

Lecture 10

Spatio-temporal evolution of instabilities II

AE209 Hydrodynamic stability

Dr Yongyun Hwang

1. Application to Ginzburg-Landau equation
2. Application to wake
3. Physical implications: oscillator vs amplifier flows

- 1. Application to Ginzburg-Landau equation**
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Complex linear Ginzburg-Landau equation

From dispersion relation, $D(k, \omega) = 0$:

$$\omega(k) = Uk - c_d k^2 + i(\mu - k^2)$$

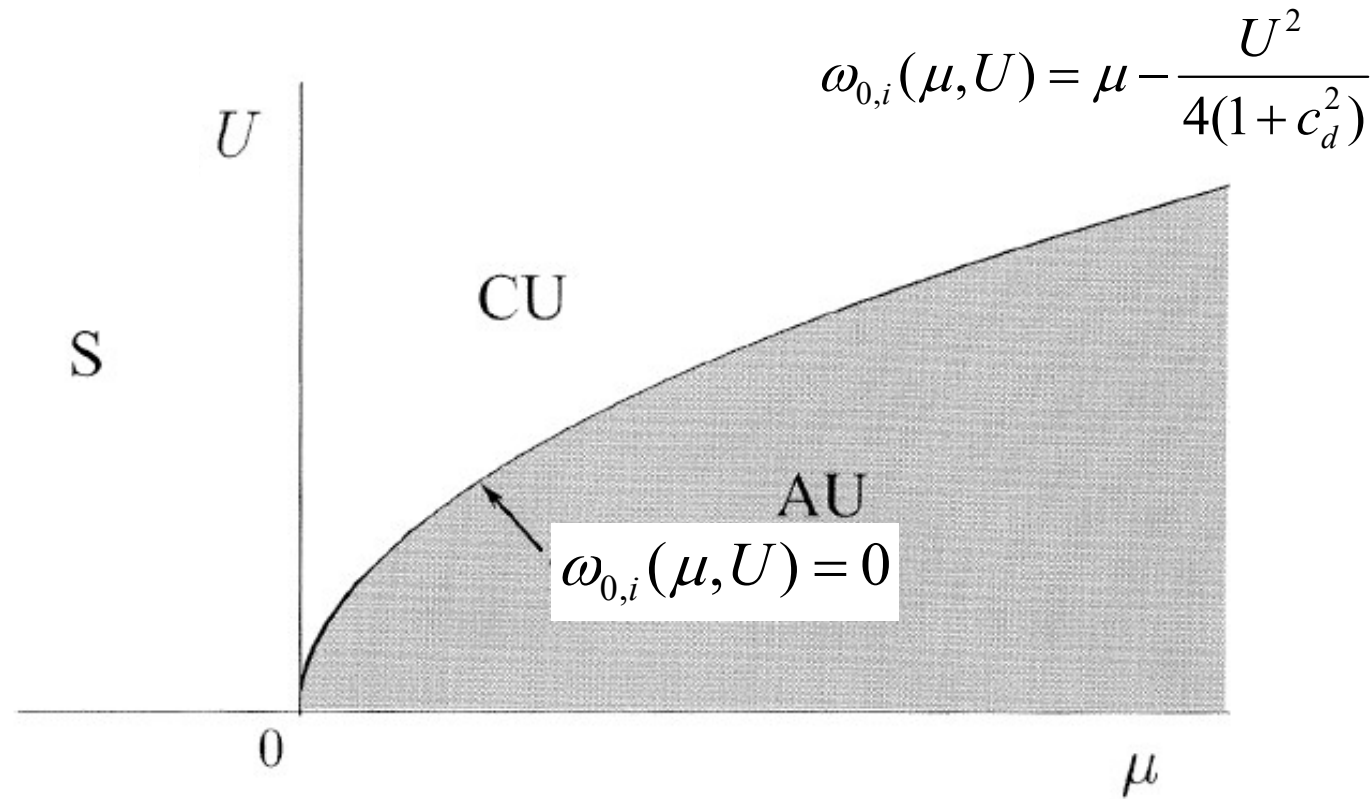
Linear stability: Calculate the maximum growth rate

$$\omega(k_{\max}) = i\omega_{i,\max} = i\mu \quad \text{with} \quad k_{\max} = 0$$

Absolute instability: Calculate the absolute growth rate

$$\omega_0 = \omega(k_0) = \frac{c_d U^2}{4(1 + c_d^2)} + i \left[\mu - \frac{U^2}{4(1 + c_d^2)} \right] \quad \text{with} \quad k_0 = \frac{U}{2(c_d + i)}$$

Absolute and convective instabilities in parametric space



Remark

This description is only important for the system with mean advection.

