ESTIMATING THE AVERAGE SIZE OF FIBER/MATRIX INTERFACE CRACKS IN UD AND CROSS-PLY LAMINATES

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Education and Culture

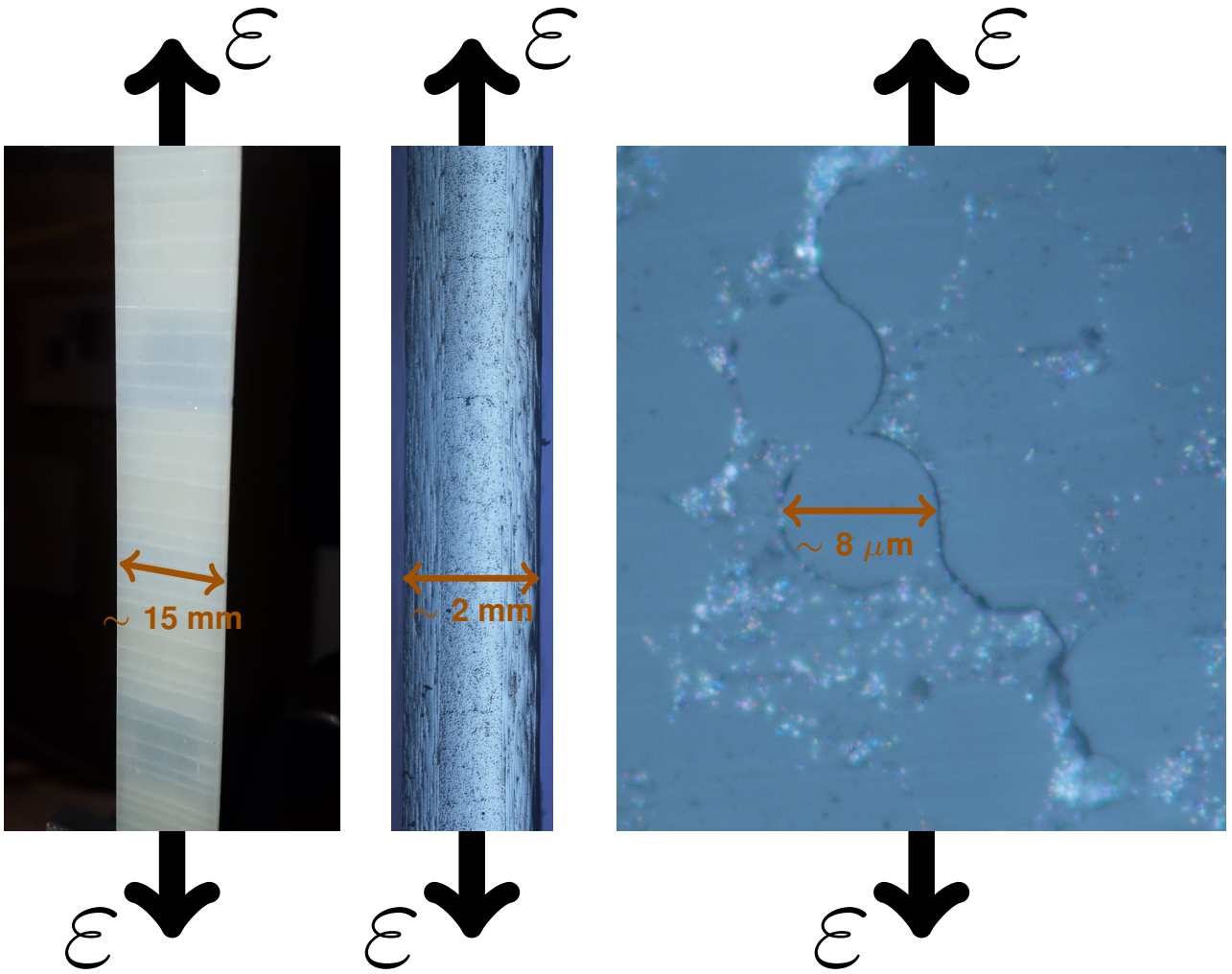
Erasmus Mundus

Outline

- Transverse Cracks Initiation in FRPC
- Modeling
- Debond Initiation
- Debond Propagation
- Conclusions

↘ TRANSVERSE CRACKS INITIATION IN FRPC

Micromechanics of Initiation: Transverse Tensile Loading



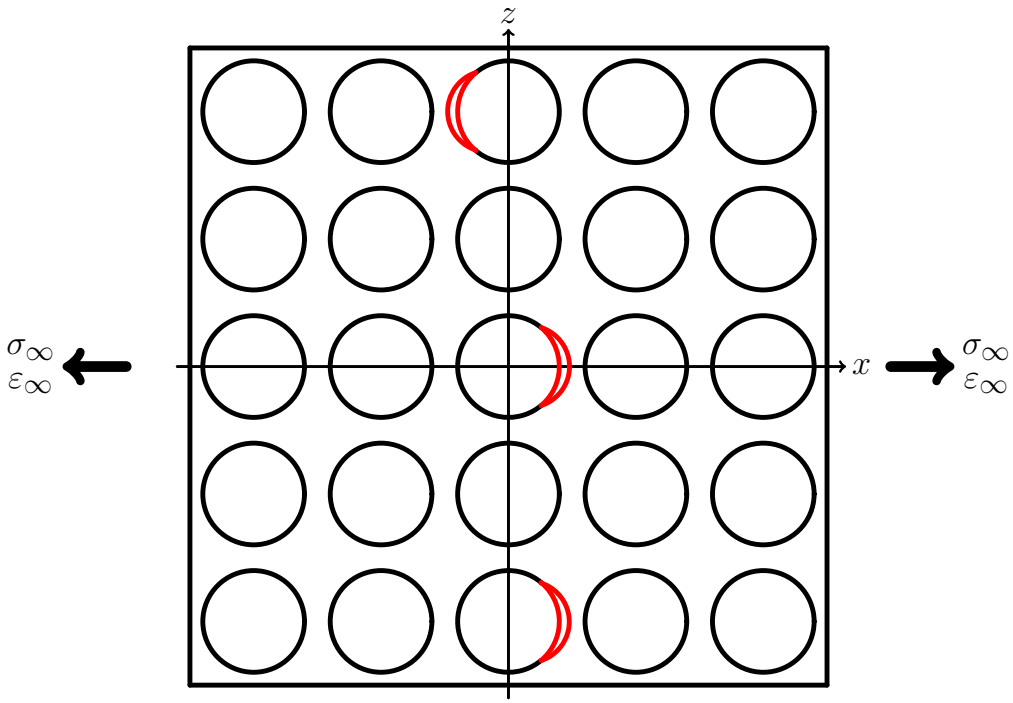
Left:
front view of $[0, 90_2]_S$, visual inspection.

Center:
edge view of $[0, 90]_S$, optical microscope.

Right:
edge view of $[0, 90]_S$, optical microscope.

Micromechanics of Initiation: Transverse Tensile Loading

Stage 1: isolated debonds



Bailey et al., P. Roy. Soc. A-Math. Phy. **366** (1727), 1979.

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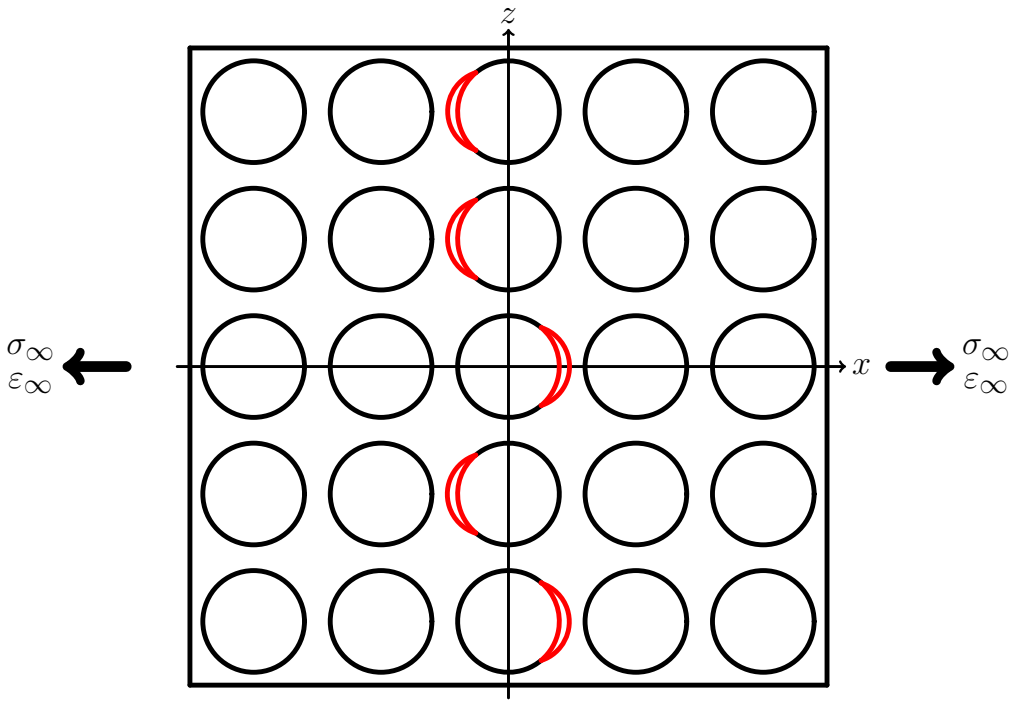
Zhang et al., Compos. Part A-Appl. S. **28** (4), 1997.

