

# Ansatz

Son Nguyen

Reference from Scalable Quantum Simulation of Molecular Energies

## Contents

<b>1</b>	<b>Decomposing the UCCSD ansatz</b>	<b>2</b>
<b>2</b>	<b>Unitary Couple Cluster Single Double (UCCSD)</b>	<b>5</b>
2.1	Coupled Cluster Theory . . . . .	5
2.2	Unitary Coupled Cluster . . . . .	6
<b>3</b>	<b>Pauli Measurement</b>	<b>6</b>
3.1	Theoretical Analysis (XX) . . . . .	9
3.2	Experimental Analysis (XX) . . . . .	10
<b>4</b>	<b>Bell Measurement (Work In Progress)</b>	<b>11</b>
<b>5</b>	<b>Cost Function</b>	<b>14</b>
<b>6</b>	<b>Experiment</b>	<b>14</b>
6.1	Mapping . . . . .	14
<b>7</b>	<b>Optical Circuit</b>	<b>16</b>
7.1	Optical Components . . . . .	16
7.2	Realization of the Circuit . . . . .	16

We start with defining the Hamiltonian of the molecular Hydrogen.

$$H = g_0\mathbb{I} + g_1Z_0 + g_2Z_1 + g_3Z_0Z_1 + g_4Y_0Y_1 + g_5X_0X_1$$

Where:  $\{X_i, Z_i, Y_i\}$  denote the Pauli matrices acting on the  $i$ -th qubit and the real scalars  $\{g_\gamma\}$  are efficiently computable functions of the hydrogen-hydrogen bond length  $R$ .

$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad \mathbb{I} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad Y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \quad Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

$$g_0\mathbb{I} = \begin{bmatrix} g_0 & 0 & 0 & 0 \\ 0 & g_0 & 0 & 0 \\ 0 & 0 & g_0 & 0 \\ 0 & 0 & 0 & g_0 \end{bmatrix}, \quad g_1Z_0 = \begin{bmatrix} g_1 & 0 & 0 & 0 \\ 0 & g_1 & 0 & 0 \\ 0 & 0 & -g_1 & 0 \\ 0 & 0 & 0 & -g_1 \end{bmatrix}, \quad g_2Z_1 = \begin{bmatrix} g_2 & 0 & 0 & 0 \\ 0 & -g_2 & 0 & 0 \\ 0 & 0 & g_2 & 0 \\ 0 & 0 & 0 & -g_2 \end{bmatrix},$$

$$g_3Z_0Z_1 = \begin{bmatrix} g_3 & 0 & 0 & 0 \\ 0 & -g_3 & 0 & 0 \\ 0 & 0 & -g_3 & 0 \\ 0 & 0 & 0 & g_3 \end{bmatrix}, \quad g_4Y_0Y_1 = \begin{bmatrix} 0 & 0 & 0 & -g_4 \\ 0 & 0 & g_4 & 0 \\ 0 & g_4 & 0 & 0 \\ -g_4 & 0 & 0 & 0 \end{bmatrix}, \quad g_5X_0X_1 = \begin{bmatrix} 0 & 0 & 0 & g_5 \\ 0 & 0 & g_5 & 0 \\ 0 & g_5 & 0 & 0 \\ g_5 & 0 & 0 & 0 \end{bmatrix}$$

$$H = \begin{bmatrix} g_0 + g_1 + g_2 + g_3 & 0 & 0 & g_5 - g_4 \\ 0 & g_0 + g_1 - g_2 - g_3 & g_5 + g_4 & 0 \\ 0 & g_5 + g_4 & g_0 - g_1 + g_2 - g_3 & 0 \\ g_5 - g_4 & 0 & 0 & g_0 - g_1 - g_2 + g_3 \end{bmatrix}$$

## 1 Decomposing the UCCSD ansatz



Figure 1: The UCCSD ansatz for the Hydrogen molecule.

Reference state  $|10\rangle$



$$(X \otimes I) \cdot (|0\rangle \otimes |0\rangle) = |10\rangle$$

$$\left( \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \cdot \left( \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right) = \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix}$$

Apply parameterized ansatz



$$\left( R_x\left(\frac{-\pi}{2}\right) \otimes R_y\left(\frac{\pi}{2}\right) \right) \cdot |10\rangle$$

$$R_x\left(\frac{-\pi}{2}\right) = e^{-iX(\frac{-\pi}{4})} = \begin{bmatrix} \cos(\frac{-\pi}{4}) & -i\sin(\frac{-\pi}{4}) \\ -i\sin(\frac{-\pi}{4}) & \cos(-\frac{\pi}{4}) \end{bmatrix} = \begin{bmatrix} \frac{\sqrt{2}}{2} & \frac{i\sqrt{2}}{2} \\ \frac{i\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix}$$

$$R_y\left(\frac{\pi}{2}\right) = e^{-iY(\frac{\pi}{4})} = \begin{bmatrix} \cos(\frac{\pi}{4}) & -\sin(\frac{\pi}{4}) \\ \sin(\frac{\pi}{4}) & \cos(\frac{\pi}{4}) \end{bmatrix} = \begin{bmatrix} \frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix}$$

$$\left( R_x\left(\frac{-\pi}{2}\right) \otimes R_y\left(\frac{\pi}{2}\right) \right) = \begin{bmatrix} \frac{\sqrt{2}}{2} & \frac{i\sqrt{2}}{2} \\ \frac{i\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix} \otimes \begin{bmatrix} \frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix} = \begin{bmatrix} \frac{1}{2} & \frac{-1}{2} & \frac{i}{2} & \frac{-i}{2} \\ \frac{1}{2} & \frac{1}{2} & \frac{i}{2} & \frac{i}{2} \\ \frac{i}{2} & \frac{-i}{2} & \frac{1}{2} & \frac{-1}{2} \\ \frac{i}{2} & \frac{i}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

$$\left( R_x\left(\frac{-\pi}{2}\right) \otimes R_y\left(\frac{\pi}{2}\right) \right) \cdot |10\rangle = \begin{bmatrix} \frac{1}{2} & \frac{-1}{2} & \frac{i}{2} & \frac{-i}{2} \\ \frac{1}{2} & \frac{1}{2} & \frac{i}{2} & \frac{i}{2} \\ \frac{i}{2} & \frac{-i}{2} & \frac{1}{2} & \frac{-1}{2} \\ \frac{i}{2} & \frac{i}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix} \cdot \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = \begin{bmatrix} \frac{i}{2} \\ \frac{i}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}$$