1. Application to Ginzburg-Landau equation
2. **Application to bluff-body wake**
3. Physical implications: oscillator vs amplifier flows

Family of wake profiles

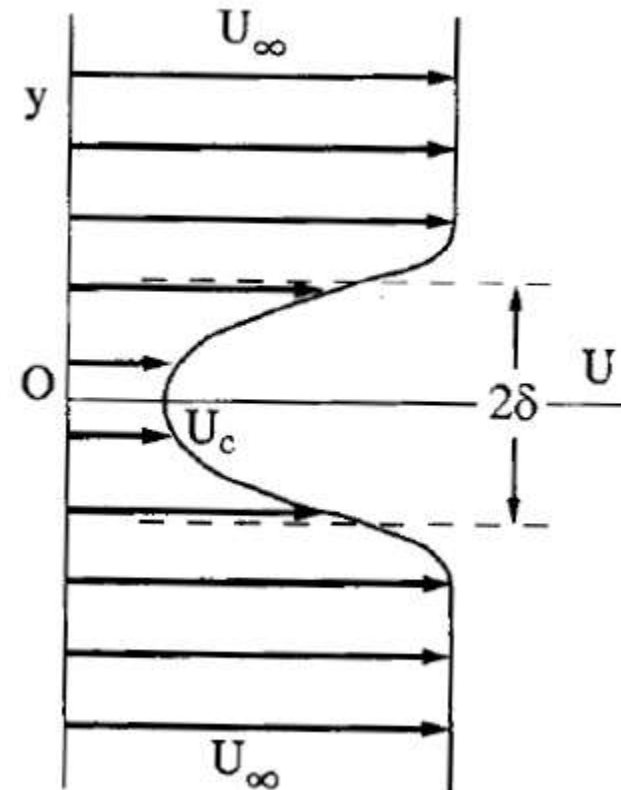
$$U(y) = U_{\infty} + (U_{\infty} - U_c)U_1\left(\frac{y}{\delta}; N\right)$$

where

$$U_1(\xi; N) = \left[1 + \sinh^{2N} \left\{ \xi \sinh^{-1}(1) \right\}\right]^{-1}$$

