

Pulse-shape Effects on the Autler-townes Doublet in Strong-field Ionization of Atomic Hydrogen

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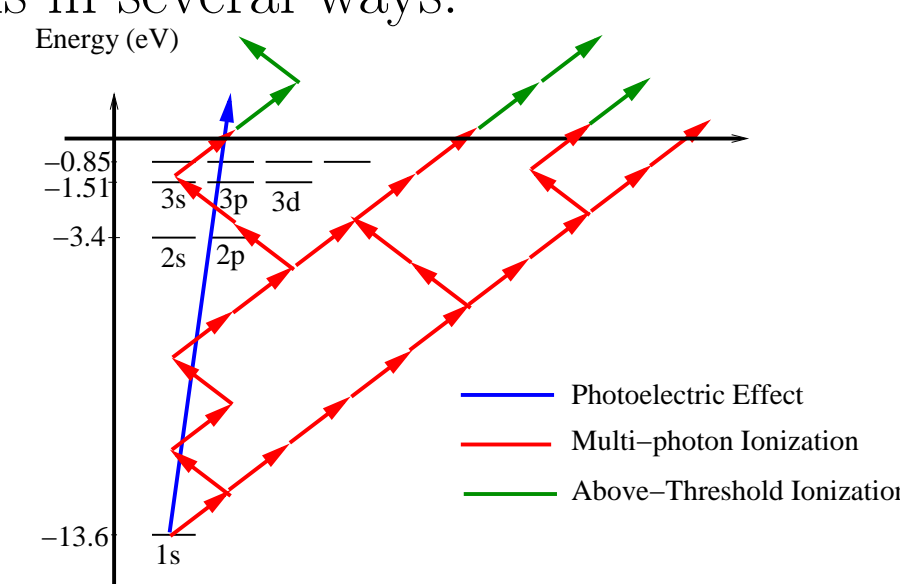
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Abstract

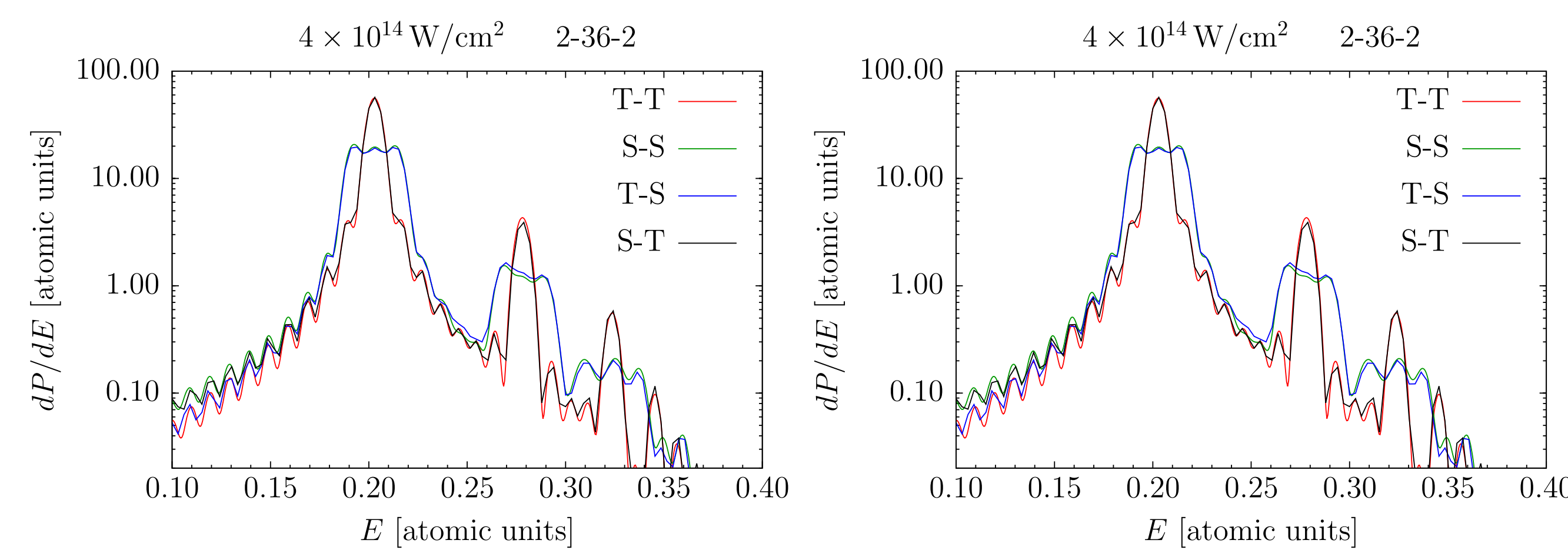
We have applied a newly developed parallelized computer code to treat the ionization of atomic hydrogen by a strong laser pulse. In particular, we studied the effect of the pulse shape, as well as the peak intensity and the central wavelength, on the theoretical results for the so-called Autler-Townes doublet. While the splitting is well known for the quasi-static case, the *dynamic (time-dependent)* Stark effect studied here is much less understood. The strong dependence on the laser pulse found in this work is not only surprising, but may also be a limiting factor for calibrating absolute laser intensities.

Introduction and Motivation

- Very short and intense laser pulses can be used to study the details of (valence) electron interactions in atoms and molecules.
- Typical laser intensities in this field range from 10^{12} to 10^{15} W/cm².
- 10^{14} W/cm² is a million billion times stronger than the radiation that the Earth receives from the Sun directly above us on a clear day.
- Such intensities can rip electrons away from atoms in several ways:
 - Multi-photon ionization
 - Above-threshold ionization
 - Field (tunnel) ionization



The Stark Effect



- The Stark effect splits up the energetically degenerate (for fixed n) energy levels in atomic hydrogen by the interaction with a strong external electric field.
- The energy splitting is proportional to the electric field strength.
- For linearly polarized light, we can “see” only the two $m = 0$ levels.
- These levels form the “Autler-Townes” doublet in the energy spectrum of the ejected electron.
- We investigate this doublet in two-photon ionization, where the central frequency of the laser is tuned in such a way that it either hits (0.375 a.u. = 10.2 eV) or just misses (0.350 a.u.) the $1s \rightarrow 2s, 2p$ resonance transition as a stepping stone.
- Also, we vary the splitting by ramping on/off the pulse.

$$\begin{array}{c}
 \text{4 degenerate states} \\
 \begin{bmatrix} |200\rangle \\ |210\rangle \\ |211\rangle \\ |21-1\rangle \end{bmatrix}
 \end{array}
 \rightarrow
 \begin{array}{c}
 m=0 \rightarrow \frac{1}{\sqrt{2}}[|200\rangle - |210\rangle] \\
 m=\pm 1 \rightarrow \begin{bmatrix} |211\rangle, |21-1\rangle \end{bmatrix} \text{ 2 degenerate states} \\
 m=0 \rightarrow \frac{1}{\sqrt{2}}[|200\rangle + |210\rangle]
 \end{array}$$

The Problem

FIG. 1: LEFT: Time to completion for analysis of a uniform filesystem tree with uniform branching factor two plotted with *depth*. RIGHT: Time to completion for analysis of a uniform filesystem with branching factor two plotted with *entry counts*.

- **Logarithmic time complexity** when sufficiently scaled. Time to completion depends on the maximum depth of the file system (i.e. most nested directory or file).
- Degenerates to linear time complexity in the worst case (same as Tree-walk).

Our Test Case

Results and Discussion

References