

# Computation and Properties of the Epstein Zeta Function

with high-performance implementation in EpsteinLib

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Joint work with Andreas A. Buchheit and Ruben Gutendorf

Python: `pip install epsteinlib`

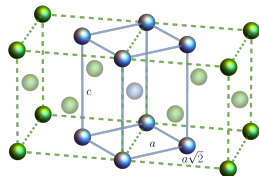
GitHub: [github.com/epsteinlib](https://github.com/epsteinlib)

Preprint: [arXiv:2412.16317](https://arxiv.org/abs/2412.16317)

20. March 2025



# Broad Applicability in Long-Range Interacting Lattices

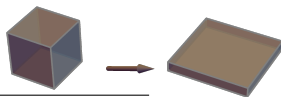


**Theoretical Chemistry:**  
Cohesive energy for Lennard-Jones<sup>1</sup>

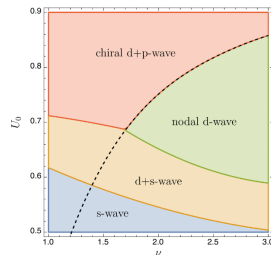
$$E_{\text{LJ}}(\mathbf{r}) \propto \frac{1}{n}|\mathbf{r}|^{-n} - \frac{1}{m}|\mathbf{r}|^{-m}$$

**High Energy Physics:**  
Casimir Energy<sup>2</sup>

$$E_{\text{cas}} = \frac{1}{2} \sum_{\mathbf{k}} \lambda_{\mathbf{k}}^{-\nu} \Big|_{\nu=-1}$$



**Spin Systems:**<sup>3</sup>  
$$H = \sum_{ij \in \Lambda} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$



**Superconductivity:**  
Long range gap equation<sup>4</sup>

$$\Delta(\mathbf{k}) = \int_{\text{BZ}} V_{\text{LR}}(\mathbf{k} - \mathbf{q}) \frac{\Delta(\mathbf{q})}{2E_{\Delta}(\mathbf{q})} d\mathbf{q}$$

<sup>1</sup>Buchheit et al., Computation and Properties of the Epstein Zeta Function (Dec. 2024)

<sup>2</sup>Schwerdtfeger et al., *Journal of Chemical Theory and Computation* 20, 3379–3405 (May 2024)

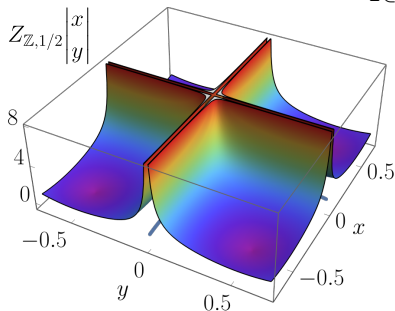
<sup>3</sup>Buchheit et al., *Phys. Rev. Res.* 5, 043065 (Oct. 2023)

<sup>4</sup>Ambjørn et al., *Annals of Physics* 147, 1–32 (Aug. 1983)

# Epstein Zeta Function

Let  $\mathbf{x}, \mathbf{y} \in \mathbb{R}^d$  and  $\Lambda = A\mathbb{Z}^d$  be a lattice. The Epstein zeta function is the meromorphic continuation of the following lattice sum to  $\nu \in \mathbb{C}$

$$Z_{\Lambda, \nu} \begin{vmatrix} \mathbf{x} \\ \mathbf{y} \end{vmatrix} = \sum'_{\mathbf{z} \in \Lambda} \frac{e^{-2\pi i \mathbf{y} \cdot \mathbf{z}}}{|\mathbf{z} - \mathbf{x}|^\nu}, \quad \operatorname{Re}(\nu) > d.$$



Numerous Applications:

- ▶ Spin systems<sup>2</sup>
- ▶ Theoretical chemistry<sup>3</sup>
- ▶ Unconventional superconductors<sup>4</sup>
- ▶ High energy physics: Casimir effect<sup>5</sup>

<sup>5</sup>Buchheit et al., Computation and Properties of the Epstein Zeta Function (Dec. 2024)

<sup>6</sup>Schwerdtfeger et al., *Journal of Chemical Theory and Computation* 20, 3379–3405 (May 2024)

<sup>7</sup>Buchheit et al., *Phys. Rev. Res.* 5, 043065 (Oct. 2023)

<sup>8</sup>Ambjørn et al., *Annals of Physics* 147, 1–32 (Aug. 1983)

