

Fundamental Concepts of Quantum Mechanics

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Aug. 18, 2016

* Roll Call

* Syllabus & HW Guideline.

Stern - Gerlach Expt. (Otto Stern & Walther Gerlach)

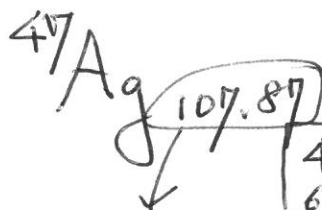
O. Stern in 1921 at Frankfurt

O. Stern & W. Gerlach in 1922

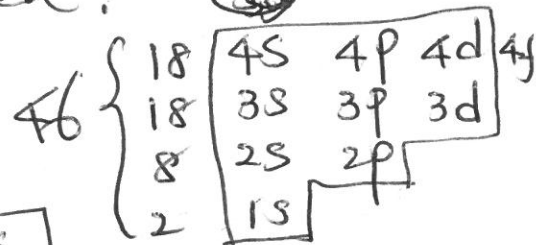
Most quantum mechanical, least classical system

- 1) Use charge zero particle because you don't want the Lorentz force to deflect the trajectory.
- 2) Atom as a whole is very heavy so that the classical trajectory can be legitimately applied.

Silver Atom



47 protons
61 neutrons



in 1 amu \approx mass of a single nucleon (proton or neutron)

$\mu_{\text{Ag}} \approx \mu_{e^-}$ since the atom is filled up to 4d and only one electron left at 5s level.

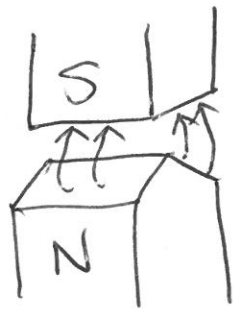
$\mu_z = \frac{e}{m_e c} S_z$ ($e < 0$ for electron) = $g \left(\frac{e}{2m_e c} \right) S_z$ ($g=2$ for electron)

of Bohr magneton $\mu_B = \frac{eh}{4\pi m_e c}$

electrons intrinsic spin magnetic moment

3) Make inhomogeneous magnetic field to split different magnetic moments. (2)

$$V(\vec{r}) = -\vec{\mu} \cdot \vec{B}(\vec{r}) \Rightarrow \vec{F} = -\vec{\nabla} V(\vec{r})$$



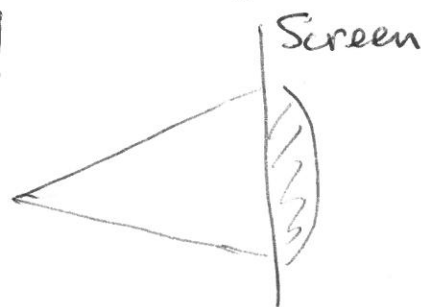
$$F_z = \frac{\partial}{\partial z} (\vec{\mu} \cdot \vec{B}) = \mu_z \frac{\partial B_z}{\partial z}$$

$$\frac{\partial B_z}{\partial z} < 0, e < 0 \Rightarrow F_z \propto S_z$$

Spin aligns to the direction of force

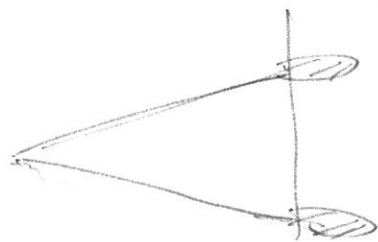
If the spin is a classical object,

then $-\mu_z < \mu_z < \mu_z$



continuous distribution
expected from
classical physics.

However, expts show "space quantization" !!



