

Calculate the CMB power spectrum: Cosmology II

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ABSTRACT

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Nomenclature

2 Constants of nature

- m_e - Mass of electron.
- $m_e = 9.10938356 \cdot 10^{-31}$ kg.
- m_H - Mass of hydrogen atom.
- $m_H = 1.6735575 \cdot 10^{-27}$ kg.
- G - Gravitational constant.
- $G = 6.67430 \cdot 10^{-11}$ m³ kg⁻¹ s⁻².
- k_B - Boltzmann constant.
- $k_B = 1.38064852 \cdot 10^{-23}$ m² kg s⁻² K⁻¹.
- \hbar - Reduced Planck constant.
- $\hbar = 1.054571817 \cdot 10^{-34}$ J s⁻¹.
- c - Speed of light in vacuum.
- $c = 2.99792458 \cdot 10^8$ m s⁻¹.
- σ_T - Thomson cross section.
- $\sigma_T = 6.6524587158 \cdot 10^{-29}$ m².
- α - Fine structure constant.
- $\alpha = \frac{m_e c}{\hbar} \sqrt{\frac{3\sigma_T}{8\pi}}$

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Cosmological parameters

- $G_{\mu\nu}$ - Einstein tensor.
- $T_{\mu\nu}$ - Stress-energy tensor.
- H - Hubble parameter.
- \mathcal{H} - Conformal Hubble parameter.
- T_{CMB0} - Temperature of CMB today.
- a - Scale factor.
- x - Logarithm of scale factor.
- t - Cosmic time.
- z - Redshift.
- η - Conformal time.
- χ - Co-moving distance.
- p - Pressure.
- ρ - Density.
- r - Radial distance.
- d_A - Angular diameter distance.
- d_L - Luminosity distance.
- n_e - Electron density.
- n_b - Baryon density.
- X_e - Free electron fraction.
- τ - Optical depth.
- \tilde{g} - Visibility function.
- s - Sound horizon.
- r_s - Sound horizon at decoupling.
- c_s - Wave propagation speed.

Density parameters

Density parameter $\Omega_X = \rho_X/\rho_c$ where ρ_X is the density and $\rho_c = 8\pi G/3H^2$ the critical density. X can take the following values:

- b - Baryons.
- CDM - Cold dark matter.
- γ - Electromagnetic radiation.
- ν - Neutrinos.
- k - Spatial curvature.
- Λ - Cosmological constant.

A 0 in the subscript indicates the present day value.

Fiducial cosmology

The fiducial cosmology used throughout this project is based on the observational data obtained by [Aghanim et al. \(2020\)](#):

$$\begin{aligned}
 h &= 0.67, \\
 T_{\text{CMB}0} &= 2.7255 \text{ K}, \\
 N_{\text{eff}} &= 3.046, \\
 \Omega_{b0} &= 0.05, \\
 \Omega_{\text{CDM}0} &= 0.267, \\
 \Omega_{k0} &= 0, \\
 \Omega_{\nu0} &= N_{\text{eff}} \cdot \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} \Omega_{\gamma0}, \\
 \Omega_{\Lambda0} &= 1 - (\Omega_{k0} + \Omega_{b0} + \Omega_{\text{CDM}0} + \Omega_{\gamma0} + \Omega_{\nu0}), \\
 \Omega_{M0} &= \Omega_{b0} + \Omega_{\text{CDM}0}, \\
 \Omega_{\text{rad}} &= \Omega_{\gamma0} + \Omega_{\nu0}, \\
 n_s &= 0.965, \\
 A_s &= 2.1 \cdot 10^{-9}.
 \end{aligned}$$

1. Introduction

Some citation [Dodelson & Schmidt \(2020\)](#) and [Weinberg \(2008\)](#)

Also write about the following:

- Cosmological principle
- Einstein field equation
- Homogeneity and isotropy
- FLRW metric

In order to explain the connection between spacetime itself and the energy distribution within it we must solve the Einstein equation:

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}, \quad (1)$$

where $G_{\mu\nu}$ is the Einstein tensor describing the geometry of spacetime, G is the gravitational constant and $T_{\mu\nu}$ is the energy and momentum tensor.

TODO: Obviously this introduction will change and amended as more milestones are completed.

2. Milestone I - Background Cosmology

The aim of this part is to set up and solve for a uniform background of the Universe. **TODO: finish this**

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2.1. Theory

2.1.1. Fundamentals

If we assume the universe to be homogeneous and isotropic, the line elements ds is given by the FLRW-metric, here in polar coordinates ([Weinberg 2008](#), eq. 1.1.11):

$$ds^2 = -dt^2 + e^{2x} \left[\frac{dr^2}{1 - kr^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right], \quad (2)$$

where we have introduced $x = \ln(a)$ which will be our primary measure of time.