The first CNOT (entanglement)



$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \cdot \begin{bmatrix} \frac{i}{2} \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} = \begin{bmatrix} \frac{i}{2} \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}$$

The  $Z_\theta$  rotation gate:

$$Z_\theta = e^{-iZ(\frac{\theta}{2})} = \begin{bmatrix} e^{-i\frac{\theta}{2}} & 0 \\ 0 & e^{i\frac{\theta}{2}} \end{bmatrix}$$

$$(Z_\theta \otimes I) \cdot \begin{bmatrix} \frac{i}{2} \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} = \begin{bmatrix} e^{-i\frac{\theta}{2}} & 0 & 0 & 0 \\ 0 & e^{-i\frac{\theta}{2}} & 0 & 0 \\ 0 & 0 & e^{i\frac{\theta}{2}} & 0 \\ 0 & 0 & 0 & e^{i\frac{\theta}{2}} \end{bmatrix} \cdot \begin{bmatrix} \frac{i}{2} \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} = \begin{bmatrix} \frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} - i\frac{\sin(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} + i\frac{\sin(\frac{\theta}{2})}{2} \\ -\frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \end{bmatrix}$$

The second CNOT (entanglement)

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \cdot \begin{bmatrix} \frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} - i\frac{\sin(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} + i\frac{\sin(\frac{\theta}{2})}{2} \\ -\frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \end{bmatrix} = \begin{bmatrix} \frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \\ -\frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} + i\frac{\sin(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} - i\frac{\sin(\frac{\theta}{2})}{2} \end{bmatrix}$$

The final rotation gates:

$$\begin{aligned} & \left( R_x\left(\frac{\pi}{2}\right) \otimes R_y\left(\frac{-\pi}{2}\right) \right) \cdot \begin{bmatrix} \frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \\ -\frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} + i\frac{\sin(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} - i\frac{\sin(\frac{\theta}{2})}{2} \end{bmatrix} \\ &= \left( \begin{bmatrix} \cos\left(\frac{\pi}{4}\right) & -i\sin\left(\frac{\pi}{4}\right) \\ -i\sin\left(\frac{\pi}{4}\right) & \cos\left(\frac{\pi}{4}\right) \end{bmatrix} \otimes \begin{bmatrix} \cos\left(\frac{-\pi}{4}\right) & -\sin\left(\frac{-\pi}{4}\right) \\ \sin\left(\frac{-\pi}{4}\right) & \cos\left(\frac{-\pi}{4}\right) \end{bmatrix} \right) \cdot \begin{bmatrix} \frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \\ -\frac{\sin(\frac{\theta}{2})}{2} + i\frac{\cos(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} + i\frac{\sin(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} - i\frac{\sin(\frac{\theta}{2})}{2} \end{bmatrix} \end{aligned}$$

$$\begin{aligned}
&= \left( \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{-i}{\sqrt{2}} \\ \frac{-i}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{bmatrix} \otimes \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{-1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{bmatrix} \right) \cdot \begin{bmatrix} \frac{\sin(\frac{\theta}{2})}{2} + i \frac{\cos(\frac{\theta}{2})}{2} \\ -\frac{\sin(\frac{\theta}{2})}{2} + i \frac{\cos(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} + i \frac{\sin(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} - i \frac{\sin(\frac{\theta}{2})}{2} \end{bmatrix} \\
&= \begin{bmatrix} \frac{1}{2} & \frac{1}{2} & \frac{-i}{2} & \frac{-i}{2} \\ \frac{-1}{2} & \frac{2}{2} & \frac{2}{2} & \frac{-i}{2} \\ \frac{2}{2} & \frac{-i}{2} & \frac{1}{2} & \frac{1}{2} \\ \frac{i}{2} & \frac{-i}{2} & \frac{-1}{2} & \frac{1}{2} \end{bmatrix} \cdot \begin{bmatrix} \frac{\sin(\frac{\theta}{2})}{2} + i \frac{\cos(\frac{\theta}{2})}{2} \\ -\frac{\sin(\frac{\theta}{2})}{2} + i \frac{\cos(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} + i \frac{\sin(\frac{\theta}{2})}{2} \\ \frac{\cos(\frac{\theta}{2})}{2} - i \frac{\sin(\frac{\theta}{2})}{2} \end{bmatrix} = \begin{bmatrix} 0 \\ -\sin(\frac{\theta}{2}) \\ \cos(\frac{\theta}{2}) \\ 0 \end{bmatrix} \quad (1)
\end{aligned}$$

Reverse to match the qiskit ordering:

$$\begin{bmatrix} 0 \\ \cos(\frac{\theta}{2}) \\ -\sin(\frac{\theta}{2}) \\ 0 \end{bmatrix} = |\phi(\vec{\theta})\rangle = \cos\left(\frac{\theta}{2}\right) |01\rangle - \sin\left(\frac{\theta}{2}\right) |10\rangle \quad (2)$$

Use  $\theta = -3.37$ :

$$\begin{bmatrix} 0 \\ \cos(\frac{-3.37}{2}) \\ -\sin(\frac{-3.37}{2}) \\ 0 \end{bmatrix} \approx \begin{bmatrix} 0 \\ -0.1139 \\ 0.9935 \\ 0 \end{bmatrix}$$

**\*Note:** This does not match the calculation, I had to switch place between the  $|01\rangle$  and  $|10\rangle$  to match the qiskit ordering

## 2 Unitary Couple Cluster Single Double (UCCSD)

### 2.1 Coupled Cluster Theory

Couple Cluster theory was introduced for calculation nuclear binding energies. It is the gold standard for the balance between accuracy and efficiency.