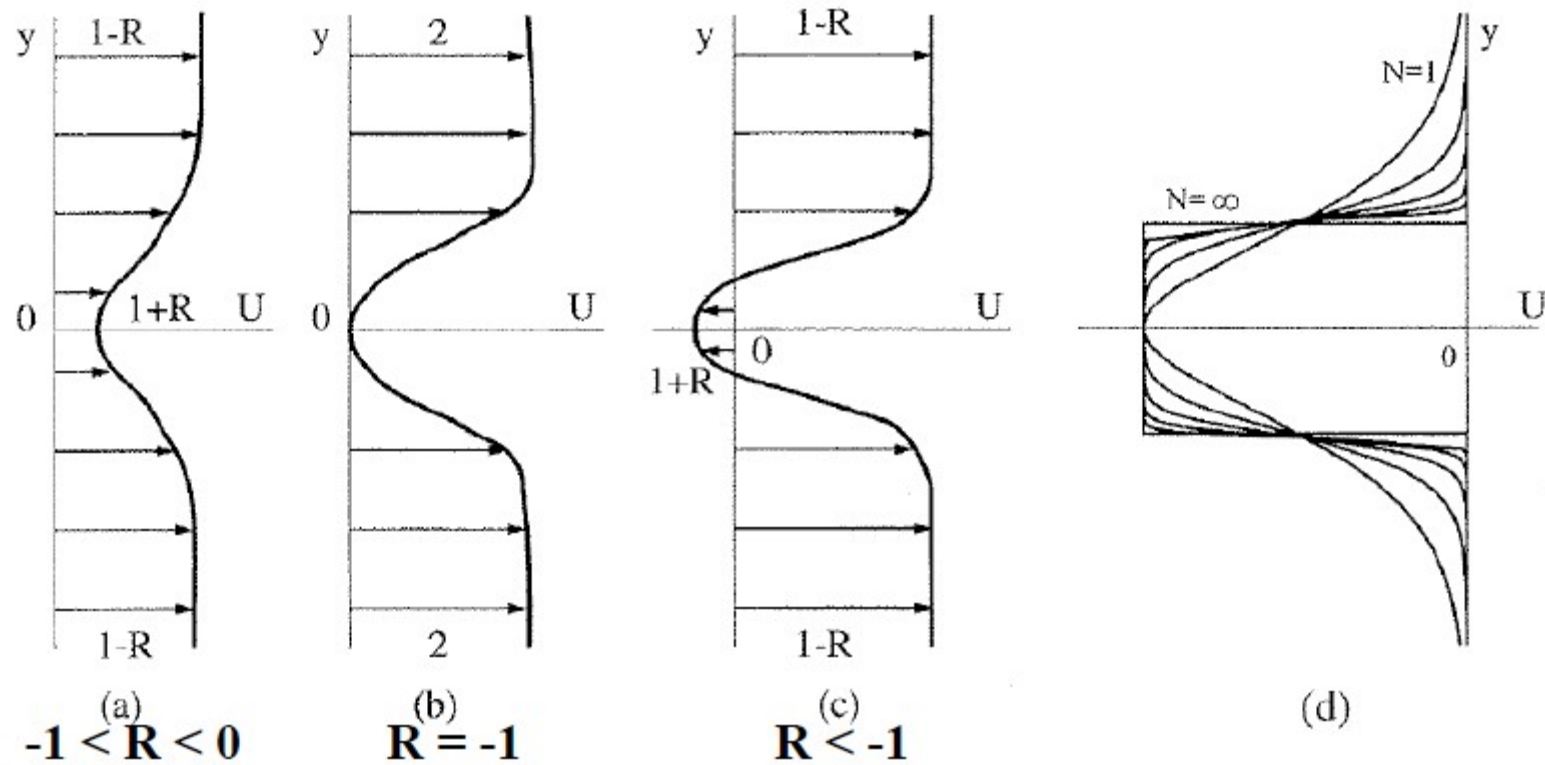
with

$$\text{Re} = \frac{\bar{U}\delta}{\nu}$$



Monkewitz (1988)

Family of wake profiles



Velocity ratio
$$R = \frac{U_c - U_\infty}{U_c + U_\infty}$$

N: stiffness

Monkewitz (1988)

Dispersion relation (obtained by solving the Orr-Sommerfeld equation)

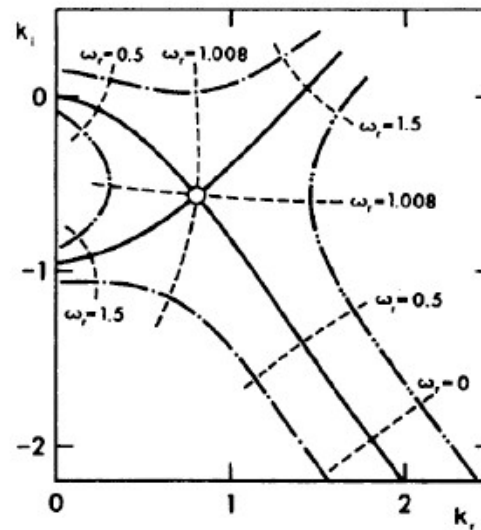
$$\left[(-i\omega + ikU)(D^2 - k^2) - ikD^2U - \frac{1}{\text{Re}}(D^2 - k^2)^2 \right] \tilde{v} = 0$$

with the **saddle point** behaviour:

