

# A study of the stochastic Saltzman-Maasch model for climate dynamics in Pleistocene

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Deterministic model

The Stochastic Model

Improved Model

## Deterministic model

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# The Saltzman-Maasch model

The model explains central features of the glacial cycles observed in the climate record of the Pleistocene Epoch.

$$\begin{cases} \dot{X} = -X - Y - vZ \\ \dot{Y} = -pZ + rY - sY^2 - Y^3 \\ \dot{Z} = -q(X + Z) \end{cases}$$

The state variables  $X, Y, Z$  represent the anomalies (deviations from long-term averages) of

- $X$ : the total continental ice mass
- $Y$ : the atmospheric  $\text{CO}_2$  concentration
- $Z$ : the mean temperature of the North Atlantic Deep Water

## Deterministic Equilibria

The model possesses the equilibrium  $E_0 = (0, 0, 0)$ , for all the parameter values. Other equilibria points  $E = (X, Y, Z)$  can be found imposing:

$$\begin{cases} -X - Y - vZ = 0 \\ -pZ + rY - sY^2 - Y^3 = 0 \\ -q(X + Z) = 0 \implies X = -Z \end{cases}$$

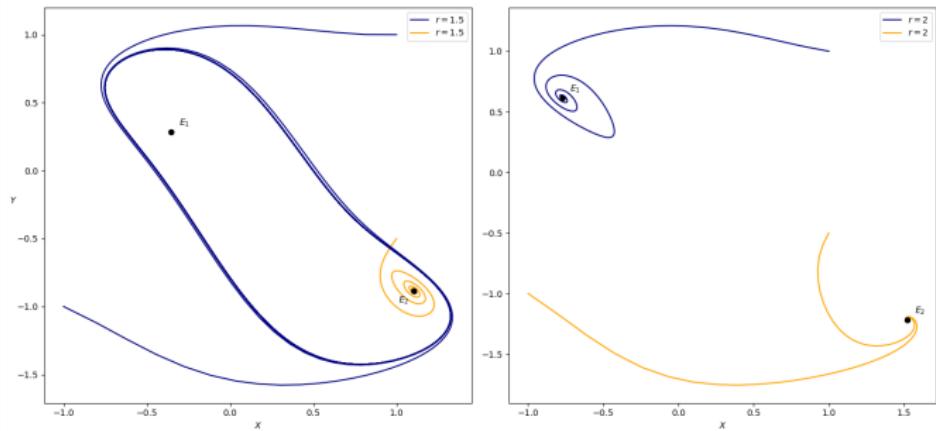
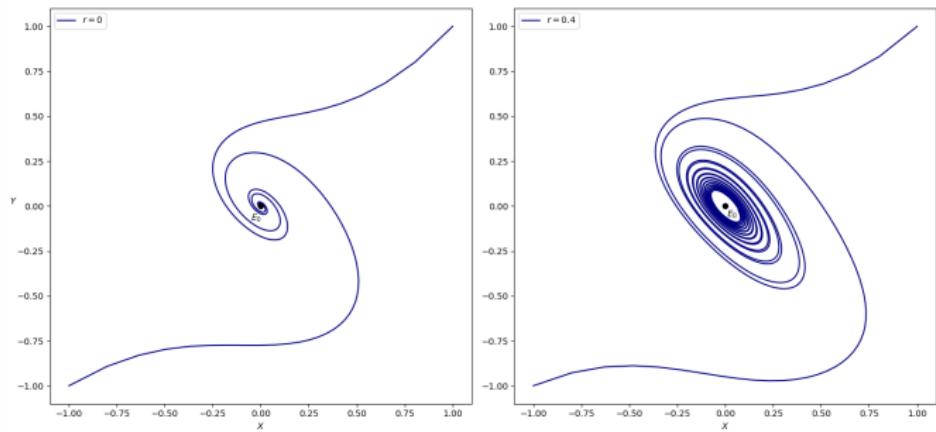
which leads to the non-trivial equilibrium points expression:

$$E_{1,2} = \left( -\frac{Y_{1,2}}{1-v}, \ Y_{1,2}, \ \frac{Y_{1,2}}{1-v} \right), \quad Y_{1,2} = -\frac{s}{2} \pm \sqrt{r + \frac{s^2}{4} - \frac{p}{1-v}}$$

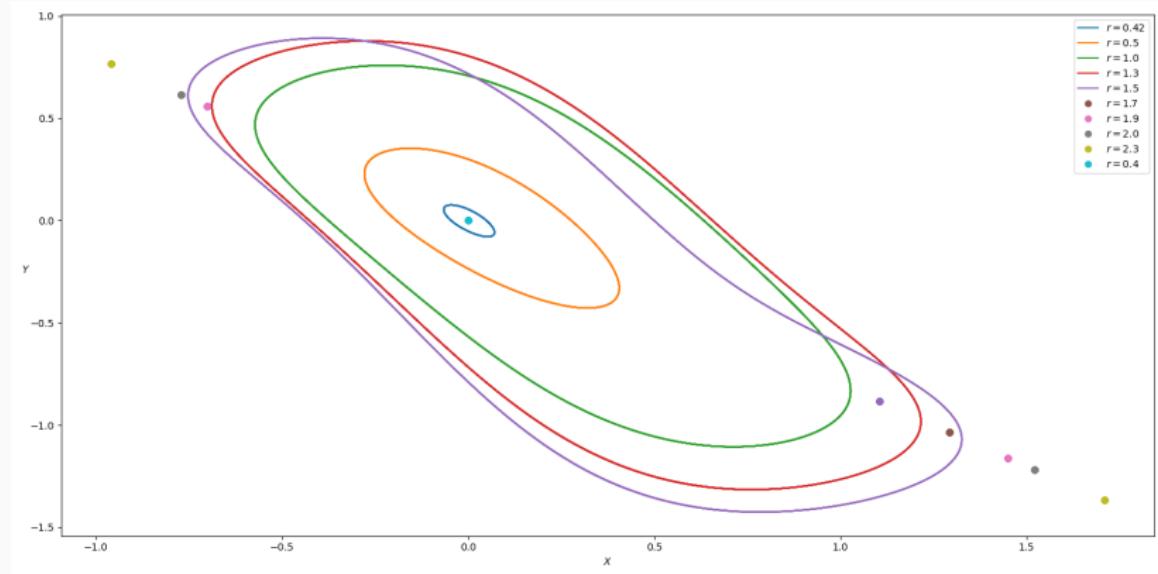
The parameter  $r$  is used as control and bifurcation parameter, fixing the others:  $p = 1$ ,  $q = 2.5$ ,  $v = 0.2$ ,  $s = 0.6$ . Within this framework:

$$E_{1,2} = (-1.25Y_{1,2}, \ Y_{1,2}, \ 1.25Y_{1,2}), \quad Y_{1,2} = -0.3 \pm \sqrt{r - 1.16} .$$

