

## LECTURE 15-1

Let's review where we were before vacation (2 weeks since last lecture!)  $w(\Gamma)$

We introduced the idea of an  $\omega$ -limit set for a trajectory  $\Gamma = \phi_t x_0$ , which is the set of points  $\tilde{x}$  s.t.  $\exists$  sequence  $t_n \rightarrow \infty$  s.t.  $\phi_{t_n} x_0 \rightarrow \tilde{x}$

We saw that this set is closed and that if  $\Gamma$  contained in a compact subset of  $\mathbb{R}^n$  then  $w(\Gamma)$  is non-empty, connected + compact

Then we defined attracting sets - a set  $S$  such that  $\exists$  open set of initial conditions  $X \supset S$  s.t.  $\phi_t x \rightarrow S$  as  $t \rightarrow \infty$   
 $\forall x \in X \quad d(\phi_t x, S) \rightarrow 0$

an attractor is an attracting set with a dense trajectory

Then we were considering fixed points that are attractors, used Lyapunov functions, LaSalle invariance principle, geometric reasoning to show regions of attraction to an asymptotically stable fixed pt.

There are (at least) 3:

There are several other types of sets that can be attractors

the SIMPLEST: Periodic orbits

We call now  $x(t)$  a periodic orbit if  $\exists T > 0$  s.t.  $x(t+T) = x(t)$  and  $T > 0$  is the minimum value s.t. this is true. (went to exclude fixed pts from definition)  
if an equation autonomous,  $\star$  uniqueness

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a periodic orbit  $\Gamma$  called stable if  $\forall \epsilon > 0$   
 $\exists$  a thd  $U \supset \Gamma$  st  
 $x \in U \Rightarrow |q_t x - \Gamma| < \epsilon \quad \forall t > 0$

unstable if not stable

asymptotically stable if stable and  $\exists$  a thd  $U \supset \Gamma$  st  
 $\forall x \in U \quad d(q_t x, \Gamma) \xrightarrow[t \rightarrow \infty]{} 0$

example (in polar coords)

$$\dot{r} = r(1-r^2)$$

$$\dot{\theta} = 1$$

then  $\begin{cases} x = \cos t \\ y = \sin t \end{cases}$

$$0 \leq t < 2\pi$$

is asymptotically stable and attracts all initial conditions except  $(0,0)$

Just like a fixed pt, a periodic orbit (in 3 or more dims) may have stable & unstable manifolds

example  $\dot{x} = -y + x(1-x^2-y^2)$

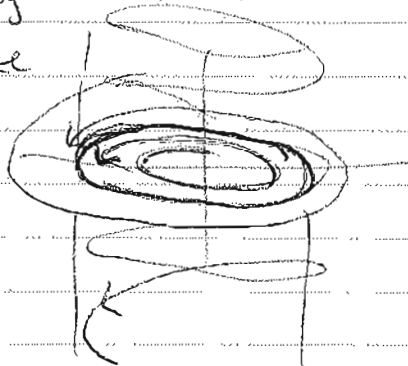
$$\dot{y} = x + y(1-x^2-y^2)$$

$$\dot{z} = z$$

fixed  $\Gamma$  given by

How should we

$$(\cos t, \sin t, 0)$$



if  $x = (x, y, 0) \neq (0,0,0)$ ,  $q_t x \rightarrow \Gamma$   
 $t \rightarrow \infty$

if  $x \neq x^2 + y^2 = 1$  but  $z \neq 0$

$$d(\Gamma, q_t x) \xrightarrow[t \rightarrow -\infty]{} 0$$

We define  $W_s(\Gamma) = \{x \mid d(q_t(x), \Gamma) \xrightarrow[t \rightarrow +\infty]{} 0\}$

$$W_u(\Gamma) = \{x \mid d(q_t(x), \Gamma) \xrightarrow[t \rightarrow -\infty]{} 0\}$$

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clearly for the SIMPCEST system

$$\dot{x} = -y$$

$$\dot{y} = x$$

periodic orbits

the ~~solution~~ are stable but not asymptotically stable

a limit cycle is a periodic orbit  $\Gamma$  which is the  $w$ -limit set (or  $\alpha$ -limit set) of some trajectory other than itself

~~a stable limit cycle~~

if a limit cycle is the  $w$ -limit set of every trajectory in a neighborhood  $U$  containing  $\Gamma$  it is called a stable limit cycle

if it is the  $\alpha$ -limit set of all trajectories in a nbhd of  $\Gamma$ , it is an unstable limit cycle

if it is the  $\alpha$ -limit set of one trajectory & the  $w$ -limit set of another,  $\Gamma$  is a semistable limit cycle

eg previous example or section 3.3 example 2

Poincaré loves to show crazy examples  $\dot{x} = f(x)$ ,  $f$  not analytic

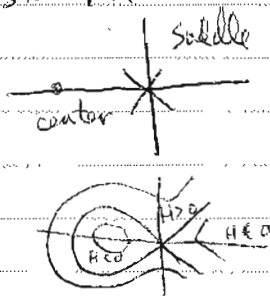
eg  $\begin{cases} \dot{r} = r^3 \sin \frac{1}{r} \\ \dot{\theta} = 1 \end{cases}$  clearly has countable sequence of limit cycles at  $r = \frac{1}{n\pi}$

