

In his theory of blackbody radiation, energy density at frequency  $\nu$  inside cavity with temperature  $T$  is given by

$$\mathcal{U}(\nu, T) = \rho u$$

where  $\rho$  is the number of radiation modes per unit volume and  $u$  is the average energy at said frequency. Both quantity respectively are given by

$$\rho = \frac{8\pi\nu^2}{c^3}, \quad u = \frac{h\nu}{\exp(h\nu/k_B T) - 1}$$

Hence

$$\mathcal{U}(\nu, T) = \frac{\nu^3}{c^3} \frac{8\pi h}{\exp(h\nu/k_B T) - 1}$$

or, we can write it in terms of  $\lambda$

$$\mathcal{U}(\lambda, T) = \frac{c}{\lambda^5} \frac{8\pi h}{\exp(hc/\lambda k_B T) - 1}$$

**Derivation.** In his work, Planck assumes that the walls of the cavity act as oscillator with energy of integer multiple  $\epsilon$ . This has some similarity with Boltzmann discrete model. We denote  $n_k$  as the number particle with  $k\epsilon$  energy, with maximum energy of  $P\epsilon$ . Therefore, we have the following constraints.

$$\sum_{k=0}^P n_k = N, \quad \sum_{k=0}^P k n_k = P$$

Planck he defined entropy as

$$S_P = k_B \ln(W) + C$$

where

$$W = \frac{\mathcal{R}}{\mathcal{J}}$$

is the probability of the  $N$  atoms have  $p\epsilon$  energy. In other hand,  $\mathcal{R}$  denotes the number of said configuration and  $\mathcal{J}$  denotes the total configuration. It then may be written as

$$S_P = k_B \ln \mathcal{R}$$

where  $\mathcal{R}$  is given by

$$\mathcal{R} = \frac{(N + L - 1)!}{L!(N - 1)!}$$

The logarithm of  $\mathcal{R}$  can be evaluated as

$$\begin{aligned} \ln \mathcal{R} = (N + L - 1)[\ln(N + L - 1) - 1] - L[\ln(L) - 1] \\ - (N - 1)[\ln(N - 1) - 1] \end{aligned}$$

and then

$$\ln \mathcal{R} = (N + L - 1) \ln(N + L - 1) - L \ln L - (N - 1) \ln(N - 1)$$

and then

$$\begin{aligned} \ln \mathcal{R} = (N + L) \ln(N + L - 1) - \ln(N + L - 1) - L \ln L \\ - N \ln(N - L) + \ln(N - 1) \end{aligned}$$

and then

$$\begin{aligned}\ln \mathcal{R} = (N + L) \left[ \ln \left( \frac{N + L - 1}{N + L} \right) + \ln (N + L) \right] - L \ln L \\ - N \left[ \ln \left( \frac{N - 1}{N} \right) + \ln(N) \right] - \ln \left( \frac{N + L - 1}{N + L} \right)\end{aligned}$$

then finally we rewrite it as

$$\begin{aligned}\ln \mathcal{R} = (N + L) \ln(N + L) - L \ln L - N \ln N \\ - N \ln \left( \frac{N - 1}{N} \right) + (N + L - 1) \ln \left( \frac{N + L - 1}{N + L} \right)\end{aligned}$$

For large  $N, L$ ; those last two terms will approach zero. Hence,

$$\ln \mathcal{R} = (N + L) \ln(N + L) - L \ln L - N \ln N$$

This equation give the same result for  $\ln(\mathcal{P}_{\max})$  in the case of Boltzmann discrete model. Therefore, we write

$$S_P = N k_B \left[ \left( 1 + \frac{u}{\epsilon} \right) \ln \left( 1 + \frac{u}{\epsilon} \right) - \frac{u}{\epsilon} \ln \left( \frac{u}{\epsilon} \right) \right]$$

Based on empirical data, Planck concludes that entropy is a function of energy and frequency

$$S_P = f \left( \frac{u}{\nu} \right)$$

On comparing those two equation, it can be seen that energy  $\epsilon$  is proportional to frequency  $\nu$ . To show this relationship, we then write  $\epsilon = h\nu$ . Solving the equation of entropy  $S_P$  above for average energy  $u$  we have derived the following result.

$$u = \frac{\epsilon}{\exp(\epsilon/k_B T) - 1}$$

Substituting  $\epsilon = h\nu$  from our empirical observation, we obtain

$$u = \frac{h\nu}{\exp(h\nu/k_B T) - 1} \quad \blacksquare$$