Micromechanics of Initiation: Transverse Tensile Loading

Stage 2: consecutive debonds



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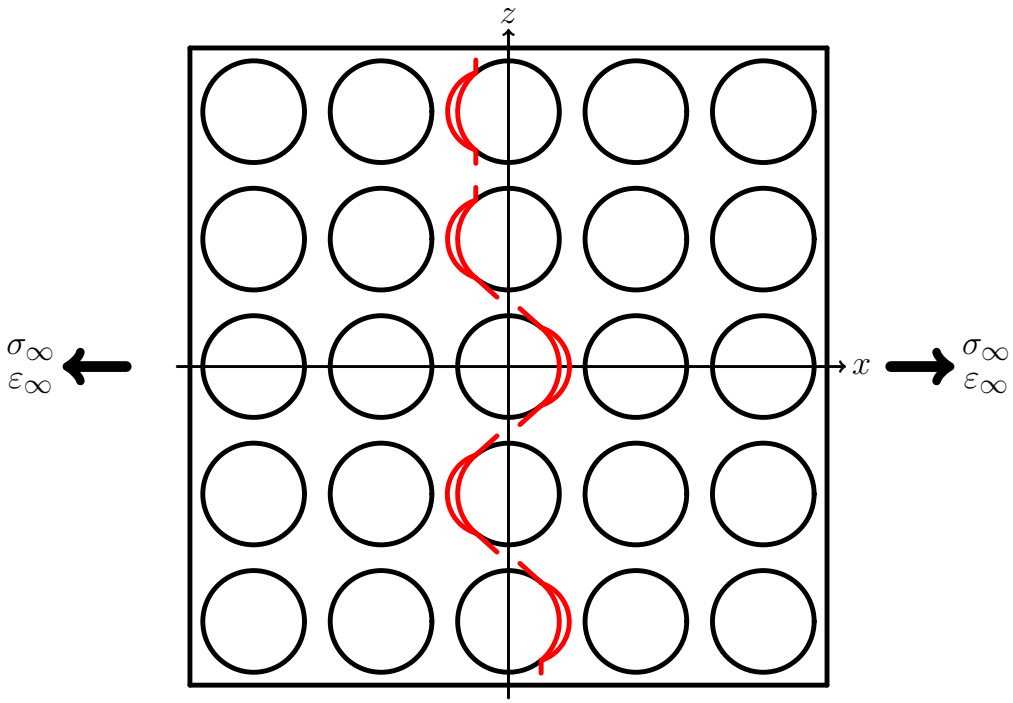
Micromechanics of Initiation: Transverse Tensile Loading

Stage 3: kinking



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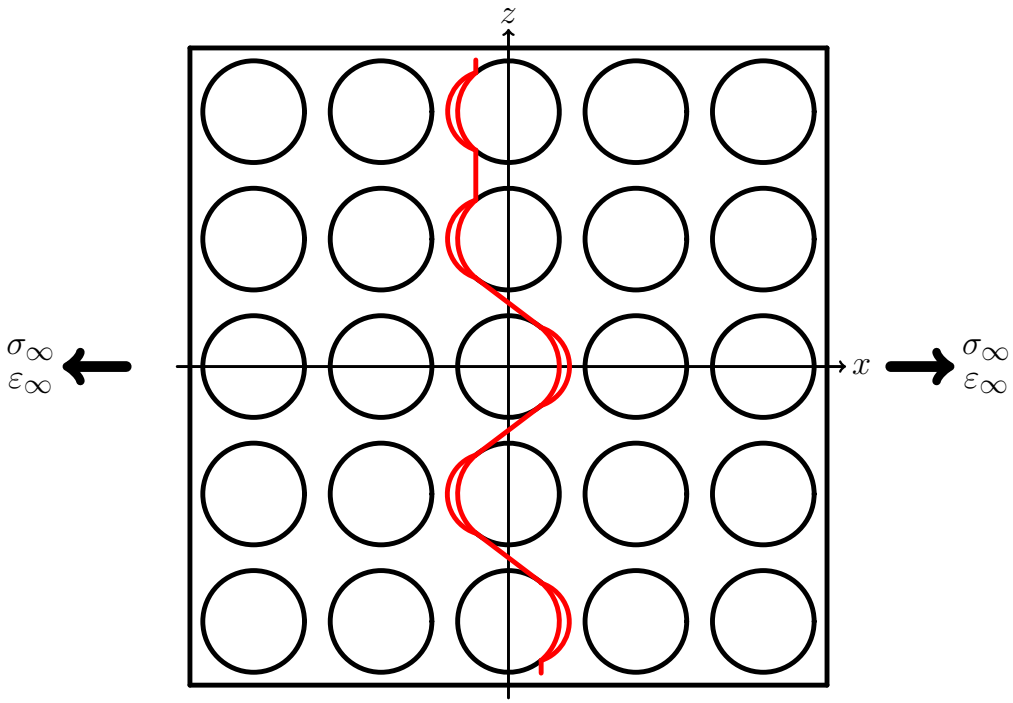
Micromechanics of Initiation: Transverse Tensile Loading

Stage 4: coalescence



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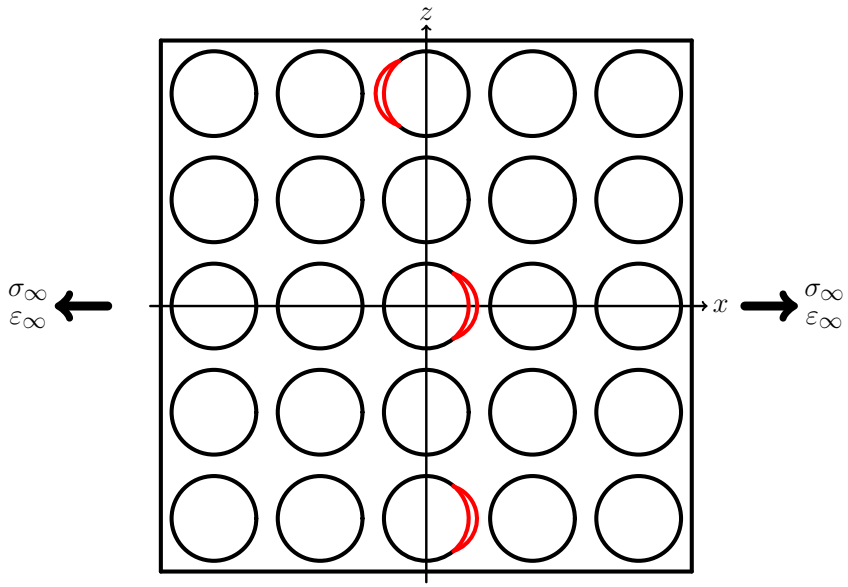
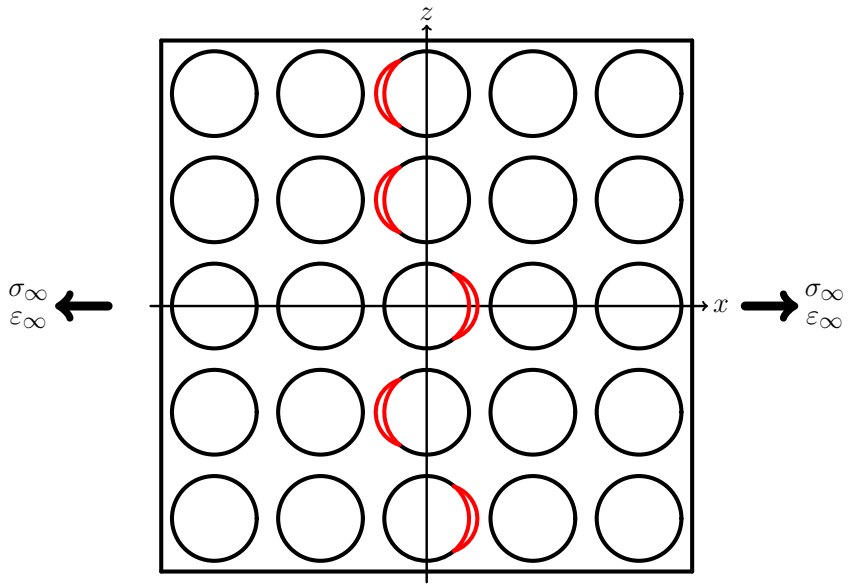
Objective of the Study

Can we talk about a ply-thickness effect for the fiber-matrix interface crack?