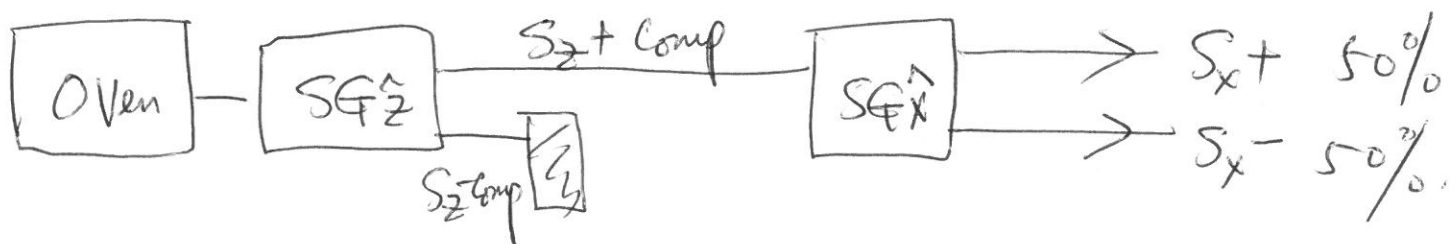
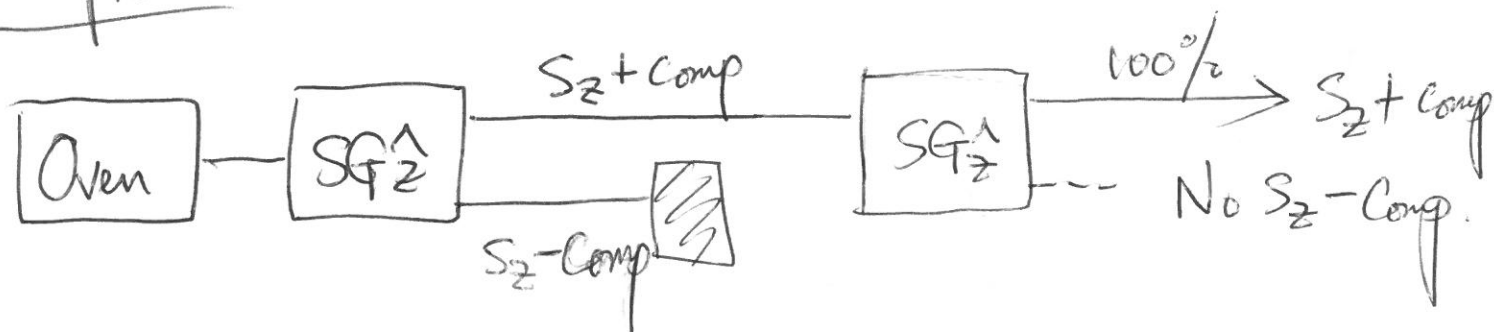
Totally different
from C.M.

$$\left. \begin{array}{l} S_z \text{ + up; } S_z = +\frac{\hbar}{2} \\ S_z \text{ - down; } S_z = -\frac{\hbar}{2} \end{array} \right\} \begin{array}{l} \text{Planck's constant} \\ \hbar = 1.0546 \times 10^{-27} \text{ erg-s} \\ = 6.5822 \times 10^{-16} \text{ eV-s} \end{array}$$

Quantization of electron spin

Expts

(3)



Does it mean that 50% of the atoms in the $S_z +$ beam coming out of the first apparatus (SG_z^1) are made up of atoms characterized by both $S_z +$ and $S_x +$, while the remaining 50% have both $S_z +$ and $S_x -$?

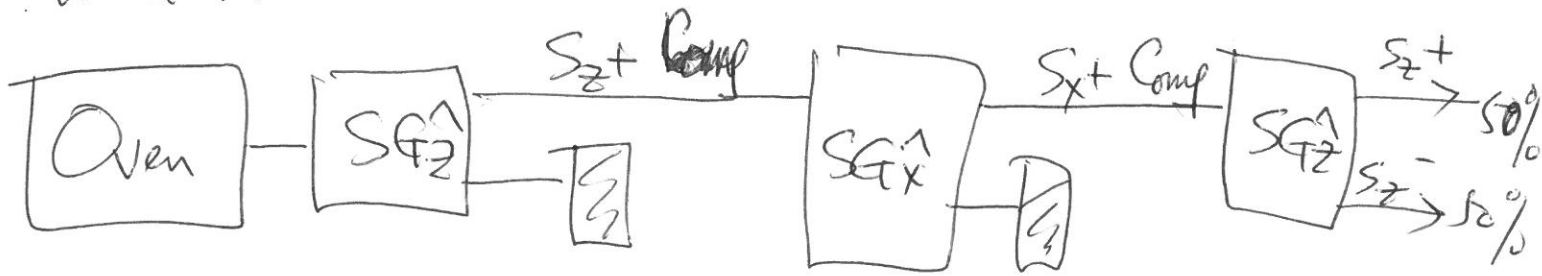
In classical picture, the spinning top with

$$\vec{L} = \vec{I} \cdot \vec{\omega}$$

has no problem in determining L_z and L_x simultaneously. In fact, no problem in classical mechanics to determine L_x , L_y and L_z once \vec{I} and $\vec{\omega}$ are known. \vec{I} is computable if we know the mass density and the geometric shape of the spinning top. Thus, there is no difficulty in specifying all the components of angular momentum simultaneously.

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Now, let's do another expt. to pin down the reality in Q.M.



The model in which the atoms ~~entering~~ the third apparatus are visualized to have both S_z+ and S_x+ is clearly unsatisfactory.

S_z and S_x cannot be determined simultaneously.

S_x+ beam from the second apparatus SG_x completely destroys any previous information about S_z (wipes out).

