

## I. INTRODUCTION

Bulk nucleonic matter is interesting for several reasons. The equation of state (EoS) of neutron matter determines properties of supernova explosions [1], and of neutron stars [2–7], and it links the latter to neutron radii in atomic nuclei [8–10] and the symmetry energy [11, 12]. Similarly, the compressibility of nuclear matter is probed in giant dipole excitations [13], and the symmetry energy of nuclear matter is related to the difference between proton and neutron radii in atomic nuclei [14–16]. The saturation point of nuclear matter determines bulk properties of atomic nuclei, and is therefore an important constraint for nuclear energy-density functionals and mass models (see, e.g., Refs. [17, 18]).

The determination and our understanding of the EoS for nuclear matter is intimately linked with our capability to solve the nuclear many-body problem. Here, correlations beyond the mean field play an important role. Theoretical studies of nuclear matter and the pertinent EoS span back to the very early days of nuclear many-body physics. Early computations are nicely described in the 1967 review by Day [19]. These early calculations were performed using Brueckner-Bethe-Goldstone theory [20, 21], see Refs.[3, 22, 23, 39] for recent reviews and developments. In these calculations, mainly particle-particle correlations were summed to infinite order. Other correlations were often included in a perturbative way. Coupled-cluster calculations of nuclear matter were performed already during the late 1970s and early 1980s [24, 25]. In recent years, there has been a considerable algorithmic development of first-principle methods for solving the nuclear many-body problem. A systematic inclusion of other correlations in a non-perturbative way are nowadays accounted for in Monte Carlo methods [26–30], self-consistent Green’s function approaches [23, 31–34] and nuclear density functional theory [16, 18].

## MANY-BODY APPROACHES TO INFINITE MATTER

For an infinite homogeneous system like nuclear matter or the electron gas, the one-particle wave functions are given by plane wave functions normalized to a volume  $\Omega$

for a quadratic box with length  $L^1$

$$\psi_{\mathbf{k}m_s}(\mathbf{r}) = \frac{1}{\sqrt{\Omega}} \exp(i\mathbf{k}\mathbf{r})\xi_{m_s}$$

where  $\mathbf{k}$  is the wave number and  $\xi_{m_s}$  is a spin function for either spin up or down

$$\xi_{m_s=+1/2} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \xi_{m_s=-1/2} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

The single-particle energies for the three-dimensional electron gas are

$$\varepsilon_{n_x, n_y, n_z} = \frac{\hbar^2}{2m} \left( \frac{2\pi}{L} \right)^2 (n_x^2 + n_y^2 + n_z^2),$$

resulting in the magic numbers 2, 14, 38, 54, etc.

In general terms, our Hamiltonian contains at most a two-body interactions. In second quantization, we can write our Hamiltonian as

$$\hat{H} = \sum_{pq} \langle p | \hat{h}_0 | q \rangle a_p^\dagger a_q + \frac{1}{4} \sum_{pqrs} \langle pq | v | rs \rangle a_p^\dagger a_q^\dagger a_s a_r, \quad (1)$$

where the operator  $\hat{h}_0$  denotes the single-particle Hamiltonian, and the elements  $\langle pq | v | rs \rangle$  are the anti-symmetrized Coulomb interaction matrix elements. Normal-ordering with respect to a reference state  $|\Phi_0\rangle$  yields

$$\hat{H} = E_0 + \sum_{pq} f_{pq} \{a_p^\dagger a_q\} + \frac{1}{4} \sum_{pqrs} \langle pq | v | rs \rangle \{a_p^\dagger a_q^\dagger a_s a_r\}, \quad (2)$$

where  $E_0 = \langle \Phi_0 | \hat{H} | \Phi_0 \rangle$  is the reference energy and we have introduced the so-called Fock matrix element defined as

$$f_{pq} = \langle p | \hat{h}_0 | q \rangle + \sum_i \langle pi | v | qi \rangle. \quad (3)$$

The curly brackets in Eq. (2) indicate that the creation and annihilation operators are normal ordered.

The unperturbed part of the Hamiltonian is defined as the sum over all the single-particle operators  $\hat{h}_0$ , resulting in

$$\hat{H}_0 = \sum_i \langle i | \hat{h}_0 | i \rangle = \sum_{\mathbf{k}_i m_{s_i}} \frac{\hbar^2 k_i^2}{2m} a_{\mathbf{k}_i m_{s_i}}^\dagger a_{\mathbf{k}_i m_{s_i}}.$$

We will throughout suppress, unless explicitly needed, all references to the explicit quantum numbers  $\mathbf{k}_i m_{s_i}$ .