We further model all forms of energy in the universe as perfect fluids, only characterised by their rest frame density ρ and isotropic pressure p , and an equation of state relating the two:

$$\omega = \frac{\rho}{p}. \quad (3)$$

By conservation of energy and momentum we must satisfy $\nabla_\mu T^{\mu\nu} = 0$, which results in the following differential equations for the density of each fluid ρ_i , from [Winther et al. \(2023\)](#):

$$\frac{d\rho_i}{dt} + 3H\rho_i(1 + \omega_i) = 0, \quad (4)$$

where we have introduced the Hubble parameter $H \equiv \dot{a}/a = dx/dt$. The solution to Eq. (4) is of the form:

$$\rho_i \propto e^{-3(1+\omega_i)x}, \quad (5)$$

where $\omega_M = 0$ (matter), $\omega_{\text{rad}} = 1/3$ (radiation), $\omega_\Lambda = -1$ (cosmological constant) and $\omega_k = -1/3$ (curvature).

With these assumptions, the solution to the Einstein equations (Eq. (1)) are the Friedmann equations, the first of which describes the expansion rate of the universe:

$$H^2 = \frac{8\pi G}{3} \sum_i \rho_i - kc^2 e^{-2x} \quad (6)$$

and the second describe how this expansion rate changes over time:

$$\frac{dH}{dt} + H^2 = -\frac{4\pi G}{3} \sum_i \left(\rho + \frac{3p}{c^2} \right). \quad (7)$$

As of now, we are primarily interested in the first Friedmann equation. By introducing the critical density, $\rho_c \equiv 2H^2/(8\pi G)$, we define the density parameters $\Omega_i = \rho_i/\rho_c$. We further define the density of the curvature $\rho_k \equiv -3kc^2 e^{-2x}/(8\pi G)$, which enables us to write Eq. (6) as simply:

$$1 = \sum_i \Omega_i, \quad (8)$$

where the curvature density Ω_k is included in the sum. From Eq. (5) we know the evolution of the densities in time, and if we assume the density values today, Ω_{i0} , are known (or are free parameters), then Eq. (6) may also be written as:

$$H = H_0 \sqrt{\sum_i \Omega_{i0} e^{-3(1+\omega_i)x}}, \quad (9)$$

which is the Hubble equation we will use further.

2.1.2. Measure of time and space

The main measure of time is usually the scale factor a , or its logarithm x . We then have the *cosmic time* t defined as:

$$t = \int_0^a \frac{da}{aH} = \int_{-\infty}^x \frac{dx}{H}. \quad (10)$$

Another temporal measure is the *conformal time* η defined as $cdt = e^x d\eta$ yielding:

$$\eta = \int_0^a \frac{cda}{a^2 H} = \int_{-\infty}^x \frac{cdx}{e^x H} \equiv \int_{-\infty}^x \frac{cdx}{\mathcal{H}}, \quad (11)$$

where $\mathcal{H} = e^x H$ is defined as the *conformal Hubble parameter*. We may also choose to measure time in terms of the *redshift* z , where $1+z = 1/a = e^{-x}$. The comoving distance is defined as follows:

$$\chi = \int_1^a \frac{cda}{a^2 H} = \int_0^x \frac{cdx}{\mathcal{H}} = \eta_0 - \eta \quad (12)$$

The radial distance is given in terms of the comoving distance and the curvature density today Ω_{k0} as:

$$r = \begin{cases} \chi \cdot \frac{\sin(\sqrt{|\Omega_{k0}|} H_0 \chi / c)}{\sqrt{|\Omega_{k0}|} H_0 \chi / c} & \Omega_{k0} < 0 \\ \chi & \Omega_{k0} = 0 \\ \chi \cdot \frac{\sinh(\sqrt{|\Omega_{k0}|} H_0 \chi / c)}{\sqrt{|\Omega_{k0}|} H_0 \chi / c} & \Omega_{k0} > 0 \end{cases} \quad (13)$$

It is then straightforward to define the angular diameter distance:

$$d_A = e^x r, \quad (14)$$

and the luminosity distance:

$$d_L = e^{-x} r, \quad (15)$$

both of which are derived in Appendix A. The temporal quantities η and t have the following evolutions with x :

$$\frac{d\eta}{dx} = \frac{c}{\mathcal{H}}. \quad (16)$$

$$\frac{dt}{dx} = \frac{1}{H}. \quad (17)$$

Both differential equations are easy to solve numerically. Their derivation may also be found in Appendix A

2.1.3. Λ CDM-model

In the Λ CDM model, the universe consists of matter in terms of baryonic matter (b) and cold dark matter (CDM), radiation in terms of photons (γ) and neutrinos (ν) and dark energy in terms of a cosmological constant (Λ). In addition, we must allow for some curvature (k). As a result, the parameters of the model will be the present values of the Hubble rate, H_0 , the baryon density Ω_{b0} , the cold dark matter density $\Omega_{\text{CDM}0}$, photon density $\Omega_{\gamma0}$, neutrino density $\Omega_{\nu0}$, dark energy density $\Omega_{\Lambda0}$, and the curvature density Ω_{k0} . The present temperature of the cosmic microwave

background radiation $T_{\text{CMB}0}$ fixes the radiation density today through:

$$\Omega_{\gamma0} = \frac{16\pi^3 G}{90} \cdot \frac{(k_b T_{\text{CMB}0})^4}{\hbar^3 c^5 H_0^2},$$

$$\Omega_{\nu0} = N_{\text{eff}} \cdot \frac{7}{8} \cdot \left(\frac{4}{3}\right)^{4/3} \cdot \Omega_{\gamma0}. \quad (18)$$

The total radiation density is $\Omega_{\text{rad}} = \Omega_{\gamma} + \Omega_{\nu}$ and the total matter density is $\Omega_{\text{M}} = \Omega_b + \Omega_{\text{CDM}}$.

