# QM HW2

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### Problem 1 (Harmonic osicllator)

(1) We've known that  $x_{nm} \neq 0 \Rightarrow m = n \pm 1$ 

$$x_{n,n}^2 = \sum_{k} x_{n,k} x_{k,n} = x_{n,n-1} x_{n-1,n} + x_{n,n+1} x_{n+1,n}$$
 (1.1)

Since

$$x_{n+1,n} = \sqrt{\frac{\hbar}{2m\omega_0}(n+1)} \tag{1.2}$$

We have

$$x_{n,n}^2 = \frac{\hbar}{m\omega_0} \left( n + \frac{1}{2} \right)$$
 (1.3)

$$(p^2)_{n,n} = m^2(\dot{x})_{n,n}^2 = \hbar\omega_0 m(n+1/2)$$
(1.4)

$$\sqrt{(x^2)_{n,n}(p^2)_{n,n}} = \hbar \left(n + \frac{1}{2}\right)$$

$$\tag{1.5}$$

When n = 0,

$$\sqrt{(x^2)_{n,n}(p^2)_{n,n}} = \frac{\hbar}{2} \tag{1.6}$$

(2)

$$x_{n,m}(t) = \sqrt{\frac{\hbar}{2m\omega_0}} \left( \sqrt{n}e^{i\omega_0 t} \delta_{m,n-1} + \sqrt{n+1}e^{-i\omega_0 t} \delta_{m,n+1} \right)$$
(1.7)

$$p_{n,m}(t) = i\sqrt{\frac{\hbar m\omega_0}{2}} \left(\sqrt{n}e^{i\omega_0 t}\delta_{m,n-1} - \sqrt{n+1}e^{-i\omega_0 t}\delta_{m,n+1}\right)$$
(1.8)

$$[x, p]_{n,m}(t) = \sum_{k} \left( x_{n,k} p_{k,m} e^{i\omega_{n,m}t} - p_{m,k} x_{k,n} e^{i\omega_{m,n}t} \right)$$
(1.9)

$$= [x_{n,n-1}p_{n-1,m} + x_{n,n+1}p_{n+1,m}]$$

$$-p_{m,n-1}x_{n-1,n} - p_{m,n+1}x_{n+1,n} ] \delta_{nm}$$
 (1.10)

$$= i\hbar \delta_{nm}$$
 (1.11)

## Problem 2 (Quantum equation of motion)

(1) By associative law,  $\forall H_2(x)$ ,

$$[x, H_2(x)] = 0 (2.1)$$

We expand  $H_1(p)$  in form of

$$H_1(p) = \sum_{n=0}^{\infty} a_n p^n \tag{2.2}$$

We have

$$[x, p^{n}] = \sum_{i=0}^{n-1} p^{i}[x, p] p^{n-i-1} = in\hbar p^{n-1} = i\hbar \frac{\partial (p^{n})}{\partial p}$$
 (2.3)

Hence,

$$\frac{\partial H}{\partial p} = \frac{1}{i\hbar} [x, H] \tag{2.4}$$

Similarly,

$$-\frac{\partial H}{\partial x} = \frac{1}{i\hbar}[p, H] \tag{2.5}$$

Therefore,

$$\dot{x} = \frac{i}{\hbar}[x, H], \quad \dot{p} = \frac{i}{\hbar}[p, H] \tag{2.6}$$

(2)  $\forall (n,m) \in \mathbb{N}^2$ 

$$[x, x^n p^m] = xx^n p^m - x^n p^m x = x^n (xp^m - p^m x) = x^n [x, p]$$
 (2.7)

So,

$$[x, x^n p^m] = i\hbar x^n \frac{\partial (p^m)}{\partial p}$$
(2.8)

Hence (2.4) and (2.5) still hold. Let  $x = x, y, p = p_x, p_y$ .

$$\dot{x} = \frac{1}{i\hbar}[x, H], \quad \dot{p}_x = \frac{1}{i\hbar}[p_x, H]$$

$$\dot{y} = \frac{1}{i\hbar}[y, H], \quad \dot{p}_y = \frac{1}{i\hbar}[p_y, H]$$
(2.9)

#### Problem 3 (De Broglie wave)

By relativity,

$$p = mv_q , E = mc^2 (3.1)$$

$$v_g = \frac{pc^2}{E} \tag{3.2}$$

In wave case,

$$v_g = \frac{\mathrm{d}\omega}{\mathrm{d}k} \tag{3.3}$$

From

$$E^2 = p^2 c^2 + m_0^2 c^4 (3.4)$$

we deduce that

$$\frac{pc^2}{E} = \frac{\mathrm{d}E}{\mathrm{d}p} = \hbar \frac{\mathrm{d}\omega}{\mathrm{d}p} \tag{3.5}$$

Thus,

$$dp = \hbar dk$$

$$p = \hbar k$$
(3.6)

### Problem 4 (Schrödinger equation)

Let

$$\Psi(r,t) = e^{iS/k} = e^{i[W(r) - Et]/k} = \psi(r)e^{-iEt/k}$$
(4.1)

Considering  $E = \hbar \omega$ , take  $k = \hbar$ . Then

$$S = -i\hbar \ln \psi(r) - Et \tag{4.2}$$

Plug in

$$\left(\nabla S\right)^2 = 2m\left(E - V(r)\right) \tag{4.3}$$

we obtain

$$\frac{\hbar^2}{2m} (\nabla \psi)^2 + (E - V(r)) \psi^2 = 0$$
 (4.4)

Now, we want to linearize the equation. Let  $\nabla$  act on (4.4) and eliminate the common factor

$$\frac{\hbar^2}{2m} \nabla^2 \psi + (E - V(r)) \psi + \frac{\nabla (E - V(r)) \psi^2}{2\nabla \psi} = 0$$
 (4.5)

God believes S, E, V are smooth at the  $\hbar$  scale. We ignore the last term.

$$\boxed{\frac{\hbar^2}{2m}\nabla^2\psi + (E - V(r))\psi = 0}$$
(4.6)

Problem 5 (Quantum Potential)

Plug

$$\frac{\partial \Psi}{\partial t} = \frac{i}{\hbar} \frac{\partial S}{\partial t} e^{iS/\hbar} \tag{5.1}$$

$$\nabla \Psi = \frac{i}{\hbar} \nabla S e^{iS/\hbar} \tag{5.2}$$

$$\nabla^2 \Psi = \frac{1}{\hbar^2} \left[ i\hbar \nabla^2 S - (\nabla S)^2 \right] e^{iS/\hbar}$$
 (5.3)

into Schrödinger equation,

$$\frac{\partial S}{\partial t} = \frac{1}{2m} \left[ i\hbar \nabla^2 S - (\nabla S)^2 \right] - V \tag{5.4}$$

When taking the limit  $\hbar \to 0$ ,

$$\frac{\partial S}{\partial t} + \frac{(\nabla S)^2}{2m} + V = 0 \tag{5.5}$$

It is exactly the Hamilton-Jacobi equation. The quantum potential is

$$V_{\text{quantum}} = -\frac{i}{2m\hbar} \nabla^2 S$$
 (5.6)