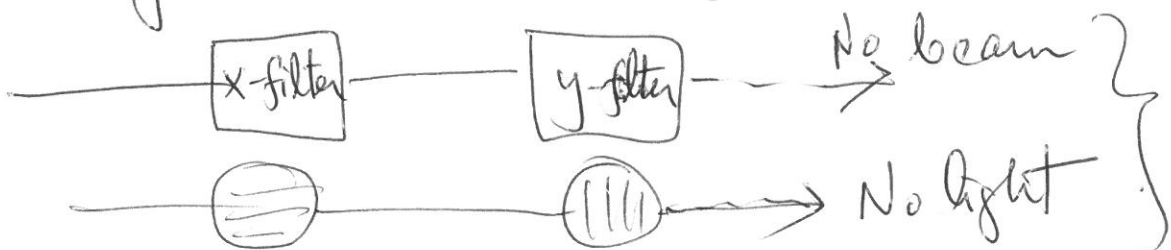
Analogy

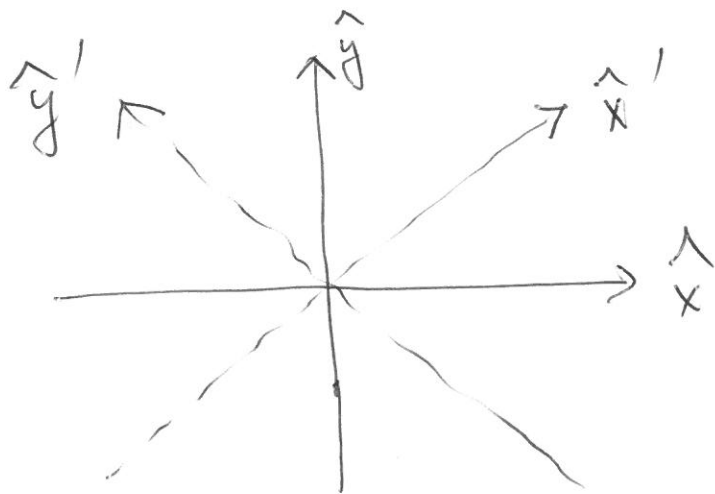
with polarization of light.

$S_z+ \leftrightarrow x\text{-polarized}$
 $S_z- \leftrightarrow y\text{-polarized}$

$$\vec{E} = \begin{cases} E_0 \hat{x} \cos(kz - \omega t) \\ E_0 \hat{y} \cos(kz - \omega t) \end{cases}$$

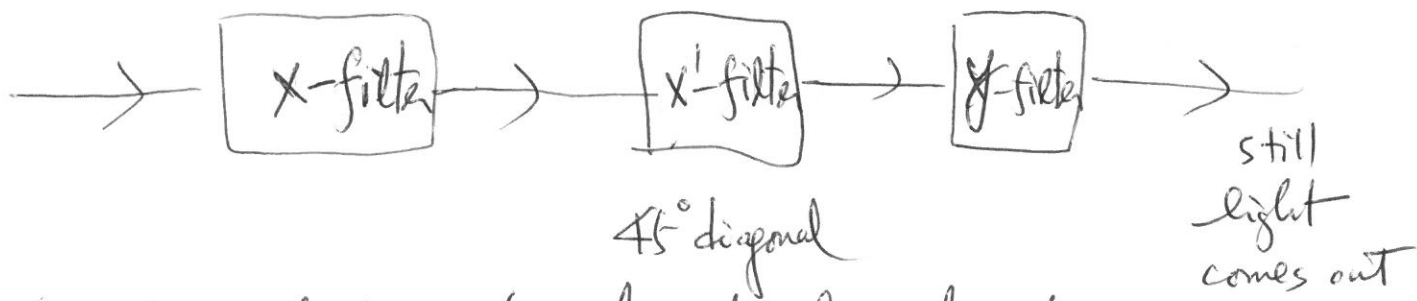
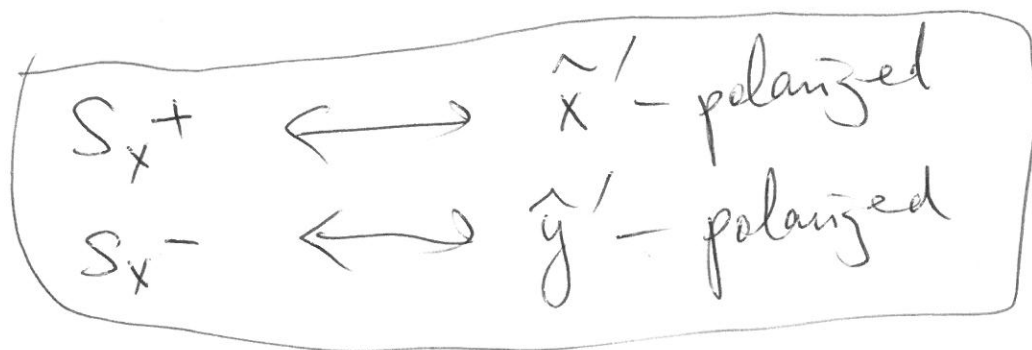
x-polarized light
 y-polarized light.





$$\hat{x}' = \frac{1}{\sqrt{2}}(\hat{x} + \hat{y})$$

$$\hat{y}' = \frac{1}{\sqrt{2}}(-\hat{x} + \hat{y})$$



The selection of the x' -polarized beam by the second polaroid destroys any previous information on light polarization
(wipes out)