Stage 1: isolated debonds

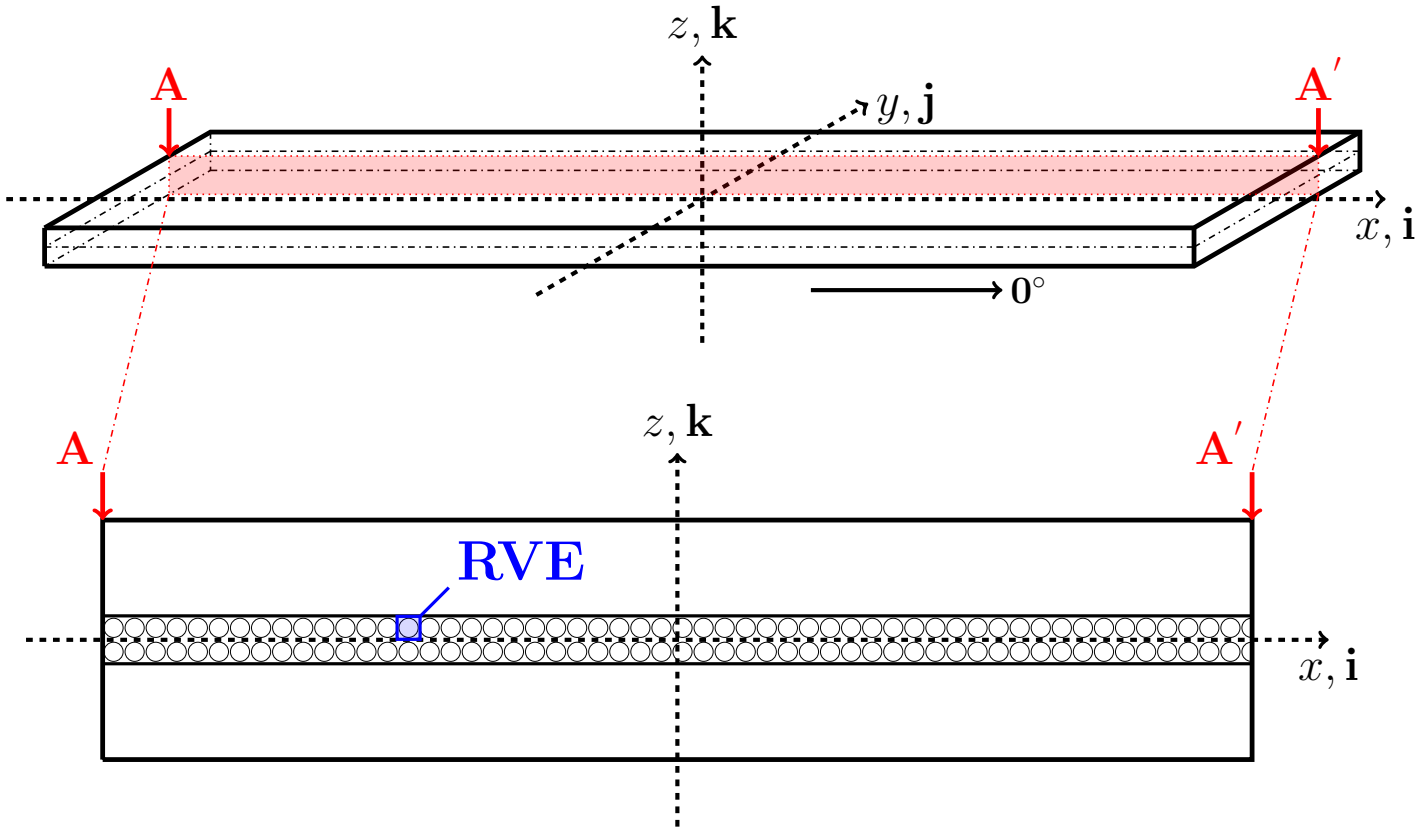


Stage 2: consecutive debonds



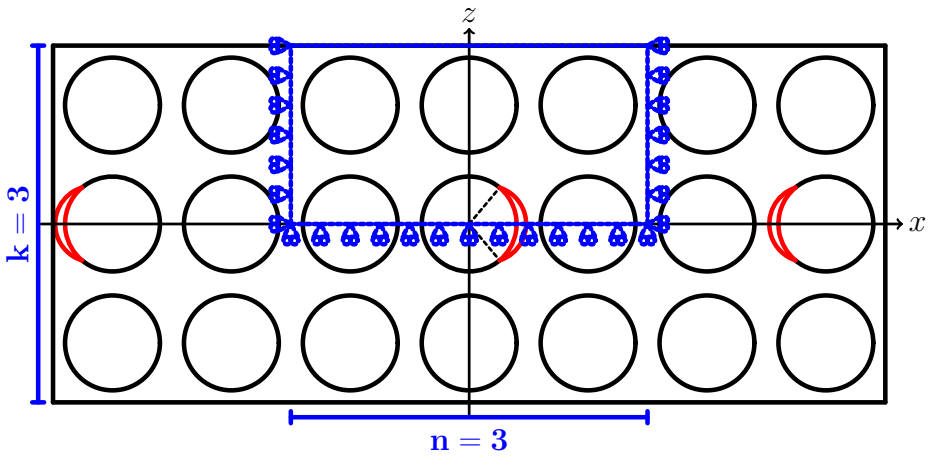
↘ MODELING

Geometry

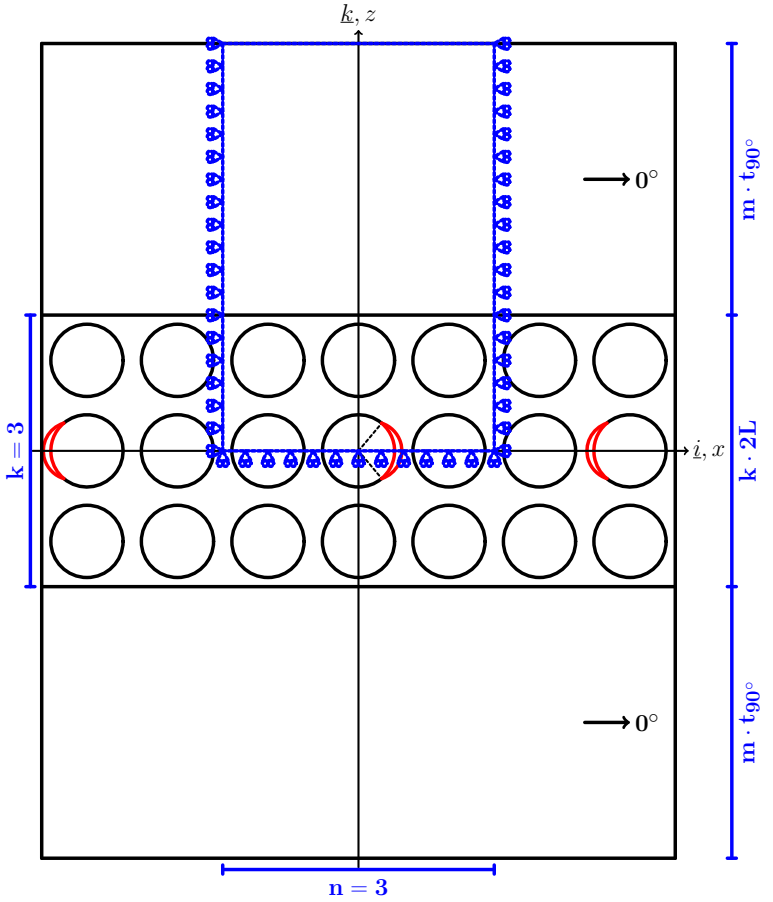


- ➔ $L, W >> t$
- ➔ $L, W \to \infty$
- ➔ Square packing
- ➔ $L_d >> \Delta\theta_d$
- ➔ 2D RVE

Representative Volume Elements

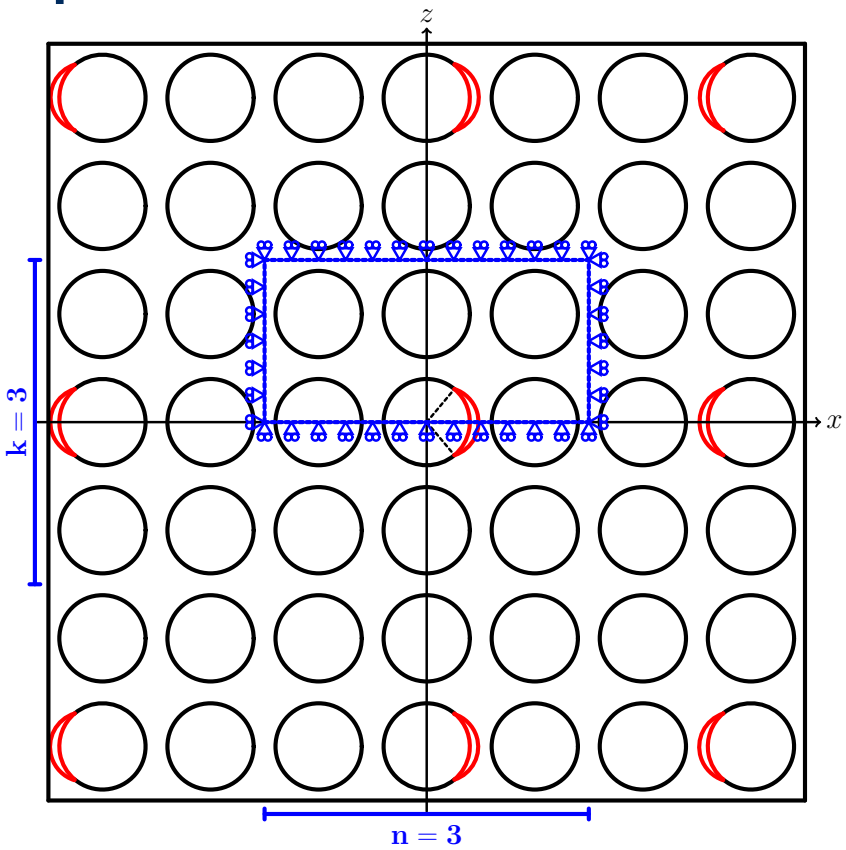


$n \times k - free$

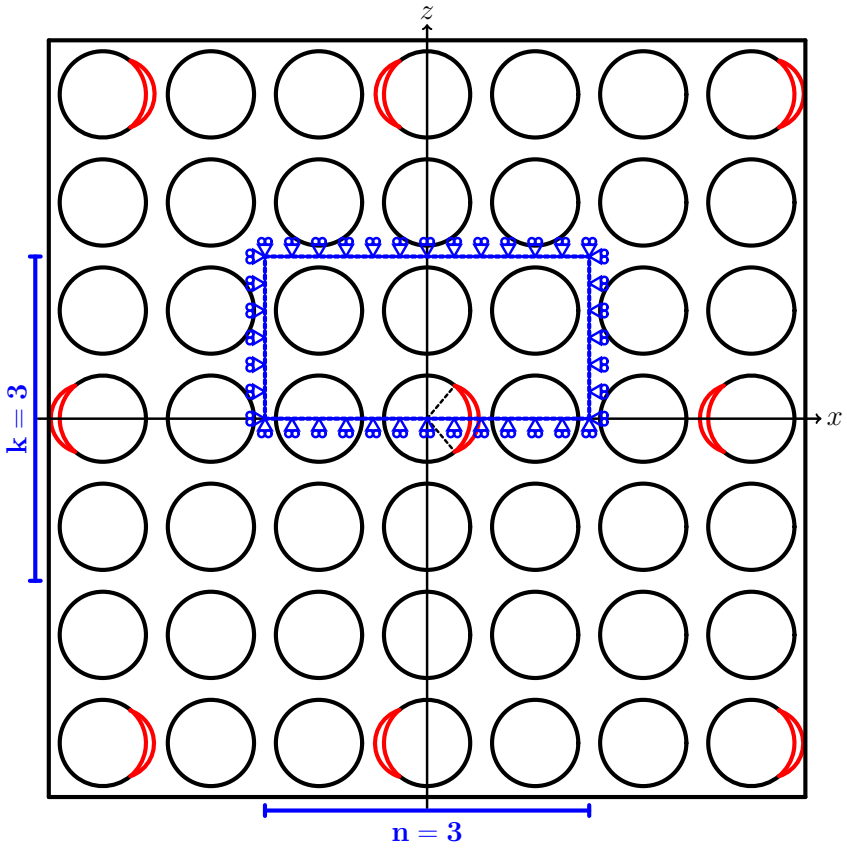


$n \times k - m \cdot t_{90°}$

Representative Volume Elements

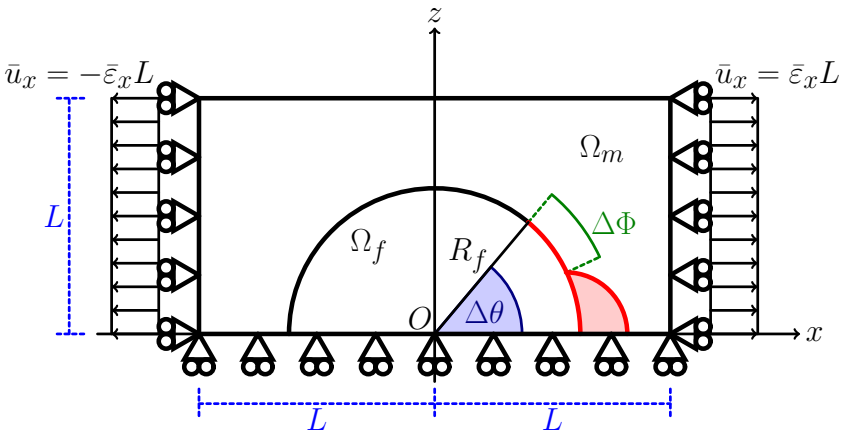


$n \times k - symm$



$n \times k - asymm$

Assumptions

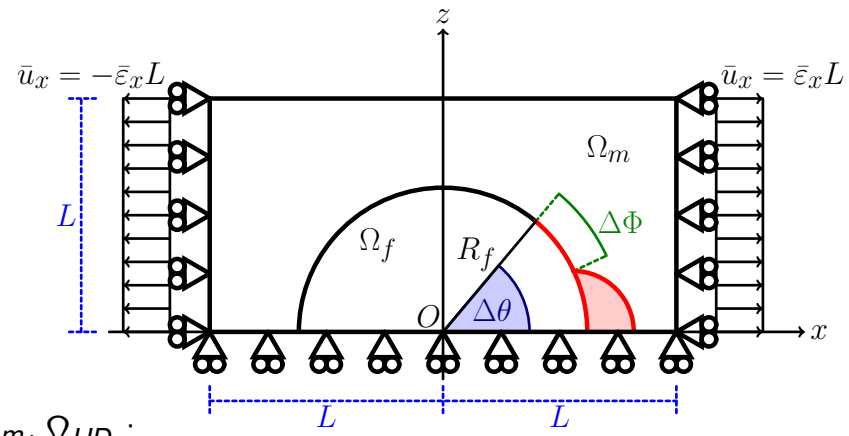


$$R_f = 1 \ [\mu m] \quad L = \frac{R_f}{2}\sqrt{\frac{\pi}{V_f}}$$

- ➔ Linear elastic, homogeneous materials
- ➔ Concentric Cylinders Assembly with Self-Consistent Shear Model for UD
- ➔ Plane strain
- ➔ Frictionless contact interaction
- ➔ Symmetric w.r.t. x-axis
- ➔ Coupling of x-displacements on left and right side (repeating unit cell)
- ➔ Applied uniaxial tensile strain $\bar{\varepsilon}_x = 1\%$
- ➔ $V_f = 60\%$