The Hubble equation from Eq. (9) may be redefined in terms of the conformal Hubble parameter \mathcal{H} as:

$$\mathcal{H} = H_0 \sqrt{U}$$

$$U \equiv \sum_i \Omega_{i0} e^{-\alpha_i x}, \quad (19)$$

where we have defined $\alpha_i \equiv (1+3\omega_i)$ and $i \in \{\text{M}, \text{rad}, \Lambda, k\}$. Since we know the values of the various ω_i it follows that:

$$\begin{aligned} \alpha_{\text{M}} &= 1 \\ \alpha_{\text{rad}} &= 2 \\ \alpha_k &= 0 \\ \alpha_{\Lambda} &= -2 \end{aligned} \quad (20)$$

2.1.4. Equalities and present day values

Given the evolution of the density parameters with time, where the proportionality constant is the present day density, we introduce the *radiation-matter equality*, i.e. the time radiation and matter densities were equal: $\rho_{\text{rad}} = \rho_{\text{M}}$. According to Eq. (5) this can be expressed as:

$$\rho_{\text{rad}0} e^{-4x} = \rho_{\text{M}0} e^{-3x}$$

$$e^x = \frac{\rho_{\text{rad}0}}{\rho_{\text{M}0}} \implies x_{\text{rM}} = \ln \left(\frac{\Omega_{\text{rad}0}}{\Omega_{\text{M}0}} \right), \quad (21)$$

where x_{rM} now denotes the time of radiation-matter equality.

Similarly, the *matter-dark energy equality*, where $\rho_{\text{M}} = \rho_{\Lambda}$ can be found to be:

$$\rho_{\Lambda} = \rho_{\text{M}0} e^{-3x}$$

$$\implies x_{\text{M}\Lambda} = \frac{1}{3} \ln \left(\frac{\Omega_{\text{M}0}}{\Omega_{\Lambda}} \right) \quad (22)$$

The time of matter-dark energy equality coincides with when the universe starts to accelerate, since this acceleration is driven by the dark energy, represented by the cosmological constant. From this time onwards, dark energy dominates the universe, and thus accelerating the expansion.

The age of the universe today, and the conformal time today can both be found by evaluating the solutions to the differential equations of t and η at the present time (where $x = 0$). This is done numerically.

2.1.5. Analytical solutions and sanity checks

There are several ways we may check that both our workings and numerical implementations are indeed correct. The simplest way is to always ensure that the sum of all density parameters add up to 1, for all times: $\sum_i \Omega_i = 1$.

If we only consider the most dominant density parameter, that is $\Omega_i \gg \sum_{j \neq i} \Omega_j$, we end up with the following analytical expressions for different temporal regimes:

$$\frac{1}{\mathcal{H}} \frac{d\mathcal{H}}{dx} \approx -\frac{\alpha_i}{2} = \begin{cases} -1 & \alpha_{\text{rad}} = 2 \\ -\frac{1}{2} & \alpha_{\text{M}} = 1 \\ 1 & \alpha_{\Lambda} = -2 \end{cases} \quad (23)$$

$$\frac{1}{\mathcal{H}} \frac{d^2\mathcal{H}}{dx^2} \approx \frac{\alpha_i^2}{4} = \begin{cases} 1 & \alpha_{\text{rad}} = 2 \\ \frac{1}{4} & \alpha_{\text{M}} = 1 \\ 1 & \alpha_{\Lambda} = -2 \end{cases} \quad (24)$$

$$\frac{\eta\mathcal{H}}{c} \approx \begin{cases} 1 & \alpha_{\text{rad}} = 2 \\ 2 & \alpha_{\text{M}} = 1 \\ \infty & \alpha_{\Lambda} = -2 \end{cases} \quad (25)$$

These equations will be useful when making sure that the implementations are correct. For a thorough derivation, see Appendix B.

2.2. Methods

2.2.1. Initial equation

We have to consider the time evolution of the density parameters, given some present value, as function of our chosen time parameter, here x . The density evolution is implemented as:

$$\Omega_n = e^{-\alpha_n x} \Omega_{n0} \mathcal{H}_{\text{rat}}^2 \quad (26)$$

where we have defined the ratio $\mathcal{H}_{\text{rat}} \equiv H_0/\mathcal{H}$, and the new index n are all the densities: $n \in \{b, \text{CDM}, \gamma, \nu, \Lambda, k\}$.

We also implement functions to solve for the luminosity distance (Eq. (15)), angular distance (Eq. (14)), and the conformal distance (Eq. (12)).