We may thus identify

$$|S_x^+\rangle = \frac{1}{\sqrt{2}}|S_z^+\rangle + \frac{1}{\sqrt{2}}|S_z^-\rangle$$

$$|S_x^-\rangle = -\frac{1}{\sqrt{2}}|S_z^+\rangle + \frac{1}{\sqrt{2}}|S_z^-\rangle$$

How about $|S_y^+\rangle = ?$ and $|S_y^-\rangle = ?$

$$\vec{E} = E_0 \left[\frac{1}{\sqrt{2}} \hat{x} \cos(kz - \omega t) + \frac{1}{\sqrt{2}} \hat{y} \cos(kz - \omega t + \frac{\pi}{2}) \right] \quad (6)$$

$$\frac{\vec{E}}{E_0} = \text{Re}(\vec{E})$$

Complex Representation

$$\vec{E} = \left[\frac{1}{\sqrt{2}} \hat{x} e^{i(kz - \omega t)} + \frac{1}{\sqrt{2}} \hat{y} e^{i(kz - \omega t + \frac{\pi}{2})} \right]$$

circularly polarized

$$\vec{E} = \left[\frac{1}{\sqrt{2}} \hat{x} e^{i(kz - \omega t)} + \frac{i}{\sqrt{2}} \hat{y} e^{i(kz - \omega t)} \right]$$

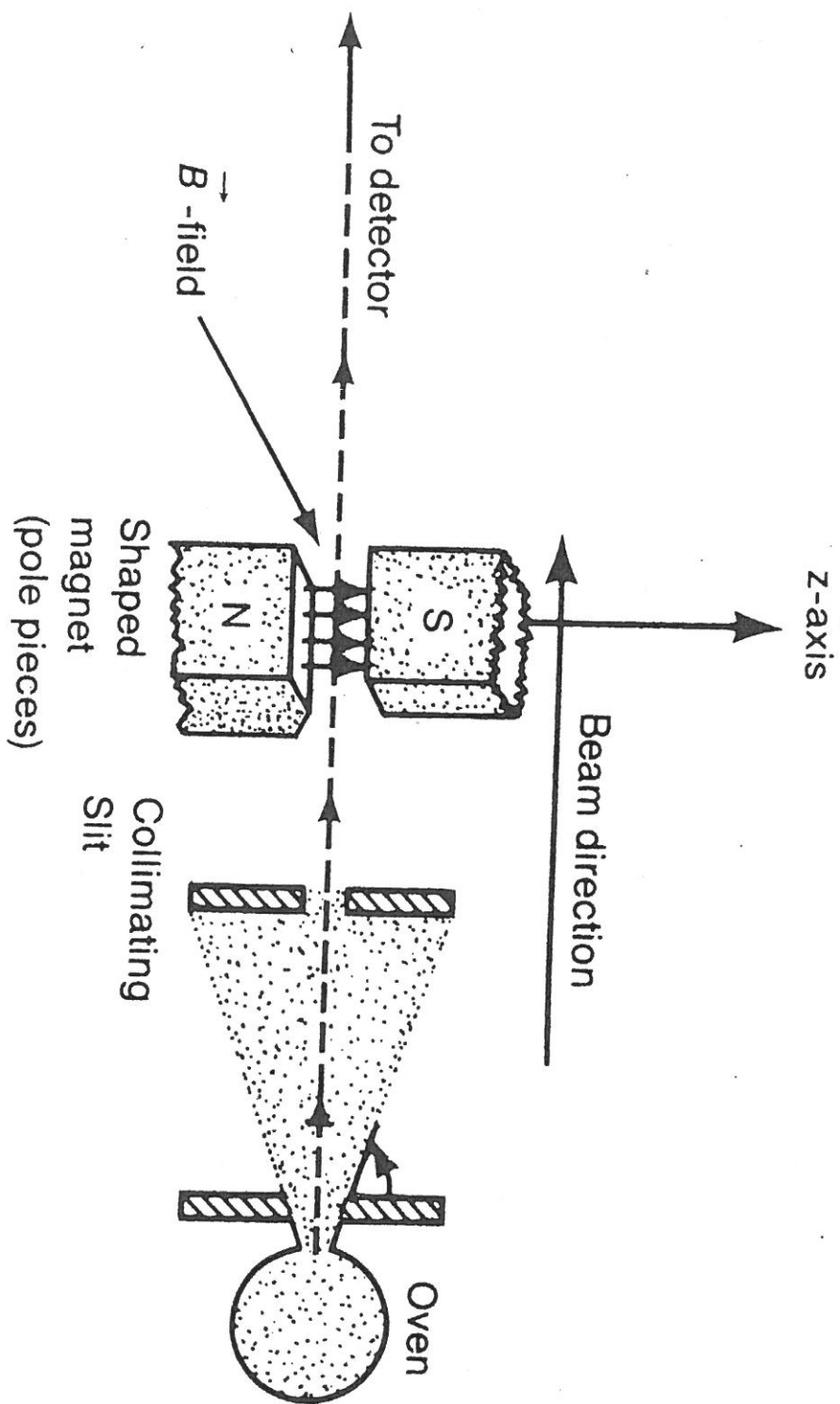
$S_y^+ \leftrightarrow$ right circularly polarized
 $S_y^- \leftrightarrow$ left "