<sup>1</sup> The limit  $L \rightarrow \infty$  is to be taken after we have computed various expectation values. In our case we will however always deal with a fixed number of particles and finite size effects become important.

The summation index  $i$  runs over all single-hole states up to the Fermi level.

The general anti-symmetrized two-body interaction matrix element

$$\langle pq|v|rs\rangle = \langle \mathbf{k}_p m_{s_p} \mathbf{k}_q m_{s_q} | v | \mathbf{k}_r m_{s_r} \mathbf{k}_s m_{s_s} \rangle,$$

is given by the following expression

$$\begin{aligned} & \langle \mathbf{k}_p m_{s_p} \mathbf{k}_q m_{s_q} | v | \mathbf{k}_r m_{s_r} \mathbf{k}_s m_{s_s} \rangle \\ &= \frac{e^2}{\Omega} \delta_{\mathbf{k}_p + \mathbf{k}_q, \mathbf{k}_r + \mathbf{k}_s} \left\{ \delta_{m_{s_p} m_{s_r}} \delta_{m_{s_q} m_{s_s}} (1 - \delta_{\mathbf{k}_p \mathbf{k}_r}) \frac{4\pi}{\mu^2 + (\mathbf{k}_r - \mathbf{k}_p)^2} \right. \\ & \quad \left. - \delta_{m_{s_p} m_{s_s}} \delta_{m_{s_q} m_{s_r}} (1 - \delta_{\mathbf{k}_p \mathbf{k}_s}) \frac{4\pi}{\mu^2 + (\mathbf{k}_s - \mathbf{k}_p)^2} \right\}, \end{aligned}$$

for the three-dimensional electron gas. The energy per electron computed with the reference Slater determinant can then be written as (using hereafter only atomic units, meaning that  $\hbar = m = e = 1$ )

$$E_0/N = \frac{1}{2} \left[ \frac{2.21}{r_s^2} - \frac{0.916}{r_s} \right],$$

for the three-dimensional electron gas. This will serve (with the addition of a Yukawa term) as the first benchmark in setting up a program for doing Coupled-cluster theories. This work will be based partly on the Master thesis project of Audun Skau Hansen. In addition to the above studies of the electron gas, it is important to study properly the boundary conditions as well. Here we will study so-called periodic boundary conditions following the work of Refs. [35–37]. See also the recent work of Hagen *et al.* and Baardsen, see Refs. [38, 39]. One aim of this of this thesis is to implement so-called twisted boundary conditions, first for the simple electron gas at the Hartree-Fock level. The next step is to perform a coupled-cluster calculations at the coupled-cluster with doubles excitations, the so-called CCD approach, see [40], for the electron gas. The next step is to include triples

correlations and perform full doubles and triples calculations (CCDT) for the electron gas. Parts of these tasks will also be studied by another Master of Science student at the computational physics group.

The main aim is to replace a Coulomb interaction with nuclear forces and study the EoS for pure neutron matter, asymmetric nuclear matter (for different proton fractions) and symmetric nuclear matter. With this one can study  $\beta$ -stable neutron star matter and extract important information about the symmetry energy in infinite matter and the composition of a neutron star.

## PROGRESS PLAN AND MILESTONES

The aims and progress plan of this thesis are as follows

- Spring 2015 and partly fall 2015: The first step is to set up a program that uses cartesian coordinates in order to compute the Hartree-Fock energy of the electron gas with a Coulomb interaction with and without a screening term (Yukawa like potential).
- Fall 2015: The next step is to include a proper treatment of periodic boundary conditions using the twisted boundary condition approach described in Refs. [35–37].
- Fall 2015: Thereafter, correlations at the electron gas level will be studied using coupled-cluster theory at the doubles level, see [40].
- Spring 2016: With a CCD code, the next step is to include triples correlations, leading to a CCDT code.
- Spring 2016: The final step is to perform systematic studies of  $\beta$ -stable nuclear matter.
- Spring 2016: The last part deals with a proper write-up of the thesis.

The thesis is expected to be handed in May/June 2016.

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