2.2.2. ODEs and Splines

The differential equations for η (Eq. (16)) and t (Eq. (17)) are solved numerically as ordinary differential equations with the Runge-Kutta 4 as advancement method. The equations are solved for $x \in (-20, 5)$. As initial condition we would like $\eta(-\infty)$ which is obviously not possible to calculate, so we pick some very early time and use the analytical approximation in the radiation dominated era (Eq. (25)), which yield:

$$\eta(x_0) = \frac{c}{\mathcal{H}(x_0)}. \quad (27)$$

Likewise for t , the initial condition is:

$$t(x_0) = \frac{1}{2H(x_0)}. \quad (28)$$

We then proceed by making splines of both η and t in order to evaluate accurately for any $x \in (-20, 5)$.

2.2.3. Supernova data

We make predictions of the luminosity distance at different redshifts z , according to the discussion in Section 2.1.2. These predictions are compared with real supernova observations, acquired by Betoule et al. (2014). In order to constrain the possible values of Ω_{M} and Ω_{Λ} we find the χ^2 -value between the luminosity distance of the supernovas and our predictions. The Ω -s are sampled with Markov-Chain Monte Carlo sampling using the Metropolis-Hastings algorithm. The χ^2 -values are given by:

$$\chi^2(h, \Omega_{m0}, \Omega_{k0}) = \sum_{i=1}^N \frac{(d_L(z, \Omega_{m0}, \Omega_{k0}) - d_L^{\text{obs}}(z_i))^2}{\sigma_i^2}. \quad (29)$$

It is easy to see from Eq. (29) that if the difference between prediction and observation matches the uncertainties in the observed data, then we have $\chi^2 \simeq N$, which is the best fit we can hope for (since we fit to data which has an intrinsic uncertainty already). If $\chi^2 \ll N$, then we *over-fit* the model, while if $\chi^2 \gg N$ then we have an increasingly bad fit. In summary, we may analyse the following quantity

$$\frac{\chi^2}{N} \begin{cases} \ll 1 & \text{over-fitting,} \\ \simeq 1 & \text{good fit,} \\ \gg 1 & \text{bad fit,} \end{cases} \quad (30)$$

in order to determine the goodness of fit.

2.3. Results and discussion

2.3.1. Tests

Fig. 1 show the evolution of the density fractions with time. They sum to one across all times which was required. At early times the radiation density dominates (orange line). The intersection between the orange and green lines mark the radiation-matter equality, after which matter is the dominating density. Likewise, the intersection between the green and purple lines mark the matter-dark energy equality, where dark energy (manifested in the cosmological constant) become the dominating density. Time can thus be divided into three regimes; radiation dominated, matter dominated, and dark energy dominated eras.

As explained in section Section 2.1.5, we have analytical solutions for constructions of η and \mathcal{H} in the different regimes. Fig. 2 is the sanity check for η , showing $\eta\mathcal{H}/c$ converging to finite values in the radiation and matter dominated eras (where $\alpha_{\text{rad}, \text{M}} > 0$), and diverging towards $+\infty$ in the dark energy dominated era ($\alpha_{\Lambda} = -2 < 0$). This is in accordance with the analytical solutions. The different regimes are shown in shaded colour. It is also worth noticing that $(d\eta/dx)\mathcal{H}/c$ is one for all regimes, as expected from equation Eq. (16).

Fig. 3 is the sanity check confirming that our constructions of \mathcal{H} and its derivatives converge to the analytical approximation in the different regimes. The second derivative, as shown in blue, takes the value of one in the radiation regime, one half in the matter regime and one in the dark energy regime. Similarly, the first derivative, as shown in green, take the value negative one in the radiation regime, negative one half in the matter regime and one in the dark energy regime. This is well in accordance with the analytical approximations put forth in section Section 2.1.5.

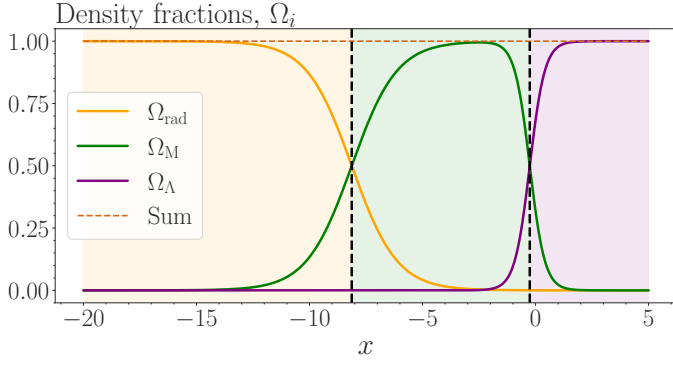


Fig. 1. Density fractions Ω_i as function of x . For low x , radiation dominates, before matter dominates and dark energy has just become the dominant energy density today $x = 0$, and will continue to dominate into the future. The sum of densities sums to one across all times, as required (white dotted line). The black dotted lines are the radiation-matter equality at $x = -8.13$ and the matter-dark energy equality at $x = -0.26$, both stated in Table 1. The domination of each regime is shown as a shaded background with similar colour as its respective graph.