$(\because e^{i\frac{\pi}{2}} = i)$
 relative phase factor

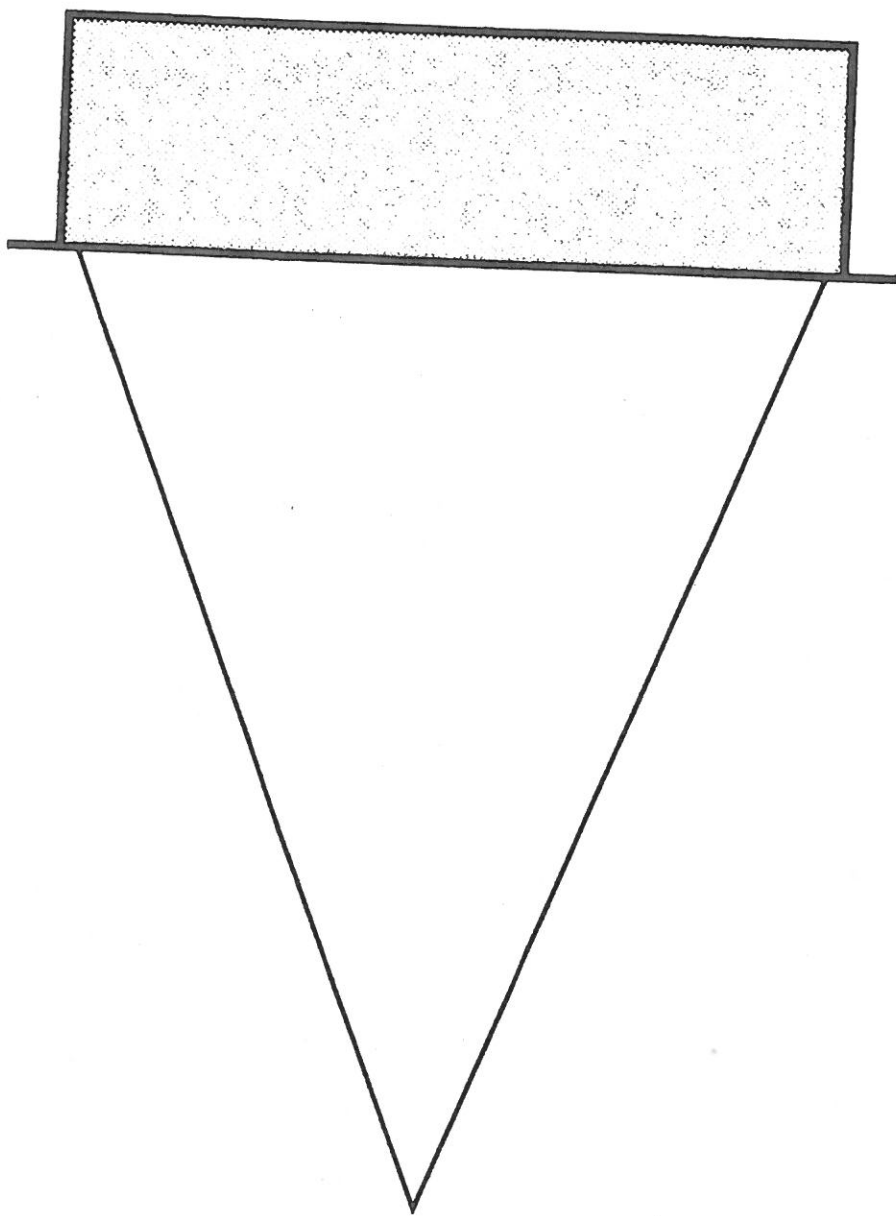
$$|S_y^+\rangle = \frac{1}{\sqrt{2}} |S_z^+\rangle + \frac{i}{\sqrt{2}} |S_z^-\rangle$$

