

# LECTURE 24 Hopf Bifurcations + Bifurcations of limit cycles

$$\dot{x} = f(x, \mu)$$

suppose at  $\mu=0$ ,  $x=0$  is a fixed pt  
and  $Df(0,0)$  has 2 imaginary eigenvalues  
Then for  $\mu$  non (complex conjugates  $\lambda = \pm i\omega$ )  
then when  $\mu \neq 0$ , these should move off imaginary  
axis. This is called a Hopf bifurcation

Canonical example

$$\begin{aligned}\dot{x} &= -y + x(\mu - x^2 - y^2) \\ \dot{y} &= x + y(\mu - x^2 - y^2)\end{aligned}$$

$$Df(0, \mu) = \begin{bmatrix} \mu & -1 \\ 1 & \mu \end{bmatrix}, \lambda = \mu \pm i$$

so  $\mu > 0 \Rightarrow$  origin unstable

of course, put in polar coordinates

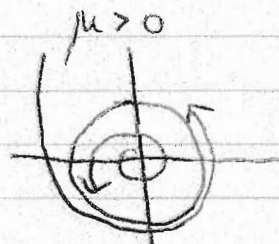
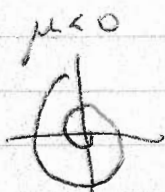
$$\dot{r} = r(\mu - r^2)$$

$$\dot{\theta} = 1$$

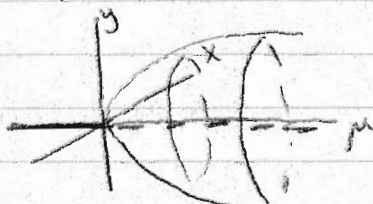
(note similarity to pitchfork)

$\mu < 0$  origin stable,  $\mu = 0$  origin nonlinear stable spiral  
 $\mu > 0$ , origin unstable and

$$r = \sqrt{\mu}, \Gamma = \sqrt{\mu}(\cos t) \text{ stable limit cycle}$$



HOPF BIFURCATION



This is a supercritical Hopf bifurcation - stable limit  
cycles beyond pt where fixed point  
loses stability

~~STABILITY~~ other case

$$\dot{x} = -y + x(\mu + x^2 + y^2)$$

$$\dot{y} = x + y(\mu + x^2 + y^2)$$



## LECTURE 24.2

Universal unfolding of the Hopf bifurcation (Theorem 2)

$$\dot{x} = \mu x - y + ax(x^2 + y^2) - by(x^2 + y^2)$$

$$\dot{y} = x + \mu y + bx(x^2 + y^2) + ay(x^2 + y^2)$$

$$r\dot{r} = x\dot{x} + y\dot{y} = \mu x^2 + \mu y^2 + a(x^2 + y^2)^2 = \mu r^2 + ar^4$$

$$\dot{r} = r(\mu + ar^2)$$

$$\theta = \frac{y\dot{x} - x\dot{y}}{x^2 + y^2} = \frac{x^2(1 + br^2) + y^2(1 + br^2)}{r^2} = 1 + br^2$$

supercriticality depends on sign  $\frac{a}{\mu}$

Text has (Theorem 1) analytic test for determining whether supercritical or subcritical

$$\dot{x} = \mu x - y + (a_{20}x^2 + a_{11}xy + a_{02}y^2) + (a_{30}x^3 + a_{21}x^2y + a_{12}xy^2 + a_{03}y^3) + \dots$$

$$\dot{y} = x + \mu y + (b_{20}x^2 + b_{11}xy + b_{02}y^2) + (b_{30}x^3 + b_{21}x^2y + b_{12}xy^2 + b_{03}y^3) + \dots$$

criticality depends on the sign of

$$\sigma = \frac{3\pi}{2} [3(a_{30} + b_{03}) + (a_{12} + b_{21}) - 2(a_{20}b_{02} - a_{02}b_{20}) + a_{11}(a_{02} + b_{20}) - b_{11}(b_{02} + b_{20})]$$

$\sigma < 0 \Rightarrow$  supercritical

example

$$\begin{aligned} \dot{x} &= \mu x - y + x^2 \\ \dot{y} &= x + \mu y + x^2 \end{aligned}$$

$$a_{20} = b_{20} = 1 \text{ all others } = 0$$

$$\text{So } \sigma = \frac{3\pi}{2} (-2 \cdot 1 \cdot 1) = -3\pi < 0$$

$\Rightarrow$  supercritical

$\Rightarrow \exists$  stable limit cycle for small  $\mu > 0$

Lacking this, one can try to use Poincaré-Bendixson theorem or other methods for showing periodic orbits do or do not exist for  $\mu > 0$  or  $\mu < 0$ .

I do not want you to learn this.