Material	V_f [%]	E_L [GPa]	E_T [GPa]	μ_{LT} [GPa]	ν_{LT} [−]	ν_{TT} [−]
Glass fiber	-	70.0	70.0	29.2	0.2	0.2
Epoxy	-	3.5	3.5	1.25	0.4	0.4
UD	60.0	43.442	13.714	4.315	0.273	0.465

Solution



in $\Omega_f, \Omega_m, \Omega_{UD}$:

$$\frac{\partial^2 \varepsilon_{xx}}{\partial z^2} + \frac{\partial^2 \varepsilon_{zz}}{\partial x^2} = \frac{\partial^2 \gamma_{zx}}{\partial x \partial z}$$

$$\varepsilon_y = \gamma_{xy} = \gamma_{yz} = 0$$

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{zx}}{\partial z} = 0$$

$$\frac{\partial \tau_{zx}}{\partial x} + \frac{\partial \sigma_{zz}}{\partial z} = 0$$

$$\sigma_{yy} = \nu \left(\sigma_{xx} + \sigma_{zz} \right)$$

for $0^\circ \leq \alpha \leq \Delta\theta, \Delta\theta \neq 0^\circ$:

$$\left(\vec{u}_m \left(R_f, \alpha \right) - \vec{u}_f \left(R_f, \alpha \right) \right) \cdot \vec{n}_\alpha \geq 0$$

for $\Delta\theta \leq \alpha \leq 180^\circ$:

$$\vec{u}_m \left(R_f, \alpha \right) - \vec{u}_f \left(R_f, \alpha \right) = 0$$

$$\sigma_{ij} = E_{ijkl} \varepsilon_{kl}$$

$+ BC$

$\forall\ \Delta\theta \neq 0^\circ$

➔ oscillating singularity

$$\sigma \sim r^{-\frac{1}{2}} \sin\left(\varepsilon \log r\right), \quad V_f \to 0$$

$$\varepsilon = \frac{1}{2\pi} \log\left(\frac{1-\beta}{1+\beta}\right)$$

$$\beta = \frac{\mu_2 \left(\kappa_1 - 1\right) - \mu_1 \left(\kappa_2 - 1\right)}{\mu_2 \left(\kappa_1 + 1\right) + \mu_1 \left(\kappa_2 + 1\right)}$$

➔ receding contact

➔ $\dfrac{G\left(R_{f,2}\right)}{G\left(R_{f,1}\right)} = \dfrac{R_{f,2}}{R_{f,1}}, \dfrac{G\left(\bar{\varepsilon}_{x,2}\right)}{G\left(\bar{\varepsilon}_{x,1}\right)} = \dfrac{\bar{\varepsilon}^2_{x,2}}{\bar{\varepsilon}^2_{x,1}}$

➔ FEM + LEFM (VCCT)

➔ regular mesh of quadrilaterals at the crack tip:

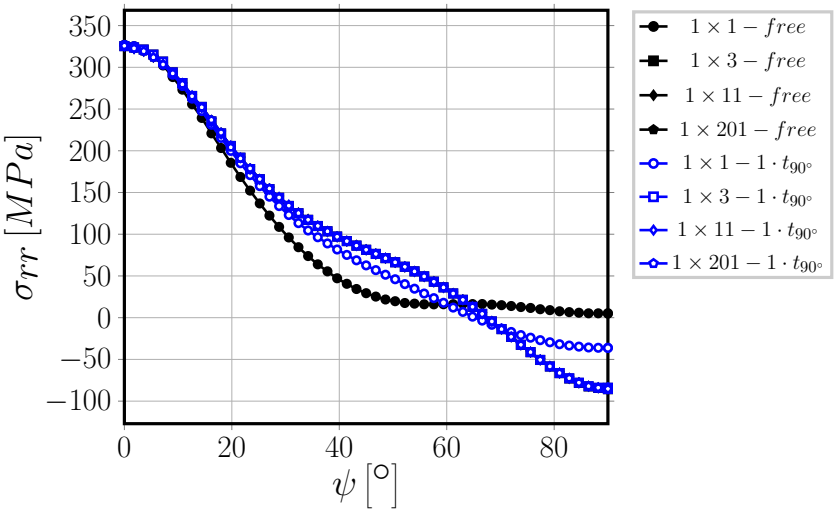
- $AR \sim 1, \quad \delta = 0.05^\circ$

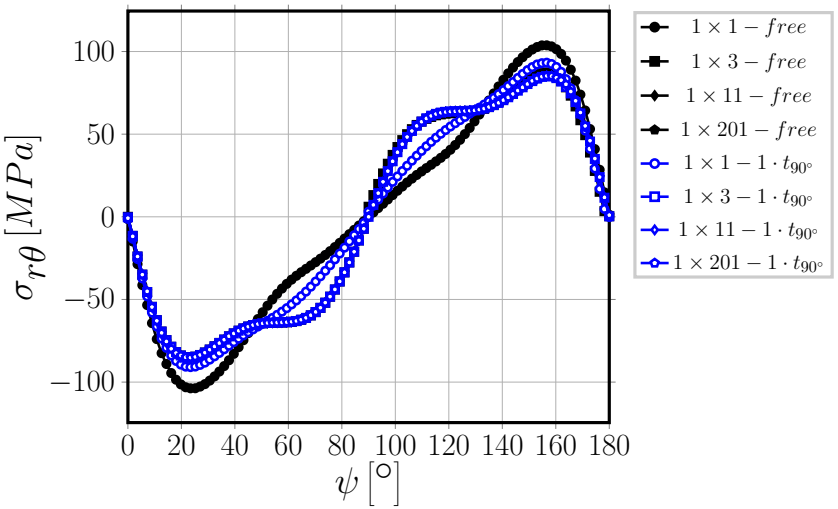
$\forall\ \Delta\theta$

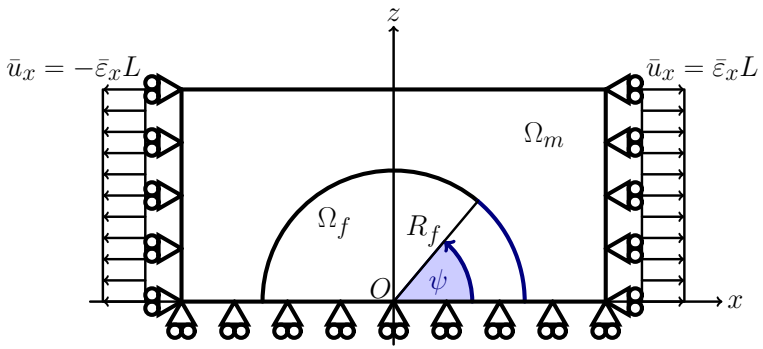
➔ 2^{nd} order shape functions

↘ Debond Initiation

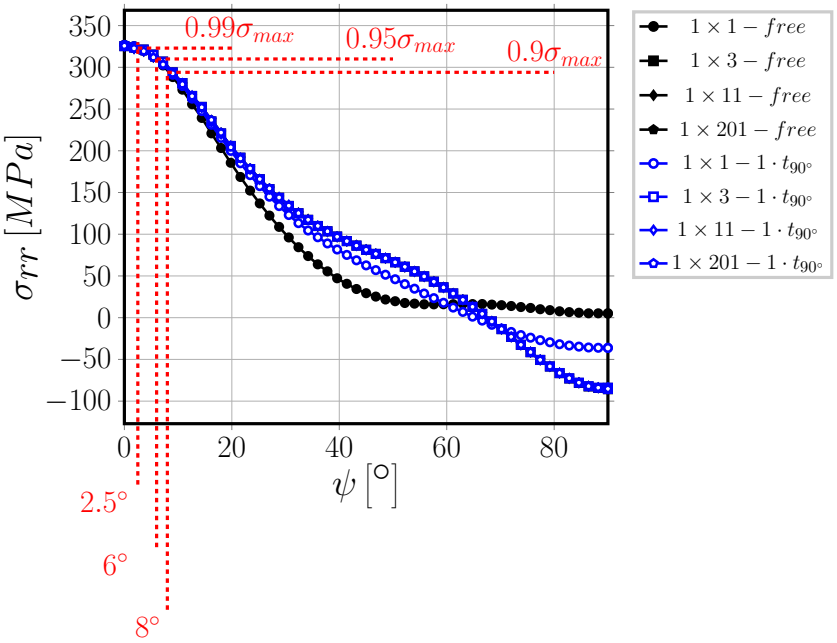
σ_{rr} vs $\tau_{r\theta}$: radial stress vs tangential shear at the interface

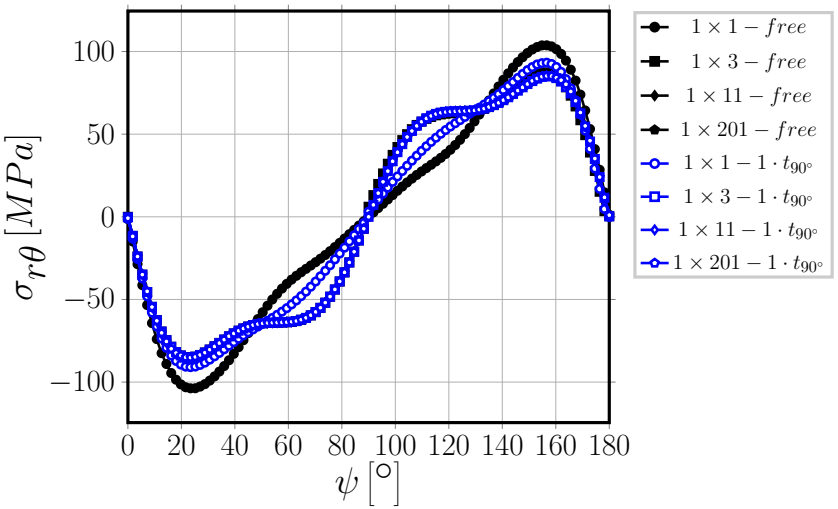


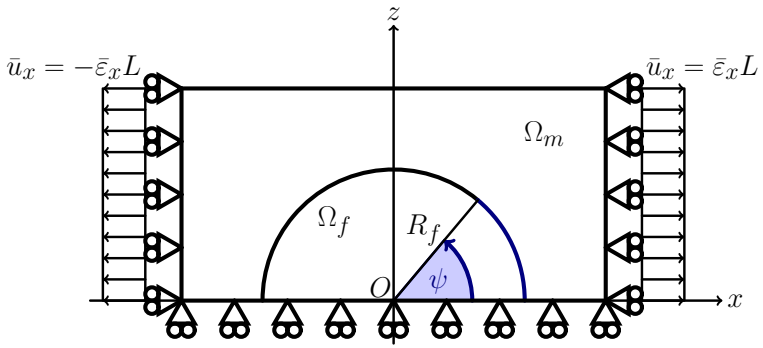




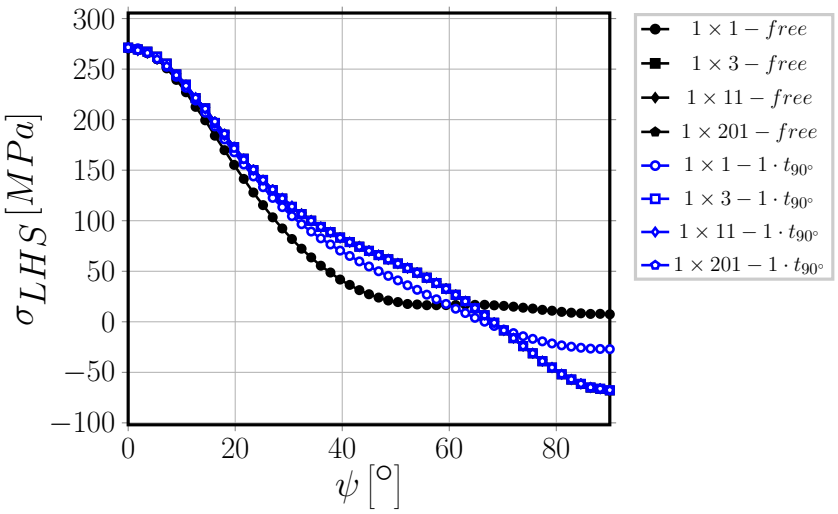
σ_{rr} vs $\tau_{r\theta}$: radial stress vs tangential shear at the interface