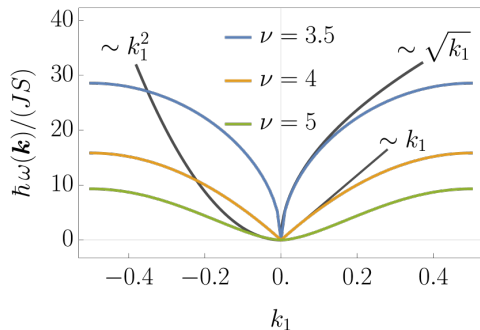
## Example: Anomalous Spin Wave Dispersion<sup>5</sup>

$$H = -c \sum'_{\mathbf{x}, \mathbf{y} \in \Lambda} \frac{\mathbf{S}_{\mathbf{x}} \cdot \mathbf{S}_{\mathbf{y}}}{|\mathbf{x} - \mathbf{y}|^{\nu}}$$

- + Holstein-Primakoff Transformation
- + Fourier Transformation

$$\omega(\mathbf{k}) \sim Z_{\Lambda, \nu} \left| \begin{pmatrix} 0 \\ 0 \end{pmatrix} \right| - Z_{\Lambda, \nu} \left| \begin{pmatrix} 0 \\ \mathbf{k} \end{pmatrix} \right|$$

- $\text{Re}(\nu) < \dim + 2$ : Anomalous scaling  $\omega(\mathbf{k}) \sim |\mathbf{k}|^{\nu-d}$
- $\text{Re}(\nu) > \dim + 2$ : typical scaling  $\omega(\mathbf{k}) \sim \mathbf{k}^2$



<sup>5</sup>Buchheit et al., Computation and Properties of the Epstein Zeta Function (Dec. 2024)

# EpsteinLib



Fast and precise evaluation of Epstein zeta function for

- ▶ every exponent  $\nu$ ,
- ▶ every lattice  $\Lambda$ ,
- ▶ every shift vector  $\mathbf{x}$ ,
- ▶ every wavevector  $\mathbf{y}$ .

```
double complex epsteinZeta(double nu, unsigned int dim, const double *A, const double *x, const double *y);
```

C

# Precision and Runtime

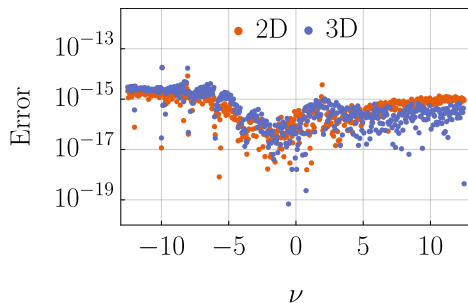
```
pip install epsteinlib
```

Shell

```
import numpy as np
from epsteinlib import epstein_zeta

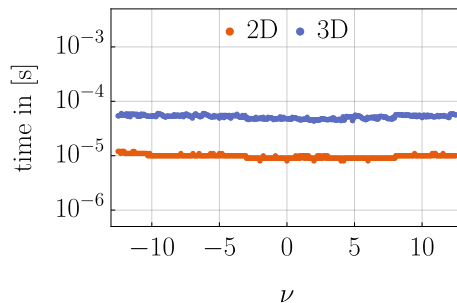
madelung = epstein_zeta(1.0, np.identity(3), np.zeros(3), np.full(3, 0.5))
```

Python



► easy to use

► precise



► ultra fast

# Crandall Representation

One algebraically decaying sum  $\Rightarrow$  two superexponentially decaying sums

$$\sum'_{\mathbf{z} \in \Lambda} \frac{e^{-2\pi i \mathbf{y} \cdot \mathbf{z}}}{|\mathbf{z} - \mathbf{x}|^\nu} = \frac{\pi^{\nu/2}}{\Gamma(\nu/2)} \left[ \sum_{\mathbf{z} \in \Lambda} G_\nu(\mathbf{z} - \mathbf{x}) e^{-2\pi i \mathbf{y} \cdot \mathbf{z}} + \frac{1}{V_\Lambda} \sum_{\mathbf{k} \in \Lambda^*} G_{d-\nu}(\mathbf{k} + \mathbf{y}) e^{-2\pi i \mathbf{x} \cdot (\mathbf{k} + \mathbf{y})} \right],$$

with  $G_\nu(\mathbf{z}) \propto \frac{\exp(-\pi \mathbf{z}^2)}{(\pi \mathbf{z}^2)^{\nu/2}}$  superexponentially decaying.

**Crandall + solving numerical challenges = EpsteinLib**

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<sup>5</sup>Crandall, Unified algorithms for polylogarithm, L-series, and zeta variants (2012)

# Epstein Zeta as a new standard function

**Now efficiently computable ✓**

**Analytical properties derived in arxiv:2412.16317 ✓**

- ▶ Functional Equation.
- ▶ Analyticity in all arguments.