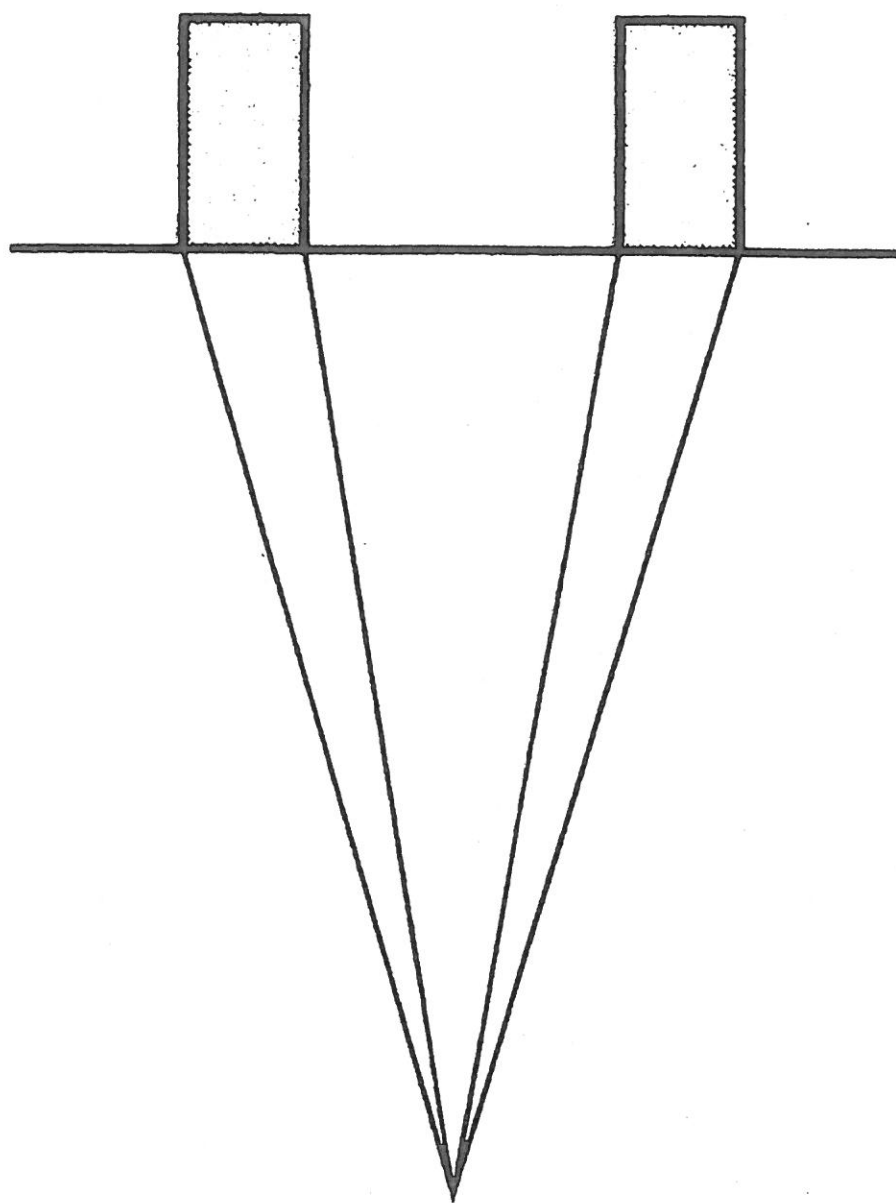
$$|S_y^-\rangle = \frac{1}{\sqrt{2}} |S_z^+\rangle - \frac{i}{\sqrt{2}} |S_z^-\rangle$$

Complex Vector space \Rightarrow Hilbert Space
(Bra & Ket space)



The Stern-Gerlach experiment.