# Phase Portraits

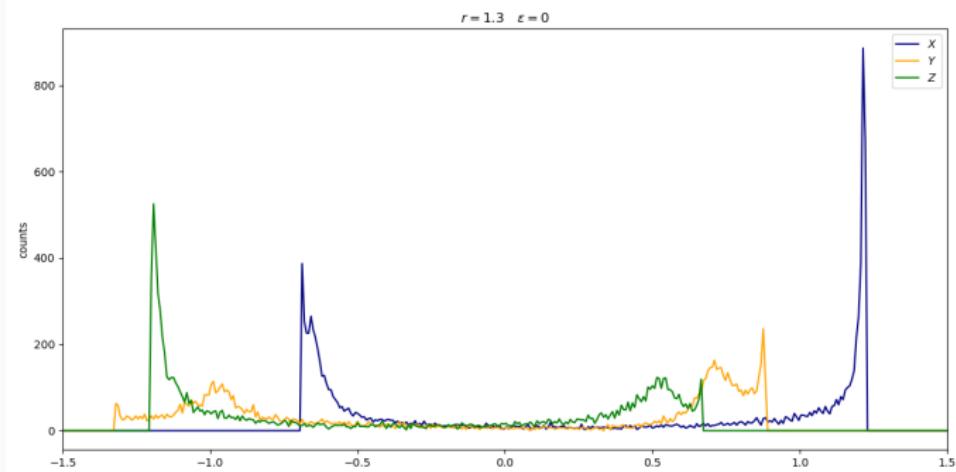
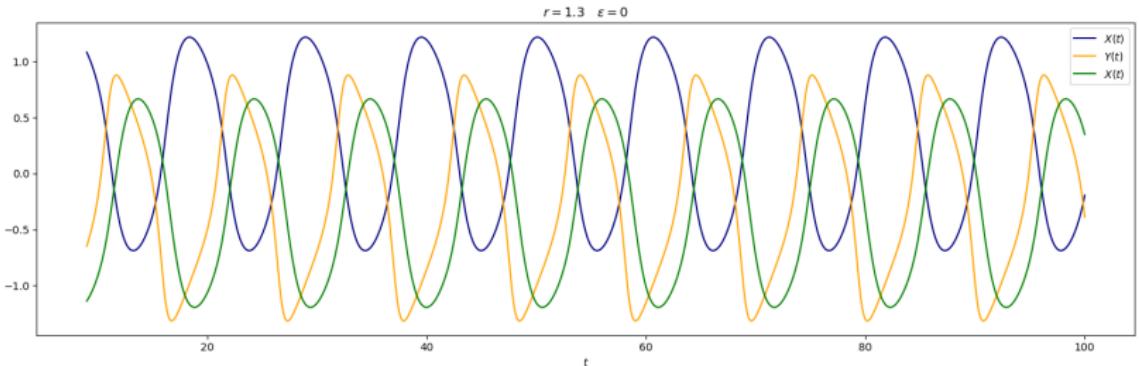


# Attractors: monostability and bistability regions

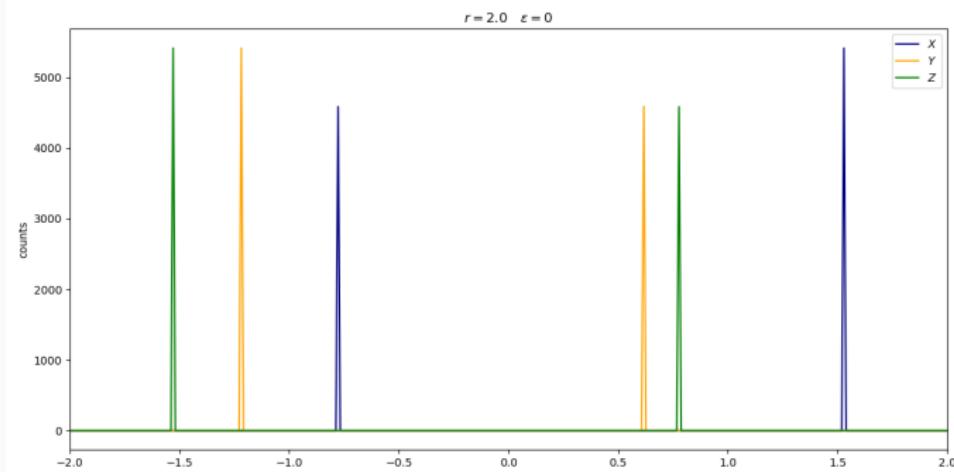
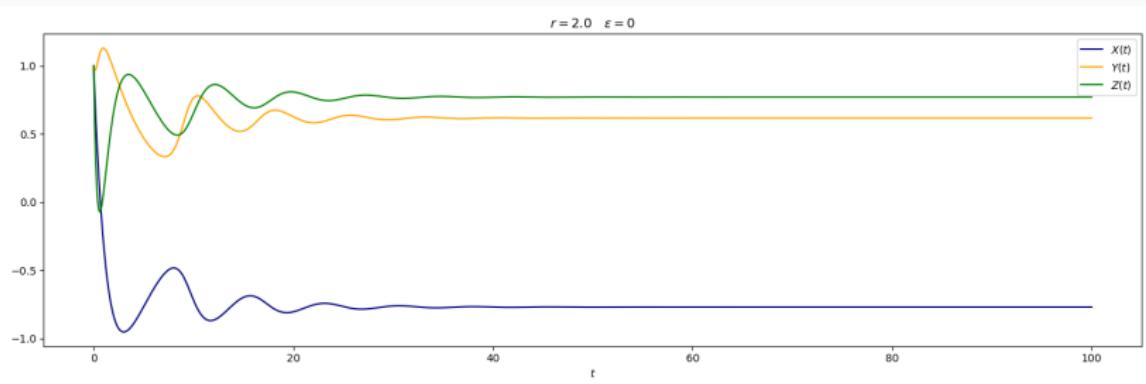


For  $r < 0.42$ : monostability in  $E_0$ . Around  $r \approx 0.42$ : bifurcation point and periodic cycles as attractors. Around  $r \approx 1.5$ , bistability with periodic cycles and  $E_2$ . Around  $r \approx 1.7$  periodic cycles loss and monostability (brown point). For  $r > 1.9$ , bistability with  $E_1$  and  $E_2$ .

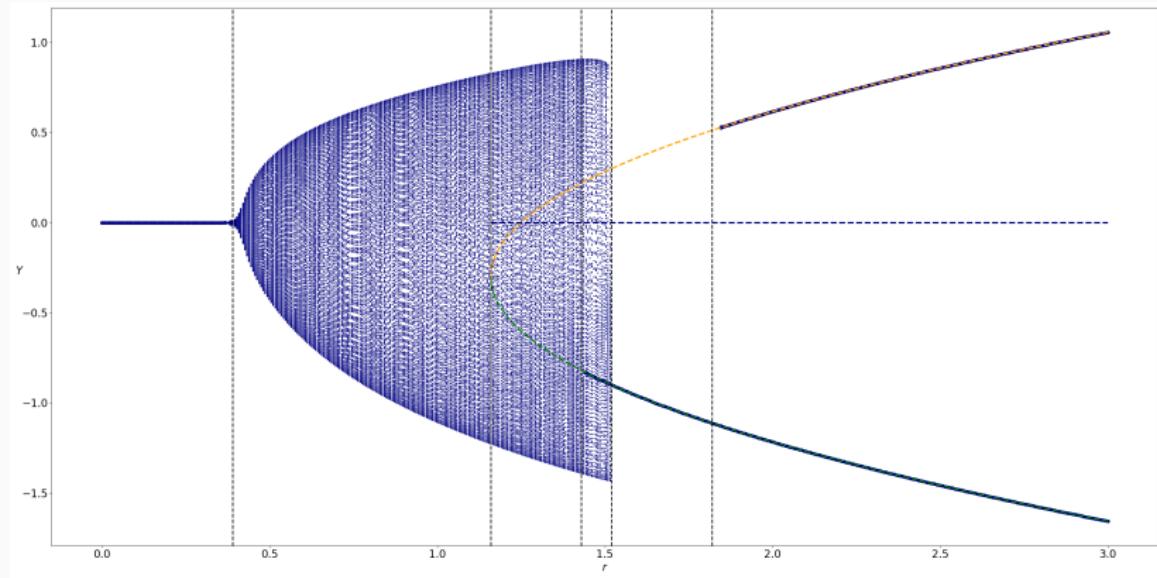
# Time Behavior and Histograms



# Time Behavior and Histograms



# Deterministic Bifurcation Diagram



$E_0 = (0, 0, 0)$  is stable for  $r < r_1 = 0.41704$ . At  $r = r_1$ , bifurcation with the birth of a stable limit cycle. This cycle loses its stability at  $r_4 = 1.517$ . The equilibria  $E_1, E_2$  appearing at  $r = r_2 = 1.16$  are unstable at first. The equilibrium  $E_2$  becomes stable at  $r_3 = 1.42271$ ,  $E_1$  at  $r_5 = 1.82025$ .