If you ever need it, look it up

# LECTURE 24-3

## Bifurcations of periodic orbits

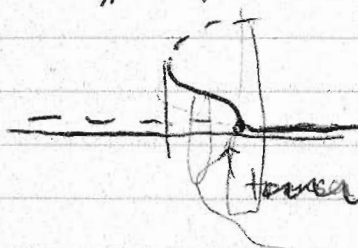
some silly examples: take your favorite bifurcation

suppose that  $\dot{\theta} = 1$

$$\text{and } \dot{r} = r(\mu - (r-1)^2)$$

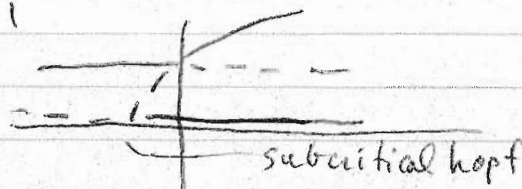
and write  $r$  in such a way that it has that bifurcation at  $r=1, \mu=0$

saddle-node  $\dot{r} = r(\mu - (r^2-1)^2)$  no ~~for~~ periodic orbits  $\mu < 0$   
fixed periodic orbits at  $r = \sqrt{1 \pm \sqrt{\mu}}$



transcritical Hopf bifurcation at  $\mu=1$

transcritical  $\dot{r} = -r(1-r^2)(1+\mu-r^2)$   
 $\dot{\theta} = 1$



$$r=0, r=1, r = \sqrt{1+\mu}$$

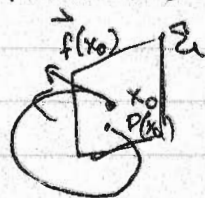
PITCHFORK  $\dot{r} = r(1-r^2)(\mu - (r^2-1)^2)$

$$r=0, r=1, r = \sqrt{1 \pm \sqrt{\mu}}$$



so what? These examples are a bit contrived

Recall The main tool for studying periodic orbits is the Poincaré map



if  $P(x_0) = x_0$  then there is a periodic orbit through  $x_0$

if for some value of  $\mu$   $f'(x^*, \mu) = 1$  we are in one of the above situations

Lecture ~~23~~ 24-25 24-4

$$x_{n+1} = f(x_n, \mu) \quad \text{s.t.} \quad x^*(\mu) = f(x^*(\mu), \mu)$$

Suppose for  $\mu < 0$ ,  $-1 < D_x f(x^*(\mu), \mu) < 1$   
and for  $\mu > 0$ ,  $D_x f(x^*(\mu), \mu) < -1$

- ~~the above bifurcations take place where  $f'(x^*(\mu)) = 1$~~
- then the system has a flip bifurcation or a period-doubling bifurcation

Notice: if  $f(0) = 0$  and  $f'(0) = -1$  were purely linear, then  
 $x_{n+1} = f(x_n) = -x_n \Rightarrow x = 0$  the only fixed pt  
all other initial conditions satisfy  $x_{n+2} = x_n$   
these are period-2 orbits of this simple map

In fact, near a flip bifurcation nearby ~~stable~~ periodic orbits always appear

Assume  $x_{n+1} = f(x_n, \mu)$ ,  $f(0,0) = 0$ ,  $f_x(0,0) = -1$   
 $= A(\mu) + (-1 + B(\mu))x + C(\mu)x^2 + D(\mu)x^3$

$$A = f_\mu \mu + \frac{f_{\mu\mu}}{2} \mu^2 + O(\mu^3) \quad \text{(note } A(0) = 0)$$

$$B = f_{\mu x} \mu + O(\mu^2)$$

$$C(\mu) = \frac{1}{2} f_{xx} + O(\mu)$$

$$D(\mu) = \frac{1}{6} f_{xxx} + O(\mu)$$

so fixed pts solve  $x^* \approx A(\mu) - x^*$   
 $x^* \approx \frac{A(\mu)}{2} = \frac{1}{2} f_\mu \mu + O(\mu^2)$

Stability Depends on  $Df(x^*)$

$$\begin{aligned} Df &= (-1 + B(\mu)) + 2C(\mu)x + O(x^2) \\ &= -1 + f_{\mu x}(\mu) + 2 \cdot \frac{1}{2} f_{xx} \cdot \frac{1}{2} f_\mu \mu \\ &= -1 + (f_{\mu x} + \frac{1}{2} f_{xx} f_\mu) \mu \end{aligned}$$

stability depends on sign  $(f_{\mu x} + \frac{1}{2} f_{xx} f_\mu)$

we need this nonzero to get a non-degenerate bifurcation (recall our calculation for transcritical)



Ex 3.24-5

How would you show this

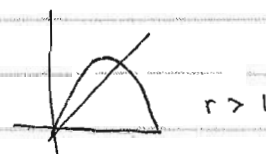
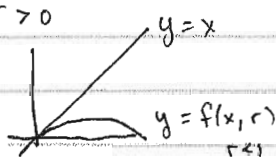
look at  $f^2(x) = f(f(x))$ , then if  $f'(x) = -1$

$$\frac{d}{dx} f(f(x)) \Big|_{x=x^*} = f'(f(x^*)) f'(x^*) = (f'(x^*))^2 = 1$$

in fact you carefully show that  $f^2(x)$  the 2nd iterate has a pitchfork bifurcation, can be supercritical or subcritical