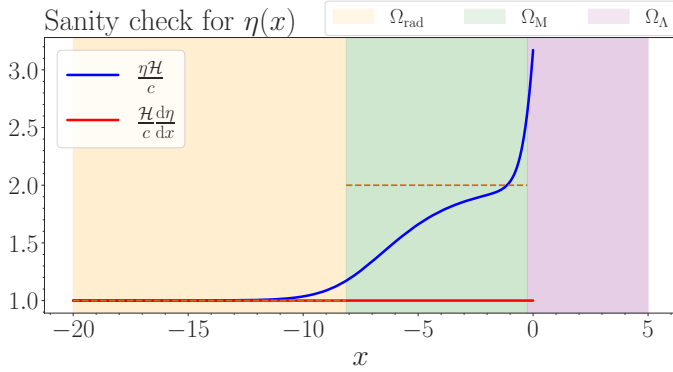


Fig. 2. Sanity check for η . $\eta H/c$, in blue, is 1 in the radiation regime, 2 in the matter regime and diverging toward $+\infty$ in the dark energy regime, as expected from the analytical approximations in each regime. Remembering that this is strictly correct only in the radiation regime explains the mismatch of the white dotted line in the matter regime. $(d\eta/dx)H/c$, in red, is 1 throughout time, as expected from the differential equation for η .

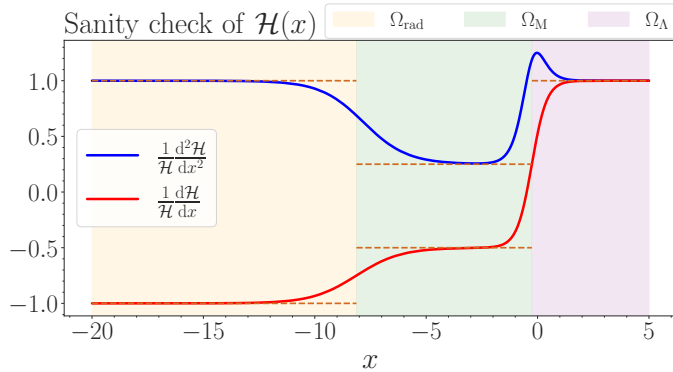


Fig. 3. Sanity check for \mathcal{H} , showing that the second derivative (blue) converge to the analytical expressions shown as white dotted lines in the different regimes. The first derivative (red) converge to its analytical values in the same regimes, which again are shown with a shaded colour.

Quantity	x	z	t
RM-equality	-8.13	3400	$51 \cdot 10^3$ yr
ML-equality	-0.26	0.29	10.378 Gyr
Accel. start	-0.49	0.63	7.752 Gyr
Age of universe	0.00	0.00	13.858 Gyr
Conformal time	0.00	0.00	46.319 Gyr

Table 1. Important quantities in the evolution of the universe.

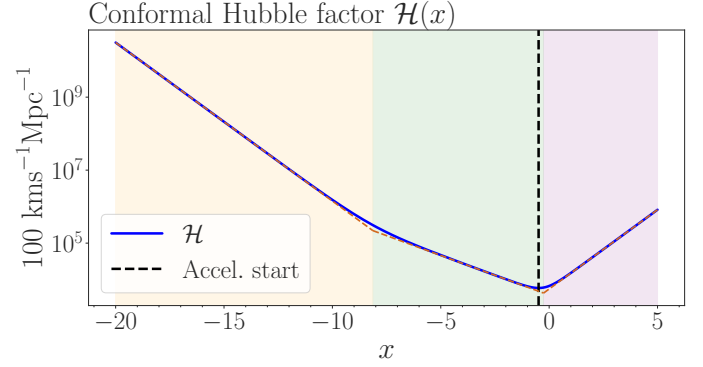


Fig. 4. \mathcal{H} as function of x . It is decreasing in the radiation and matter regimes, and increasing in the dark energy regime, tightly following its analytical approximation in each regime.

These sanity checks are confirmations that the implementation of the model yields the same result as the analytical approximation in the different regimes for various constructions of η and \mathcal{H} and their derivatives.

2.3.2. Analysis

Section 2.1.4 indicate how we can calculate the radiation-matter equality (RM-equality), matter-dark energy equality (ML-equality), when the acceleration of the universe started, the age of the universe and the conformal time today. The result is shown in Table 1. We note that the equalities is in accordance with the sanity checks, and the age of the universe today (in cosmic time) is about 13.9 Gyr.

The conformal Hubble factor, \mathcal{H} , is plotted against time, x , in Fig. 4. It is decreasing in the radiation and matter regimes and increasing in the dark energy regime. Since this is a measure of the expansion of the universe, the acceleration seem to coincide with the matter-dark energy equality.

Fig. 5 show the cosmic time t and the conformal time η/c . We observe that the cosmic time accelerates around matter-dark energy equality, and seems to diverge, whereas the conformal time increases in the matter regime and seem to converge in the dark energy regime.

Fig. 6 shows the supernova data as blue error bars, with the theoretical prediction plotted above it. The quantity plotted is the luminosity distance divided by redshift, d_L/z for better comparison. We notice the accordance between the two, and also note that the x -axis in this plot is the redshift z instead of the logarithm of the scale factor. This means that earlier times are to the right in the plot (high redshift), as opposed to the other plots.