# Deterministic Bifurcation Diagram

- In the parametric region  $0 < r < r_1$ , the system is monostable with the trivial equilibrium  $E_0$  as a single attractor.
- In the zone  $r_1 < r < r_3$ , the climate system is monostable with the stable limit cycle as a single attractor.
- For  $r_3 < r < r_4$ , the system is bistable with two coexisting attractors, namely the equilibrium  $E_2$  and the limit cycle.
- For  $r_4 < r < r_5$ , the deterministic system is monostable with the equilibrium  $E_2$  representing a single attractor.
- If  $r > r_5$ , the system is bistable with two coexisting stable equilibria  $E_1$  and  $E_2$ .

## The Stochastic Model

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# The Stochastic Saltzman-Maasch model

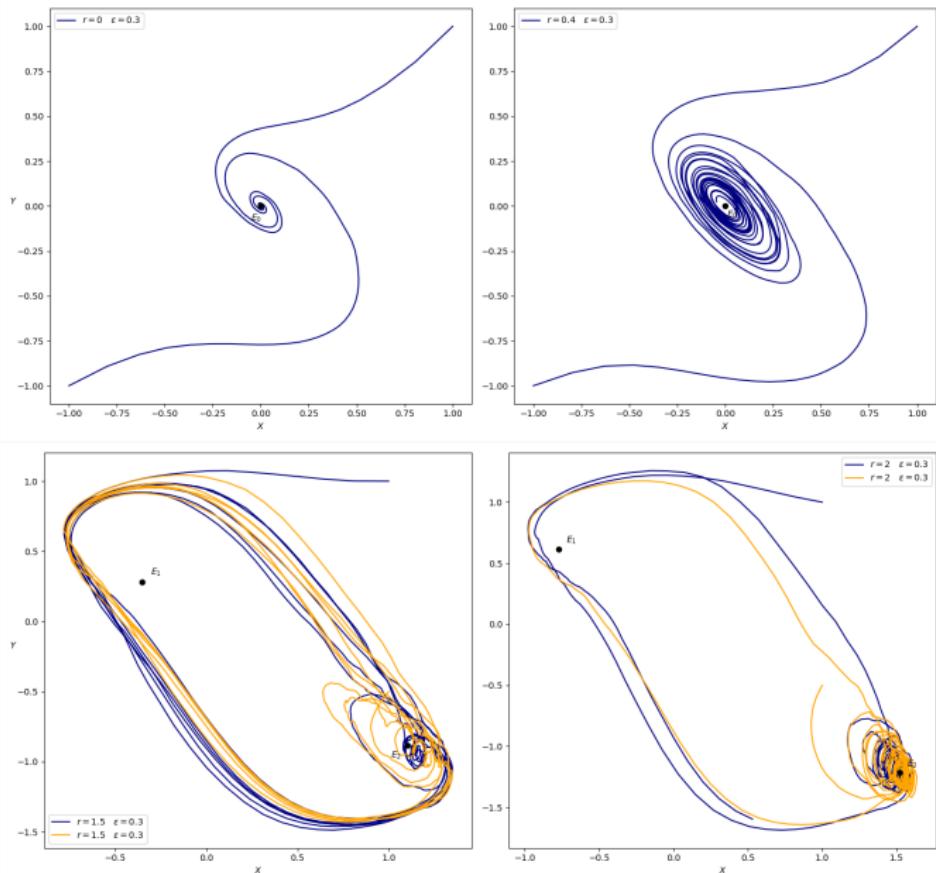
Now a stochastic dynamics is introduced, in order to reveal new features of noise-induced climate variability. Introduce in the model a **multiplicative white gaussian noise** term in the third equation.

$$\begin{cases} \dot{X} = -X - Y - vZ \\ \dot{Y} = -pZ + rY - sY^2 - Y^3 \\ \dot{Z} = -q(X + Z) + \epsilon Z\xi(t) \end{cases}$$

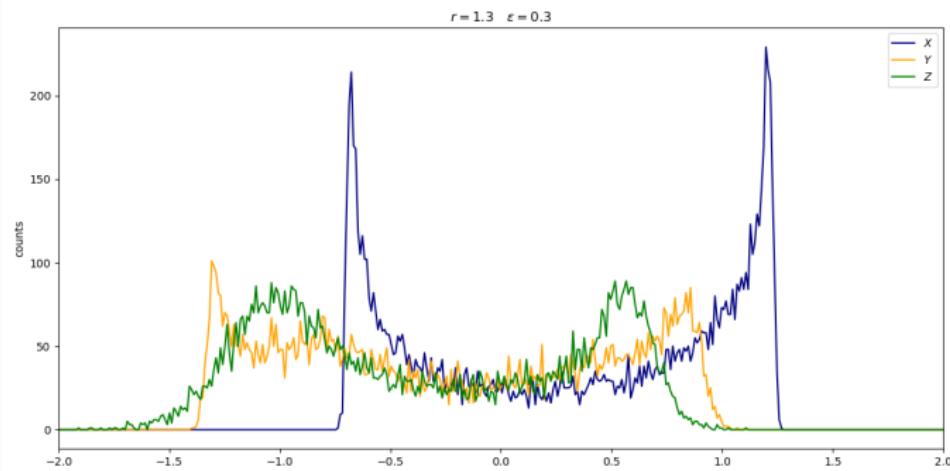
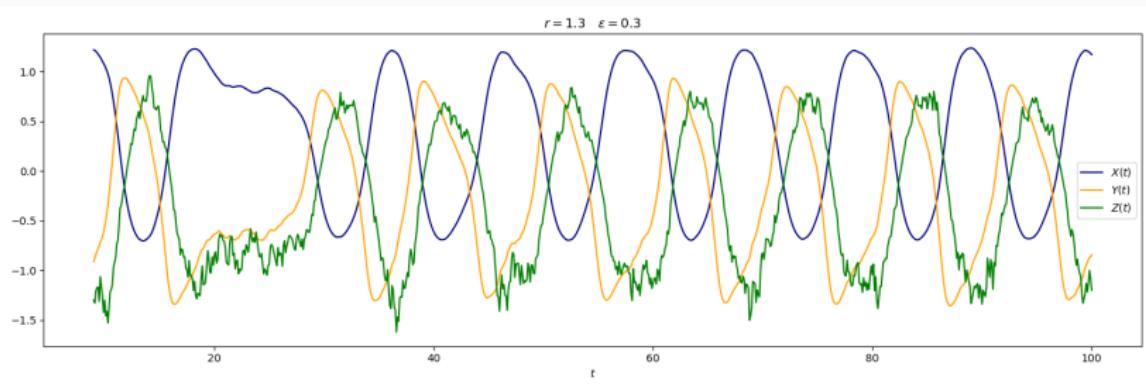
where  $\xi(t)$  is the multiplicative white noise.

