

Lecture 9

Spatio-temporal evolution of instabilities I

AE209 Hydrodynamic stability

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1. Ginzburg-Landau equation
2. Absolute and convective instabilities in parallel flow
3. Criterion for absolute instability

1. **Ginzburg-Landau equation**
2. Absolute and convective instabilities in parallel flow
3. Criterion for absolute instability

Complex linear Ginzburg-Landau equation

$$\frac{\partial u}{\partial t} + U \frac{\partial u}{\partial x} - \mu u - (1 - ic_d) \frac{\partial^2 u}{\partial x^2} = 0$$

with boundary condition

$$u(x = \pm\infty) = 0$$

and initial condition

$$u(x, t = 0) = u_0(x)$$

Complex linear Ginzburg-Landau equation

$$\frac{\partial u}{\partial t} + U \frac{\partial u}{\partial x} - \mu u - (1 - ic_d) \frac{\partial^2 u}{\partial x^2} = 0$$

Normal mode solution

$$u = Ae^{i(kx - \omega t)}$$

Dispersion relation

$$\omega - Uk + c_d k^2 - i(\mu - k^2) = 0$$

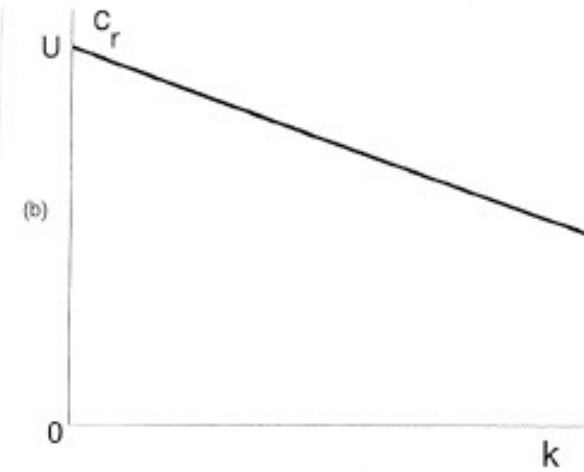
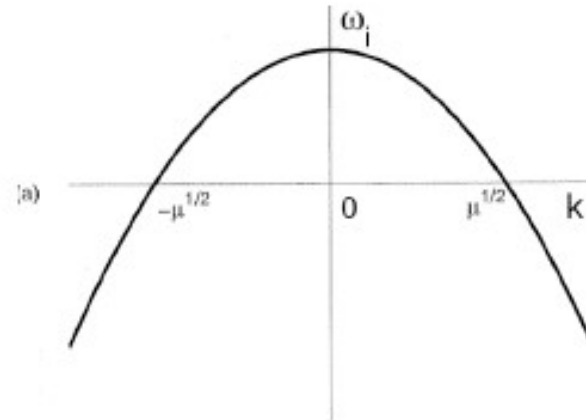
Temporal stability

$$\omega(k) = Uk - c_d k^2 + i(\mu - k^2)$$

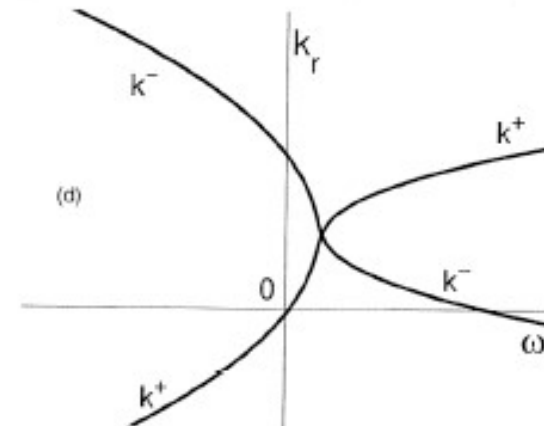
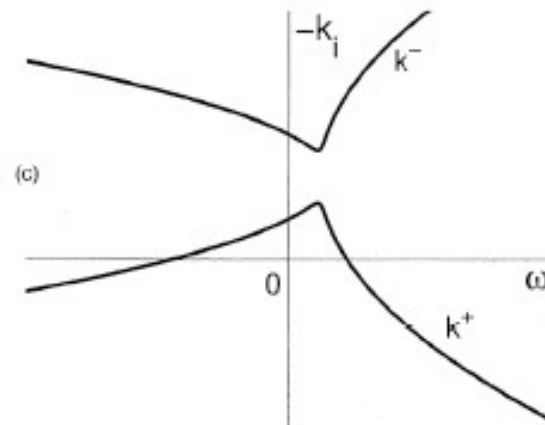
Spatial stability

$$k^\pm(\omega) = \frac{U}{2(c_d + i)} \pm \left(\frac{-1}{c_d + i} \right)^{1/2} \left[\omega - \frac{c_d U^2}{4(1 + c_d^2)} - i \left\{ \mu - \frac{U^2}{4(1 + c_d^2)} \right\} \right]$$