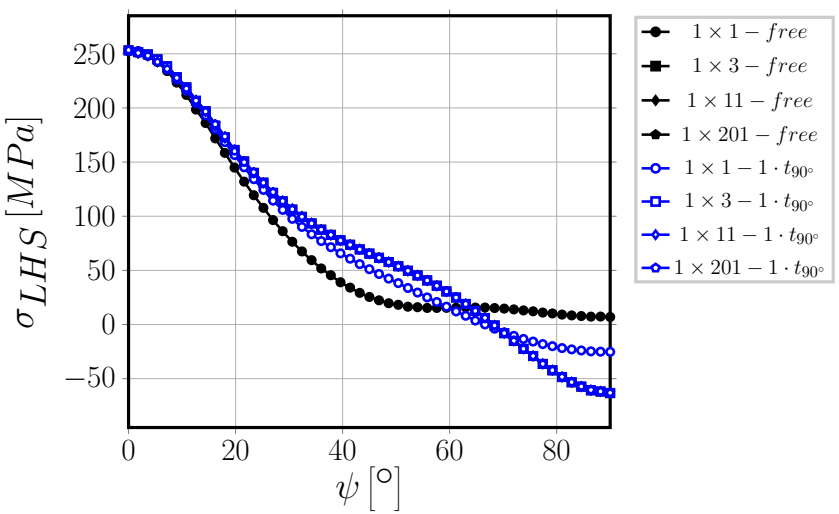






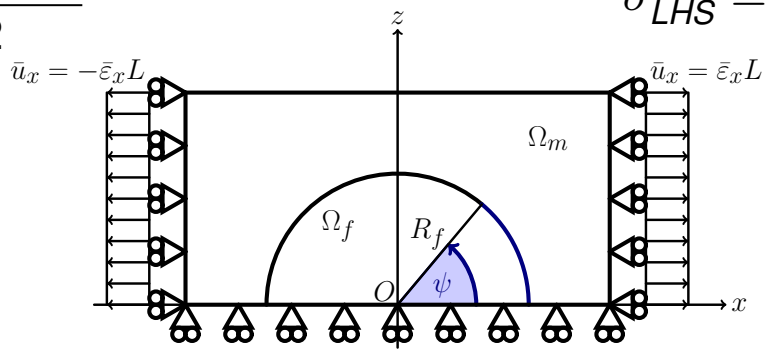
σ_{LHS}: local hydrostatic stress at the interface





$$\sigma^{2D}_{LHS} = \frac{\sigma_{rr} + \sigma_{\theta\theta}}{2}$$

$$\sigma^{3D}_{LHS} = \frac{\sigma_{rr} + \sigma_{\theta\theta} + \sigma_{yy}}{3}$$



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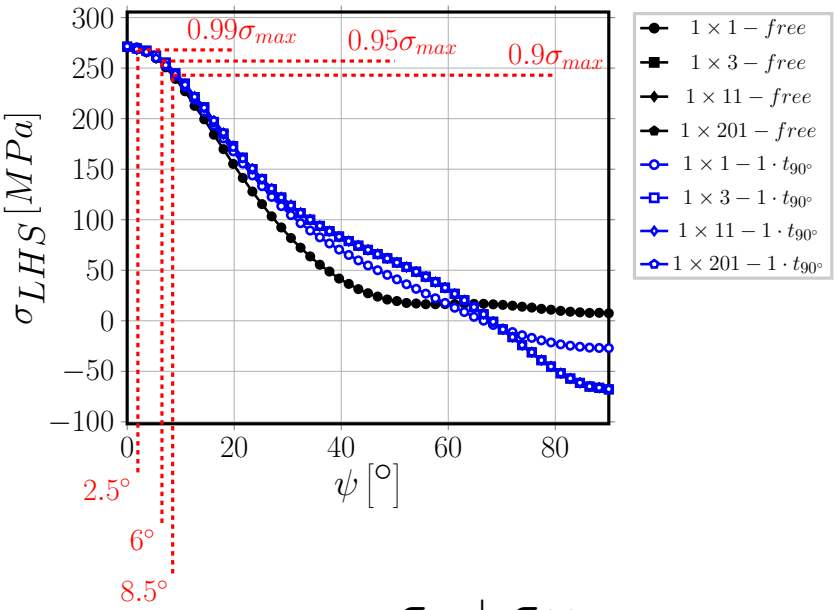

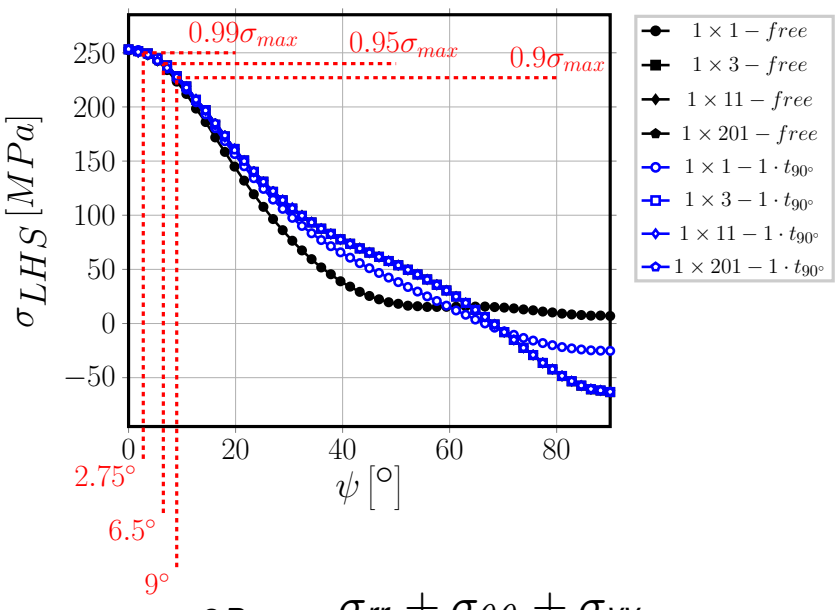

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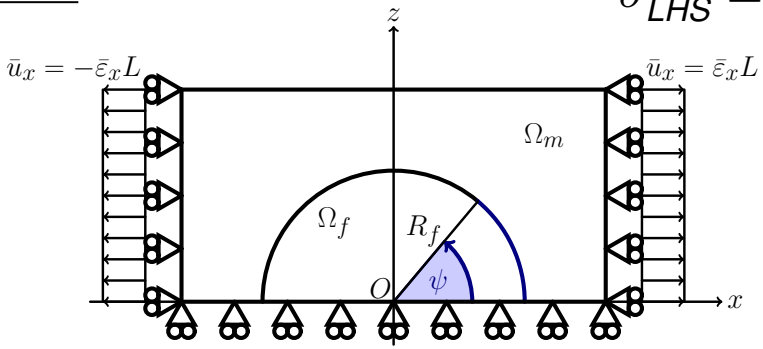
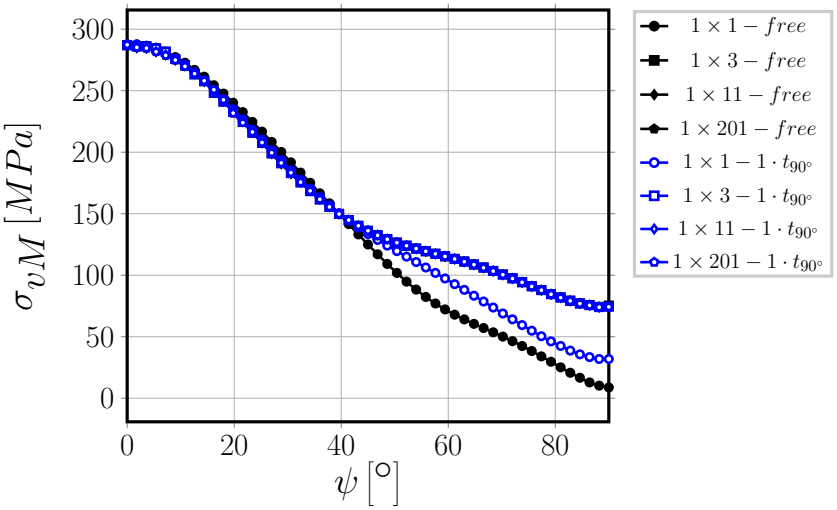
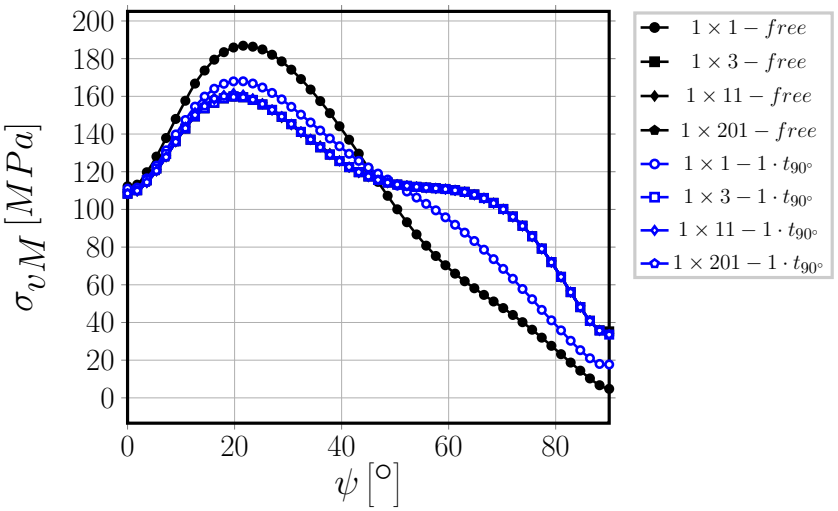