It will be shown that the climate system can be highly noise-excitable and it possesses large-amplitude fluctuations even in those regions where the deterministic model does not contain any self-sustained oscillations.

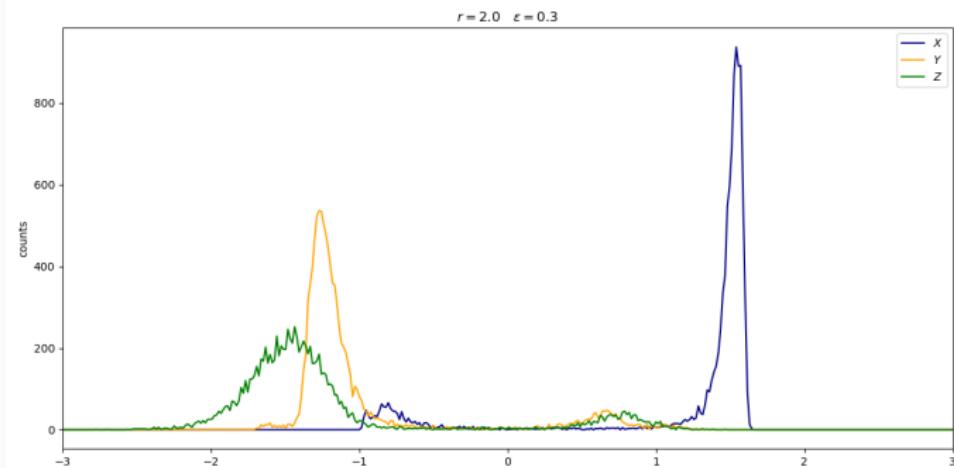
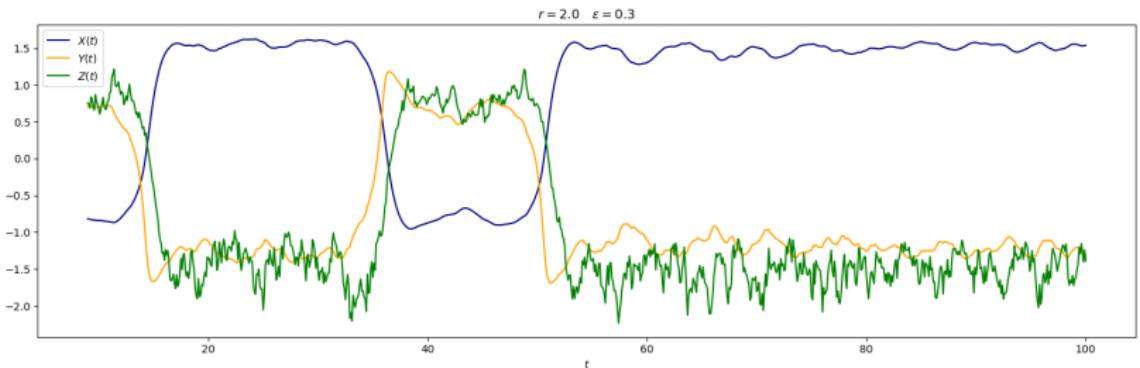
# Noisy Trajectories: Noisy Phase Portraits