$$R = -1$$

$$N = 2$$

$$\text{Re} = 11.3$$



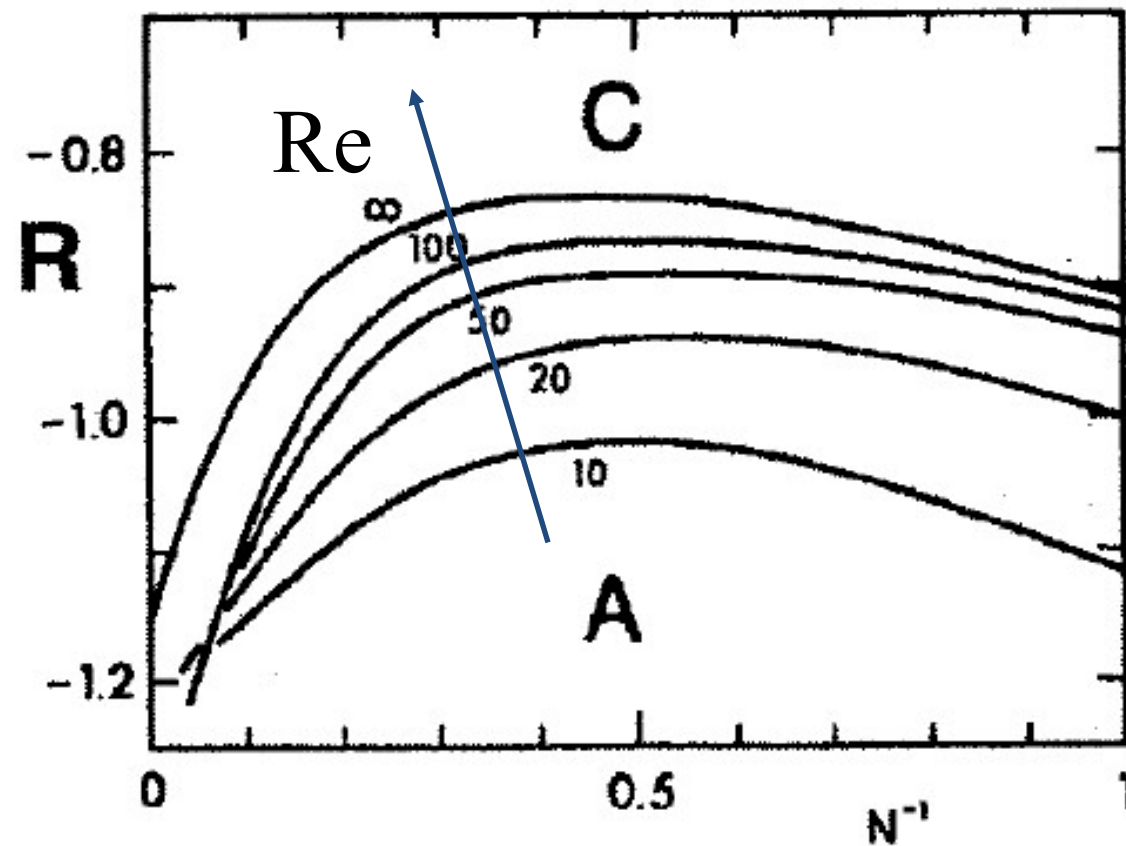
$$\omega_0 = 1.008 + 0i$$

$$\omega - \omega_0 \sim \frac{1}{2} \frac{\partial^2 \omega}{\partial k^2}(k_0)(k - k_0)^2$$

Monkewitz (1988)