Key concepts:

- First quantization: individual particles are described by wavefunction  $\psi(x)$  that satisfies the Schrodinger equation.
- Second quantization: instead of describing each particle separately, we define creation and annihilation operators that act on quantum states of an entire system.

The fundamental operators:

- Creation operator:  $a_i^\dagger$  creates a particle in state  $i$ .
- Annihilation operator:  $a_i$  removes a particle from state  $i$ .

Fermionic Second Quantization: We described electron using second quantization.

$$|\Psi\rangle = a_1^\dagger a_3^\dagger |0\rangle$$

Which means we have occupied states 1 and 3 in the vacuum state  $|0\rangle$ .

## 2.2 Unitary Coupled Cluster

The UCC ansatz  $|\phi(\vec{\theta})\rangle$  is constructed from the reference state (Hartree-Fock state  $|\varphi\rangle$ )

$$|\phi(\vec{\theta})\rangle = e^{T(\vec{\theta}) - T(\vec{\theta})^\dagger} |\varphi\rangle \quad (3)$$

Where  $T(\vec{\theta})$  is the anti-Hermitian cluster operator.

## 3 Pauli Measurement

The expectation value (Pauli Measurement):

In this case, our state we want to reconstruct is  $|\phi(\vec{\theta})\rangle$ .

Starting with the density matrix:

$$\rho = |\phi(\vec{\theta})\rangle\langle\phi(\vec{\theta})|$$

The general two qubits wavefunction can be written as:

$$|\phi\rangle = a_{00}|00\rangle + a_{01}|01\rangle + a_{10}|10\rangle + a_{11}|11\rangle$$

Where  $a_{ij} \in \mathbb{C}$ , and  $\sum_{i,j} |a_{ij}|^2 = 1$ . For our case, we have:

$$|\phi(\vec{\theta})\rangle = \cos\left(\frac{\theta}{2}\right) |01\rangle - \sin\left(\frac{\theta}{2}\right) |10\rangle$$

Where:  $a_{00} = 0, a_{01} = \cos\left(\frac{\theta}{2}\right), a_{10} = -\sin\left(\frac{\theta}{2}\right), a_{11} = 0$ , the goal is to reconstruct  $a_{01}$  and  $a_{10}$ . To achieve this, we need to make measurement in different basis. (X, Y, Z, ...).

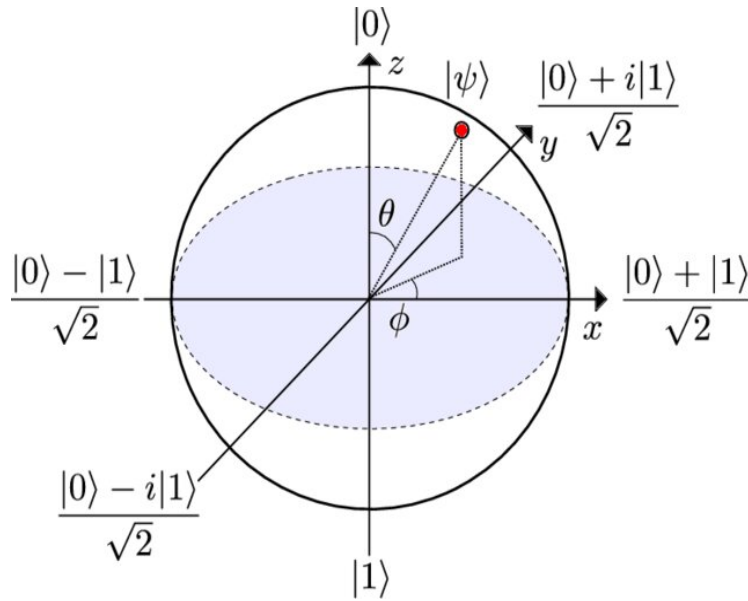


Figure 2: [Reference](#)

Let say we have a 1-qubit state  $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$ . This state is a superposition of  $|0\rangle$  and  $|1\rangle$ . This is also called the **Z basis (computational basis)**.

Measurement in the X basis - Diagonal basis/ Hadamard basis: superposition collapses the quantum state of the qubit  $|\psi\rangle$  to either  $|+\rangle$  or  $|-\rangle$ .

$$\begin{aligned}
 |+\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \\
 |-\rangle &= \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) \\
 |0\rangle &= \frac{1}{\sqrt{2}}(|+\rangle + |-\rangle) \\
 |1\rangle &= \frac{1}{\sqrt{2}}(|+\rangle - |-\rangle)
 \end{aligned}$$

$$H|\psi\rangle = H(\alpha|0\rangle + \beta|1\rangle) \quad (4)$$

$$= \alpha H|0\rangle + \beta H|1\rangle \quad (5)$$

$$= \alpha \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \beta \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \quad (6)$$

$$= \alpha \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{bmatrix} + \beta \begin{bmatrix} \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \end{bmatrix} \quad (7)$$

$$= \alpha|+\rangle + \beta|-\rangle \quad \text{X basis} \quad (8)$$

Measurement in the Y basis (Imaginary basis):

$$\begin{aligned} (S^\dagger \cdot H)|\psi\rangle &= (S^\dagger \cdot H)(\alpha|0\rangle + \beta|1\rangle) \\ &= \alpha(S^\dagger \cdot H)|0\rangle + \beta(S^\dagger \cdot H)|1\rangle \\ &= \alpha \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{-i}{\sqrt{2}} & \frac{i}{\sqrt{2}} \end{bmatrix} \cdot \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \beta \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{-i}{\sqrt{2}} & \frac{i}{\sqrt{2}} \end{bmatrix} \cdot \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\ &= \alpha \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{-i}{\sqrt{2}} \end{bmatrix} + \beta \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{i}{\sqrt{2}} \end{bmatrix} \\ &= \alpha \frac{1}{\sqrt{2}}(|0\rangle - i|1\rangle) + \beta \frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle) \quad \text{Y basis} \end{aligned}$$