# Epstein Zeta as a new standard function

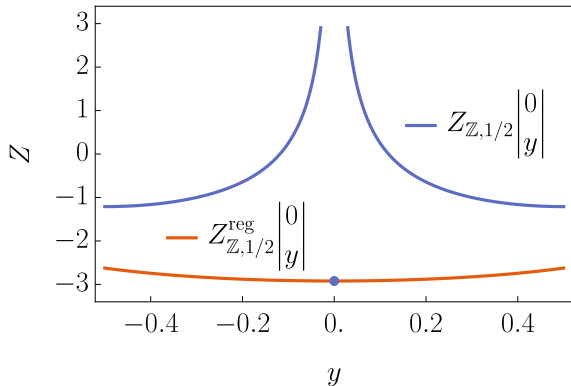
Now efficiently computable ✓

Analytical properties derived in arxiv:2412.16317 ✓

- ▶ Functional Equation.
- ▶ Analyticity in all arguments.

Physically relevant singularities

$$\underbrace{Z_{\Lambda,\nu}\left|\begin{smallmatrix} 0 \\ \mathbf{y} \end{smallmatrix}\right|}_{\text{singularity in } \mathbf{y}=0} = \underbrace{Z_{\Lambda,\nu}^{\text{reg}}\left|\begin{smallmatrix} 0 \\ \mathbf{y} \end{smallmatrix}\right|}_{\text{smooth in } \mathbf{y}\in\text{BZ}} + c_{\Lambda}(\nu)|\mathbf{y}|^{\nu-d}$$



# Conclusions

## **First high-performance library for Epstein zeta function.**

- ▶ General: Any lattice, any interaction exponent, any shift, any wavevector.
- ▶ Fast and precise: exact 3D lattice sums in  $10^{-4}$  seconds, on your laptop.
- ▶ Easy-to-use: `pip install epsteinlib`

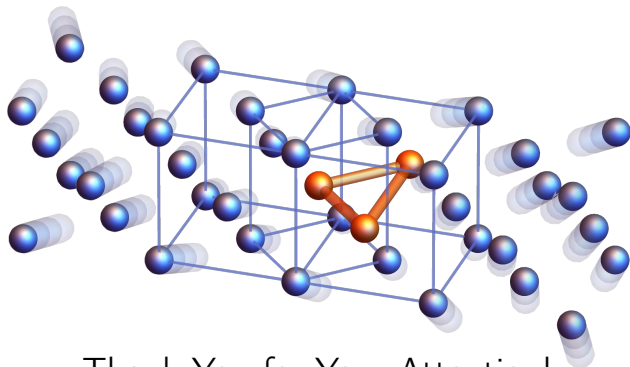
## **Actively used by the community:**

- ▶ Long-range interacting hard-core bosons: SciPost Physics 17.4 (2024): 111,
- ▶ Long-range Dicke-Ising model: arXiv:2503.02734 (2025).

# Outlook

## Next steps:

- ▶ Multi-body interactions.
- ▶ Investigation into stability of matter in theoretical chemistry.



Thank You for Your Attention!

# Unconventional Superconductors

Tight-binding Hamiltonian with power-law long-range interaction<sup>6</sup>

$$H = H_0 + H_{\text{int}}$$
$$H_{\text{int}} = \frac{1}{2} \sum_{\sigma, \sigma'} c_{\sigma, \mathbf{x}}^\dagger c_{\sigma', \mathbf{x} - \mathbf{y}}^\dagger V_{\sigma\sigma'}(\mathbf{y}) c_{\sigma', \mathbf{x} - \mathbf{y}} c_{\sigma, \mathbf{x}}, \quad V_{\sigma\sigma'}(\mathbf{y}) = \begin{cases} -C_{\sigma\sigma'} \leq 0 & \mathbf{y} = 0 \\ -\frac{U_{\sigma\sigma'}}{|\mathbf{y}|^\nu} \leq 0 & \mathbf{y} \neq 0 \end{cases}$$

Fouriertransform:

$$H_{\text{int}} = -\frac{V_\Lambda}{2} \sum_{\sigma, \sigma'} \int_{E^*} \int_{E^*} \oint_{E^*} \left( c_{\sigma\sigma'} + U_{\sigma\sigma'} Z_{\Lambda, \nu} \left| \frac{0}{\mathbf{q}} \right| \right),$$
$$\times c_\sigma^\dagger(\mathbf{k} + \mathbf{q}) c_{\sigma'}^\dagger(\mathbf{k}' - \mathbf{q}) c_{\sigma'}(\mathbf{k}') c_\sigma(\mathbf{k}) d\mathbf{q} d\mathbf{k} d\mathbf{k}'$$

