

Lecture 4

Linear stability of parallel flows I

AE209 Hydrodynamic stability

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1. **Concept of parallel flows**
2. **Linearised equation for inviscid parallel shear flows**
3. **Normal mode solution**
4. **When does an instability occur?**

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Parallel shear flows

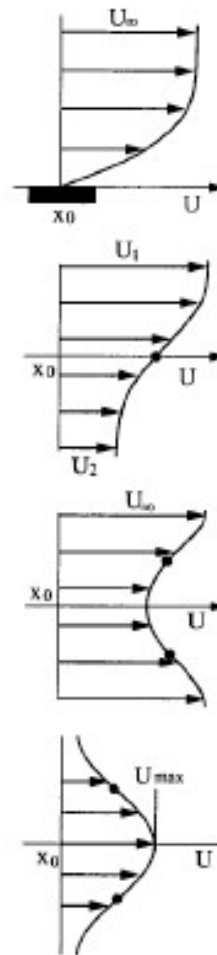
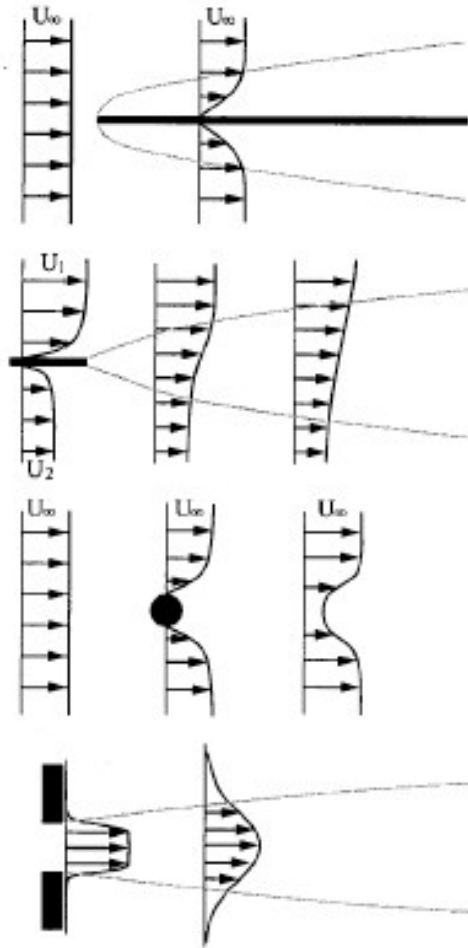
The flow configuration, the base flow of which is given by

$$\mathbf{U} = (U(y), 0, 0)$$

Examples

Plane Couette flow, Poiseuille flow, Pipe flow, and etc.

Weakly non-parallel flows (or spatially developing flows)



Boundary layer

Mixing layer

Cylinder wake

Jet

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Euler equation

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\nabla p$$

$$\nabla \cdot \mathbf{u} = 0$$

Consider $\mathbf{u}(x, y, z, t) = (U(y), 0, 0) + \varepsilon \mathbf{u}'(x, y, z, t),$

$$p(x, y, z, t) = P(y) + \varepsilon p'(x, y, z, t)$$

and neglect the terms at $O(\varepsilon^2)$. Then,

Linearised Euler equation around parallel base flow

$$\frac{\partial u'}{\partial t} + U \frac{\partial u'}{\partial x} + v' \frac{dU}{dy} = -\frac{\partial p'}{\partial x}$$

$$\frac{\partial v'}{\partial t} + U \frac{\partial v'}{\partial x} = -\frac{\partial p'}{\partial y} \quad \Rightarrow \quad \nabla^2 p = -2 \frac{dU}{dy} \frac{\partial v}{\partial x}$$

$$\frac{\partial w'}{\partial t} + U \frac{\partial w'}{\partial x} = -\frac{\partial p'}{\partial z}$$

$$\frac{\partial u'}{\partial x} + \frac{\partial v'}{\partial y} + \frac{\partial w'}{\partial z} = 0$$

1) Equation for wall-normal velocity

$$\frac{\partial v'}{\partial t} + U \frac{\partial v'}{\partial x} = -\frac{\partial p'}{\partial y} \qquad \nabla^2 p = -2 \frac{dU}{dy} \frac{\partial v}{\partial x}$$

$$\left[\left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right) \nabla^2 - \frac{d^2 U}{dy^2} \frac{\partial}{\partial x} \right] v' = 0$$

2) Equation for wall-normal vorticity

$$\frac{\partial u'}{\partial t} + U \frac{\partial u'}{\partial x} + v' \frac{dU}{dy} = - \frac{\partial p'}{\partial x}$$

$$\frac{\partial w'}{\partial t} + U \frac{\partial w'}{\partial x} = - \frac{\partial p'}{\partial z}$$

Wall-normal vorticity

$$\eta' = \frac{\partial u'}{\partial z} - \frac{\partial w'}{\partial x}$$

$$\left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right) \eta' = - \frac{dU}{dy} \frac{\partial v'}{\partial z}$$

