# Homework 6

Son Nguyen

## Problem 3.5

(a)

$$\langle f|xf\rangle = \int_{-\infty}^{\infty} f^*xfdx$$

$$= \int_{-\infty}^{\infty} xf^*fdx \quad \text{(since } x \text{ is real: } x = x^*\text{)}$$

$$= \int_{-\infty}^{\infty} (xf)^*fdx$$

$$= \langle xf|f\rangle \Rightarrow x^{\dagger} = x$$

$$\langle f|if\rangle = \int_{-\infty}^{\infty} f^*ifdx$$

$$= \int_{-\infty}^{\infty} if^*fdx$$

$$= \int_{-\infty}^{\infty} (-if)^*fdx$$

$$= \langle -if|f\rangle \Rightarrow i^{\dagger} = -i$$

$$\langle f | \frac{d}{dx} f \rangle = \int_{-\infty}^{\infty} f^* \frac{df}{dx} dx$$

$$= f^* f \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{df^*}{dx} f dx$$

$$= -\int_{-\infty}^{\infty} \frac{df^*}{dx} f dx = \langle -\frac{d}{dx} f | f \rangle \Rightarrow -\frac{d}{dx} = \frac{d}{dx}^{\dagger}$$

(b)

$$\begin{split} \langle f | \hat{Q} \hat{R} f \rangle &= \int_{-\infty}^{\infty} f^* \hat{Q} \hat{R} f dx \\ &= \int_{-\infty}^{\infty} f^* \hat{Q} \underbrace{\left(\hat{R} f\right)}_{g} dx \\ &= \langle f | \hat{Q} g \rangle = \langle \hat{Q}^{\dagger} f | g \rangle \\ &= \int_{-\infty}^{\infty} \underbrace{\left(\hat{Q}^{\dagger} f\right)^*}_{h} (\hat{R} f) dx \\ &= \langle \hat{R}^{\dagger} h | f \rangle = \langle \hat{R}^{\dagger} \hat{Q}^{\dagger} f | f \rangle \Rightarrow (\hat{Q} \hat{R})^{\dagger} = \hat{R}^{\dagger} \hat{Q}^{\dagger} \end{split}$$

Next, we need to prove  $(\hat{Q} + \hat{R})^{\dagger} = \hat{Q}^{\dagger} + \hat{R}^{\dagger}$ 

$$\begin{split} \langle f(\hat{Q}+\hat{R})|f\rangle &= \int_{-\infty}^{\infty} f^*(\hat{Q}+\hat{R})fdx \\ &= \int_{-\infty}^{\infty} f^*\left(\hat{Q}f+\hat{R}f\right)dx \\ &= \langle f|\hat{Q}f\rangle + \langle f|\hat{R}f\rangle \\ &= \langle \hat{Q}^{\dagger}f|f\rangle + \langle \hat{R}^{\dagger}f|f\rangle \\ &= \int_{-\infty}^{\infty} (\hat{Q}^{\dagger}f)^*f + (\hat{R}^{\dagger}f)^*fdx \\ &= \int_{-\infty}^{\infty} \left(\hat{Q}^{\dagger}f + \hat{R}^{\dagger}f\right)^*fdx \\ &= \int_{-\infty}^{\infty} \left[\left(\hat{Q}^{\dagger}+\hat{R}^{\dagger}\right)f\right]^*fdx \\ &= \langle \left(\hat{Q}^{\dagger}+\hat{R}^{\dagger}\right)f|f\rangle \Rightarrow (\hat{Q}+\hat{R})^{\dagger} = \hat{Q}^{\dagger}+\hat{R}^{\dagger} \end{split}$$

$$\begin{split} \langle f|c\hat{Q}f\rangle &= \int_{-\infty}^{\infty} f^*c\hat{Q}fdx \\ &= c\langle f|\hat{Q}f\rangle = c\langle \hat{Q}^{\dagger}f|f\rangle \\ &= \int_{-\infty}^{\infty} c(\hat{Q}^{\dagger}f)^*fdx \\ &= \int_{-\infty}^{\infty} (c^*\hat{Q}^{\dagger}f)^*fdx \\ &= \langle c^*\hat{Q}^{\dagger}f|f\rangle \Rightarrow (c\hat{Q})^{\dagger} = c^*\hat{Q}^{\dagger} \end{split}$$