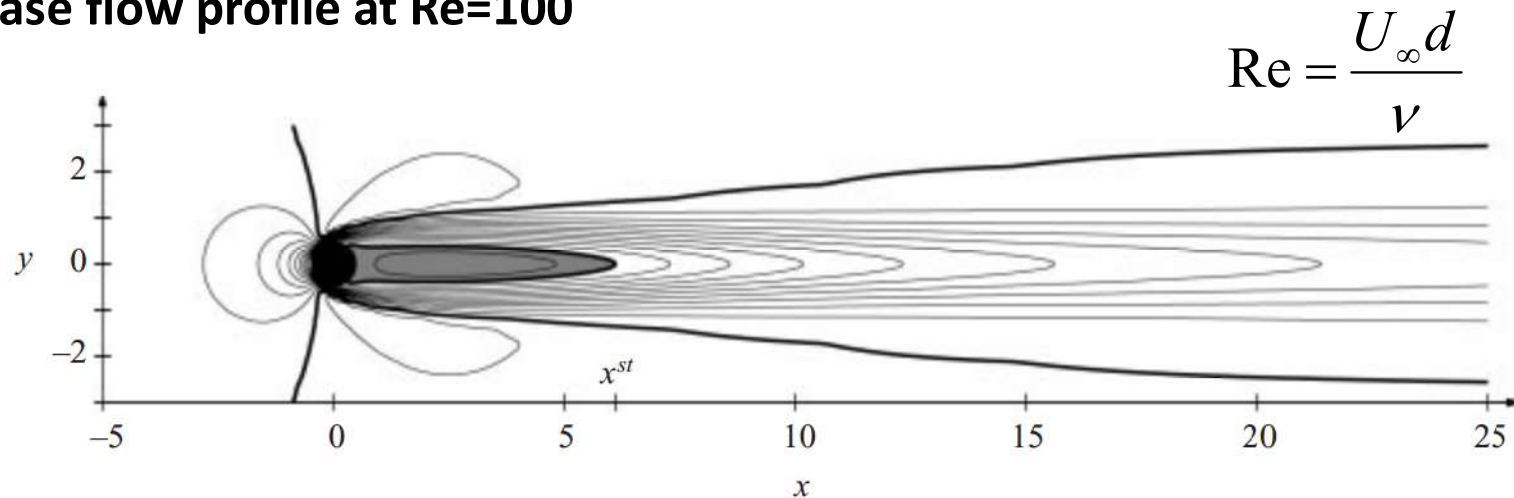
—, curve $\omega_i = 0$; - · - · - , $\omega_i = 0.25$; - · · - · - , $\omega_i = -0.25$

Effect of velocity ratio, stiffness and Reynolds number

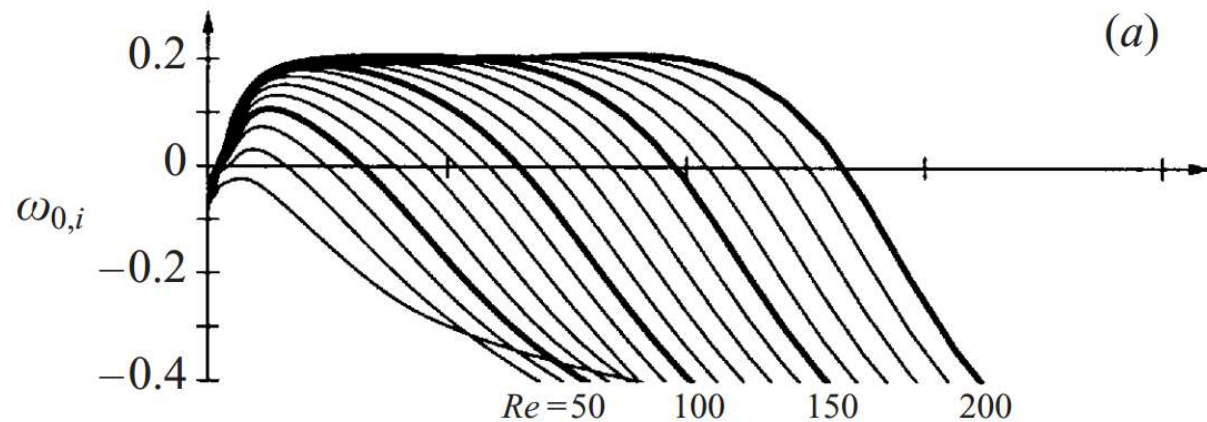


Monkewitz (1988)

Base flow profile at $Re=100$



Absolute growth rate of velocity profile at each streamwise location



Pier (2002)

Emergence of vortex shedding

$$5 < \text{Re} < 25$$

Some regions are convectively unstable

$$25 < \text{Re} < 47$$

Some regions are absolutely unstable

$$\text{Re} \approx 47$$

Strong local absolute instability leads to a **global instability** in the form of