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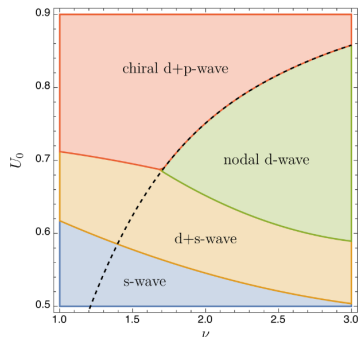
<sup>6</sup>Buchheit et al., *Phys. Rev. Res.* 5, 043065 (Oct. 2023)

# Gap Equation for Unconventional Superconductors

Gap Equation:

$$\Delta(\mathbf{k}) = c \int_{\text{BZ}} \left( C_0 + U_0 Z_{\Lambda, \nu} \left| \begin{smallmatrix} 0 \\ \mathbf{k} - \mathbf{q} \end{smallmatrix} \right| \right) \times \left( \frac{1}{2} \Delta E[\Delta]^{-1} \right) (\mathbf{k} - \mathbf{q}) \, d\mathbf{q}$$

- + fast and efficient computation
- + regularization  $\Rightarrow$



<sup>6</sup>Buchheit et al., *Phys. Rev. Res.* 5, 043065 (Oct. 2023)

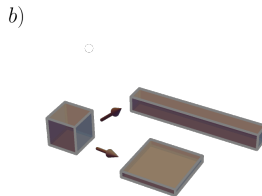
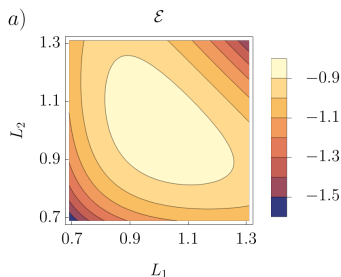
# Casimir Effect

Consider Field modes of the klein gordon wave equation

$$\phi(t, \mathbf{x}) = e^{-i\lambda_{\mathbf{k}}t} e^{i\mathbf{k}\cdot\mathbf{x}}, \quad \lambda_{\mathbf{k}} = \sqrt{m^2 + \mathbf{k}^2}$$

where  $\mathbf{k} = 2\pi(z_1/L_1, \dots, z_d/L_d)^T$  for  $\mathbf{z} \in \mathbb{Z}^d$  and edge-lengths  $0 < L_1, \dots, L_d < \infty$ . In the massless case  $m = 0$  for  $\Lambda^* = \text{diag}(1/L_1, \dots, 1/L_d)$ , the total energy  $\mathcal{E}$  is given by

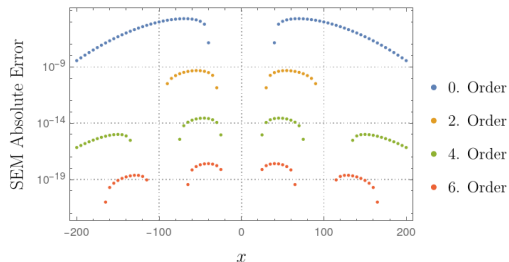
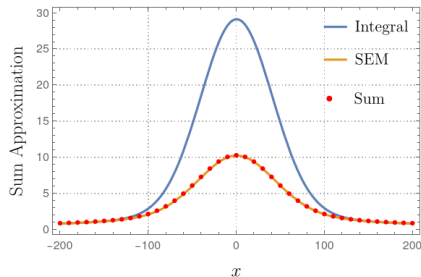
$$\mathcal{E}(\Lambda) = \frac{1}{2} \sum_{\mathbf{k}} \lambda_{\mathbf{k}}^{-\nu} \Big|_{\nu=-1} = \pi Z_{\Lambda^*, -1} \begin{vmatrix} 0 \\ 0 \end{vmatrix}.$$



# Singular Euler-Maclaurin expansion (SEM) Example

$$\overbrace{\sum'_{y=-\infty}^{\infty} \frac{g(y)}{|y-x|^\nu}}^{\text{sum (exact)}} = \underbrace{\frac{1}{V_\Lambda} \int_{-\infty}^{\infty} \frac{g(y)}{|y-x|^\nu} dy}_{\text{integral}} + \underbrace{\sum_{n=0}^{\text{order}} \frac{1}{n} \left( \frac{\partial_y}{-2\pi i} \right)^n Z_{\mathbb{Z},\nu}^{\text{reg}} \Big|_x}_{\text{lattice contribution}} \partial_x^n g(x)$$

SEM (exact for order =  $\infty$ )



<sup>6</sup>Buchheit et al., *Nonlinearity* 35, 3706 (2022)

# Computational Challenges

- ▶ Stable evaluation close to singularities
- ▶ Round-off error in the sum
- ▶ Choice of adequate cutoff
- ▶ Stable evaluation of incomplete gamma function for negative  $\nu$