Pauli Measurement	Unitary Transformation
$Z \otimes 1$	$1 \otimes 1$
$X \otimes 1$	$H \otimes 1$
$Y \otimes 1$	$HS^\dagger \otimes 1$
$1 \otimes Z$	SWAP
$1 \otimes X$	$(H \otimes 1)\text{SWAP}$
$1 \otimes Y$	$(HS^\dagger \otimes 1)\text{SWAP}$
$Z \otimes Z$	$\text{CNOT}_{10}$
$X \otimes Z$	$\text{CNOT}_{10}(H \otimes 1)$
$Y \otimes Z$	$\text{CNOT}_{10}(HS^\dagger \otimes 1)$
$Z \otimes X$	$\text{CNOT}_{10}(1 \otimes H)$
$X \otimes X$	$\text{CNOT}_{10}(H \otimes H)$
$Y \otimes X$	$\text{CNOT}_{10}(HS^\dagger \otimes H)$
$Z \otimes Y$	$\text{CNOT}_{10}(1 \otimes HS^\dagger)$
$X \otimes Y$	$\text{CNOT}_{10}(H \otimes HS^\dagger)$
$Y \otimes Y$	$\text{CNOT}_{10}(HS^\dagger \otimes HS^\dagger)$

$$\langle H \rangle = g_0 \mathbb{I} + g_1 \langle Z_0 \rangle + g_2 \langle Z_1 \rangle + g_3 \langle Z_0 Z_1 \rangle + g_4 \langle Y_0 Y_1 \rangle + g_5 \langle X_0 X_1 \rangle$$

$$\langle P \rangle = \sum_{\text{bitstring}} (-1)^{\text{parity}} \left( \frac{\text{count}}{\text{shots}} \right)$$



Where:

- parity = `bitstring.count('1')` mod 2
- count = number of times the bitstring was measured
- shots = total number of measurements

### 3.1 Theoretical Analysis (XX)

- For the  $XX$  operator

$$XX = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \otimes \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix}$$

The expectation value of an observable  $O$  is given by:

$$\langle O \rangle = \sum_i \lambda_i p_i$$

Where  $p_i$  is the probability of measuring the state in the  $i^{th}$  eigenstate.  $\lambda_i$  is the expectation value corresponding to eigenstate.

From equation (4) we can see that  $H|0\rangle = |+\rangle, H|1\rangle = |-\rangle$ . The corresponding eigenvalues for the operator  $X$  are

$$\langle +|X|+ \rangle = 1 \text{ and } \langle -|X|-\rangle = -1$$

For the 2 qubit states, we have:

$$\begin{aligned} (H \otimes H)|00\rangle &= \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} = |++\rangle = \frac{1}{2}(|00\rangle + |01\rangle + |10\rangle + |11\rangle) \\ (H \otimes H)|01\rangle &= |+-\rangle = \frac{1}{2}(|00\rangle - |01\rangle + |10\rangle - |11\rangle) \\ (H \otimes H)|10\rangle &= |-+\rangle = \frac{1}{2}(|00\rangle + |01\rangle - |10\rangle - |11\rangle) \\ (H \otimes H)|11\rangle &= |--\rangle \end{aligned}$$

Apply the operator  $XX$  on each state:

$$\begin{aligned}
XX|++\rangle &= \frac{1}{2}(XX|00\rangle + XX|01\rangle + XX|10\rangle + XX|11\rangle) \\
&= \frac{1}{2}(|11\rangle + |10\rangle + |01\rangle + |00\rangle) = |++\rangle \\
&\Rightarrow \langle ++ | XX | ++ \rangle = \langle ++ | ++ \rangle = 1 \\
XX|+-\rangle &= \frac{1}{2}(XX|00\rangle - XX|01\rangle + XX|10\rangle - XX|11\rangle) \\
&= \frac{1}{2}(|11\rangle - |10\rangle + |01\rangle - |00\rangle) = \frac{1}{2} \begin{bmatrix} -1 \\ 1 \\ -1 \\ 1 \end{bmatrix} \\
&\Rightarrow \langle +- | XX | +- \rangle = \begin{bmatrix} \frac{1}{2} & -\frac{1}{2} & \frac{1}{2} & -\frac{1}{2} \end{bmatrix} \cdot \begin{bmatrix} -\frac{1}{2} \\ \frac{1}{2} \\ -\frac{1}{2} \\ \frac{1}{2} \end{bmatrix} = -1
\end{aligned}$$

and so we can get the remaining expectation values  $\langle -+ | XX | -+ \rangle = -1, \langle -- | XX | -- \rangle = 1$

Measurement	XX basis equivalent	XX Expectation Value
$ 00\rangle$	$ ++\rangle$	1
$ 01\rangle$	$ +-\rangle$	-1
$ 10\rangle$	$  - + \rangle$	-1
$ 11\rangle$	$ --\rangle$	1

### 3.2 Experimental Analysis (XX)

From our circuit that produce the wavefunction  $|\phi(\vec{\theta})\rangle$  (see Figure 1). We just need to apply two Hadamard gates to the qubits and measure the expectation value of the  $XX$  operator.