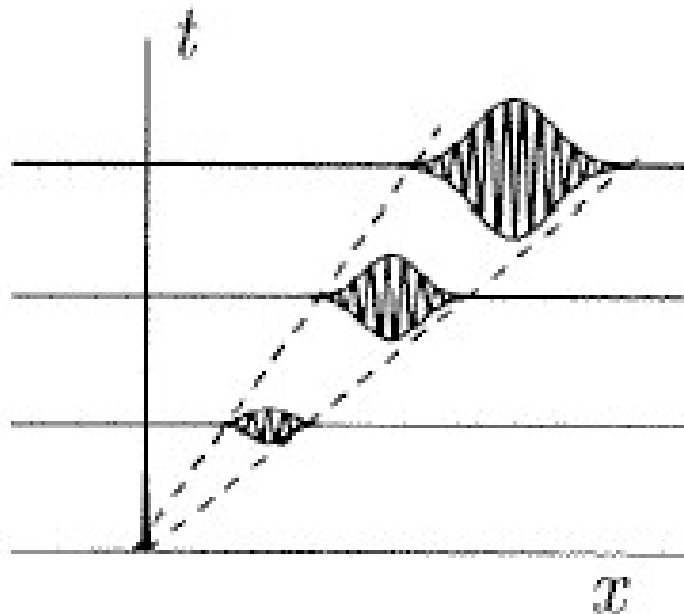
$$\mathbf{u}'(x, y, t) = \hat{\mathbf{u}}(x, y)e^{-i\omega_G t} \quad \text{with} \quad \omega_{G,i} > 0$$

Remark

Local absolute instability is a **necessary condition** for the onset of a **global instability of a fully non-parallel open flow**

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Remarks



Convective instability

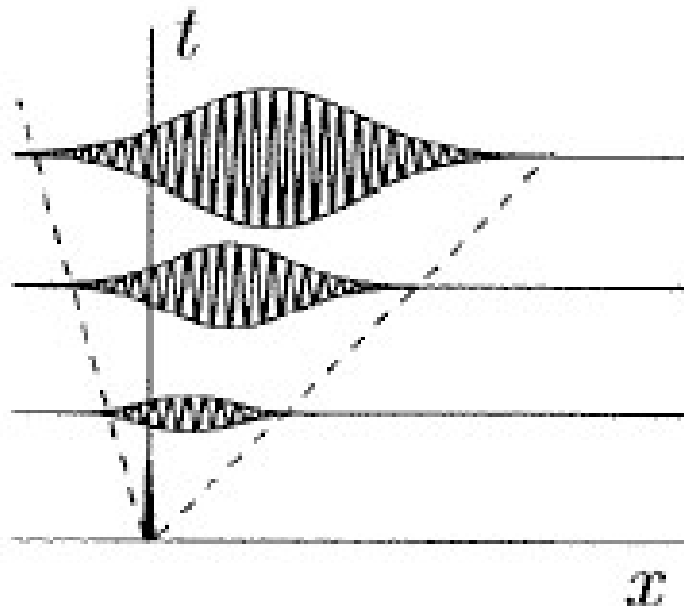
Ex)

1. Reference control volume returns to the original state after the impulse moves away downstream.

2. **Spatial stability analysis** becomes **meaningful** in this situation.

3. Instability dynamics is **driven by upstream noise**

Remarks



Absolute instability

Ex)

1. Reference control volume never return to the original state.

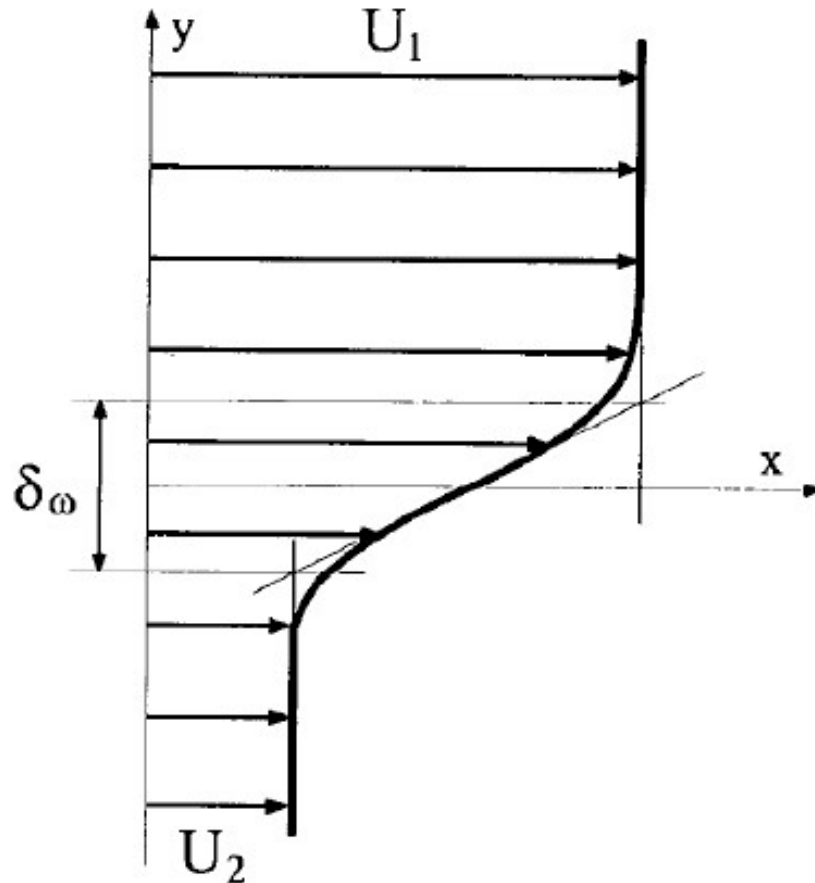
2. **Spatial stability analysis** becomes **meaningless** in this situation.

