

# Quantum Mechanics 2025 HW11

No need to turn in

December 22, 2025

## Problem 1. Resonance scattering

Consider a 3D attractive spherical  $\delta$ -potential well with  $\frac{2m}{\hbar^2}V(r) = \gamma\delta(r - R)$  where  $\gamma < 0$  and  $R$  is the radius of the shell of the  $\delta$ -potential. In the following, we only consider the  $s$ -wave channel, i.e.,  $l = 0$ , and define  $u(r) = rR_0(r)$  where  $R_0(r)$  is the radial solution in the  $s$ -wave channel.

1) Prove that there exists a value  $|\gamma_c|$ , such that, when  $|\gamma| > |\gamma_c|$  there is one bound state, and when  $|\gamma| < |\gamma_c|$  there is no bound state. Determine the value of the dimensionless parameter  $|\gamma_c|R$ , and the localization length of the bound state in the limit of  $|\gamma| \rightarrow |\gamma_c|$ .

2) Find the equation of phase shift that  $\tan \delta_0(k)$  satisfies, and check that this expression is consistent with the fact  $\delta_0(k)$  is an odd function of  $k$ .

It is not required here, but if you can justify why  $\delta_0(k)$  is an odd function, you can get some extra credits.

Expand the expression  $k \cot \delta_0(k)$  around  $k = 0$  and keep the first two leading terms. Show that it can be expanded as

$$k \cot \delta_0(k) = -\frac{1}{a_0} + r_0 k^2, \quad (1)$$

and determine the values of scattering length  $a_0$  and the interaction range  $r_0$ . Show that in the case that  $|\gamma| \rightarrow |\gamma_c| + 0^+$ ,  $a_0$  is the same as the localization length of the bound state at the leading order. In this case, between  $|a_0|$  and  $r_0$ , which can be much larger than  $R$ , and which is at the same order of  $R$ ?

Now let us consider the locations of resonance scattering.

3) Show that at  $|\gamma| < |\gamma_c|$ , there is no resonance. Sketch the plots of  $\tan \delta_0(k)$  v.s.  $kR$ , and  $\delta_0(k)$  v.s.  $kR$ . What is the value of  $\delta_0(k = 0)$ ? If right at  $|\gamma| = |\gamma_c|$ , again sketch the plots of  $\tan \delta_0(k)$  v.s.  $kR$ , and  $\delta_0(k)$  v.s.  $kR$  at  $|\gamma| = |\gamma_c|$ . What is the value of  $\delta_0(k = 0)$  in this case?

4) Let us consider the case that  $|\gamma|$  slightly larger than  $|\gamma_c|$ . Show that there is only one solution for  $\tan \delta_0(k) = \pm\infty$  at  $0 < kR < \frac{\pi}{2}$ . Again sketch the plots of  $\tan \delta_0(k)$  v.s.  $kR$ , and  $\delta_0(k)$  v.s.  $kR$ . What is the value of  $\delta_0(k = 0)$  in this case?

If  $|\gamma|$  further goes large will  $\delta_0(k = 0)$  change or not? You further check the case at which there are three solutions  $\tan \delta_0(k) = \pm\infty$ , and use plots to support your arguments.

Combining all the cases in 4) and 5), can you relate  $\delta_0(k = 0)$  to the number of bound states?

5) Consider the limit of  $|\gamma| \rightarrow +\infty$ , show that the resonance solutions, i.e.,  $\tan \delta_0(k) = \pm\infty$ , occur at  $kR \rightarrow (n + \frac{1}{2})\pi + 0^-$ , and at  $kR \rightarrow n\pi + 0^+$ . For which case, the radial

wavefunction behaves like a meta-stable quasi-bound state, i.e., the magnitude of  $u(r)$  is much stronger at  $r < R$  than that at  $r > R$ ?