Nice Theorem if  $f \in C^1(\mathbb{R}^2)$  and  $\Gamma$  is a periodic orbit and  $\exists$  trajectory exterior to  $\Gamma$  s.t.  $\Gamma$  is its  $w$ -limit set then  $\exists$  nbhd  $U$  exterior to  $\Gamma$  with  $w$ -limit set  $\Gamma$ , further every trajectory in  $U$  spirals around  $\Gamma$   $\infty$ -often

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## Heteroclinic orbits

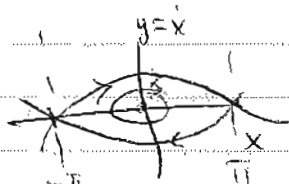
eg  $\ddot{x} = x + x^2$  fixed pts  $x=0, x=-1$   
 $\frac{1}{2}\dot{x}^2 - \frac{1}{2}x^2 + \frac{1}{3}x^3 = \phi H$



The curves  $H=0$  are called separatrices since they separate solutions of different asymptotic behavior

Note these are the stable & unstable manifolds of the origin

Heteroclinic orbits  
 pendulum  $\ddot{x} + \sin x = 0$



Note  $\frac{1}{2}\dot{x}^2 + 1 - \cos x = E$

also note sometimes it is easier to find heteroclinic orbits than general periodic orbits

~~at~~  $\cos \pi = -1$

so  $E=2$  on the heteroclinic orbits

$$\frac{1}{2}\dot{x}^2 + 1 - \cos x = 2$$

$$\frac{1}{2}\dot{x}^2 = 1 + \cos x$$

$$\frac{dx}{\sqrt{1 + \cos x}} = \pm \sqrt{2} dt$$

use  $\cos x = \frac{1}{2} + \frac{1}{2}\cos 2x$

$$\cos 2x =$$

$$\cos^2 x = \frac{1}{2} + \frac{1}{2}\cos 2x$$

$$\cos^2 \frac{x}{2} = \frac{1}{2} + \frac{1}{2}\cos x$$

$$2\cos^2 \frac{x}{2} = 1 + \cos x$$

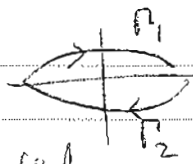
$$\frac{dx}{\sqrt{2\cos^2 \frac{x}{2}}} = \pm \sqrt{2} dt$$

$$\frac{dx}{2\cos \frac{x}{2}} = \pm \sqrt{2} dt$$

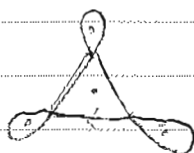
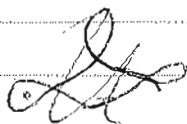
etc

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separatrix cycles -  
 & compound separatrix cycles



above  
 is a SIMPLE CLOSED  
 CURVE



union of a finite # of separatrices +  
 their limit points

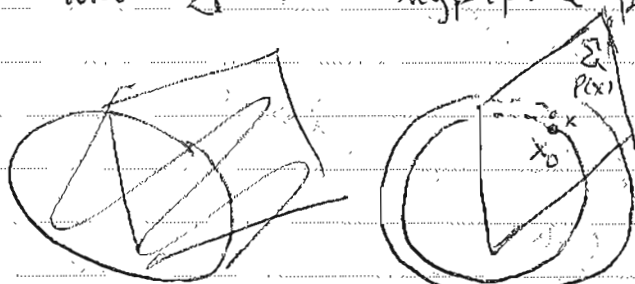
The Poincaré

The goal of the next few lectures: methods for determining  
 existence, non-existence, stability etc of periodic  
 orbits, homoclinic, heteroclinic orbits

Important tool: the Poincaré map,  $x \in \mathbb{R}^n$

$$x \mapsto f(x), \quad f \in C^1(E), \quad E \subset \mathbb{R}^n \text{ open}$$

Suppose  $\Gamma = \{x, f(x), \dots\}$  is a periodic orbit to  $x = f(x)$  through  $x_0$   
 and  $\Sigma$  is a hyperplane perpendicular to  $\Gamma$  at  $x_0$



if  $x$  is a point on  $\Sigma$  near  $x_0$ , then the trajectory will  
 cross the plane  $\Sigma$  in another point  $p(x)$

In fact, as long as  $\Sigma$  is any  $(n-1)$ -dimensional hypersurface  
 such that  $\Gamma$  is not tangent to  $\Sigma$  such that  $\Gamma$  is  
 not tangent to  $\Sigma$ , (the trajectory is said to  
 intersect  $\Sigma$  transversally) we can make this argument