$$

$$
= \begin{cases} \underline{\underline{I}} \; for \; p = q < 3 \\ \underline{\underline{0}} \; otherwise \end{cases} \quad \text{for } 2^{nd} \text{ order quadrilateral elements}
$$

$$
= \begin{cases} \underline{\underline{I}} \; for \; p = q < 4 \\ \underline{\underline{0}} \; otherwise \end{cases} \quad \text{for } 3^{rd} \text{ order quadrilateral elements}
$$

$$
= \begin{cases} \left(14 - \frac{33\pi}{8}\right) \underline{\underline{I}} \; for \; p = 1, q = 1 \\ \left(-52 + \frac{33\pi}{2}\right) \underline{\underline{I}} \; for \; p = 1, q = 2 \\ \left(17 - \frac{21\pi}{4}\right) \underline{\underline{I}} \; for \; p = 2, q = 1 \\ \left(-\frac{7}{2} + \frac{21\pi}{16}\right) \underline{\underline{I}} \; for \; p = 2, q = 2 \\ \left(8 - \frac{21\pi}{8}\right) \underline{\underline{I}} \; for \; p = 1, q = 3 \\ \left(-32 + \frac{21\pi}{2}\right) \underline{\underline{I}} \; for \; p = 2, q = 3 \\ \underline{\underline{0}} \; otherwise \end{cases} \quad \text{for } 2^{nd} \text{ order quarter-point quadrilateral elements}
$$

$$
= \begin{cases} \left(-11187 + \frac{7155\pi}{2}\right) \underline{\underline{I}} \; for \; p = 1, q = 1 \\ \left(38556 - \frac{24543\pi}{2}\right) \underline{\underline{I}} \; for \; p = 1, q = 2 \\ \left(-53055 + \frac{33777\pi}{2}\right) \underline{\underline{I}} \; for \; p = 1, q = 3 \\ \left(\frac{11396}{3} - \frac{9575\pi}{8}\right) \underline{\underline{I}} \; for \; p = 2, q = 1 \\ \left(-12936 + \frac{33003\pi}{8}\right) \underline{\underline{I}} \; for \; p = 2, q = 2 \\ \left(17988 - \frac{45837\pi}{8}\right) \underline{\underline{I}} \; for \; p = 2, q = 3 \\ \left(-\frac{8453}{3} + \frac{3595\pi}{4}\right) \underline{\underline{I}} \; for \; p = 3, q = 1 \\ \left(9804 - \frac{12411\pi}{4}\right) \underline{\underline{I}} \; for \; p = 3, q = 2 \\ \left(-13587 + \frac{17289\pi}{4}\right) \underline{\underline{I}} \; for \; p = 3, q = 3 \\ \left(6948 - \frac{17685\pi}{8}\right) \underline{\underline{I}} \; for \; p = 1, q = 4 \\ \left(-23976 + \frac{60993\pi}{8}\right) \underline{\underline{I}} \; for \; p = 2, q = 4 \\ \left(33372 - \frac{84807\pi}{8}\right) \underline{\underline{I}} \; for \; p = 3, q = 4 \\ \underline{\underline{0}} \; otherwise \end{cases} \quad \text{for } 3^{rd} \text{ order quarter-point quadrilateral elements}
$$

(B.1)

where $\underline{\underline{I}}$ is the identity matrix.