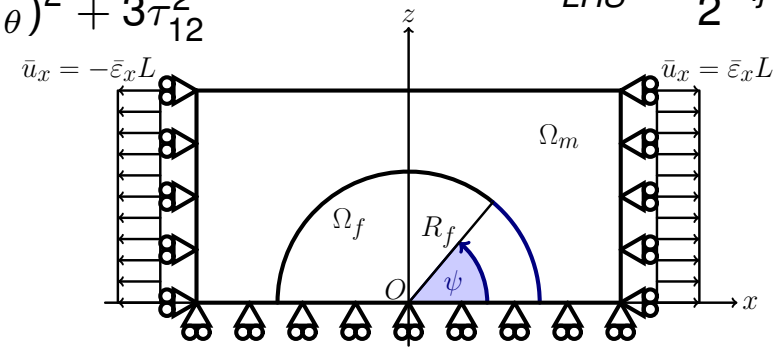
σ_{vM}: von Mises stress at the interface





$$\sigma^{2D}_{vM} = \sqrt{(\sigma_{rr}-\sigma_{\theta\theta})^2 + 3\tau^2_{12}}$$

$$\sigma^{3D}_{LHS} = \frac{3}{2} s_{ij} s_{ij} \quad s_{ij} = \sigma_{ij} - \frac{1}{3}\sigma_{kk}\delta_{ij}$$



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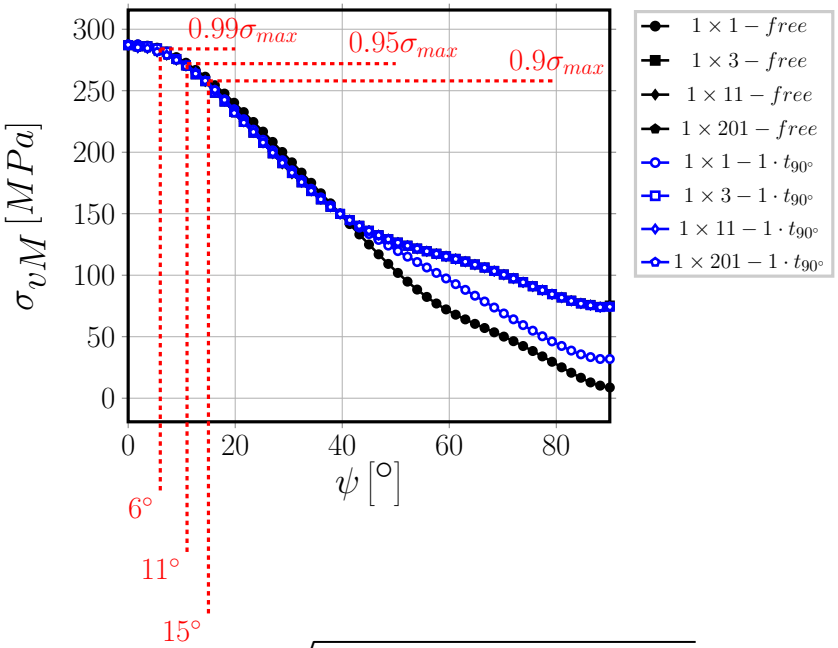

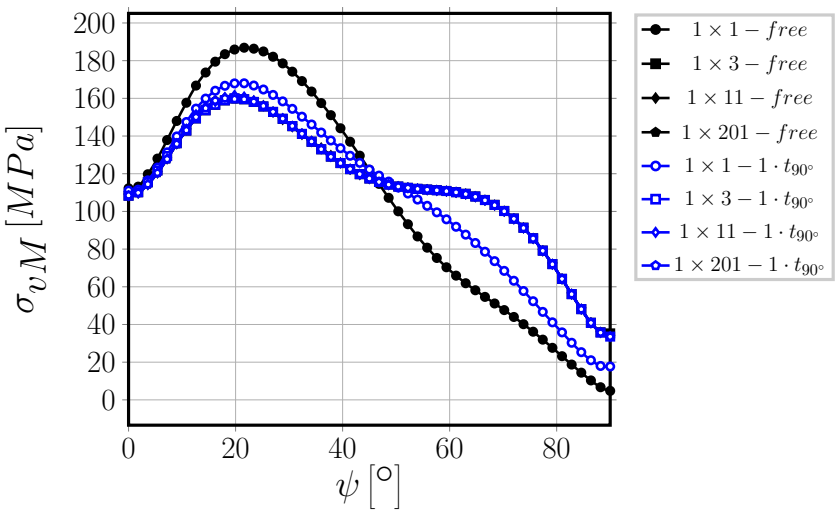

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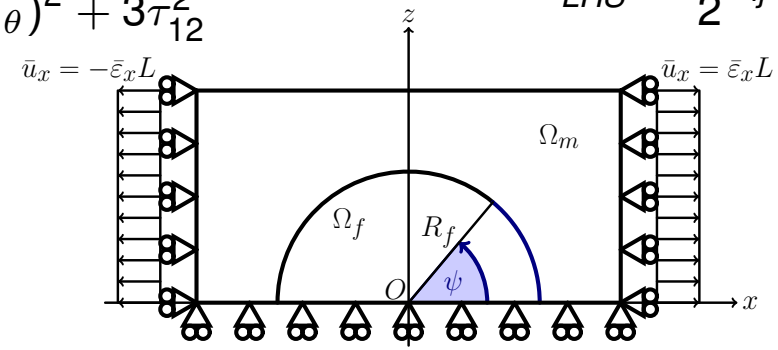
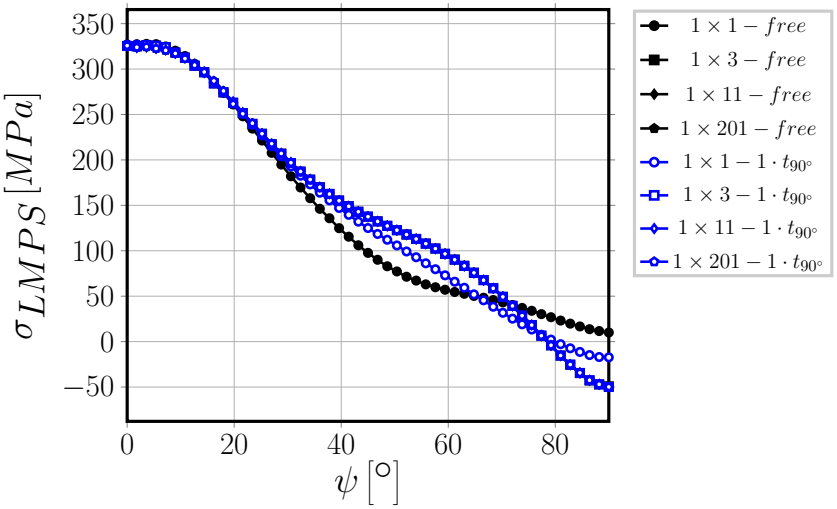
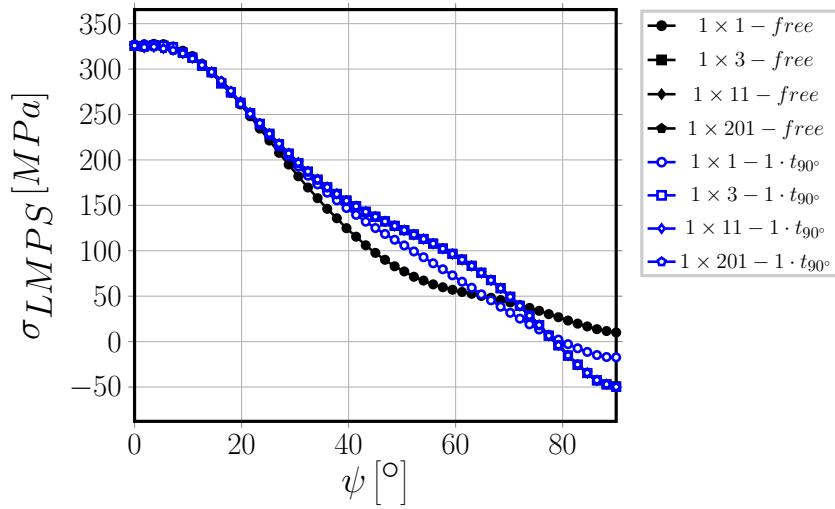
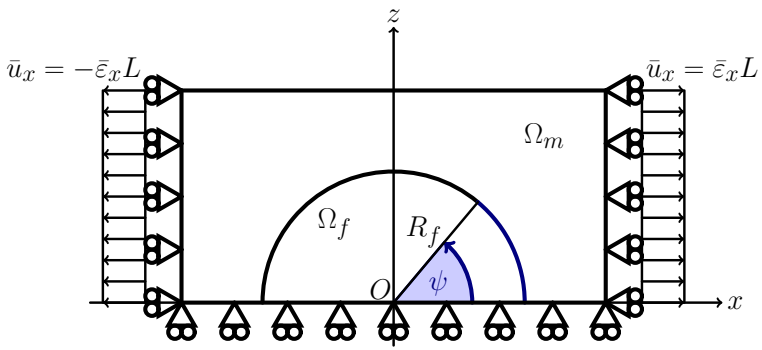