(c) Equation(2.48):

$$\hat{a}_{\pm} = \frac{1}{\sqrt{2\hbar m\omega}} (\mp i\hat{p} + m\omega x)$$

$$\hat{a}_{+}^{\dagger} = \left[ \frac{1}{\sqrt{2\hbar m\omega}} (-i\hat{p} + m\omega x) \right]^{\dagger}$$

$$= \left( \frac{1}{\sqrt{2\hbar m\omega}} \right)^{*} \left[ (-i\hat{p})^{\dagger} + (m\omega x)^{\dagger} \right]$$

$$= \frac{1}{\sqrt{2\hbar m\omega}} \left[ i\hat{p}^{\dagger} + m\omega x \right]$$

$$= \frac{1}{\sqrt{2\hbar m\omega}} \left[ i \left( -i\hbar \frac{d}{dx} \right)^{\dagger} + m\omega x \right]$$

$$= \frac{1}{\sqrt{2\hbar m\omega}} \left[ i \left( i\hbar \left( -\frac{d}{dx} \right) \right) + m\omega x \right] = \frac{1}{\sqrt{2\hbar m\omega}} (i\hat{p} + m\omega x) = \hat{a}_{-}$$

$$= \frac{1}{\sqrt{2\hbar m\omega}} \left[ \hbar \frac{d}{dx} + m\omega x \right]$$

#### Problem 3.14

(a) Prove equation (3.64):  $[\hat{A} + \hat{B}, \hat{C}] = [\hat{A}, \hat{C}] + [\hat{B}, \hat{C}]$ 

$$\begin{split} [\hat{A} + \hat{B}, \hat{C}] f(x) &= [(\hat{A} + \hat{B})\hat{C} - \hat{C}(\hat{A} + \hat{B})] f(x) \\ &= \hat{A}\hat{C}f(x) + \hat{B}\hat{C}f(x) - \hat{C}\hat{A}f(x) - \hat{C}\hat{B}f(x) \\ &= \hat{A}\hat{C}f(x) - \hat{C}\hat{A}f(x) + \hat{B}\hat{C}f(x) - \hat{C}\hat{B}f(x) \\ &= [\hat{A}, \hat{C}]f(x) + [\hat{B}, \hat{C}]f(x) \\ &\Rightarrow [\hat{A}, \hat{C}] + [\hat{B}, \hat{C}] \end{split}$$

Prove equation (3.65):  $[\hat{A}\hat{B}, \hat{C}] = \hat{A}[\hat{B}, \hat{C}] + [\hat{A}, \hat{C}]\hat{B}$ 

$$\begin{split} [\hat{A}\hat{B},\hat{C}]f(x) &= \hat{A}\hat{B}\hat{C}f(x) - \hat{C}\hat{A}\hat{B}f(x) \\ [\hat{A}[\hat{B},\hat{C}] + [\hat{A},\hat{C}]\hat{B}]f(x) &= \hat{A}\hat{B}\hat{C}f(x)\underbrace{-\hat{A}\hat{C}\hat{B}f(x) + \hat{A}\hat{C}\hat{B}f(x)}_{=0} - \hat{C}\hat{A}\hat{B}f(x) \\ &= [\hat{A}\hat{B},\hat{C}]f(x) \Rightarrow [\hat{A}\hat{B},\hat{C}] = \hat{A}[\hat{B},\hat{C}] + [\hat{A},\hat{C}]\hat{B} \end{split}$$

(b) Prove  $[x^n, \hat{p}] = i\hbar nx^{n-1}$ 

$$\begin{split} [x^n,\hat{p}]f(x) &= x^n\hat{p}f(x) - \hat{p}x^nf(x) \\ &= x^n \left( -i\hbar \frac{d}{dx} \right) f(x) - \left( -i\hbar \frac{d}{dx} \right) x^n f(x) \\ &= -x^n i\hbar \frac{d}{dx} f(x) + i\hbar \frac{d}{dx} (x^n f(x)) \\ &= -x^n i\hbar \frac{df}{dx} + i\hbar (nx^{n-1}f(x) + x^n \frac{df}{dx}) \\ &= -x^n i\hbar \frac{df}{dx} + i\hbar nx^{n-1}f(x) + i\hbar x^n \frac{df}{dx} \\ &= i\hbar nx^{n-1}f(x) \end{split}$$

(c)

$$\begin{split} [f(x),\hat{p}]F &= f(x)\hat{p}F - \hat{p}f(x)F \\ &= f(x)\left(-i\hbar\frac{d}{dx}\right)F - \left(-i\hbar\frac{d}{dx}\right)f(x)F \\ &= f(x)(-i\hbar\frac{dF}{dx}) + i\hbar\frac{df}{dx}F + i\hbar f(x)\frac{dF}{dx} \\ &= i\hbar\frac{df}{dx}F \Rightarrow [f(x),\hat{p}] = i\hbar\frac{df}{dx} \end{split}$$