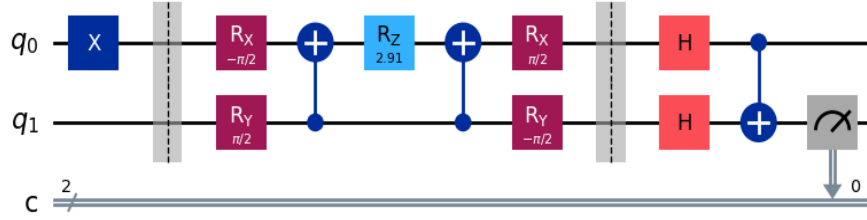


Figure 3: Circuit for the  $XX$  operator

Give the trial wavefunction  $|\phi(\vec{\theta})\rangle = \cos\left(\frac{\theta}{2}\right)|01\rangle - \sin\left(\frac{\theta}{2}\right)|10\rangle$ . Applying the Hadamard gates given:

$$\begin{aligned}
& (H \otimes H)|\phi(\vec{\theta})\rangle \\
&= (H \otimes H)\left(\cos\left(\frac{\theta}{2}\right)|01\rangle - \sin\left(\frac{\theta}{2}\right)|10\rangle\right) \\
&= \cos\left(\frac{\theta}{2}\right)(H|0\rangle \otimes H|1\rangle) - \sin\left(\frac{\theta}{2}\right)(H|1\rangle \otimes H|0\rangle) \\
&= \cos\left(\frac{\theta}{2}\right)\left[\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)\right] - \sin\left(\frac{\theta}{2}\right)\left[\frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)\right] \\
&= \cos\left(\frac{\theta}{2}\right)\left[\frac{1}{2}(|00\rangle - |01\rangle + |10\rangle - |11\rangle)\right] - \sin\left(\frac{\theta}{2}\right)\left[\frac{1}{2}(|00\rangle + |01\rangle - |10\rangle - |11\rangle)\right] \\
&= \frac{1}{2}\left\{\left[\cos\left(\frac{\theta}{2}\right) - \sin\left(\frac{\theta}{2}\right)\right]|00\rangle - \left[\cos\left(\frac{\theta}{2}\right) + \sin\left(\frac{\theta}{2}\right)\right]|01\rangle + \left[\cos\left(\frac{\theta}{2}\right) + \sin\left(\frac{\theta}{2}\right)\right]|10\rangle + \left[\sin\left(\frac{\theta}{2}\right) - \cos\left(\frac{\theta}{2}\right)\right]|11\rangle\right\}
\end{aligned}$$

Then we apply the  $\text{CNOT}_{01}$  gate to the state:

$$\begin{aligned}
& \text{CNOT}_{01}\left\{\cos\left(\frac{\theta}{2}\right)\left[\frac{1}{2}(|00\rangle - |01\rangle + |10\rangle - |11\rangle)\right] - \sin\left(\frac{\theta}{2}\right)\left[\frac{1}{2}(|00\rangle + |01\rangle - |10\rangle - |11\rangle)\right]\right\} \\
&= \cos\left(\frac{\theta}{2}\right)\left[\frac{1}{2}\text{CNOT}_{01}(|00\rangle - |01\rangle + |10\rangle - |11\rangle)\right] - \sin\left(\frac{\theta}{2}\right)\left[\frac{1}{2}\text{CNOT}_{01}(|00\rangle + |01\rangle - |10\rangle - |11\rangle)\right] \\
&= \cos\left(\frac{\theta}{2}\right)\left[\frac{1}{2}(|00\rangle - |01\rangle + |11\rangle - |10\rangle)\right] - \sin\left(\frac{\theta}{2}\right)\left[\frac{1}{2}(|00\rangle + |01\rangle - |11\rangle - |10\rangle)\right] \\
&= \frac{1}{2}\left\{\left[\cos\left(\frac{\theta}{2}\right) - \sin\left(\frac{\theta}{2}\right)\right]|00\rangle - \left[\cos\left(\frac{\theta}{2}\right) + \sin\left(\frac{\theta}{2}\right)\right]|01\rangle + \left[\sin\left(\frac{\theta}{2}\right) - \cos\left(\frac{\theta}{2}\right)\right]|10\rangle + \left[\sin\left(\frac{\theta}{2}\right) + \cos\left(\frac{\theta}{2}\right)\right]|11\rangle\right\}
\end{aligned}$$

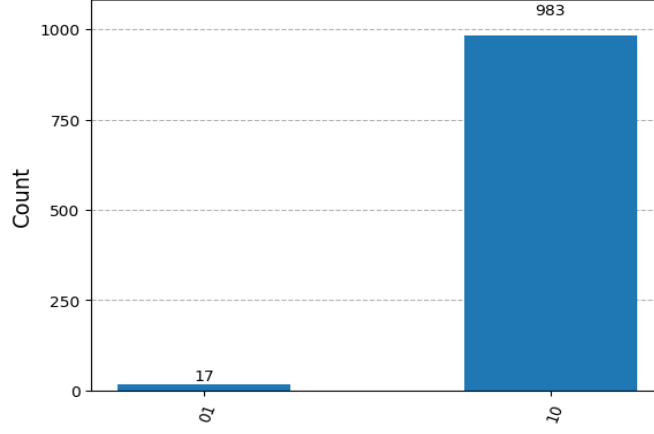
After applying  $\text{CNOT}_{01}$  gate.

$$\underbrace{|00\rangle}_{\text{even}} \rightarrow |0\mathbf{0}\rangle; \underbrace{|01\rangle}_{\text{odd}} \rightarrow |0\mathbf{1}\rangle; \underbrace{|10\rangle}_{\text{odd}} \rightarrow |\mathbf{1}1\rangle; \underbrace{|11\rangle}_{\text{even}} \rightarrow |\mathbf{1}\mathbf{0}\rangle$$

We can see that after applying the  $\text{CNOT}_{01}$  gate, the second qubit is flipped (the colored). Now measuring just the second qubit give us the parity of the state. If the parity is even we will see 0 and if the parity is odd we will see 1. A CNOT gate is used to compute the parity of two qubits and store it in one qubit without fully collapsing the state.