σ_I: maximum principal stress at the interface







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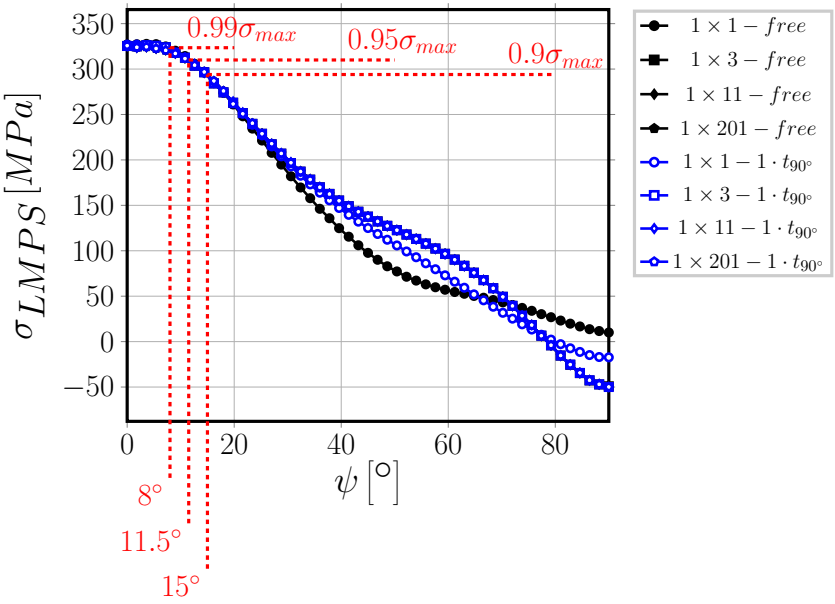
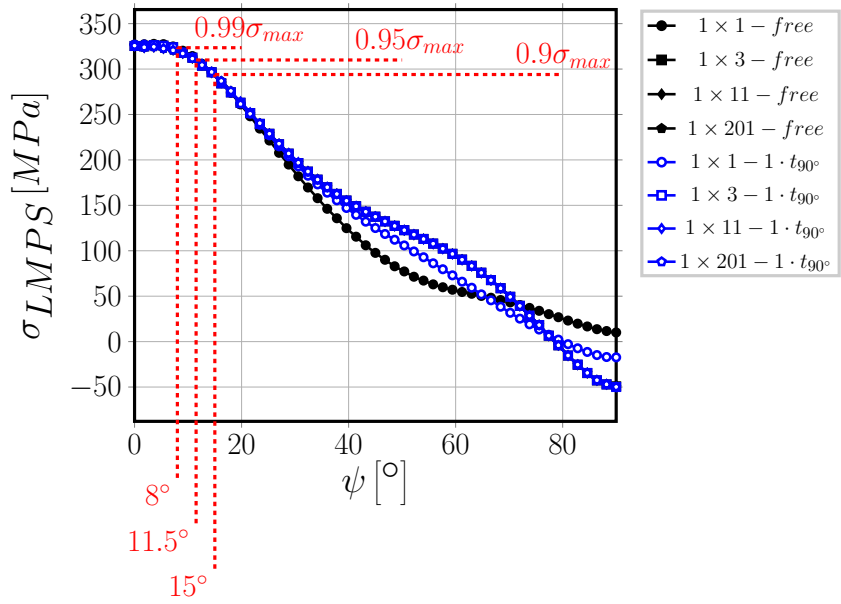
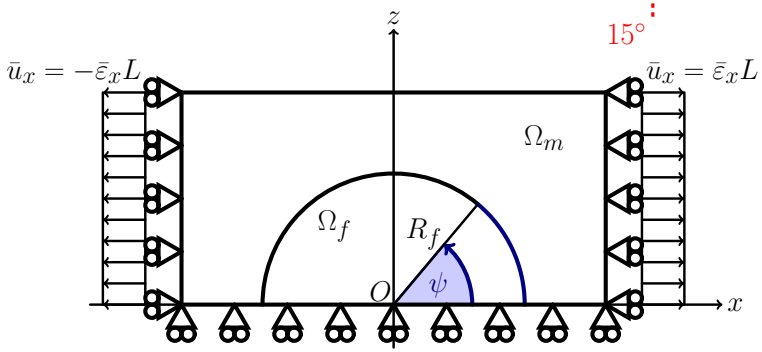






Observations & Conclusions

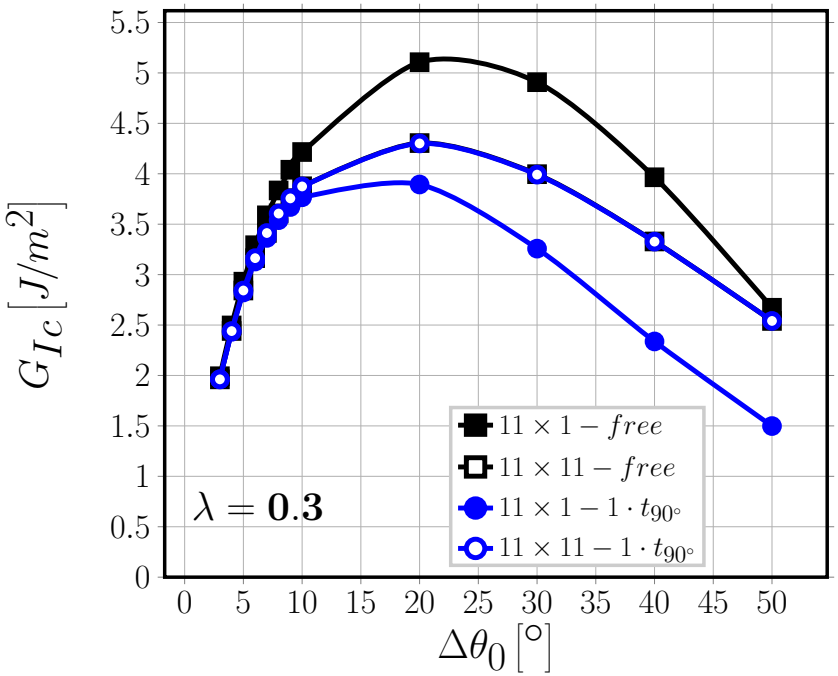
- ➔ For all stresses analyzed, no significant difference is present between the different RUCs for $\psi \leq 10°$;
- ➔ for all stresses analyzed, no difference can be observed by increasing k when $k \geq 3$;
- ➔ for all stresses analyzed, no difference can be observed between $1 \times k - \textit{free}$ and $1 \times k - 1 \cdot t_{90°}$ for $k \geq 3$;
- ➔ σ_{rr}, $\sigma_{LHS,2D}$, $\sigma_{LHS,3D}$, $\sigma_{vM,2D}$, $\sigma_{LMPS,2D}$ and $\sigma_{LMPS,3D}$ all reach their peak value at 0° and 180° and decrease to 99% the peak value between 2° and 8°, to 95% the peak value between 6° and 12° and to 90% the peak value between 8° and 15° from the occurrence of the maximum.

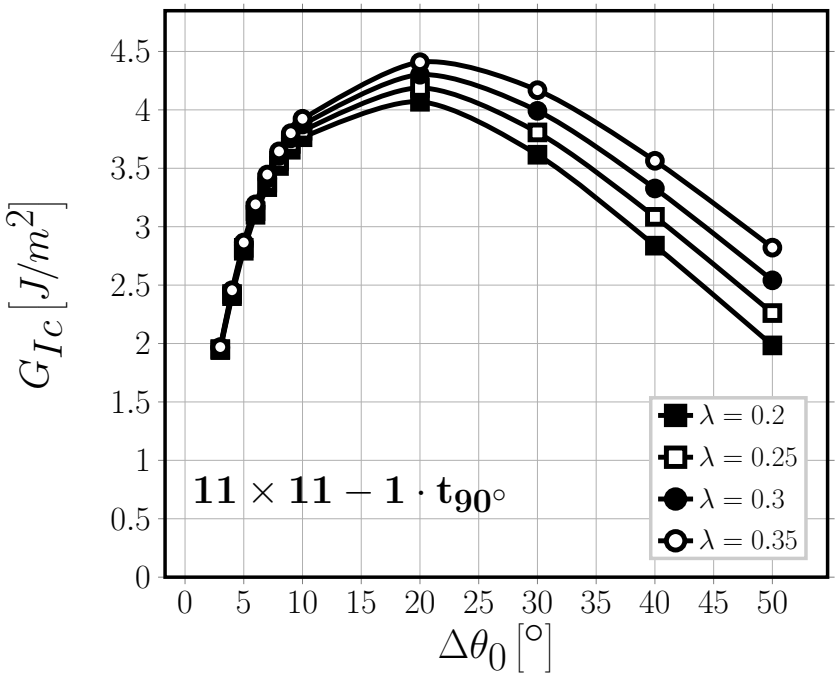
It seems reasonable to conclude that...

...a stress-based criterion would predict, irrespectively of the specific criterion chosen, the onset of an interface crack at 0° or 180° with an initial size at least comprised in the range $2° - 8°$ (1% margin) and likely in the range $6° - 12°$ (5% margin).

↘ Debond Propagation

Estimation of G_{Ic}

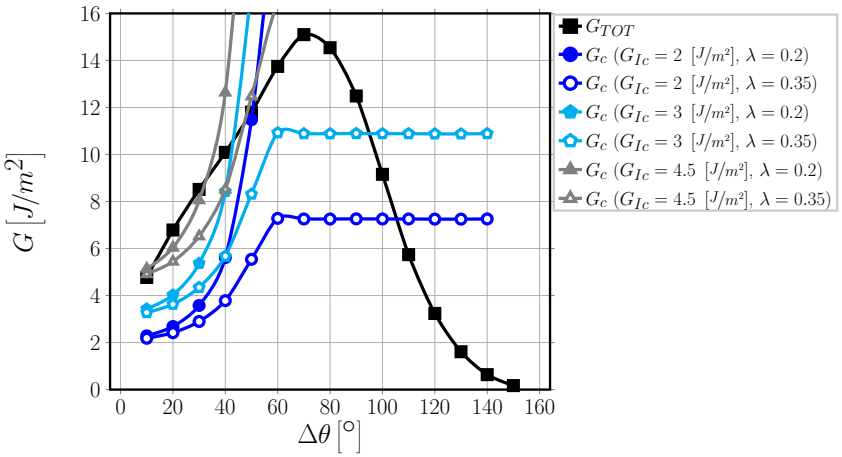




$$G_{Ic} = \left.\frac{G_c}{1 + \tan^2\left((1-\lambda)\,\Psi_G\right)}\right|_{G_c = G_{TOT}(\Delta\theta_0)}, \qquad \Psi_G = tan^{-1}\left(\left.\sqrt{\frac{G_{II}}{G_I}}\right|_{\Delta\theta_0}\right)$$

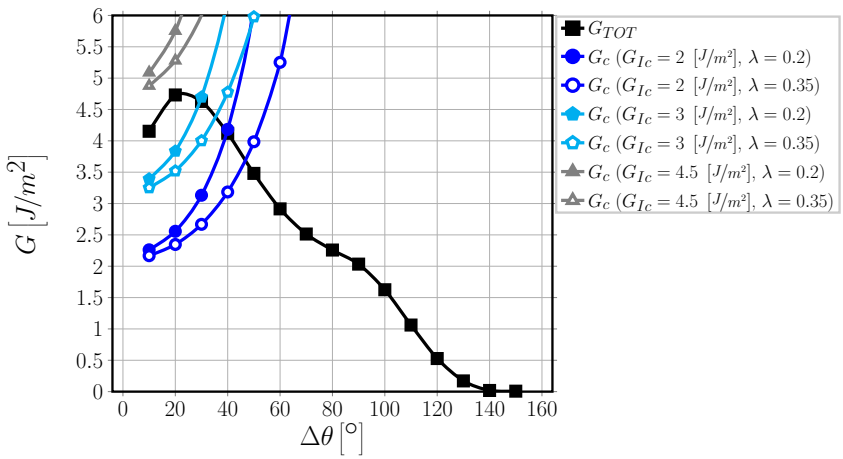
Estimation of $\Delta\theta_{max}$

$21 \times 1 - free$



$\Delta\theta_{max} \in (30^\circ - 105^\circ)$

$21 \times 1 - 1 \cdot t_{90^\circ}$

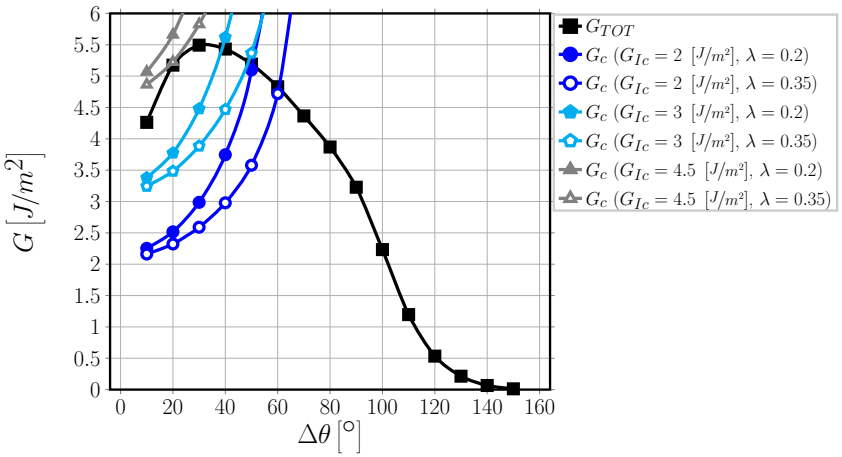


$\Delta\theta_{max} \in (30^\circ - 50^\circ)$

$$G_{TOT}(\Delta\theta) > G_c = G_{Ic}\left(1 + \tan^2\left((1-\lambda)\,\Psi_G\right)\right), \quad \Psi_G = tan^{-1}\left(\sqrt{\frac{G_{II}}{G_I}}\right)\Bigg|_{\Delta\theta}$$