(d) Equation (2.54):  $\hat{H} = \hbar\omega(\hat{a}_{-}\hat{a}_{+} - \frac{1}{2})$ 

$$\begin{split} [\hat{H}, \hat{a}_{\pm}] &= \hat{H} \hat{a}_{\pm} - \hat{a}_{\pm} \hat{H} \\ &= \hbar \omega (\hat{a}_{-} \hat{a}_{+} - \frac{1}{2}) \hat{a}_{\pm} - \hat{a}_{\pm} \hbar \omega (\hat{a}_{-} \hat{a}_{+} - \frac{1}{2}) \\ &= \hbar \omega (\hat{a}_{-} \hat{a}_{+} \hat{a}_{\pm} - \frac{1}{2} \hat{a}_{\pm}) - \hbar \omega (\hat{a}_{\pm} \hat{a}_{-} \hat{a}_{+} - \hat{a}_{\pm} \frac{1}{2}) \\ &= \hbar \omega (\hat{a}_{-} \hat{a}_{+} \hat{a}_{\pm} - \frac{1}{2} \hat{a}_{\pm} - \hat{a}_{\pm} \hat{a}_{-} \hat{a}_{+} + \hat{a}_{\pm} \frac{1}{2}) \\ &= \hbar \omega (\hat{a}_{-} \hat{a}_{+} \hat{a}_{\pm} - \hat{a}_{\pm} \hat{a}_{-} \hat{a}_{+}) \\ &= \hbar \omega (\hat{a}_{-} \hat{a}_{+} \hat{a}_{-} - \hat{a}_{-} \hat{a}_{-} \hat{a}_{+}) \quad \text{or} \quad \hbar \omega (\hat{a}_{-} \hat{a}_{+} \hat{a}_{+} - \hat{a}_{+} \hat{a}_{-} \hat{a}_{+}) \end{split}$$

We know that  $[\hat{a}_{-}, \hat{a}_{+}] = 1$  and  $[\hat{a}_{+}, \hat{a}_{-}] = -1$ 

$$\begin{split} [\hat{H}, \hat{a}_{-}] &= \hbar \omega (\hat{a}_{-} \hat{a}_{+} \hat{a}_{-} - \hat{a}_{-} \hat{a}_{-} \hat{a}_{+}) \\ &= \hbar \omega \hat{a}_{-} [\hat{a}_{+}, \hat{a}_{-}] = -\hbar \omega \hat{a}_{-} \\ [\hat{H}, \hat{a}_{+}] &= \hbar \omega (\hat{a}_{-} \hat{a}_{+} \hat{a}_{+} - \hat{a}_{+} \hat{a}_{-} \hat{a}_{+}) \\ &= \hbar \omega [\hat{a}_{-}, \hat{a}_{+}] \hat{a}_{+} = \hbar \omega \hat{a}_{+} \\ \\ &\Rightarrow [\hat{H}, \hat{a}_{\pm}] = \pm \hbar \omega \hat{a}_{\pm} \end{split}$$

#### Problem 3.20

The energy-time uncertainty principle is given by:

$$\Delta E \Delta t \ge \frac{\hbar}{2}$$

From equation (3.75) we have  $\Delta E = \sigma_H$  and  $\Delta_t = \frac{\sigma_x}{\left|\frac{d(x)}{dt}\right|}$ 

$$\Rightarrow \sigma_H \frac{\sigma_x}{\left|\frac{d\langle x \rangle}{dt}\right|} \ge \frac{\hbar}{2}$$

$$\Psi(x,0) = A[\psi_1(x) + \psi_2(x)]$$

For this equation, we have:

$$A = \frac{1}{\sqrt{2}}$$

$$\langle H \rangle = \frac{1}{2}E_1 + \frac{1}{2}E_2 = \frac{5\pi^2\hbar^2}{4ma^2}$$

$$\langle H^2 \rangle = \frac{1}{2} E_1^2 + \frac{1}{2} E_2^2 = \frac{17 \hbar^4 \pi^4}{8 m^2 a^4}$$

$$\langle x \rangle = \frac{a}{2} - \frac{16a}{9\pi^2} \cos\left(\frac{3\hbar\pi^2}{2ma^2}t\right)$$

$$\langle x^2 \rangle = \frac{a^2}{3} - \frac{5a^2}{16\pi^2} - \frac{16a^2}{9\pi^2} \cos\left(\frac{3\hbar\pi^2}{2ma^2}t\right)$$