Temporal



Spatial

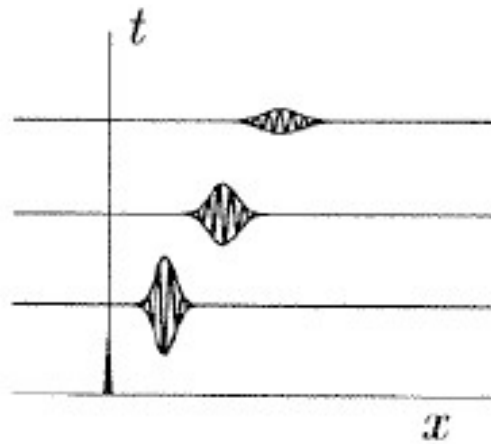


1. Ginzburg-Landau equation
2. **Absolute and convective instabilities in parallel flow**
3. Criterion for absolute instability

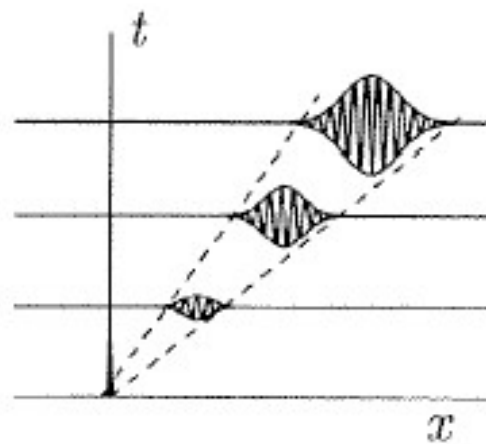
Evolution of localised disturbances

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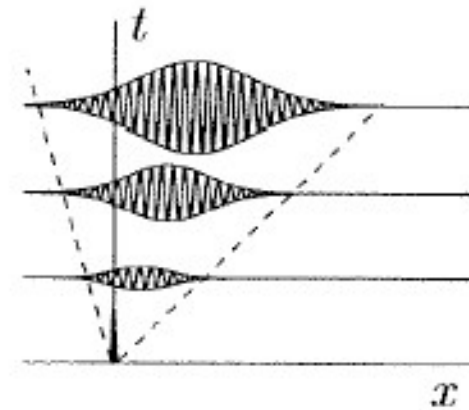
Spatio-temporal evolution of a wave packet (i.e. impulse) in parallel flow



**Linearly
Stable**



Linearly unstable



Impulse response of Ginzburg-Landau equation

$$\left[\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} - \mu - (1 - ic_d) \frac{\partial^2}{\partial x^2} \right] G(x, t) = \delta(x) \delta(t)$$

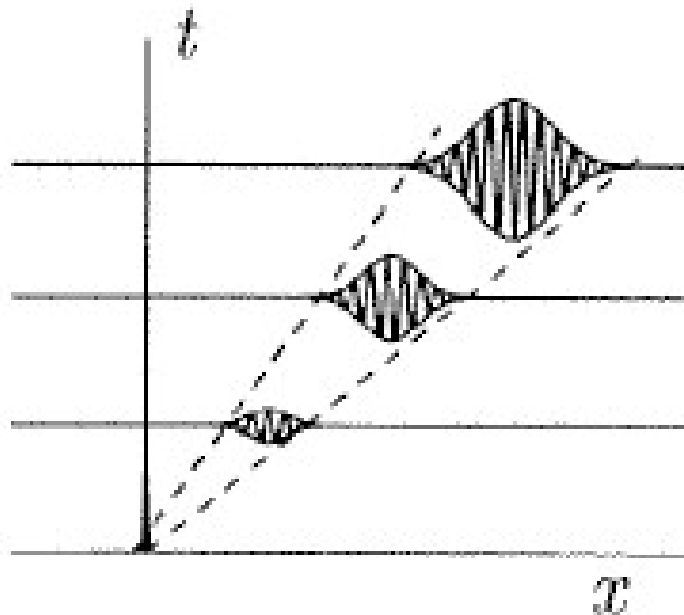
Linearly stable:

$$\lim_{t \rightarrow \infty} G(x, t) = 0 \quad \text{for all rays } x/t = \textit{const}$$

