

Deriving the Lagrange Equation from the Principle of the Least Action

The **action**, denoted by S , of a physical system, is a **functional** of $x(t)$:

$$S = \int_{t_1}^{t_2} L(x, \dot{x}) dt.$$

Mathematically, the **Principle of the Least Action** is $\delta S = 0$, where δ means a *infinitesimal* change of S .

We now look into a possible path near the actual path $x(t)$:

$$\tilde{x}(t) = x(t) + \delta x(t).$$

The two paths above take the same positions at t_1 and t_2 . Hence,

$$\tilde{S} - S = \int_{t_1}^{t_2} (L(\tilde{x}, \dot{\tilde{x}}; t) - L(x, \dot{x}; t)) dt, \quad (1)$$

$$\delta x(t_1) = \delta x(t_2) = 0. \quad (2)$$

As the difference δx between $\tilde{x}(t)$ and $x(t)$ is infinitesimal, we can do the **Taylor expansion** as follows:

$$\begin{aligned} L(\tilde{x}, \dot{\tilde{x}}; t) &= L(x, \dot{x}; t) + \left(\frac{\partial L}{\partial x} \delta x + \frac{\partial L}{\partial \dot{x}} \delta \dot{x} \right) + O(\delta^2) \\ &= L(x, \dot{x}; t) + \left[\frac{\partial L}{\partial x} \delta x + \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \delta x \right) - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) \delta x \right] + O(\delta^2). \end{aligned} \quad (3)$$

With Eq. (1), (2) and (3), we could have:

$$\begin{aligned} \tilde{S} - S &= \delta S = \int_{t_1}^{t_2} \left[\frac{\partial L}{\partial x} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) \right] \delta x dt + \left(\frac{\partial L}{\partial \dot{x}} \delta x \right) \Big|_{t_1}^{t_2} + O(\delta^2) \\ &= \int_{t_1}^{t_2} \left[\frac{\partial L}{\partial x} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) \right] \delta x dt. \end{aligned}$$

According to the Principle of the Least Action, $\delta S = 0$, then there must be

$$\frac{\partial L}{\partial x} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) = 0, \text{ which is so-called the } \mathbf{Lagrange \text{ equation}.}$$