3. Instability dynamics is **intrinsically** driven by the given system and often results in a **nonlinear oscillation** with a **distinct frequency**.

Example: Mixing layer

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Hyperbolic tangent mixing layer



Base flow profile

$$U(y) = \bar{U} + \frac{\Delta U}{2} \tanh\left(\frac{2y}{\delta_w}\right)$$

Velocity ratio

$$R = \frac{U_1 - U_2}{U_1 + U_2} = \frac{\Delta U}{2\bar{U}}$$

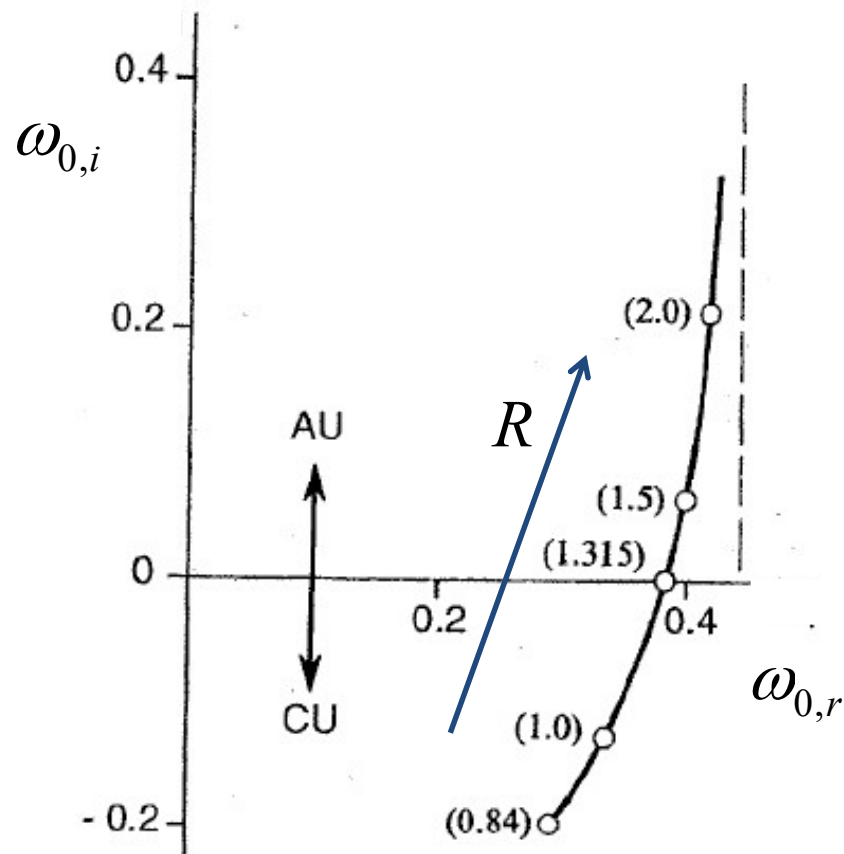
$$\delta_w = \frac{\Delta U}{(dU/dy)_{\max}}$$