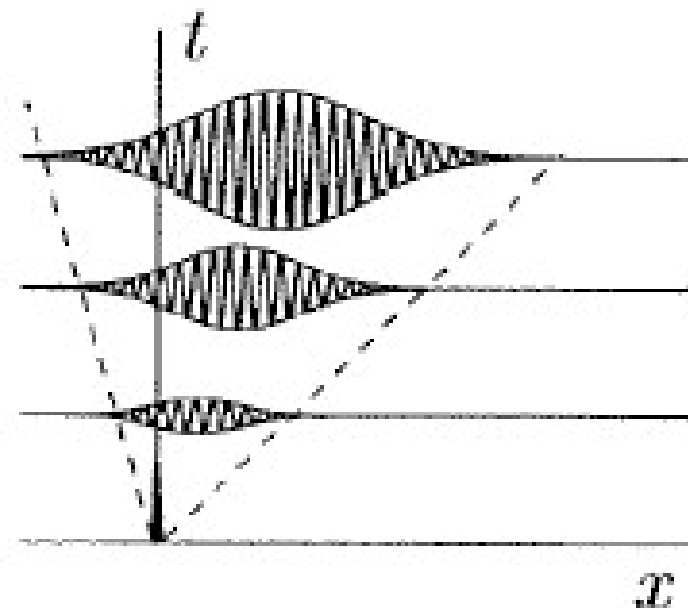
Linearly unstable:

$$\lim_{t \rightarrow \infty} G(x, t) = \infty \quad \text{for at least one ray } x/t = \textit{const}$$

Spatio-temporal evolution of **unstable wavepackets** in parallel flow



Convectively unstable



Absolutely unstable

Impulse response of Ginzburg-Landau equation

$$\left[\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} - \mu - (1 - ic_d) \frac{\partial^2}{\partial x^2} \right] G(x, t) = \delta(x) \delta(t)$$

Convectively unstable

$$\lim_{t \rightarrow \infty} G(x, t) = 0 \quad \text{along the ray of } x/t = 0$$

Absolutely unstable

$$\lim_{t \rightarrow \infty} G(x, t) = \infty \quad \text{along the ray of } x/t = 0$$

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Impulse response of parallel flow

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Impulse response of Ginzburg-Landau equation

$$\left[\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} - \mu - (1 - ic_d) \frac{\partial^2}{\partial x^2} \right] G(x, t) = \delta(x) \delta(t)$$

Solution)

Step 1) Perform Fourier transform in x and Laplace transform in t : i.e.

$$\tilde{G}(k, \omega) == \int_0^\infty \int_{-\infty}^\infty G(x, t) e^{-i(kx - \omega t)} dx dt$$

Step 2) Construct solution in the wavenumber space

Step 3) Invert the Fourier-Laplace transform

$$G(x, t) = \frac{1}{(2\pi)^2} \int_{-\infty}^\infty \int_{-\infty}^\infty \tilde{G}(k, \omega) e^{i(kx - \omega t)} dk d\omega$$

Step 1) Perform Fourier transform in x and Laplace transform in t

i) Fourier transform in x

$$\hat{G}(k, t) = \int_{-\infty}^{\infty} G(x, t) e^{-ikx} dx$$

Examples:

$$\int_{-\infty}^{\infty} \frac{\partial G(x, t)}{\partial t} e^{-ikx} dx = \frac{\partial \hat{G}(k, t)}{\partial t}$$

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{\partial G(x, t)}{\partial x} e^{-ikx} dx &= \left[G(x, t) e^{-ikx} \right]_{-\infty}^{\infty} - \int_{-\infty}^{\infty} G(x, t) \frac{\partial e^{-ikx}}{\partial x} dx \\ &= ik \hat{G}(k, t) \end{aligned}$$

$$\int_{-\infty}^{\infty} \delta(x) \delta(t) e^{-ikx} dx = \delta(t)$$

Step 1) Perform Fourier transform in x and Laplace transform in t

ii) Laplace transform in t

$$\tilde{G}(k, \omega) = \int_0^{\infty} \hat{G}(k, t) e^{i\omega t} dt$$

Examples:

$$\begin{aligned} \int_0^{\infty} \frac{\partial \hat{G}(k, t)}{\partial t} e^{i\omega t} dt &= [\hat{G}(k, t) e^{i\omega t}]_0^{\infty} - \int_0^{\infty} \hat{G}(k, t) \frac{\partial e^{i\omega t}}{\partial t} dt \\ &= -i\omega \tilde{G}(k, \omega) \end{aligned}$$