Fig. 8 shows the χ^2 -values found from Eq. (29), to 1σ accuracy. The black dotted line represent a flat universe, where the matter and dark energy are the main constituents

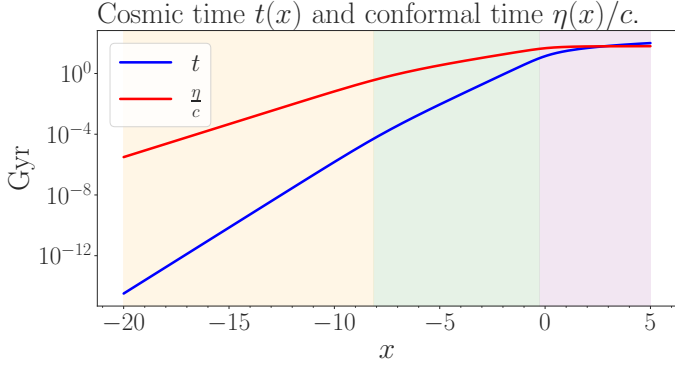


Fig. 5. Cosmic time (in blue) and conformal time (red). **TODO:** find meaning behind this plot.

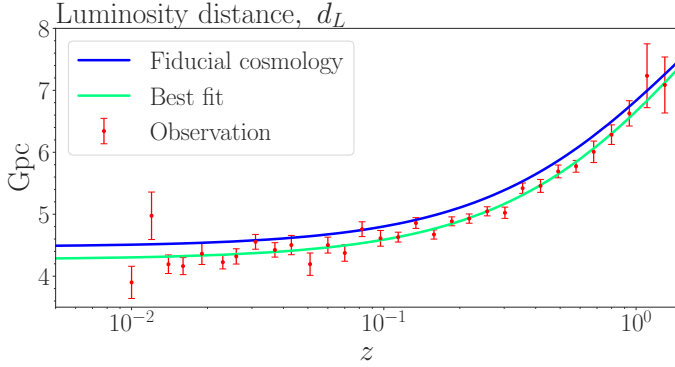


Fig. 6. The luminosity distance predicted using the fiducial cosmology in blue, against observations of actual supernovas in red (or rather the confidence interval of the observations). The green line is found from computing the luminosity distance using a cosmology of the best fit parameters from the supernova fitting; $h = 0.702, \Omega_{b0} = 0.05, \Omega_{CDM0} = 0.209, \Omega_k = 0.067, N_{\text{eff}} = 3.046, T_{\text{CMB}} = 2.7255\text{K}$. Notice the x -axis is now the redshift $z = e^x - 1$.

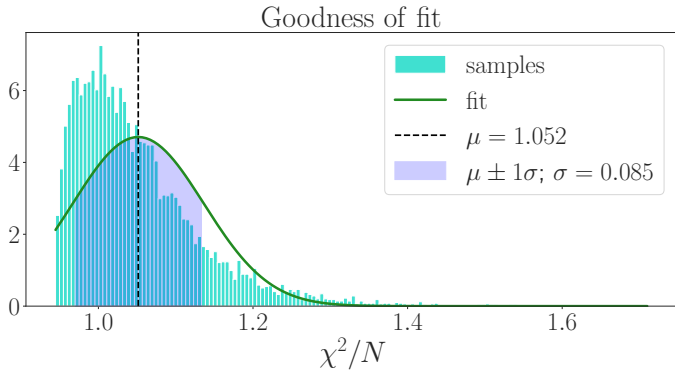


Fig. 7. Goodness of fit

of the universe. The supernova data originate in close temporal proximity to us, we thus assume that the contribution from the radiation density is negligible for making constraints on Ω_M and Ω_Λ today.

The MCMC method allow us to make a posterior probability distribution of the parameters in question. Fig. 10 the distribution of the Hubble factor H_0 constructed from the sampled values of h . The actual samples are shown in green and the corresponding fitted probability distribution in blue.

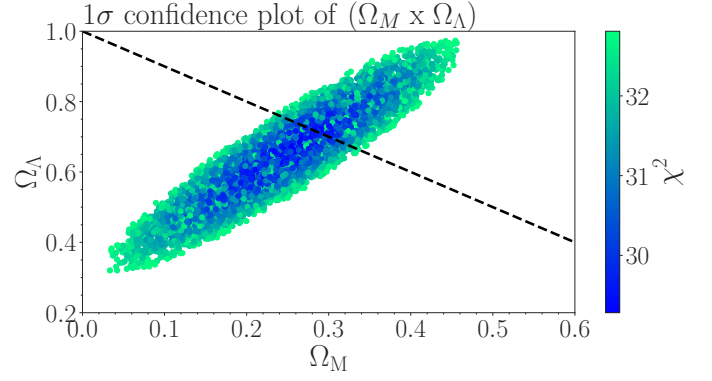


Fig. 8. Scatter plot showing the χ^2 -error of the luminosity distance d_L between the observed values and the prediction, as function of Ω_M and Ω_Λ . The data shown is within 1σ (standard deviation). The black dotted line signifies a flat universe.

