## In-class Exercise 8.2

Error analysis for time-stepping routines AMATH 301 University of Washington Jakob Kotas

Truncation error for time-stepping methods

| Method           | Local error       | Global error      |
|------------------|-------------------|-------------------|
| Euler's          | $O((\Delta t)^2)$ | $O(\Delta t)$     |
| Backward Euler's | $O((\Delta t)^2)$ | $O(\Delta t)$     |
| Heun's           | $O((\Delta t)^3)$ | $O((\Delta t)^2)$ |
| Runge-Kutta      | $O((\Delta t)^5)$ | $O((\Delta t)^4)$ |

## 1. Conceptual questions

- (a) What are the two types of error in time-stepping methods?
  - Truncation error, due to only keeping a finite number of terms in Taylor series expansion. Typically written as  $O((\Delta t)^k)$  for some k, where higher k means better accuracy.
  - Roundoff error, due to only keeping a finite number of digits on the floating-point numbers while doing calculations. Always present regardless of method.

Even if truncation error was made very small by using a very sophisticated method, roundoff error will still be present. The sweet spot for many problems is Runge-Kutta.

(b) Why is global error one power lower than local error?

Intuitively, because errors accumulate as you take more and more steps.

Mathematically, there are  $N = (b - a)/(\Delta t)$  steps, where a and b are the starting and ending times. If you sum up N local error terms which are  $O((\Delta t)^k)$ , the global error is

$$O((\Delta t)^k) \cdot \frac{b-a}{\Delta t} = O((\Delta t)^{k-1})$$

(c) In the 8.1 in-class exercises we looked at the IVP:

$$y' = -y, \ y(0) = 1$$

and found that Euler's method was unstable with  $\Delta t = 2$  but stable with  $\Delta t = 0.5$ . We also found that Backward Euler's method was stable with  $\Delta t = 2$ . What are the largest  $\Delta t$  we would expect to be stable for this IVP using Euler's and Backward Euler's methods?

From the Kutz video, Euler's method is unstable (roundoff error grows) for  $|1 + \lambda \Delta t| > 1$  and Backward Euler's method is unstable for  $\left| \frac{1}{1 - \lambda \Delta t} \right| > 1$ . The video looked at the IVP  $y' = \lambda y$  so here  $\lambda = -1$ .

Thus for Euler's method we will have stability if:

$$|1 - \Delta t| < 1$$
$$-1 < 1 - \Delta t < 1$$

1

$$\Delta t < 2$$

where we drop the left inequality as it is not physically meaningful. We avoid the edge-case exactly at  $|1 + \lambda \Delta t| = 1$ . Similarly, for Backward Euler's method we will have stability if:

$$\left| \frac{1}{1 + \Delta t} \right| < 1$$

$$|1 + \Delta t| > 1$$

$$1 + \Delta t > 1$$

(again, ignoring the case where  $\Delta t$  is negative).

$$\Delta t > 0$$

Meaning any  $\Delta t$  will be stable for Backward Euler's method on this IVP.

2. Consider the initial value problem:

$$\frac{dy}{dt} = \frac{3(t-1)^2}{y}, \ y(0) = 2.$$

The analytical solution is:

$$y(t) = \sqrt{2x^3 - 6x^2 + 6x + 4}.$$

(a) Approximate the solution on  $0 \le t \le t_{\text{fin}} = 4$  using Euler's method with a step size of  $\Delta t = 0.5$ . Save the local error after one step:

absolute local error = 
$$|y_{\text{Euler}}(\Delta t) - y_{\text{true}}(\Delta t)|$$

and global error after all steps:

absolute global error = 
$$|y_{\text{Euler}}(t_{\text{fin}}) - y_{\text{true}}(t_{\text{fin}})|$$

- (b) Reduce  $\Delta t$  to half its original value and repeat, saving both errors again.
- (c) Repeat until  $\Delta t = 2^{-8}$ . Also plot the analytical solution. Do the curves approach the analytical solution as  $\Delta t \to 0$ ?
- (d) Make a plot of  $\ln(\Delta t)$  vs.  $\ln(\text{error})$  for both the local and global errors. Do both decrease as  $\Delta t \to 0$ ?
- (e) Preview of chapter 4: find the slope of the best-fit line through these data sets and print the values. Use np.polyfit. Do we see evidence of second-order local error, first-order global error?
- 3. The Blasius equation is used in fluid dynamics to describe the laminar boundary layer on a semi-infinite plate under constant fluid flow. The stream function f(x) is defined by the differential equation:

$$2\frac{d^3f}{dx^3} + \frac{d^2f}{dx^2}f = 0$$

with initial/boundary conditions:

$$f'(0) = 0$$
,  $f(0) = 0$ ,  $f'(\infty) = 1$ 

(a) Use the shooting method to find a value of  $f''(0) = \alpha$  such that  $f'(100) \approx 1$  when solving with solve\_ivp. (We assume that 100 is a large enough time that  $f'(100) \approx f'(\infty)$ .)

Let f' = g, g' = h. Then the system of equations is:

$$\frac{dv}{dt} = \frac{d}{dt} \begin{bmatrix} f \\ g \\ h \end{bmatrix} = \begin{bmatrix} g \\ h \\ -fh/2 \end{bmatrix}$$

with initial conditions:

$$\left[\begin{array}{c} f(0) \\ g(0) \\ h(0) \end{array}\right] = \left[\begin{array}{c} 0 \\ 0 \\ \alpha \end{array}\right]$$

and we wish to find  $\alpha$  such that  $f''(0) = \alpha$  leads to  $f'(100) \approx 1$ .

4. There are some differential equations where time-stepping methods perform poorly because the number of time-steps required to avoid instability is extremely large. These are called **stiff equations**.

As an example, solve the IVP below:

$$\begin{cases} x'(t) & = -56x + 55y \\ y'(t) & = 44x - 45y \end{cases}, \quad x(0) = 1, \ y(0) = 0.$$

on the time interval  $0 \le t \le 1000$ .

- (a) Use solve\_ivp with the RK45 method; use time.time() to print the elapsed time. You'll need to import time first.
- (b) Repeat part (a) with the Radau method, which is designed to perform better on stiff equations. Print the elapsed time.