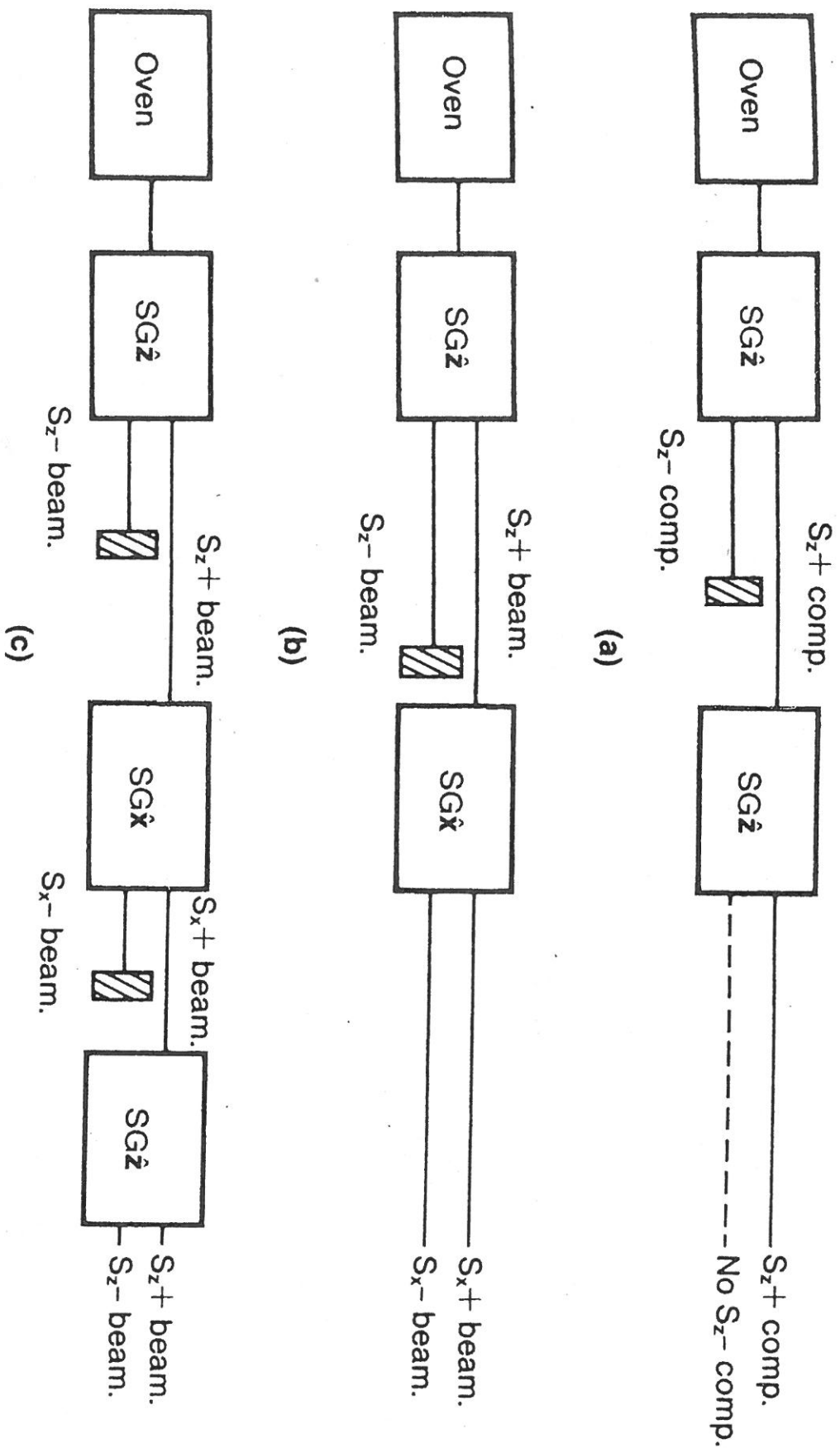
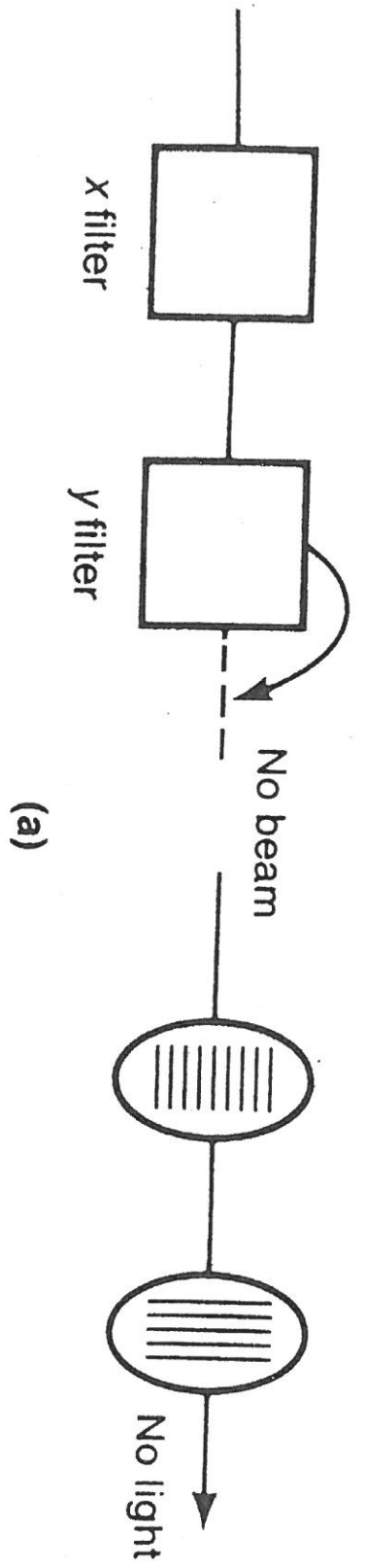
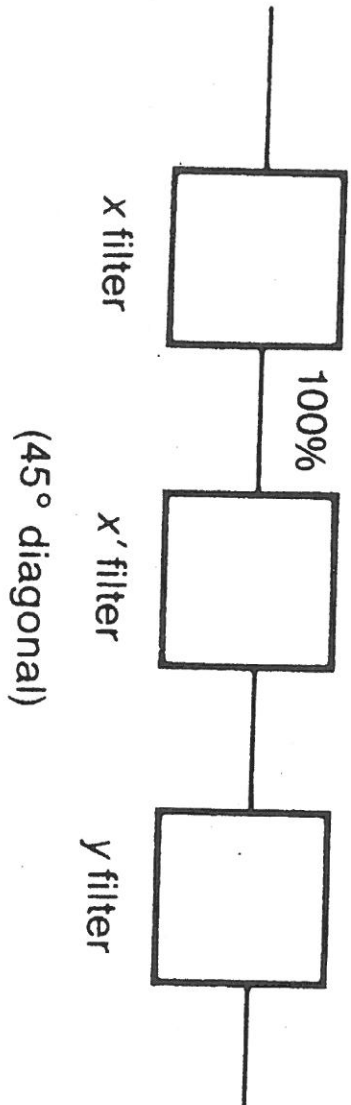


FIGURE 1.3. Sequential Stern-Gerlach experiments.