Huerre & Monkewitz (1985)

Example: Mixing layer

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Theory: transition from convective to absolute instability with R



Velocity ratio

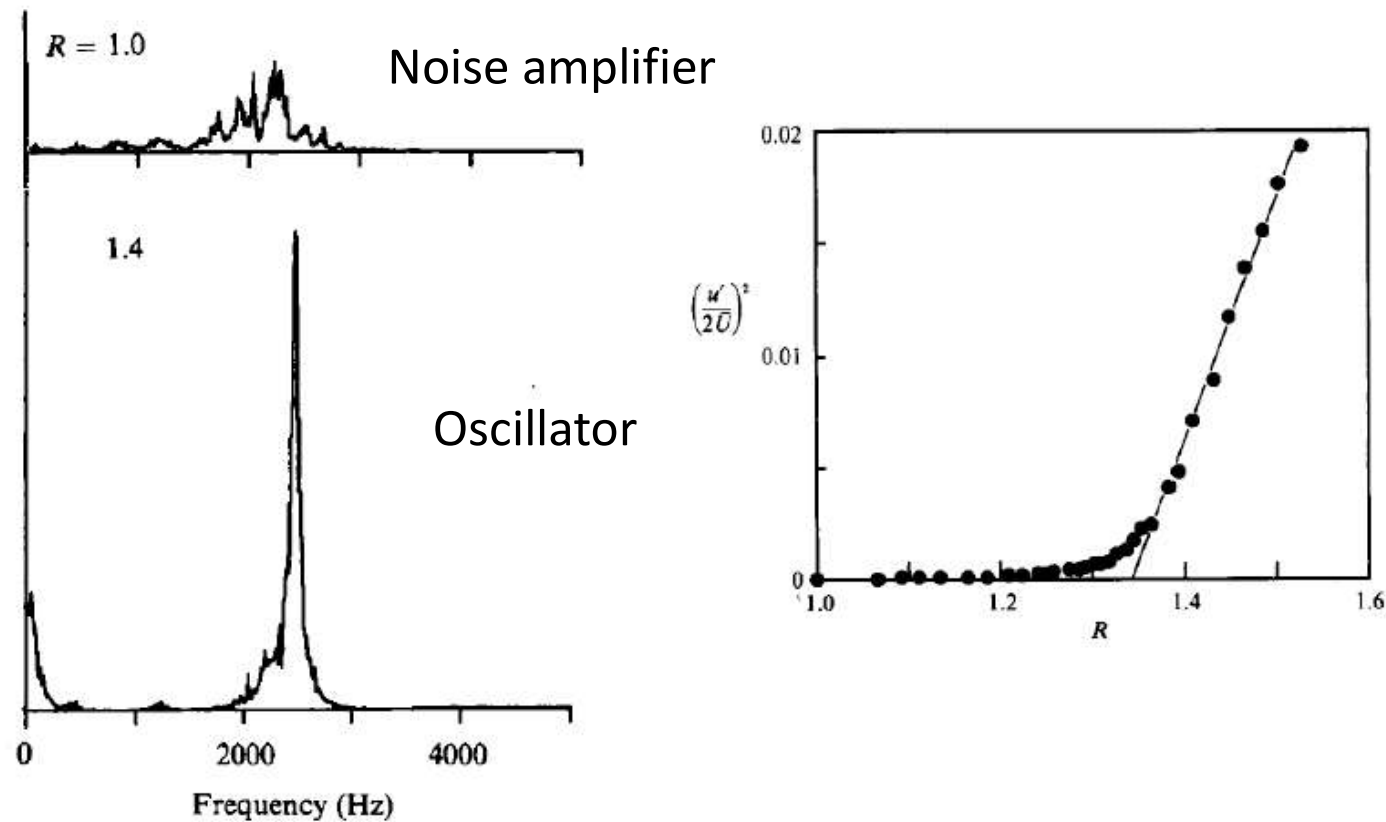
$$R = \frac{U_1 - U_2}{U_1 + U_2} = \frac{\Delta U}{2\bar{U}}$$

Huerre & Monkewitz (1985)

Example: Mixing layer

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Experiment: shift from noise amplifier to oscillator with R



Strykowski & Niccum (1991)

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