$$\frac{d\langle x\rangle}{dt} = \frac{8\hbar}{3ma} \sin\left(\frac{3\hbar\pi^2}{2ma^2}t\right)$$

$$\Delta E \Delta t = \sqrt{\left(\frac{17\hbar^4\pi^4}{8m^2a^4}\right) - \left(\frac{5\pi^2\hbar^2}{4ma^2}\right)^2} \frac{\sqrt{\left(\frac{a^2}{3} - \frac{5a^2}{16\pi^2} - \frac{16a^2}{9\pi^2}\cos\left(\frac{3\hbar\pi^2}{2ma^2}t\right)\right) - \left(\frac{a}{2} - \frac{16a}{9\pi^2}\cos\left(\frac{3\hbar\pi^2}{2ma^2}t\right)\right)^2}}{\left|\frac{8\hbar}{3ma}\sin\left(\frac{3\hbar\pi^2}{2ma^2}t\right)\right|} \geq \frac{\hbar}{2}$$

### Problem 3.24

Let operator  $\hat{Q}$  has the othonormal basis  $|v_1\rangle, |v_2\rangle, \dots, |v_n\rangle$ , the matrix representing the operator can be written as:

$$\hat{Q}|v_1\rangle = \hat{Q}_{11}|v_1\rangle + \hat{Q}_{21}|v_2\rangle + \dots + \hat{Q}_{n1}|v_n\rangle$$

$$\hat{Q}|v_2\rangle = \hat{Q}_{12}|v_1\rangle + \hat{Q}_{22}|v_2\rangle + \dots + \hat{Q}_{n2}|v_n\rangle$$

$$\vdots$$

$$\hat{Q}|v_n\rangle = \hat{Q}_{1n}|v_1\rangle + \hat{Q}_{2n}|v_2\rangle + \dots + \hat{Q}_{nn}|v_n\rangle$$

or

$$\hat{Q}|v_i\rangle = \sum_{j=1}^n \hat{Q}_{ji}|v_j\rangle$$

The goal is to find the matrix element  $\hat{Q}_{ji}$ , and we can do that by taking the inner product of both sides:

$$\langle v_j | \hat{Q} | v_i \rangle = \hat{Q}_{ji}$$

Take the complex of both sides:

$$\langle v_j | \hat{Q} | v_i \rangle^* = \hat{Q}_{ji}^* = \langle v_i | \underbrace{\hat{Q}^{\dagger}}_{\hat{Q}} | v_j \rangle = \hat{Q}_{ij}$$

#### Problem 3.25

Since  $|1\rangle, |2\rangle$  are orthonormal basis, we have:

$$\begin{split} \hat{H}|1\rangle &= H_{11}|1\rangle + H_{21}|2\rangle \\ &= \epsilon(|1\rangle\langle 1| - |2\rangle\langle 2| + |1\rangle\langle 2| + |2\rangle\langle 1|)|1\rangle \\ &= \epsilon(|1\rangle + |2\rangle) = \underbrace{\epsilon}_{H_{11}}|1\rangle + \underbrace{\epsilon}_{H_{21}}|2\rangle \end{split}$$

$$\begin{split} \hat{H}|2\rangle &= H_{12}|1\rangle + H_{22}|2\rangle \\ &= \epsilon(|1\rangle\langle 1| - |2\rangle\langle 2| + |1\rangle\langle 2| + |2\rangle\langle 1|)|2\rangle \\ &= \epsilon(-|2\rangle + |1\rangle) = \underbrace{\epsilon}_{H_{12}}|1\rangle + \underbrace{-\epsilon}_{H_{22}}|2\rangle \end{split}$$

So the matrix representation of  $\hat{H}$  is:

$$\hat{H} = \begin{pmatrix} \epsilon & \epsilon \\ \epsilon & -\epsilon \end{pmatrix}$$

To find the eigenvalues and eigenvectors of  $\hat{H}$ , we need to solve the equation:

$$\begin{split} H|\psi\rangle &= E|\psi\rangle \\ H|\psi\rangle - EI|\psi\rangle &= 0 \\ (H-EI)|\psi\rangle &= 0 \\ \Rightarrow \det(H-EI) &= 0 \Rightarrow E = \pm \epsilon \sqrt{2} \end{split}$$

For 
$$E_1 = \epsilon \sqrt{2}$$
, its eigenvector is  $\begin{pmatrix} 1 \\ \sqrt{2} - 1 \end{pmatrix}$   
For  $E_2 = -\epsilon \sqrt{2}$ , its eigenvector is  $\begin{pmatrix} 1 \\ -\sqrt{2} - 1 \end{pmatrix}$