$$\int_0^{\infty} \delta(t) e^{i\omega t} dt = 1$$

Step 2) Construct solution in the wavenumber space

$$D(k, \omega) \tilde{G}(k, \omega) = 1$$

where $D(k, \omega) = -i\omega + iUk - ic_d k^2 - (\mu - k^2)$

Step 3) Construct solution in the wavenumber space

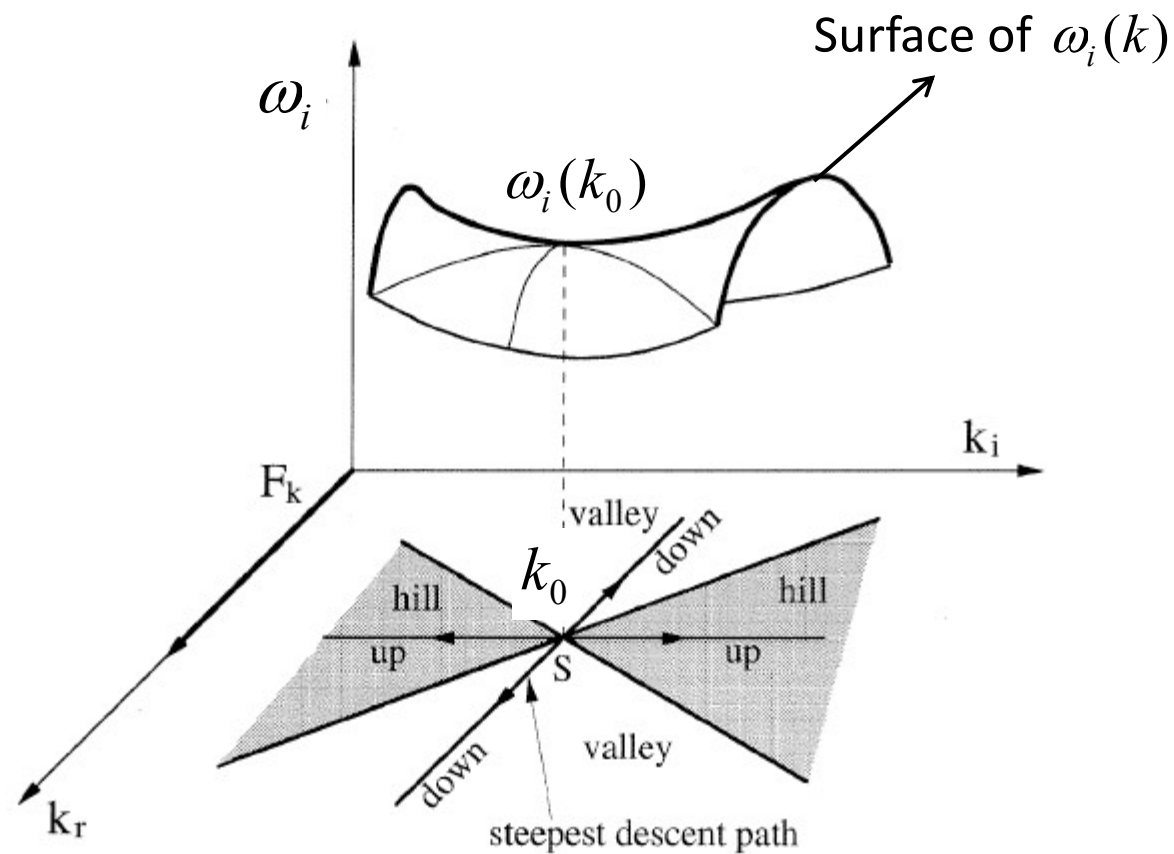
$$G(x, t) = \frac{1}{(2\pi)^2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{1}{D(k, \omega)} e^{i(kx - \omega t)} dk d\omega$$
$$\sim \frac{e^{i(k_0 x - \omega_0 t)}}{\frac{\partial D}{\partial \omega}(k_0, \omega_0) \left[\frac{\partial^2 \omega}{\partial k^2}(k_0) t \right]^{1/2}} \quad \text{as } t \rightarrow \infty$$

where

$$\frac{\partial \omega(k)}{\partial k} = 0 \quad \text{at } k = k_0 \quad \text{and} \quad \omega_0 = \omega(k_0)$$

Remark

The point $\omega_0 = \omega(k_0)$ forms a saddle point over complex k plane



Criterion of absolute instability

$$G(x, t) \sim e^{i(k_0 x - \omega_0 t)}$$

From the definition of absolute instability, the growth rate along $x/t = 0$ is given by $\omega_{0,i}$ called absolute growth rate. In general,

$$\omega_i(k_{\max}) < 0 \quad : \text{Linearly stable}$$

$$\omega_i(k_{\max}) > 0 \quad \text{and} \quad \omega_{0,i} < 0 \quad : \text{Convectively unstable}$$

$$\omega_i(k_{\max}) > 0 \quad \text{and} \quad \omega_{0,i} > 0 \quad : \text{Absolutely unstable}$$

- 1. Ginzburg-Landau equation: a toy model of NS equation**
- 2. Absolute and convective instabilities in parallel flow**
- 3. Criterion for absolute instability**