## 4 Bell Measurement (Work In Progress)

Alternatively, using Bell Measurement to reconstruct the trial wavefunction with the parameter  $\theta \approx -3.37$ , getting the expectation after 1000 measurements:



From the figure, we have see there is 1.7% of  $|01\rangle$  and 98.3% of  $|10\rangle$ .

$$\begin{aligned}\sqrt{1.7\%}|01\rangle + \sqrt{98.3\%}|10\rangle &= |\phi(\vec{\theta})\rangle \\ \pm 0.13|01\rangle \pm 0.99|10\rangle &= |\phi(\vec{\theta})\rangle\end{aligned}$$

To determine the sign of our trial wavefunction, we can use Bell measurements. We can measures any state which is an superposition of  $|00\rangle, |01\rangle, |10\rangle, |11\rangle$  in the Bell basis.

$$\begin{aligned}|\Phi^+\rangle &= \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \\ |\Phi^-\rangle &= \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle) \\ |\Psi^+\rangle &= \frac{1}{\sqrt{2}}(|01\rangle + |10\rangle) \\ |\Psi^-\rangle &= \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)\end{aligned}$$

By combining a CNOT gate followed by a Hadamard gate, we can measure the state in the Bell basis.

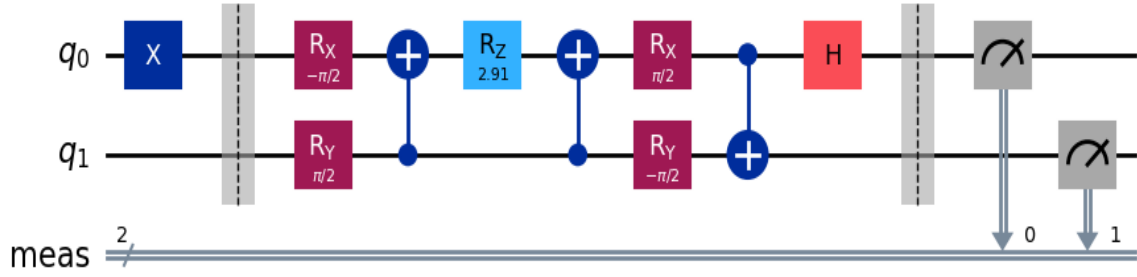
$$\begin{aligned}U|\Phi^+\rangle &= |00\rangle \\ U|\Phi^-\rangle &= |01\rangle \\ U|\Psi^+\rangle &= |10\rangle \\ U|\Psi^-\rangle &= |11\rangle\end{aligned}$$

Where  $U_{Bell} = (H \otimes I) \cdot \text{CNOT}(0,1)$

$$U_{Bell} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 0 & 1 \\ 1 & 0 & 0 & -1 \\ 0 & 1 & 1 & 0 \\ 0 & -1 & 1 & 0 \end{bmatrix}$$

Applying the  $U_{Bell}$  on  $\begin{bmatrix} A \\ B \\ C \\ D \end{bmatrix}$

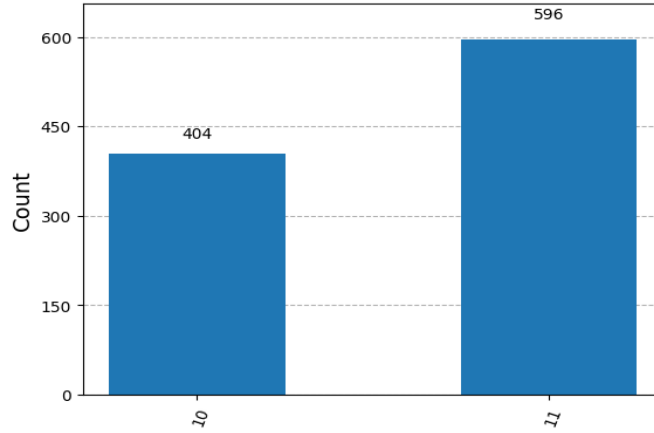
$$U_{Bell} \cdot \begin{bmatrix} A \\ B \\ C \\ D \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} A + D \\ A - D \\ B + C \\ B - C \end{bmatrix}$$



The trial wavefunction after applying the  $U_{Bell}$  unitary gate:

$$U_{Bell} \cdot |\phi(\vec{\theta})\rangle = U_{Bell} \cdot \begin{bmatrix} 0 \\ \cos\left(\frac{\theta}{2}\right) \\ -\sin\left(\frac{\theta}{2}\right) \\ 0 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ 0 \\ \cos\left(\frac{\theta}{2}\right) - \sin\left(\frac{\theta}{2}\right) \\ \cos\left(\frac{\theta}{2}\right) + \sin\left(\frac{\theta}{2}\right) \end{bmatrix} \begin{matrix} |00\rangle \\ |01\rangle \\ |10\rangle \\ |11\rangle \end{matrix} \begin{matrix} \\ \\ \approx 0.39\% \\ \approx 0.61\% \end{matrix}$$

Using  $\theta \approx -3.37$  we have:



We can see the counts of  $|11\rangle$  is dominant, which means the state is  $|\Psi^-\rangle$ . Therefore, the sign between  $|01\rangle$  and  $|10\rangle$  is negative.

$$0.13|01\rangle - 0.99|10\rangle = |\phi(\vec{\theta})\rangle \quad (9)$$

#### Reference.

Now we plug in the  $\theta$  to equation (1) to compare with equation (2), we have:

$$\begin{aligned} -\sin\left(\frac{-3.37}{2}\right)|01\rangle + \cos\left(\frac{-3.37}{2}\right)|10\rangle &= |\phi(\vec{\theta})\rangle \\ 0.993|01\rangle - 0.11|10\rangle &= |\phi(\vec{\theta})\rangle \end{aligned}$$