Example: logistic map  $x > 0, r > 0$

$$x_{n+1} = r x_n (1 - x_n)$$



$f(x)$  fixed pts  $x = rx(1-x)$

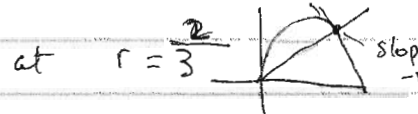
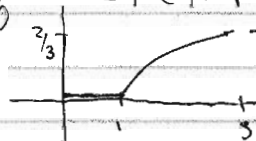
$$x=0 \text{ or } 1 = r(1-x) \Rightarrow x = 1 - \frac{1}{r}$$

$$f = r(x - x^2)$$

$$f' = r(1 - 2x), \quad f'(0) = r \quad \text{stable for } 0 < r < 1$$

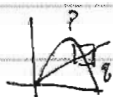
$$f'(1 - \frac{1}{r}) = r(1 - 2(1 - \frac{1}{r})) = r(-1 + \frac{2}{r}) = 1 - r$$

$$\text{stable if } -1 < 1 - r < 1 \Rightarrow 1 < r < 3$$



period 2

$r > 2$



$$\begin{aligned} f(p) &= q \\ f(q) &= p \\ f^2(q) &= q \end{aligned}$$

$$f(f(q)) = r^2 q(1-q)(1-rq+rq^2)$$

know  $q=0$ ,  $q=1-\frac{1}{r}$  are solutions, so we can factor out  $q(1-\frac{1}{r})$

after a bunch of algebra

$$\text{find } \left( q = \frac{r+1 + \sqrt{(r-3)(r+1)}}{2r} \right) \left( q = \frac{r+1 - \sqrt{(r-3)(r+1)}}{2r} \right) = 0$$

$$q = \frac{1}{2} + \frac{1 + \sqrt{(r-3)(r+1)}}{2r}, \quad p = \frac{1}{2} + \frac{1 - \sqrt{(r-3)(r+1)}}{2r}$$

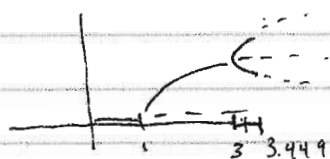
LEC 24/4/21-6

$$\begin{aligned}
 \frac{d}{dg} f(f(g)) &= f'(f(g)) f'(g) \\
 &= f'(p) f'(g) \\
 &= r(1-2p)(1-2g) \\
 &= r^2(1-2(p+g)+4pg) \\
 &= r^2(1 - \frac{2(r+1)}{4r} + \frac{4(r+1)}{r^2}) \\
 &= r^2 - 2(r+1)r + 4(r+1) \\
 &= r^2 - 2r^2 - 2r + 4r + 4 \\
 &= 4 + 2r - r^2
 \end{aligned}$$

$$\begin{aligned}
 pg &= \frac{1}{4r^2} ((r+1)^2 - (r+1)(r-3)) \\
 &= \frac{1}{4r^2} (r+1)(r+1-r+3) \\
 &= \frac{1}{4r^2} (r+1)(4) = \frac{r+1}{r^2}
 \end{aligned}$$

Stable if  $|4+2r-r^2| < 1$

stable for  $3 < r < 1+\sqrt{6} = 3.449...$



then  $\frac{d}{dg} f^2(g) = -1$  another flip

this is the start of a period-doubling cascade  
show logistic

~~pp~~ bifurcation

lorenz attractor

$$\dot{x} = 10(y-x)$$

$$\dot{y} = \mu x - y - xz$$

$$\dot{z} = xy - \frac{8}{3}z$$

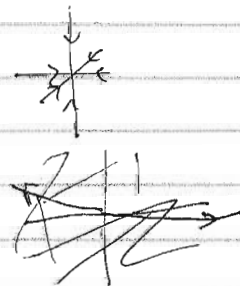
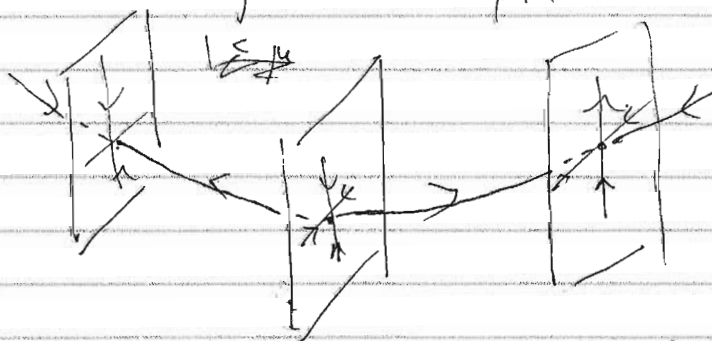
$$0 < \mu < 1$$

origin stable

$$\mu = 1$$

pitchfork

to



$$1.34 < \mu < 24.74$$

still stable but oscillatory