(a)



(45° diagonal)

(b)

FIGURE 1.4. Light beams subjected to Polaroid filters.

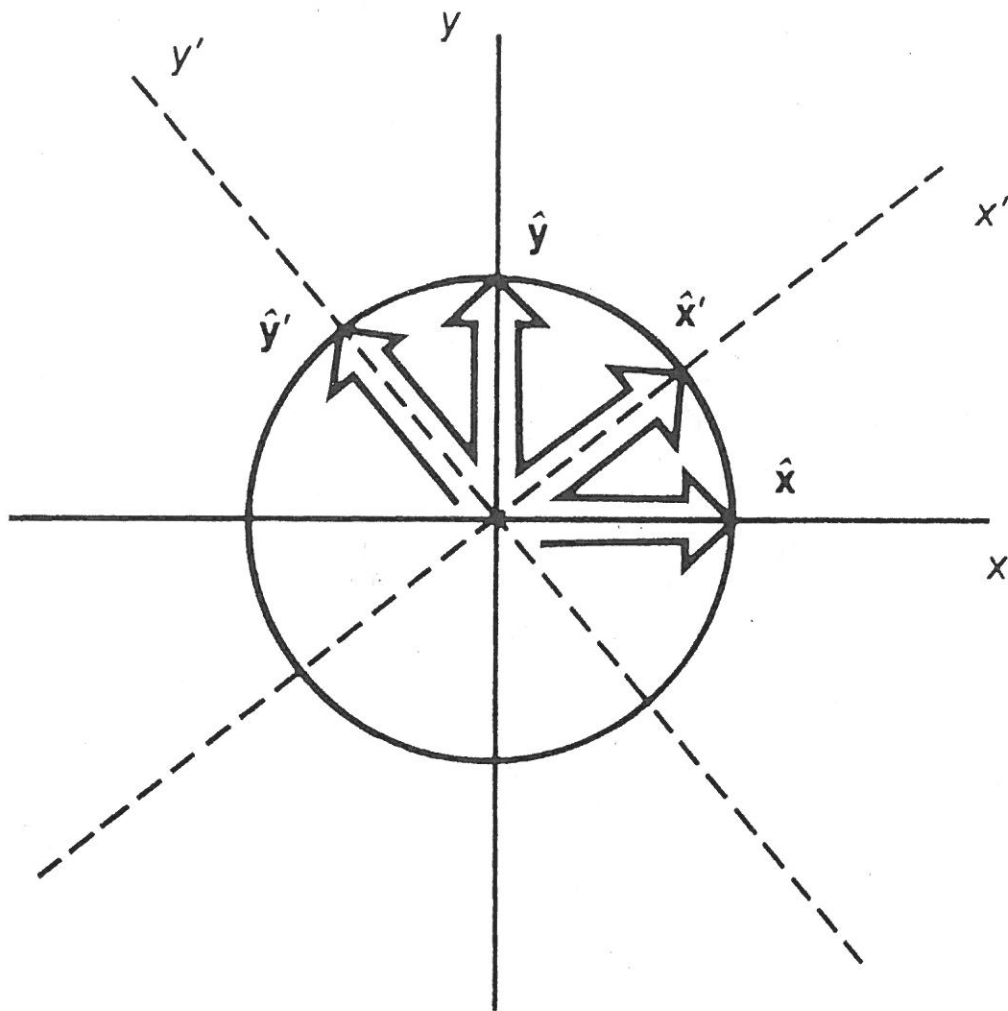


FIGURE 1.5. Orientations of the x' - and y' -axes.

assical electrodynamics. Using Figure 1.5 we obtain

$$\begin{aligned}
 E_0 \hat{x}' \cos(kz - \omega t) &= E_0 \left[\frac{1}{\sqrt{2}} \hat{x} \cos(kz - \omega t) + \frac{1}{\sqrt{2}} \hat{y} \cos(kz - \omega t) \right], \\
 E_0 \hat{y}' \cos(kz - \omega t) &= E_0 \left[-\frac{1}{\sqrt{2}} \hat{x} \cos(kz - \omega t) + \frac{1}{\sqrt{2}} \hat{y} \cos(kz - \omega t) \right].
 \end{aligned}
 \tag{1.1.8}$$