**There is a mistake for my bell measurement, I will correct it later.**

## 5 Cost Function

Mathematically we can use the Hamiltonian and the trial wavefunction, we can get our cost function (energy) as:

$$E = \langle \phi(\vec{\theta}) | H | \phi(\vec{\theta}) \rangle$$

$$\begin{bmatrix} 0 & \cos(\frac{\theta}{2}) & -\sin(\frac{\theta}{2}) & 0 \end{bmatrix} \cdot \begin{bmatrix} g_0 + g_1 + g_2 + g_3 & 0 & 0 & g_5 - g_4 \\ 0 & g_0 + g_1 - g_2 - g_3 & g_5 + g_4 & 0 \\ 0 & g_5 + g_4 & g_0 - g_1 + g_2 - g_3 & 0 \\ g_5 - g_4 & 0 & 0 & g_0 - g_1 - g_2 + g_3 \end{bmatrix} \cdot \begin{bmatrix} 0 \\ \cos(\frac{\theta}{2}) \\ -\sin(\frac{\theta}{2}) \\ 0 \end{bmatrix}$$

Plug in  $g_0 = -0.4804, g_1 = 0.3435, g_2 = -0.4347, g_3 = 0.5716, g_4 = 0.091, g_5 = 0.091$  we have:

$$\begin{bmatrix} 0 & \cos(\frac{\theta}{2}) & -\sin(\frac{\theta}{2}) & 0 \end{bmatrix} \cdot \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & -0.2738 & 0.182 & 0 \\ 0 & 0.182 & -1.8302 & 0 \\ 0 & 0 & 0 & 0.1824 \end{bmatrix} \cdot \begin{bmatrix} 0 \\ \cos(\frac{\theta}{2}) \\ -\sin(\frac{\theta}{2}) \\ 0 \end{bmatrix} \approx -1.851 \quad (\text{with } \theta = -3.37)$$

The minimum energy can be found using classical optimization techniques.

$$E_{min} = \langle \phi_{min}(\vec{\theta}) | H | \phi_{min}(\vec{\theta}) \rangle$$

## 6 Experiment

### 6.1 Mapping

From equation (2), we have a trial wavefunction of:

$$|\phi(\vec{\theta})\rangle = 0|00\rangle + \cos\left(\frac{\theta}{2}\right)|01\rangle - \sin\left(\frac{\theta}{2}\right)|10\rangle + 0|11\rangle \quad (10)$$

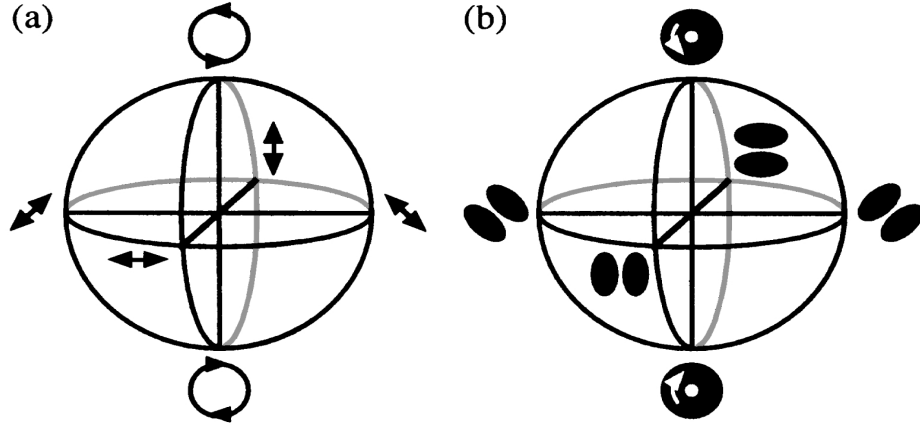


Figure 4: Poincare

For this experiment, we will use "01"/"10" Spatial Mode as our first qubit and Polarization as our second qubit.

$$|0\rangle = \text{Spatial Mode 10 and Horizontal Polarization or } |0\rangle_s, |0\rangle_p$$

$$|1\rangle = \text{Spatial Mode 01 and Vertical Polarization or } |1\rangle_s, |1\rangle_p$$

$$\frac{|0\rangle - i|1\rangle}{\sqrt{2}} = \text{South Pole of the Poincare Sphere or Left-circular Polarization } \frac{|L\rangle_s}{\sqrt{2}}, \frac{|L\rangle_p}{\sqrt{2}}$$

$$\frac{|0\rangle + i|1\rangle}{\sqrt{2}} = \text{North Pole of the Poincare Sphere or Right-circular Polarization } \frac{|R\rangle_s}{\sqrt{2}}, \frac{|R\rangle_p}{\sqrt{2}}$$

$$|+\rangle = \text{Diagonal Polarization } \frac{|D\rangle_s}{\sqrt{2}}, \frac{|D\rangle_p}{\sqrt{2}}$$

$$|-\rangle = \text{Anti-diagonal Polarization } \frac{|A\rangle_s}{\sqrt{2}}, \frac{|A\rangle_p}{\sqrt{2}}$$

Therefore our initial state is  $|0\rangle_p \otimes |0\rangle_s \equiv |0_p 0_s\rangle$