# Time Behavior and Histograms

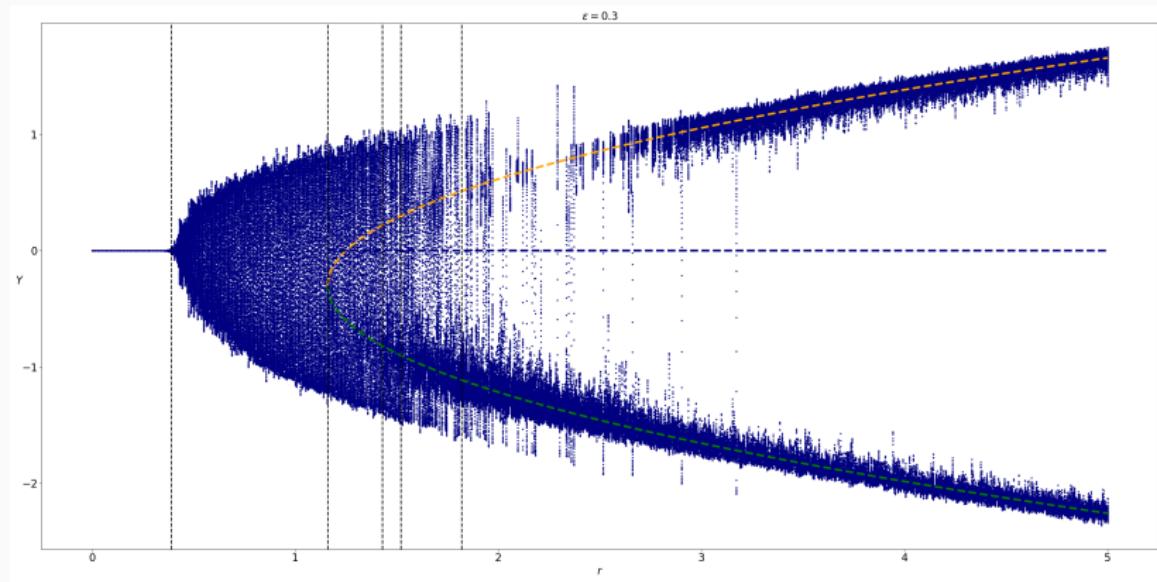


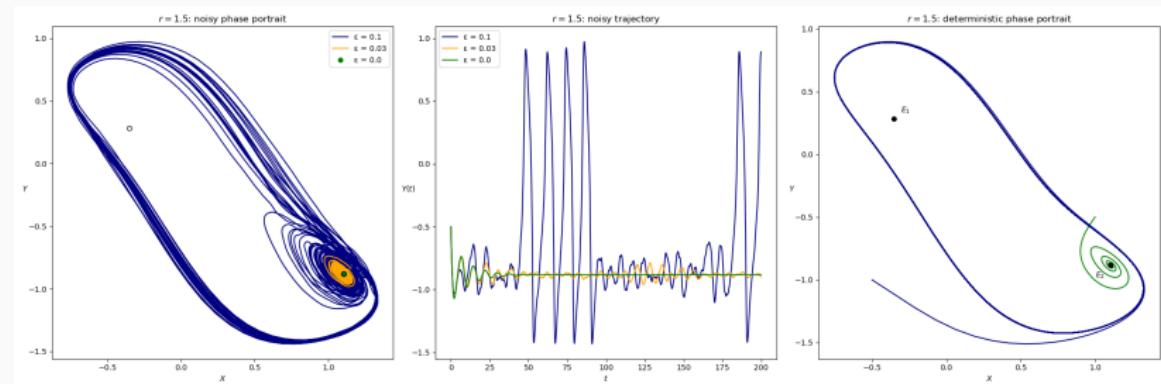
# Time Behavior and Histograms



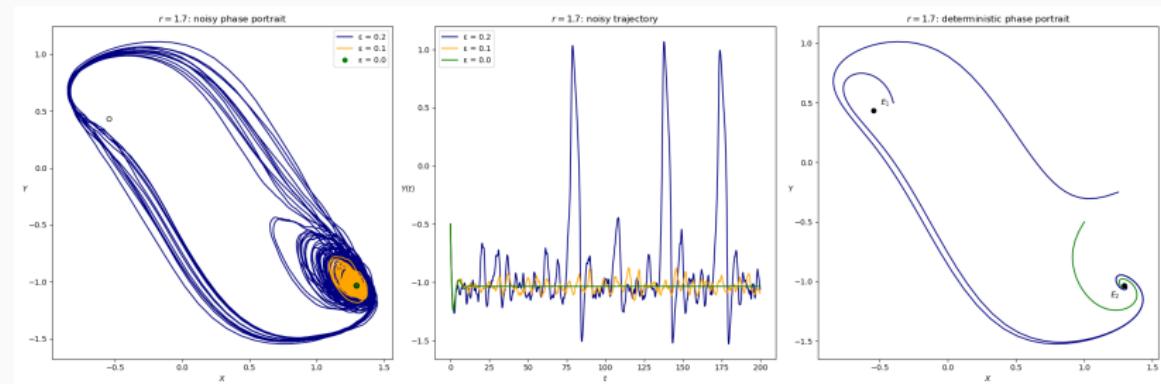
# Bifurcation Diagram with Noise

How the Phase Portrait has changed under the effect of noise: presence of noise-induced phenomena of transition.

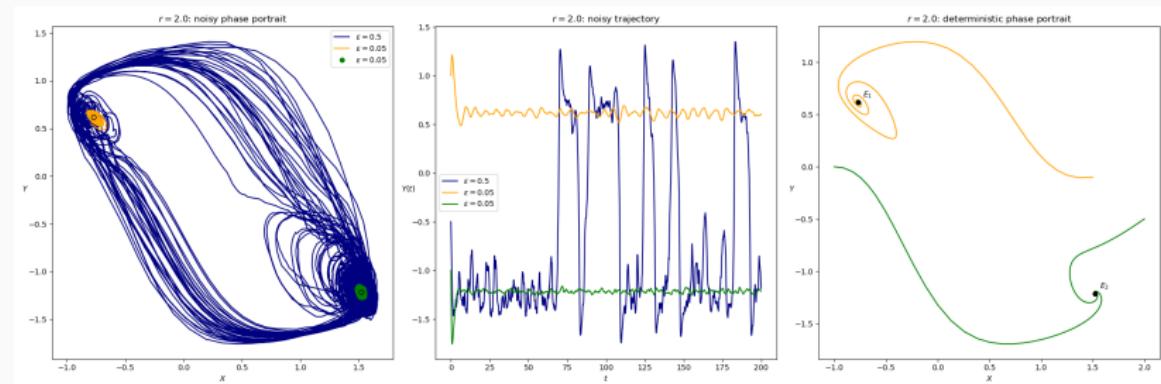




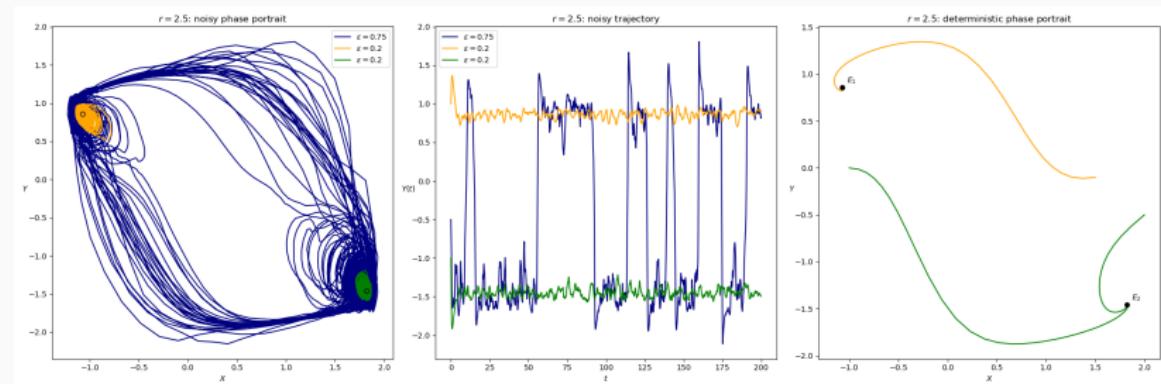
- $r = 1.5$ : bistability region  $r_3 < r < r_4$ .
- For weak noise,  $\epsilon = 0.03$ , the solutions exhibit small amplitude stochastically-induced oscillations (SASIO) near  $E_2$
- With increasing noise, the system starts to demonstrate large amplitude stochastically-induced oscillations (LASIO)
- This change is due to noise-induced transitions between the basins of attraction of the equilibrium  $E_2$  and the periodic cycle



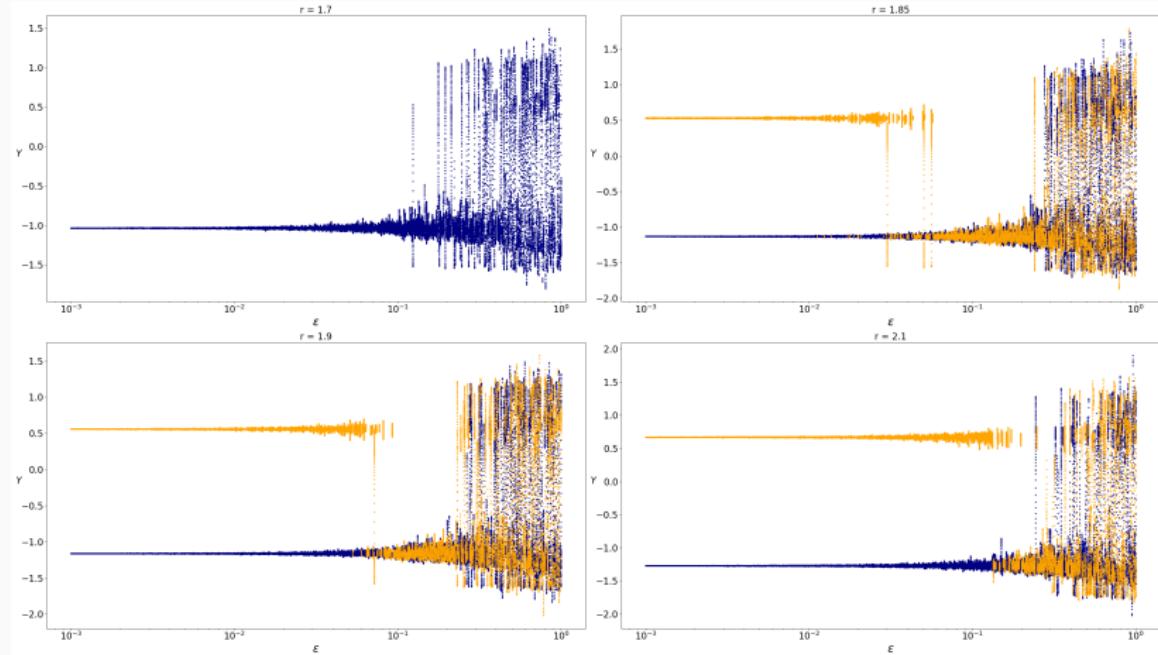
- Monostability region with  $r = 1.7$ :  $E_2$  as single attractor and no stable cycles
- For small deviations of initial states (subthreshold zone), the trajectories tend to  $E_2$ . For larger deviations (superthreshold zone), the trajectory exhibits a large-amplitude loop
- For  $\epsilon = 0.1$ , SASIO is observed near  $E_2$ , while for  $\epsilon = 0.2$ , a complex regime with alternating SASIO and LASIO is evident



