

A set of criteria for the prediction of initiation and propagation of transverse cracks

Luca Di Stasio^{a,b}, Janis Varna^b, Zoubir Ayadi^a

^a *Université de Lorraine, EEIGM, IJL, 6 Rue Bastien Lepage, F-54010 Nancy, France*

^b *Luleå University of Technology, University Campus, SE-97187 Luleå, Sweden*

Abstract

1. Normalization function

$$G_0 = G_0(\varepsilon_0, V_f, E_{1f}, E_{2f}, E_m, \nu_{12f}, \nu_{23f}, \nu_m, G_{12f}, G_{23f}) \quad (1)$$

Given the elastic properties of the transversely isotropic UD ply $E_1, E_2, \nu_{12}, \nu_{23}$, for a 90° ply under transverse tension the cross section along the direction of the load coincides with the plane of transversal isotropy. It is thus possible, for
5 a system in plane strain, to define equivalent isotropic Young's modulus and Poisson's ratio as follows. The effective Young's modulus and Poisson's ratio in plane strain in the plane of isotropy are defined as

$$E^* = \frac{E_2}{1 - \nu_{21}\nu_{12}} \quad \nu^* = \frac{\nu_{23} + \nu_{21}\nu_{12}}{1 + \nu_{23}} \quad (2)$$

$$G_0 = \frac{\sigma_0^2}{E^*} \pi R_f \quad \text{for a stress or force controlled test} \quad (3)$$

$$G_0 = E^* \varepsilon_0^2 \pi R_f \quad \text{for a strain or displacement controlled test}$$

2. Boundary conditions

The ratio of maximum radial and tangential crack displacements with respect
10 to the free case (single repeating element or single fiber layer ply?) can be considered as proxies for the effect of boundary conditions

$$\frac{u_{r,max}^{BC}}{u_{r,max}^{free}}, \frac{u_{\theta,max}^{BC}}{u_{\theta,max}^{free}} \quad (4)$$

3. Initiation of fiber-matrix debonds

4. Propagation of fiber-matrix debonds

$$\frac{G_I}{G_0} = \begin{cases} A_\delta (V_f) \log(\delta) + A_{\Delta\theta} \left(V_f, \frac{u_{r,max}^{BC}}{u_{r,max}^{free}}, \frac{u_{\theta,max}^{BC}}{u_{\theta,max}^{free}} \right) \sin(B_{\Delta\theta} \Delta\theta + C_{\Delta\theta}) + D \\ B_{\Delta\theta} \Delta\theta_{max} \left(V_f, \frac{u_{r,max}^{BC}}{u_{r,max}^{free}}, \frac{u_{\theta,max}^{BC}}{u_{\theta,max}^{free}} \right) + C_{\Delta\theta} = \frac{\pi}{2} \\ B_{\Delta\theta} \Delta\theta_{CZ} \left(\frac{u_{r,max}^{BC}}{u_{r,max}^{free}}, \frac{u_{\theta,max}^{BC}}{u_{\theta,max}^{free}} \right) + C_{\Delta\theta} = \pi \\ \text{for } \Delta\theta < \Delta\theta_{CZ} \left(\frac{u_{r,max}^{BC}}{u_{r,max}^{free}}, \frac{u_{\theta,max}^{BC}}{u_{\theta,max}^{free}} \right) \\ 0 & \text{otherwise} \end{cases} \quad (5)$$

$$\frac{G_{II}}{G_0} = \begin{cases} E_\delta (V_f) \log(\delta) + F_{\Delta\theta} \left(V_f, \frac{u_{r,max}^{BC}}{u_{r,max}^{free}}, \frac{u_{\theta,max}^{BC}}{u_{\theta,max}^{free}} \right) \sin(G_{\Delta\theta} \Delta\theta + H_{\Delta\theta}) + \\ + I_{\Delta\theta} \left(V_f, \frac{u_{r,max}^{BC}}{u_{r,max}^{free}}, \frac{u_{\theta,max}^{BC}}{u_{\theta,max}^{free}} \right) \sin(2G_{\Delta\theta} \Delta\theta + H_{\Delta\theta}) + L \\ G_{\Delta\theta} \Delta\theta_{max} \left(V_f, \frac{u_{r,max}^{BC}}{u_{r,max}^{free}}, \frac{u_{\theta,max}^{BC}}{u_{\theta,max}^{free}} \right) + H_{\Delta\theta} = \frac{\pi}{2} \\ G_{\Delta\theta} \Delta\theta_{CZ} \left(\frac{u_{r,max}^{BC}}{u_{r,max}^{free}}, \frac{u_{\theta,max}^{BC}}{u_{\theta,max}^{free}} \right) + H_{\Delta\theta} = \pi \end{cases} \quad (6)$$

5. Transition to collective mesoscopic behavior