Velocity and vorticity form of linearised Euler equation

$$\left[\left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right) \nabla^2 - \frac{d^2 U}{dy^2} \frac{\partial}{\partial x} \right] v' = 0$$

$$\left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right) \eta' + \frac{dU}{dy} \frac{\partial v'}{\partial z} = 0$$

with the boundary condition,

$$v' = \eta' = 0 \quad \text{at solid boundary and/or the far field}$$

and the initial condition,

$$v'(x, y, z, t = 0) = v'_0(x, y, z)$$

$$\eta'(x, y, z, t = 0) = \eta'_0(x, y, z)$$

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Consider two-dimensional case such that

$$\left[\left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right) \nabla^2 - \frac{d^2 U}{dy^2} \frac{\partial}{\partial x} \right] v' = 0$$

Then, the normal mode solution takes the following form:

$$v'(x, y, t) = \tilde{v}(y) e^{i(\alpha x - \omega t)} + c.c$$

where $\alpha \in R$ and $\omega \in C$.

For example,

$$\frac{\partial v'}{\partial t} =$$

$$\frac{d^2 U}{dy^2} \frac{\partial v'}{\partial x} =$$

$$\nabla^2 v' =$$

Rayleigh equation

Let $\omega = \alpha c$ and $D \equiv d / dy$. Then, we get

$$(U - c)(D^2 - \alpha^2)\tilde{v} - D^2U \tilde{v} = 0$$

with the boundary condition,

$$\tilde{v} = 0 \quad \text{at solid boundary and/or the far field}$$

Remark

If $\alpha \in R$ is given, then $c \in C$ becomes unknown with \tilde{v} , resulting in **eigenvalue problem**.

Normal mode solution

$$\begin{aligned}v'(x, y, t) &= \tilde{v}(y)e^{i\alpha(x-ct)} + c.c \\&= \text{Real}\left\{ \tilde{v}(y) \mid e^{i\phi(y)} e^{i\alpha(x-(c_r+ic_i)t)} \right\} \\&= |\tilde{v}(y)| e^{\alpha c_i t} \cos[\alpha(x - c_r t) + \phi(y)]\end{aligned}$$

c_r : Phase speed

$c_i > 0$ Linearly unstable

$c_i = 0$ Marginally stable (or neutral)

$c_i < 0$ Linearly stable

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Theorem: Rayleigh inflection point criterion

If there exist perturbations with $c_i > 0$, then d^2U / dy^2 must be zero at some $y \in \Omega$ ($\Omega = [a, b]$ is the flow domain in y).

Proof

$$(U - c)(D^2 - \alpha^2)\tilde{v} - D^2U \tilde{v} = 0$$

$$\int_a^b \tilde{v}^* (D^2 - \alpha^2) \tilde{v} dy - \int_a^b \frac{D^2U}{(U - c)} \tilde{v}^* \tilde{v} dy = 0$$

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$$\int_a^b |D\tilde{v}|^2 + \alpha^2 |\tilde{v}|^2 dy + \int_a^b \frac{d^2 U / dy^2}{(U - c)} |\tilde{v}|^2 dy = 0$$

Imaginary part:

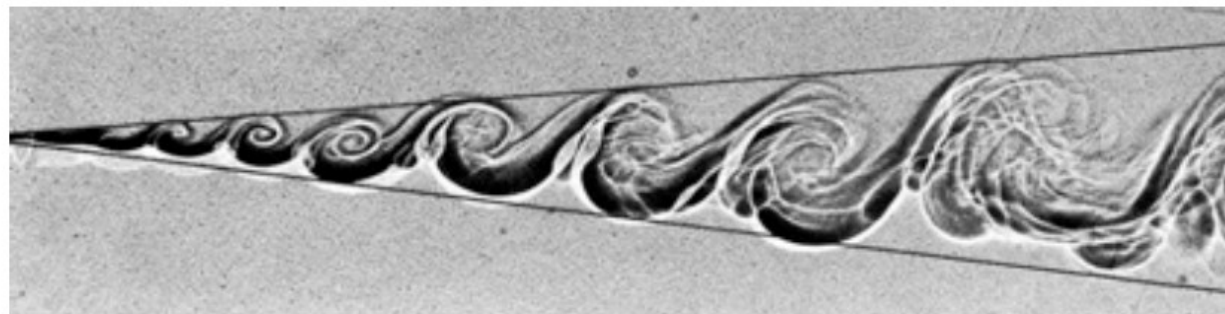
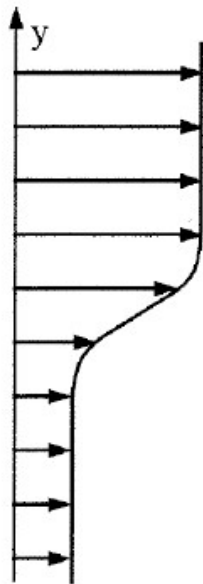
$$\text{Im} \left(\int_a^b \frac{d^2 U / dy^2}{(U - c)} |\tilde{v}|^2 dy \right) = \int_a^b \frac{c_i d^2 U / dy^2 |\tilde{v}|^2}{|U - c|^2} dy = 0$$

Remark 1

The presence of some **inflection points** is a **necessary condition** for **linear instability**.

Remark 2

The presence of an **inflection point** is often a **sign of instability**.



Instability in a mixing layer
(Brown & Roshko 1974)

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4. **Rayleigh thereom**