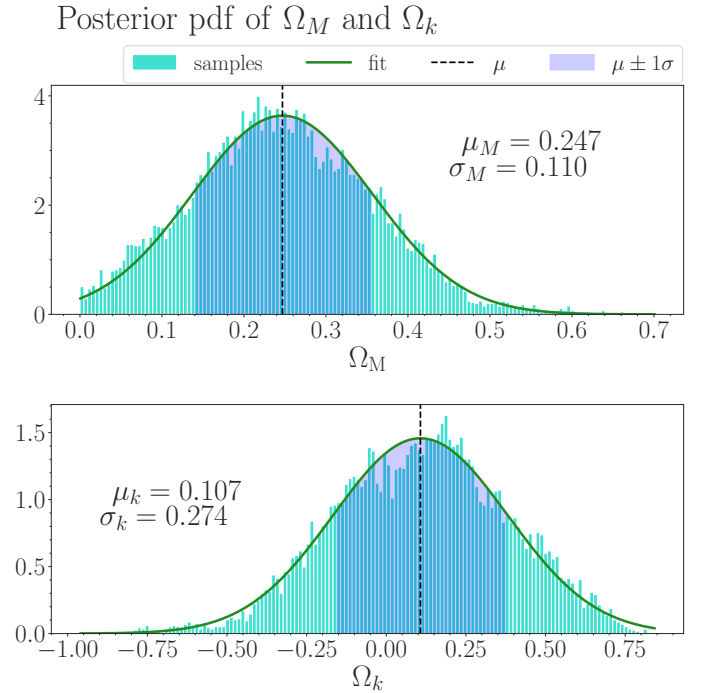


Fig. 9. Posterior probability distributions (pdfs) of Ω_M and Ω_k as result of the MCMC sampling. The samples as shown in turquoise and the constructed pdfs in green. The mean μ is shown as a black dotted line, with the 1σ confidence interval in shaded blue ($\mu \pm 1\sigma$)

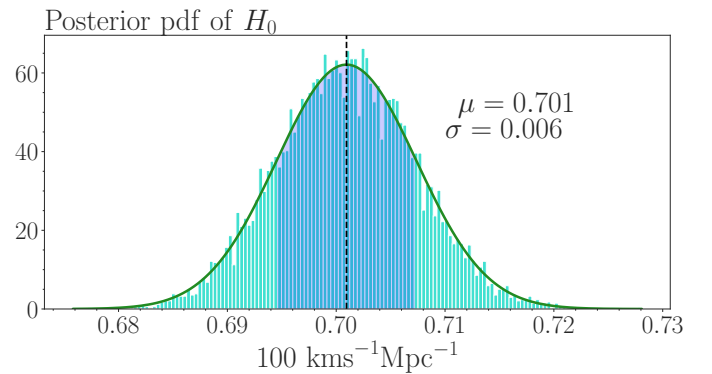


Fig. 10. Posterior probability distribution (pdf) of H_0 as result of the MCMC sampling.

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Appendix A: Useful derivations

A.1. Angular diameter distance

This is related to the physical distance of say, an object, whose extent is small compared to the distance at which we observe is. If the extension of the object is Δs , and we measure an angular size of $\Delta\theta$, then the angular distance to the object is:

$$d_A = \frac{\Delta s}{\Delta\theta} = \frac{ds}{d\theta} = \sqrt{e^{2x} r^2} = e^x r, \quad (\text{A.1})$$

where we inserted for the line element ds as given in equation Eq. (2), and used the fact that $dt/d\theta = dr/d\theta = d\phi/d\theta = 0$ in polar coordinates.

A.2. Luminosity distance

If the intrinsic luminosity, L of an object is known, we can calculate the flux as: $F = L/(4\pi d_L^2)$, where d_L is the luminosity distance. It is a measure of how much the light has dimmed when travelling from the source to the observer. For further analysis we observe that the luminosity of objects moving away from us is changing by a factor a^{-4} due to the energy loss of electromagnetic radiation, and the observed flux is changed by a factor $1/(4\pi d_A^2)$. From this we draw the conclusion that the luminosity distance may be written as:

$$d_L = \sqrt{\frac{L}{4\pi F}} = \sqrt{\frac{d_A^2}{a^4}} = e^{-x} r \quad (\text{A.2})$$

A.3. Differential equations

From the definition of $e^x d\eta = c dt$ we have the following:

$$\begin{aligned} \frac{d\eta}{dt} &= \frac{d\eta}{dx} \frac{dx}{dt} = \frac{d\eta}{dx} H = e^{-x} c \\ \Rightarrow \frac{d\eta}{dx} &= \frac{c}{H}. \end{aligned} \quad (\text{A.3})$$

Likewise, for t we have:

$$\begin{aligned} \frac{d\eta}{dt} &= \frac{d\eta}{dx} \frac{dx}{dt} = \frac{dx}{dt} \frac{c}{H} = e^{-x} c \\ \Rightarrow \frac{dt}{dx} &= \frac{e^x}{H} = \frac{1}{H}. \end{aligned} \quad (\text{A.4})$$

Appendix B: Sanity checks

B.1. For \mathcal{H}

We start with the Hubble equation from Eq. (19) and realize that we may write any derivative of U as

$$\frac{d^n U}{dx^n} = \sum_i (-\alpha_i)^n \Omega_{i0} e^{-\alpha_i x}. \quad (\text{B.1})$$

We further have:

$$\frac{d\mathcal{H}}{dx} = \frac{H_0}{2} U^{-\frac{1}{2}} \frac{dU}{dx}, \quad (\text{B.2})$$

and

$$\begin{aligned} \frac{d^2 \mathcal{H}}{dx^2} &= \frac{d}{dx} \frac{d\mathcal{H}}{dx} \\ &= \frac{H_0}{2} \left[\frac{dU}{dx} \left(\frac{d}{dx} U^{-\frac{1}{2}} \right) + U^{-\frac{1}{2}} \left(\frac{d}{dx} \frac{dU}{dx} \right) \right] \\ &= H_0 \left[\frac{1}{2U^{\frac{1}{2}}} \frac{d^2 U}{dx^2} - \frac{1}{4U^{\frac{3}{2}}} \left(\frac{dU}{dx} \right)^2 \right] \end{aligned} \quad (\text{B.3})$$

Multiplying both equations with $\mathcal{H}^{-1} = 1/(H_0 U^{\frac{1}{2}})$ yield the following:

$$\frac{1}{\mathcal{H}} \frac{d\mathcal{H}}{dx} = \frac{1}{2U} \frac{dU}{dx}, \quad (\text{B.4})$$

and

$$\begin{aligned} \frac{1}{\mathcal{H}} \frac{d^2 \mathcal{H}}{dx^2} &= \frac{1}{2U} \frac{d^2 U}{dx^2} - \frac{1}{4U^2} \left(\frac{dU}{dx} \right)^2 \\ &= \frac{1}{2U} \frac{d^2 U}{dx^2} - \left(\frac{1}{\mathcal{H}} \frac{dU}{dx} \right)^2 \end{aligned} \quad (\text{B.5})$$

We now make the assumption that one of the density parameters dominate $\Omega_i \gg \sum_{j \neq i} \Omega_j$, enabling the following approximation:

$$U \approx \Omega_{i0} e^{-\alpha_i x}$$

$$\frac{d^n U}{dx^n} \approx (-\alpha_i)^n \Omega_{i0} e^{-\alpha_i x}, \quad (\text{B.6})$$

from which we are able to construct:

$$\frac{1}{\mathcal{H}} \frac{d\mathcal{H}}{dx} \approx \frac{-\alpha_i \Omega_{i0} e^{-\alpha_i x}}{2\Omega_{i0} e^{-\alpha_i x}} = -\frac{\alpha_i}{2}, \quad (\text{B.7})$$

and

$$\begin{aligned} \frac{1}{\mathcal{H}} \frac{d^2 \mathcal{H}}{dx^2} &\approx \frac{\alpha_i^2 \Omega_{i0} e^{-\alpha_i x}}{2\Omega_{i0} e^{-\alpha_i x}} - \left(\frac{\alpha_i}{2} \right)^2 \\ &= \frac{\alpha_i^2}{2} - \frac{\alpha_i^2}{4} = \frac{\alpha_i^2}{4} \end{aligned} \quad (\text{B.8})$$

which are quantities which should be constant in different regimes and we can easily check if our implementation of \mathcal{H} is correct, which is exactly what we sought.

B.2. For η

In order to test η we consider the definition, solve the integral and consider the same regimes as above, where one density parameter dominates:

$$\begin{aligned} \eta &= \int_{-\infty}^x \frac{cdx}{\mathcal{H}} = \frac{-2c}{\alpha_i} \int_{x=-\infty}^{x=x} \frac{d\mathcal{H}}{\mathcal{H}^2} \\ &= \frac{2c}{\alpha_i} \left(\frac{1}{\mathcal{H}(x)} - \frac{1}{\mathcal{H}(-\infty)} \right), \end{aligned} \quad (\text{B.9})$$

where we have used that:

$$\begin{aligned} \frac{d\mathcal{H}}{dx} &= -\frac{\alpha_i}{2} \mathcal{H} \\ \Rightarrow dx &= -\frac{2}{\alpha_i \mathcal{H}} d\mathcal{H}. \end{aligned} \quad (\text{B.10})$$

Since we consider regimes where one density parameter dominates, we have that $\mathcal{H}(x) \propto \sqrt{e^{-\alpha_i x}}$, meaning that we have:

$$\left(\frac{1}{\mathcal{H}(x)} - \frac{1}{\mathcal{H}(-\infty)} \right) \approx \begin{cases} \frac{1}{\mathcal{H}} & \alpha_i > 0 \\ -\infty & \alpha_i < 0. \end{cases} \quad (\text{B.11})$$

Combining the above yields:

$$\frac{\eta \mathcal{H}}{c} \approx \begin{cases} \frac{2}{\alpha_i} & \alpha_i > 0 \\ \infty & \alpha_i < 0. \end{cases} \quad (\text{B.12})$$

Notice the positive sign before ∞ . This is due to α_i now being negative.