- For  $r = 1.9$ , the deterministic system has two stable equilibria,  $E_1$  (warm climate state, in orange) and  $E_2$  (cold state, in green)
- For  $\epsilon = 0.05$ , the solutions starting near  $E_1$  and  $E_2$  are concentrated near their corresponding equilibria (subthreshold)
- With increasing noise, mutual transitions between the basins of attraction of these equilibria are revealed with complex oscillatory regime with alternating SASIO and LASIO



- For  $\epsilon = 0.2$ , solutions starting near  $E_1$ , result in small-amplitude oscillations
- As noise increases, trajectories fall into the superthreshold zone, then to the subthreshold zone of equilibrium  $E_2$  exhibiting SASIO
- With further noise increase, trajectories fall into the superthreshold zone of equilibrium  $E_2$ , generating large-amplitude oscillations around both  $E_1$  and  $E_2$



Interesting multiphase process of stochastic noise-induced transformations in the bistability zone  $r > r_5 \approx 1.82$

- If a trajectory begins near  $E_2$  (in blue), increasing noise reveals only two phases: small amplitude stochastic oscillations (SASIO) near  $E_1$  and large amplitude stochastic oscillations (LASIO) between  $E_1$  and  $E_2$
- If the trajectory starts from  $E_1$  (in orange), three phases are observed: SASIO near  $E_1$ , a transition to SASIO near  $E_2$ , and LASIO between  $E_1$  and  $E_2$ . This reveals a sharp transition from warm to cold climate state

## Improved Model

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## Improved Saltzman-Maasch model

1. Generalize the deterministic part of the model, adding nonlinearities and coupling terms
2. Generalize the multiplicative stochastic term: use a more sophisticated noise term, such as an Ornstein-Uhlenbeck process  $\eta(t)$ , and incorporating a different dependency on the variables
3. Including external stochastic forcing terms

$$\begin{cases} \dot{X} = -a(X + Y) - vZ + \omega_X \xi(t) \\ \dot{Y} = -pZ + rY - sY^2 - Y^3 + bYZ + cZY^2 + \omega_Y \xi(t) \\ \dot{Z} = -q(X + Z) + \epsilon Z(1 + dX)\eta(t) \end{cases}$$

where:

- $a, b, c, d \geq 0$
- $\xi(t)$  is a white noise term
- $\eta(t)$  can be a white noise term, or an Ornstein-Uhlenbeck process

# Equilibrium Points

In addition to the equilibrium  $E_0 = (0, 0, 0)$ , the new model presents equilibrium points corresponding to the values of  $Y$ :

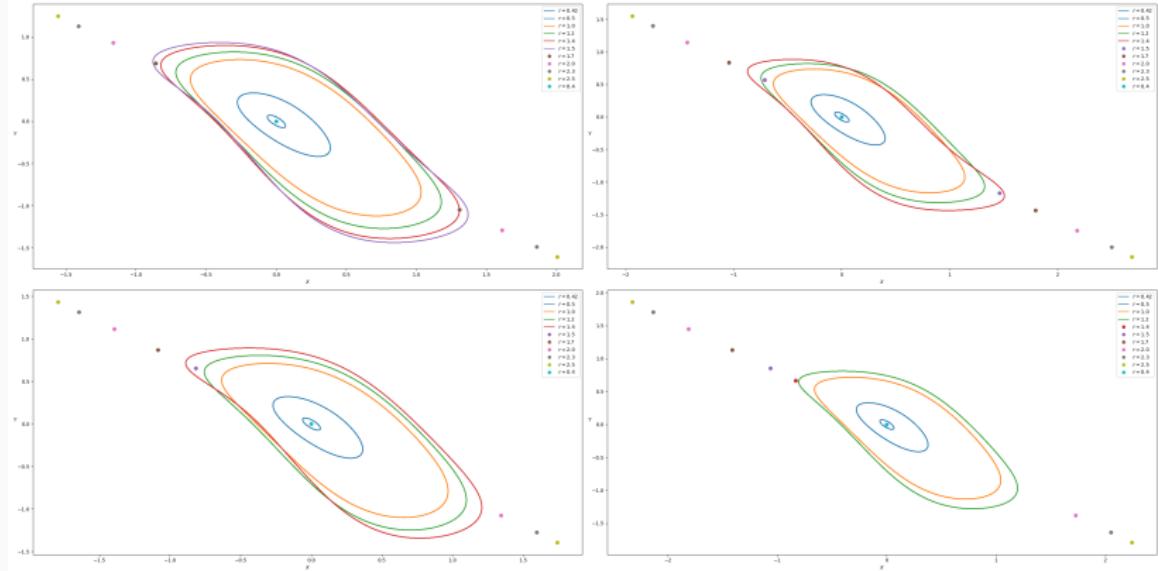
$$Y_{1,2} = \frac{\frac{ba}{a-v} - s \pm \sqrt{\left(\frac{ba}{a-v} - s\right)^2 - 4 \left(\frac{ca}{a-v} - 1\right) \left(r - \frac{pa}{a-v}\right)}}{2 \left(1 - \frac{ca}{a-v}\right)}$$

which leads to the non-trivial equilibrium points expression:

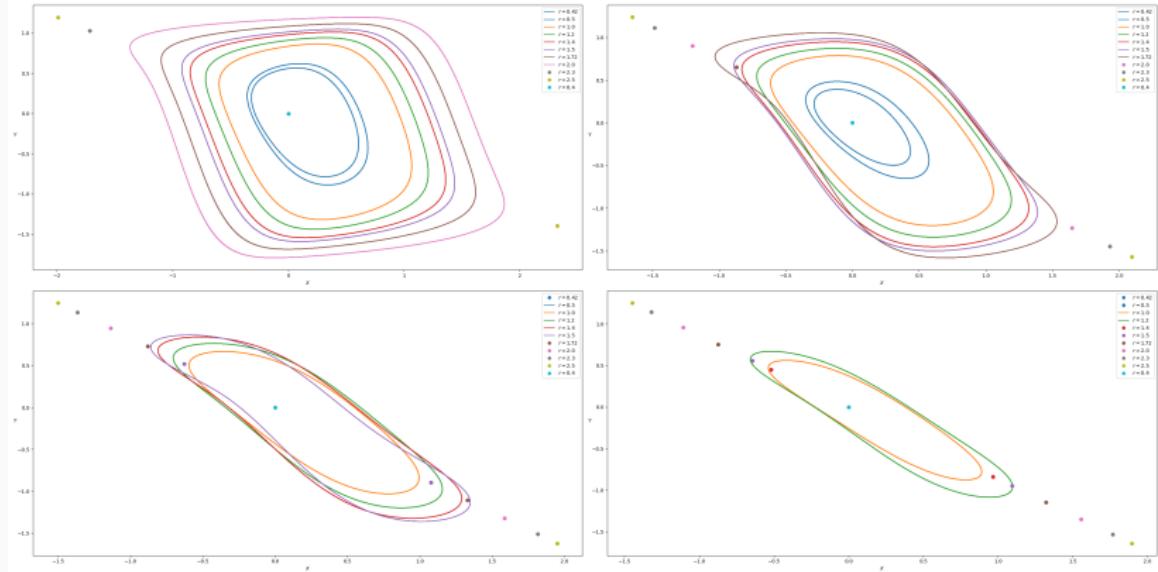
$$E_{1,2} = \left(-\frac{aY_{1,2}}{a-v}, Y_{1,2}, \frac{aY_{1,2}}{a-v}\right)$$

For  $b = c = 0$  and  $a = 1$ , we obtain the previous deterministic model.

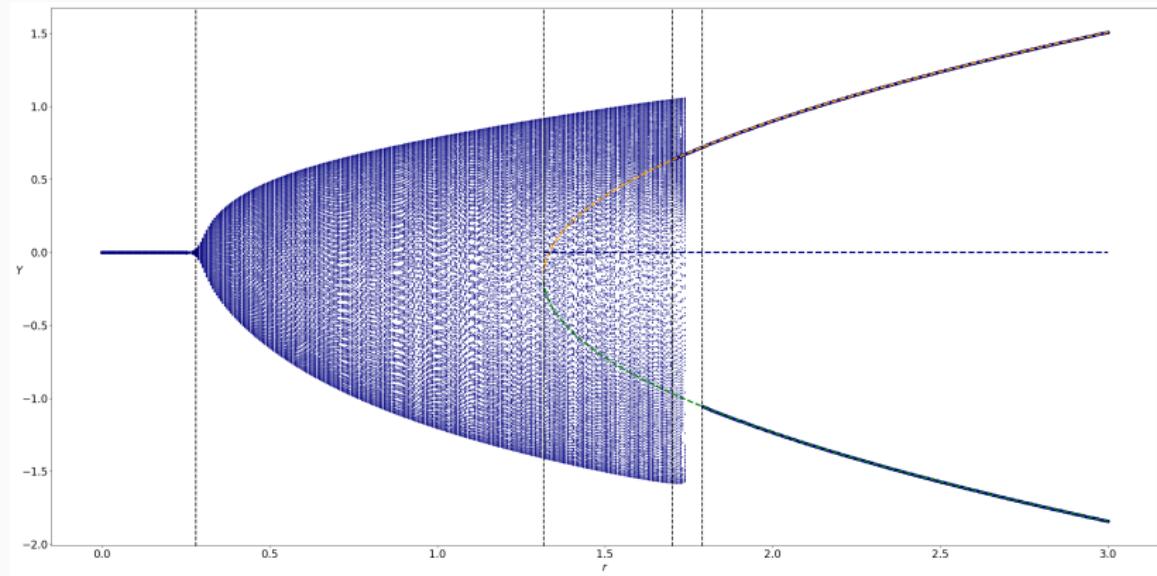
# Effects of the parameters $b$ and $c$



# Effects of the parameter $a$



# Deterministic Bifurcation Diagram

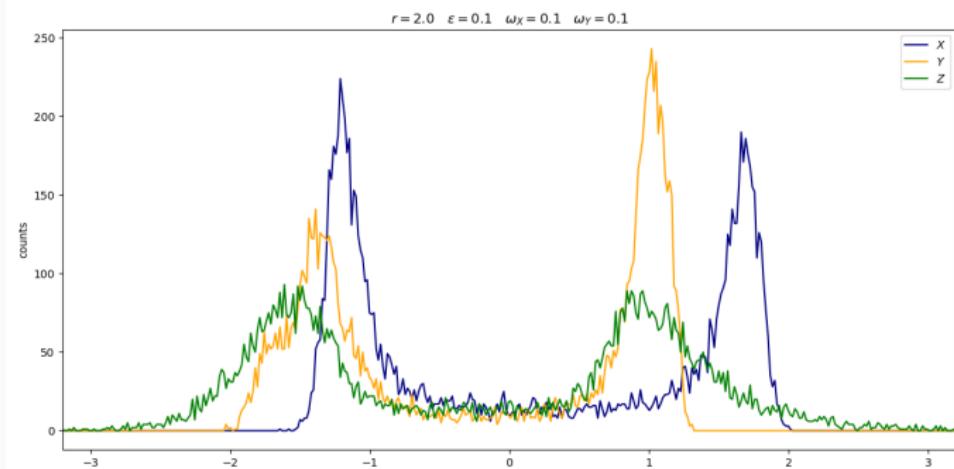
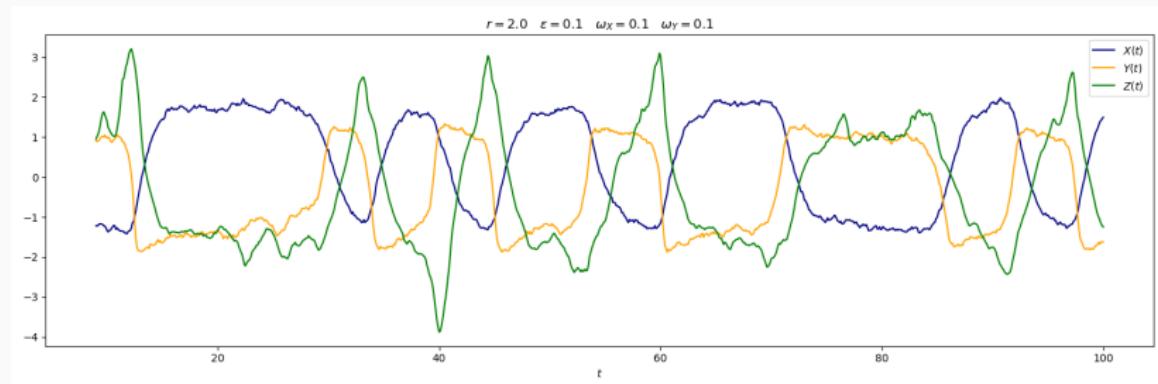


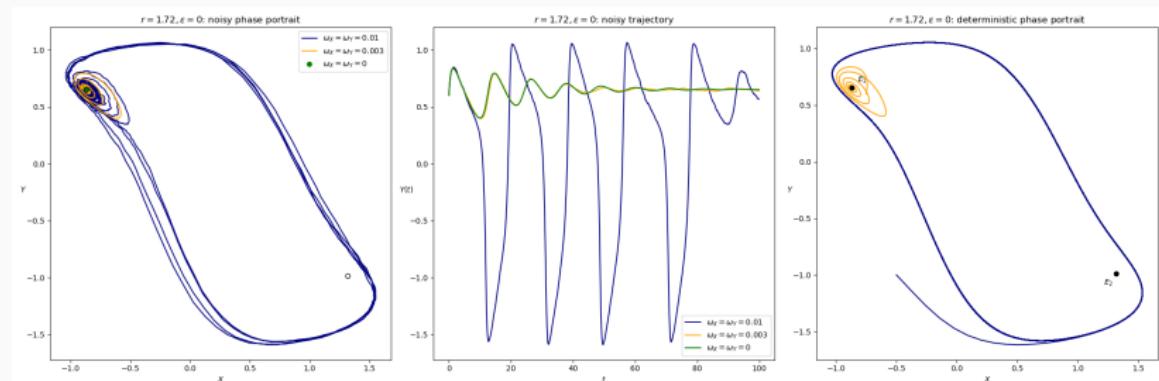
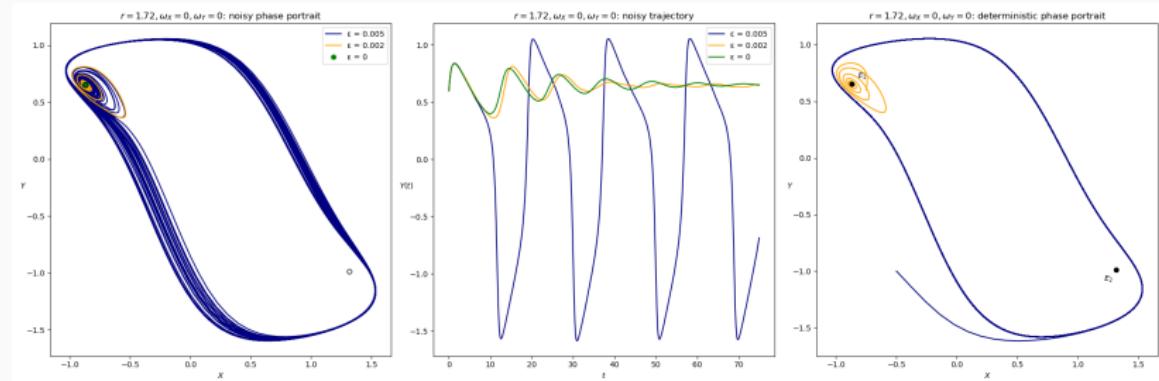
Fixed for the following analysis the values of the parameters:

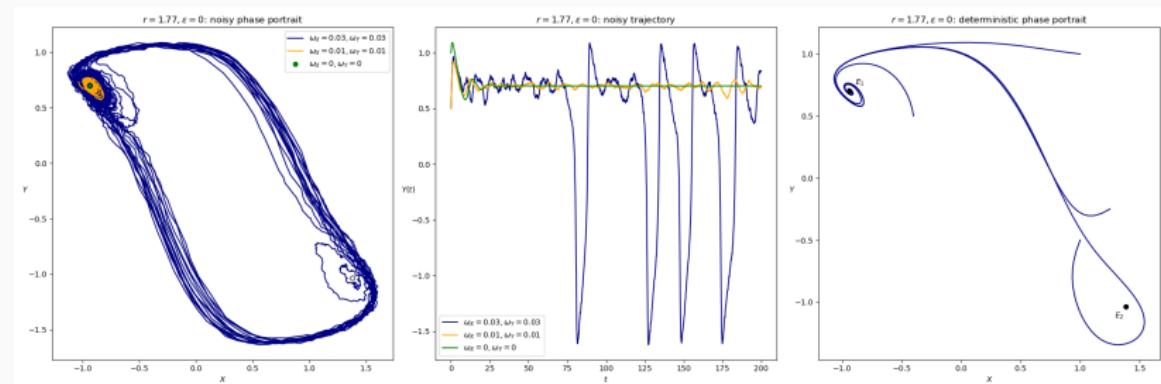
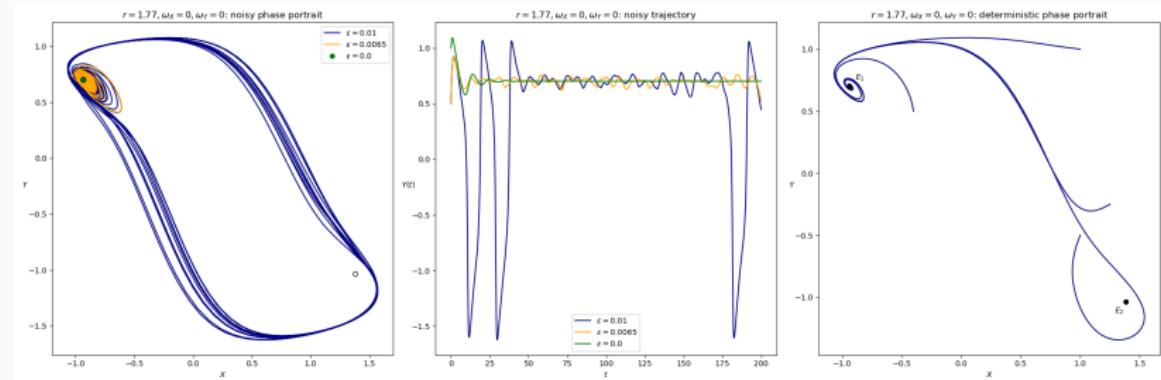
$$b = 0.3, c = 0.3, d = 0.2, a = 0.8.$$

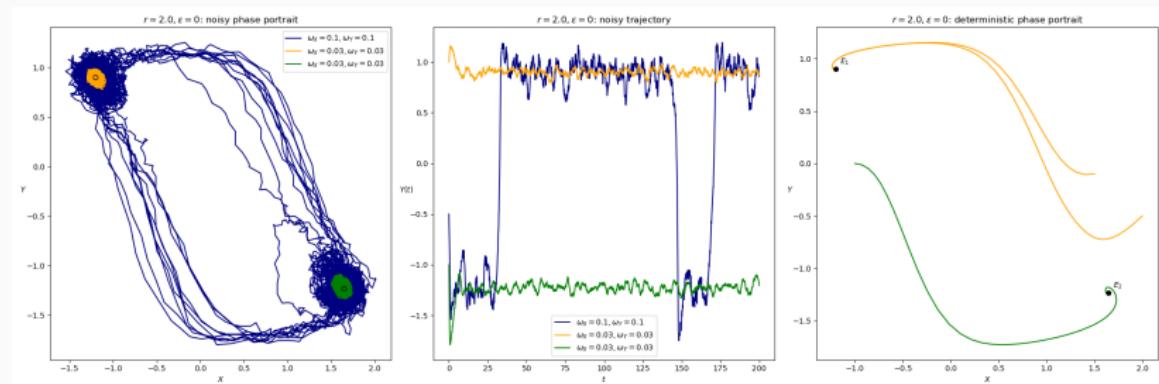
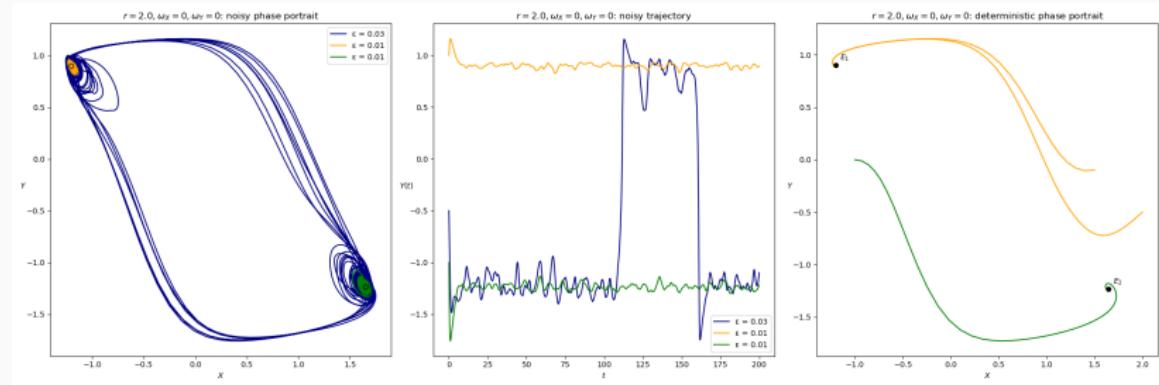
Interesting thing: with this combination of parameters, the equilibrium points  $E_1$  and  $E_2$  have the inverse stability.

# Stochastic system: time Behavior and Histograms









Some considerations on NITs in this case....