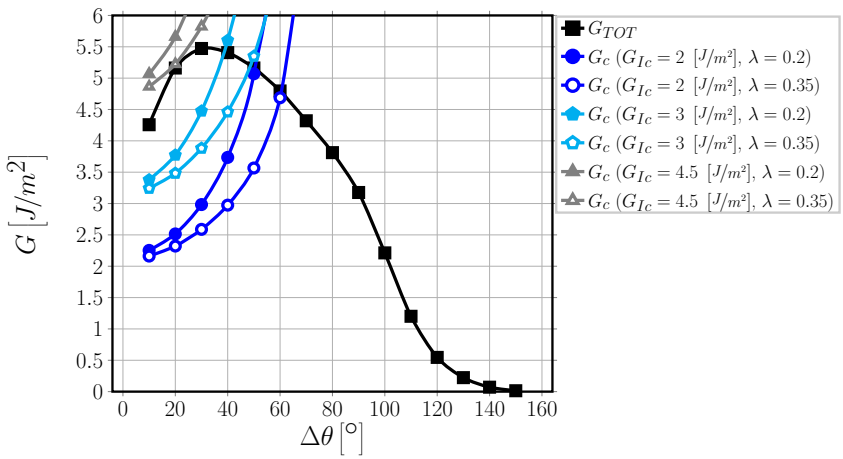
Estimation of $\Delta\theta_{max}$

$21 \times 3 - free$



$\Delta\theta_{max} \in (40° - 60°)$

$21 \times 3 - 1 \cdot t_{90°}$

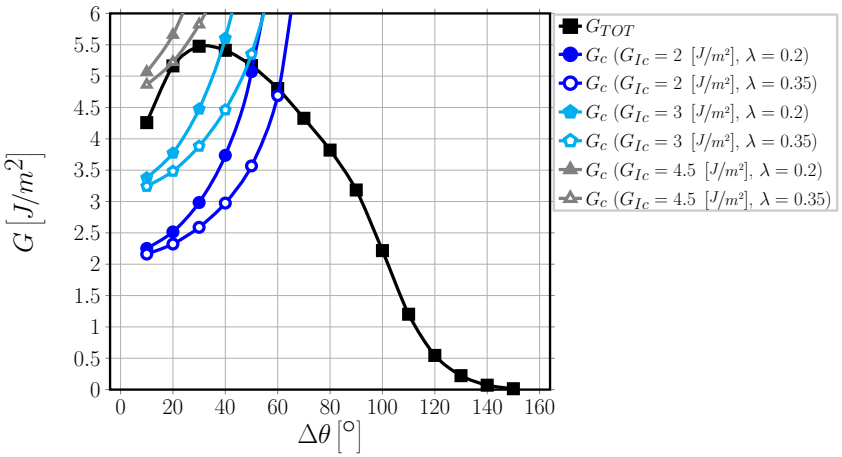


$\Delta\theta_{max} \in (40° - 60°)$

$$G_{TOT}(\Delta\theta) > G_c = G_{Ic}\left(1 + \tan^2\left((1-\lambda)\,\Psi_G\right)\right), \quad \Psi_G = tan^{-1}\left(\sqrt{\frac{G_{II}}{G_I}}\right)\Bigg|_{\Delta\theta}$$

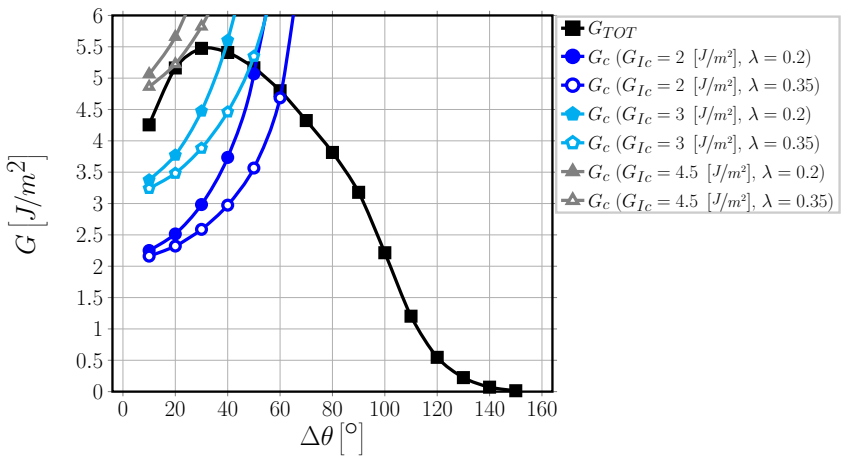
Estimation of $\Delta\theta_{max}$

$21 \times 21 - free$



$\Delta\theta_{max} \in (40° - 60°)$

$21 \times 21 - 1 \cdot t_{90°}$

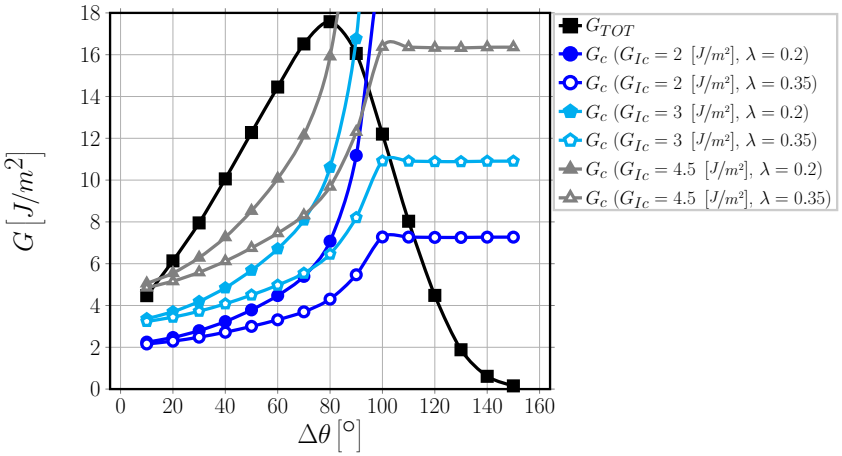


$\Delta\theta_{max} \in (40° - 60°)$

$$G_{TOT}(\Delta\theta) > G_c = G_{Ic}\left(1 + \tan^2\left((1-\lambda)\,\Psi_G\right)\right), \quad \Psi_G = tan^{-1}\left(\sqrt{\frac{G_{II}}{G_I}}\right)\Bigg|_{\Delta\theta}$$

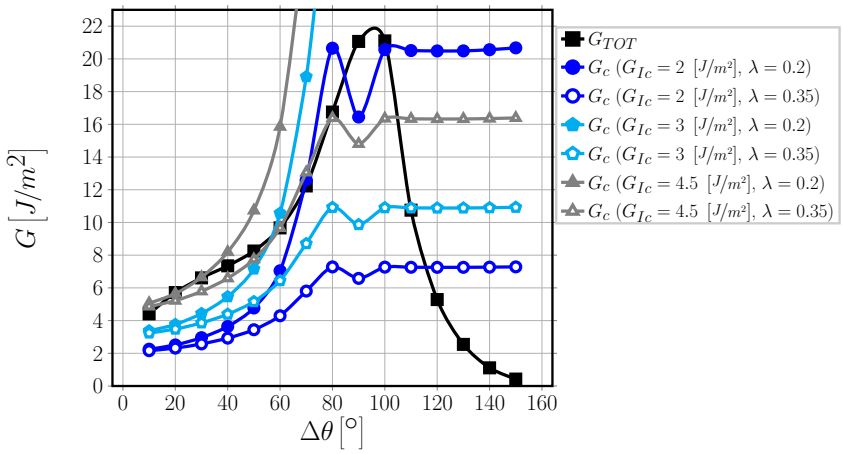
Estimation of $\Delta\theta_{max}$

$21 \times 21 - symm$



$\Delta\theta_{max} \in (80° - 110°)$

$21 \times 21 - asymm$



$\Delta\theta_{max} \in (55° - 115°)$

$$G_{TOT}(\Delta\theta) > G_c = G_{Ic}\left(1 + \tan^2\left((1-\lambda)\,\Psi_G\right)\right), \quad \Psi_G = tan^{-1}\left(\sqrt{\frac{G_{II}}{G_I}}\right)\Bigg|_{\Delta\theta}$$

Observations & Conclusions

- ➔ The effect of the presence of the 0° layer is to reduce the maximum size of debonds;
- ➔ however, the estimated debond size is the same for $n \times k - free$ and $n \times k - 1 \cdot t_{90°}$ for $k \geq 3$;
- ➔ the presence of a debond on the neighboring fiber in the vertical direction ($21 \times 1 - coupling$ and $21 \times 1 - asymm$) favors instead the growth of larger debonds;
- ➔ the largest size is achieved when debonds are on opposite sides of consecutive fibers.

Estimated debond size range in cross-ply ($n \times k - 1 \cdot t_{90°}$)

$40° - 60°$

Measured debond size range in cross-ply (Correa et al., Compos. Sci. Technol. 155 (213-220), 2018)

$21.4° - 89.2°$, average $49.3°$, standard deviation of $11.7°$
63% of measurements in $40° - 60°$ range

↘ CONCLUSIONS

Conclusions

- ➔ A stress criterion for initiation would likely predict, irrespectively of which criterion from those proposed in the literature is chosen, the onset of a debond at 0° or 180° with a semi-aperture $\Delta\theta_0$ in the range $2^\circ - 12^\circ$, corresponding to a margin of 5% on the satisfaction of the criterion.
- ➔ Assuming that debond propagation occurs unstably immediately after debond onset at the same level of global applied strain $\bar{\varepsilon}_x$, it is possible to evaluate the parameter G_{Ic} in the expression of the critical ERR and with it to estimate the range of expected maximum debond size.
- ➔ The prediction for a cross-ply laminate (models $n \times k - 1 \cdot t_{90^\circ}$, $k \geq 3$) agrees well with the debond size distribution in $\left[0_3^\circ, 90_3^\circ\right]_S$ estimated in *Correa et al., Compos. Sci. Technol.* ***155*** *(213-220), 2018* through microscopic observations.

Thank you for listening today!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