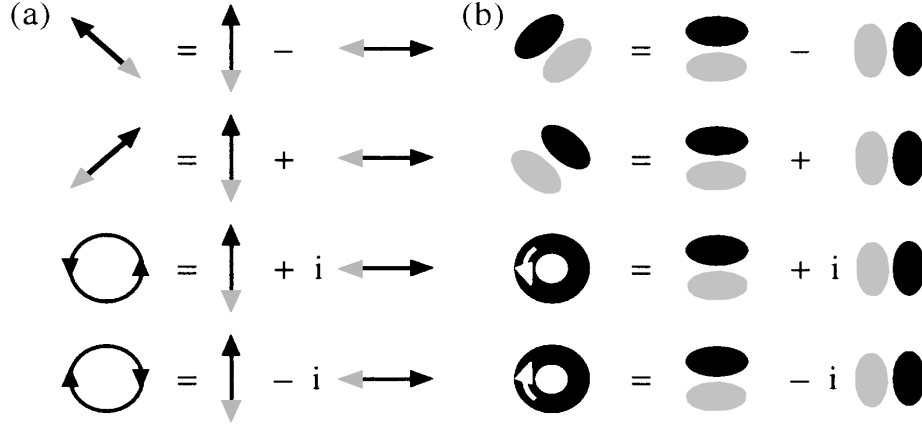


Figure 5: **Geometric Phase Shift**

## 7 Optical Circuit

### 7.1 Optical Components

#### Reference

Quarter-wave plate (QWP):

$$e^{-\frac{i\pi}{4}} \begin{bmatrix} \cos^2(\theta) + i \sin^2(\theta) & (1-i) \sin(\theta) \cos(\theta) \\ (1-i) \sin(\theta) \cos(\theta) & \sin^2(\theta) + i \cos^2(\theta) \end{bmatrix}$$

Quarter-wave plate with fast axis vertical (QWPv):

$$e^{\frac{i\pi}{4}} \begin{bmatrix} 1 & 0 \\ 0 & -i \end{bmatrix}$$

Quarter-wave plate with fast axis horizontal (QWPh):

$$e^{-\frac{i\pi}{4}} \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$$

Half-wave plate (HWP):

$$\begin{bmatrix} \cos(2\phi) & \sin(2\phi) \\ \sin(2\phi) & -\cos(2\phi) \end{bmatrix}$$

**Dove Prism** (DP) using Rotation matrix:

$$\begin{bmatrix} \cos(\omega) & -\sin(\omega) \\ \sin(\omega) & \cos(\omega) \end{bmatrix}$$

### 7.2 Realization of the Circuit

$$R_x\left(\frac{\pi}{2}\right) = \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{-i}{\sqrt{2}} \\ \frac{-i}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & -i \\ -i & 1 \end{bmatrix}$$



$$\begin{aligned}
R_y\left(\frac{\pi}{2}\right) &= \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{-1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & -1 \\ 1 & 1 \end{bmatrix} \\
R_z(\theta) = e^{-iZ(\frac{\theta}{2})} &= \begin{bmatrix} e^{-i\frac{\theta}{2}} & 0 \\ 0 & e^{i\frac{\theta}{2}} \end{bmatrix} = \begin{bmatrix} \cos(\frac{\theta}{2}) - i\sin(\frac{\theta}{2}) & 0 \\ 0 & \cos(\frac{\theta}{2}) + i\sin(\frac{\theta}{2}) \end{bmatrix} \\
H &= \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \\
S^\dagger &= \begin{bmatrix} 1 & 0 \\ 0 & -i \end{bmatrix}
\end{aligned}$$

$R_x(\frac{\pi}{2})$  realization in the optical circuit for Polarization:

$$\begin{aligned}
\text{QWP}\left(\frac{\pi}{4}\right) &= \begin{bmatrix} \cos^2(\frac{\pi}{4}) + i\sin^2(\frac{\pi}{4}) & (1-i)\sin(\frac{\pi}{4})\cos(\frac{\pi}{4}) \\ (1-i)\sin(\frac{\pi}{4})\cos(\frac{\pi}{4}) & \sin^2(\frac{\pi}{4}) + i\cos^2(\frac{\pi}{4}) \end{bmatrix} \\
&= \begin{bmatrix} \frac{1}{2} + i\frac{1}{2} & (1-i)\frac{1}{2} \\ (1-i)\frac{1}{2} & \frac{1}{2} + i\frac{1}{2} \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1+i & (1-i) \\ (1-i) & 1+i \end{bmatrix} = \frac{1+i}{2} \begin{bmatrix} 1 & \frac{1-i}{1+i} \\ \frac{1-i}{1+i} & 1 \end{bmatrix} \\
&= \frac{1+i}{2} \begin{bmatrix} 1 & -i \\ -i & 1 \end{bmatrix}
\end{aligned}$$

$R_y(\frac{\pi}{2})$  realization in the optical circuit for Spatial Mode:

$$\begin{aligned}
\text{DP}\left(\frac{\pi}{4}\right) &= \begin{bmatrix} \cos(\frac{\pi}{4}) & -\sin(\frac{\pi}{4}) \\ \sin(\frac{\pi}{4}) & \cos(\frac{\pi}{4}) \end{bmatrix} \\
&= \begin{bmatrix} \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{bmatrix}
\end{aligned}$$

$R_z(\theta)$  realization in the optical circuit for Polarization:

$$\text{QWPh} \cdot \text{QWP}\left(\frac{-\pi}{4}\right) \cdot \text{HWP}\left(\frac{\pi}{2} - \frac{\phi}{4}\right) \cdot \text{QWP}\left(\frac{-\pi}{4}\right) \cdot \text{QWPh} = \begin{bmatrix} -ie^{-\frac{i\phi}{2}} & 0 \\ 0 & -ie^{\frac{i\phi}{2}} \end{bmatrix}$$

Hadamard gate ( $H$ ) realization in the optical circuit for Polarization:

$$\begin{aligned}
\text{HWP}\left(\frac{\pi}{8}\right) &= \begin{bmatrix} \cos(\frac{\pi}{8}) & \sin(\frac{\pi}{8}) \\ \sin(\frac{\pi}{8}) & -\cos(\frac{\pi}{8}) \end{bmatrix} \\
&= \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix}
\end{aligned}$$

Mach-Zehnder Interferometer (MZI) realization in the optical circuit for Spatial Mode: