

A symmetry-based structural SI analysis of the SEI model

Johannes G. Borgqvist

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1 Automated notes generated by the SymPy-based script

ODE for S :

$$\frac{d}{dt}S(t) = -\beta I(t)S(t) + c - \mu_S S(t) \quad (1)$$

ODE for E :

$$\frac{d}{dt}E(t) = \beta(1-v)I(t)S(t) - \delta E(t) - \mu_E E(t) \quad (2)$$

ODE for I :

$$\frac{d}{dt}I(t) = \beta v I(t)S(t) + \delta E(t) - \mu_I I(t) \quad (3)$$

2 Symmetry-based local SI analysis of the SEI model

We consider the following SEI model of epidemiological transmission of tuberculosis:

$$\dot{S} = -\beta IS + c - \mu_S S, \quad (4)$$

$$\dot{E} = \beta(1-v)IS - \delta E - \mu_E E, \quad (5)$$

$$\dot{I} = \beta v IS + \delta E - \mu_I I. \quad (6)$$

We also observe the following two outputs

$$y_E = k_E E, \quad (7)$$

$$y_I = k_I I, \quad (8)$$

and their interpretation is that we observe proportions k_E and k_I of the exposed and infected populations, respectively. In total, we have nine parameters collected in the vector $\boldsymbol{\theta} \in \mathbb{R}^9$

which are given by

$$\boldsymbol{\theta} = \begin{pmatrix} c \\ \beta \\ \mu_S \\ \mu_E \\ \mu_I \\ \delta \\ v \\ k_E \\ k_I \end{pmatrix}. \quad (9)$$

We are looking for a family of infinitesimal generators of the Lie group given by

$$\begin{aligned} X = & \xi(t, S, E, I, \boldsymbol{\theta}) \partial_t + \eta_S(t, S, E, I, \boldsymbol{\theta}) \partial_t + \eta_E(t, S, E, I, \boldsymbol{\theta}) \partial_t + \eta_I(t, S, E, I, \boldsymbol{\theta}) \partial_t \\ & + \chi_c(\boldsymbol{\theta}) \partial_c + \chi_\beta(\boldsymbol{\theta}) \partial_\beta + \chi_{\mu_S}(\boldsymbol{\theta}) \partial_{\mu_S} + \chi_{\mu_E}(\boldsymbol{\theta}) \partial_{\mu_E} + \chi_{\mu_I}(\boldsymbol{\theta}) \partial_{\mu_I} + \chi_\delta(\boldsymbol{\theta}) \partial_\delta \\ & + \chi_v(\boldsymbol{\theta}) \partial_v + \chi_{k_E}(\boldsymbol{\theta}) \partial_{k_E} + \chi_{k_I}(\boldsymbol{\theta}) \partial_{k_I}. \end{aligned} \quad (10)$$

The three first prolongations are given by

$$\eta_S^{(1)}(t, S, E, I, \dot{S}, \boldsymbol{\theta}) = D_t \eta_S(t, S, E, I, \boldsymbol{\theta}) - \dot{S} D_t \xi(t, S, E, I, \boldsymbol{\theta}), \quad (11)$$

$$\eta_E^{(1)}(t, S, E, I, \dot{E}, \boldsymbol{\theta}) = D_t \eta_E(t, S, E, I, \boldsymbol{\theta}) - \dot{E} D_t \xi(t, S, E, I, \boldsymbol{\theta}), \quad (12)$$

$$\eta_I^{(1)}(t, S, E, I, \dot{I}, \boldsymbol{\theta}) = D_t \eta_I(t, S, E, I, \boldsymbol{\theta}) - \dot{I} D_t \xi(t, S, E, I, \boldsymbol{\theta}), \quad (13)$$

where the total derivative is defined by: $D_t = \partial_t + \dot{S} \partial_S + \dot{E} \partial_E + \dot{I} \partial_I$. These prolongations define the first prolongation of the infinitesimal generator $X^{(1)}$ according to

$$X^{(1)} = X + \eta_S^{(1)} \partial_{\dot{S}} + \eta_E^{(1)} \partial_{\dot{E}} + \eta_I^{(1)} \partial_{\dot{I}}. \quad (14)$$

Before, we define the linearised symmetry conditions, we make two critical simplifications. First, the model of interest is autonomous which implies that the time infinitesimal is a constant, i.e. $\xi(t, S, E, I) = K$ for some $K \in \mathbb{R}$ and thus $D_t \xi = 0$.

Second, the fact that the observed outputs are differential invariants of our generator yields equations for the infinitesimals η_E and η_I , respectively. Starting with the infinitesimal for E , we have that

$$X(y_E) = 0 \implies k_E \eta_E + \chi_{K_E} E = 0,$$

which gives us the following equation for η_E

$$\eta_E = - \left(\frac{\chi_{k_E}}{k_E} \right) E. \quad (15)$$

Analogously, the equation for η_I is given by

$$\eta_I = - \left(\frac{\chi_{k_I}}{k_I} \right) I. \quad (16)$$

Thus the corresponding prolongations simplify to

$$\eta_E^{(1)} = - \left(\frac{\chi_{k_E}}{k_E} \right) \dot{E}. \quad (17)$$

$$\eta_I^{(1)} = - \left(\frac{\chi_{k_I}}{k_I} \right) \dot{I}. \quad (18)$$